

Lecture 10  
Thur. 9.20.2012

*Quadratic form geometry and development of mechanics  
of Lagrange and Hamilton*

*(Ch. 12 of Unit 1 and Ch. 4-5 of Unit 7)*

*Scaling transformation between Lagrangian and Hamiltonian views of KE (Review of Lecture 9)  
Introducing 1st Lagrange and Hamilton differential equations of mechanics (Review Of Lecture 9)*

*Introducing the Poincare' and Legendre contact transformations*

*Geometry of Legendre contact transformation*


*Example from thermodynamics*

*Legendre transform: special case of General Contact Transformation (lights, camera, **ACTION!**)*

*A general contact transformation from sophomore physics*

*Algebra-calculus development of "The Volcanoes of Io" and "The Atoms of NIST"*

*Intuitive-geometric development of " " " and " " "*

*Scaling transformation between Lagrangian and Hamiltonian views of KE (Review of Lecture 9)*  
*Introducing the (partial) differential equations of mechanics (Review Of Lecture 9)*  
 *1st equations of Lagrange and Hamilton*

# Introducing the (partial $\frac{\partial}{\partial}$ ) differential equations of mechanics

Starts out with simple demands for explicit-dependence, “loyalty” or “fealty to the colors”

*Lagrangian and Estrangian* have no explicit dependence on **momentum  $\mathbf{p}$**

$$\frac{\partial L}{\partial \mathbf{p}_k} \equiv 0 \equiv \frac{\partial E}{\partial \mathbf{p}_k}$$

*Hamiltonian and Estrangian* have no explicit dependence on **velocity  $\mathbf{v}$**

$$\frac{\partial H}{\partial \mathbf{v}_k} \equiv 0 \equiv \frac{\partial E}{\partial \mathbf{v}_k}$$

*Lagrangian and Hamiltonian* have no explicit dependence on **speedinum  $\mathbf{V}$**

$$\frac{\partial L}{\partial \mathbf{V}_k} \equiv 0 \equiv \frac{\partial H}{\partial \mathbf{V}_k}$$

Such non-dependencies hold in spite of “under-the-table” matrix and partial-differential connections

$$\begin{aligned} \nabla_{\mathbf{v}} L &= \frac{\partial L}{\partial \mathbf{v}} = \frac{\partial}{\partial \mathbf{v}} \frac{\mathbf{v} \cdot \mathbf{M} \cdot \mathbf{v}}{2} \\ &= \mathbf{M} \cdot \mathbf{v} = \mathbf{p} \end{aligned}$$

$$\begin{aligned} \nabla_{\mathbf{p}} H &= \mathbf{v} = \frac{\partial H}{\partial \mathbf{p}} = \frac{\partial}{\partial \mathbf{p}} \frac{\mathbf{p} \cdot \mathbf{M}^{-1} \cdot \mathbf{p}}{2} \\ &= \mathbf{M}^{-1} \cdot \mathbf{p} = \mathbf{v} \end{aligned}$$

(Forget Estrangian for now)

$$\begin{pmatrix} \frac{\partial L}{\partial v_1} \\ \frac{\partial L}{\partial v_2} \end{pmatrix} = \begin{pmatrix} m_1 & 0 \\ 0 & m_2 \end{pmatrix} \begin{pmatrix} v_1 \\ v_2 \end{pmatrix} = \begin{pmatrix} p_1 \\ p_2 \end{pmatrix}$$

Lagrange's 1<sup>st</sup> equation(s)

$$\frac{\partial L}{\partial v_k} = p_k \quad \text{or:} \quad \frac{\partial L}{\partial \mathbf{v}} = \mathbf{p}$$

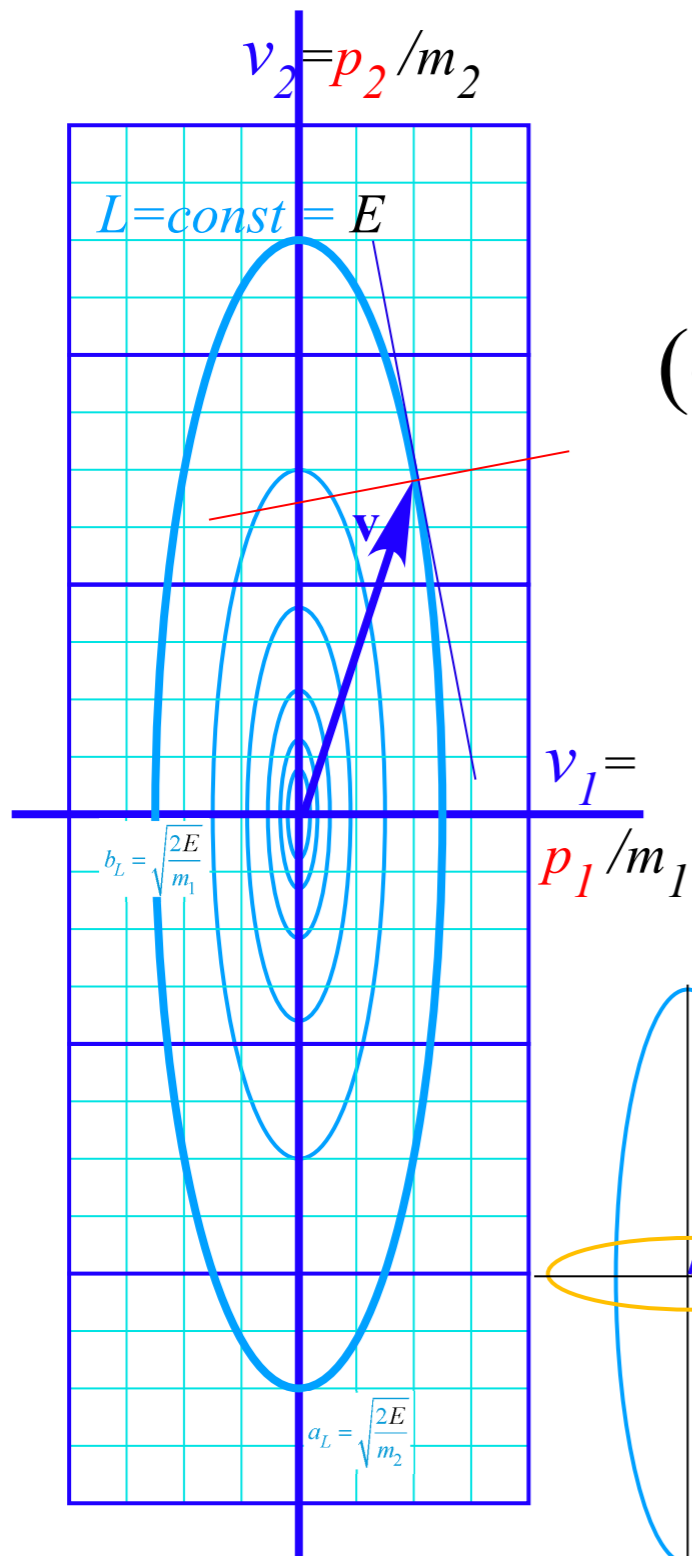
$$\begin{pmatrix} \frac{\partial H}{\partial p_1} \\ \frac{\partial H}{\partial p_2} \end{pmatrix} = \begin{pmatrix} m_1^{-1} & 0 \\ 0 & m_2^{-1} \end{pmatrix} \begin{pmatrix} p_1 \\ p_2 \end{pmatrix} = \begin{pmatrix} v_1 \\ v_2 \end{pmatrix}$$

Hamilton's 1<sup>st</sup> equation(s)

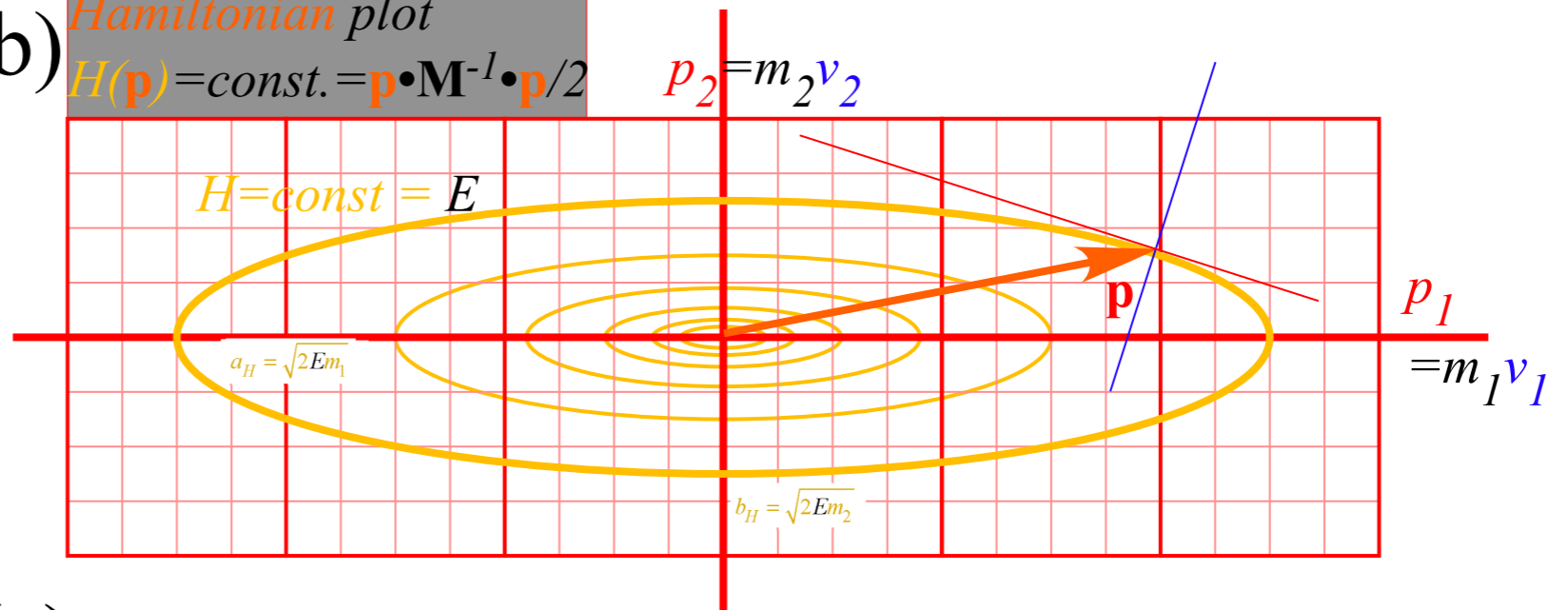
$$\frac{\partial H}{\partial p_k} = v_k \quad \text{or:} \quad \frac{\partial H}{\partial \mathbf{p}} = \mathbf{v}$$

Unit 1  
Fig. 12.2

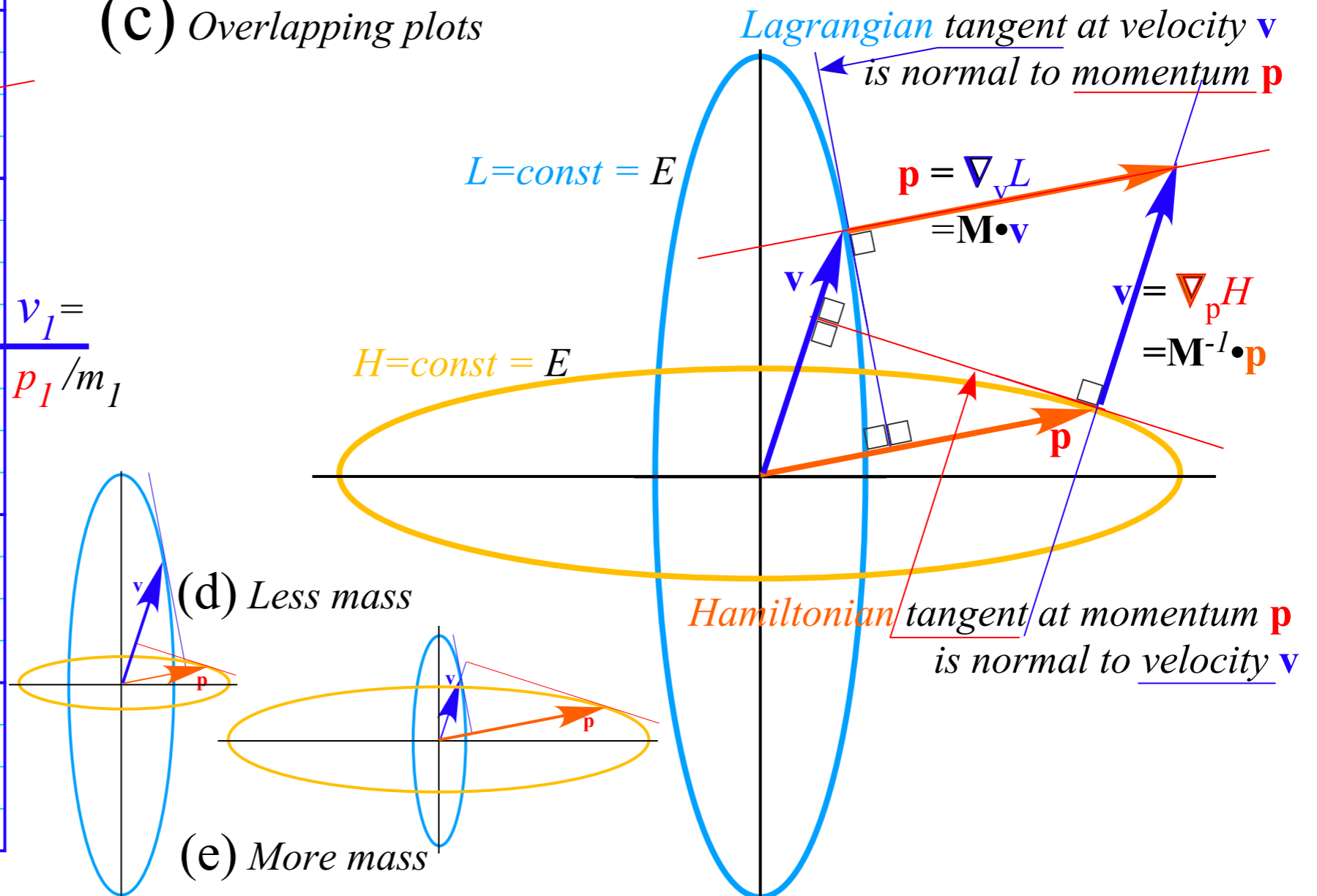
(a) *Lagrangian plot*  
 $L(\mathbf{v}) = \text{const.} = \mathbf{v} \cdot \mathbf{M} \cdot \mathbf{v} / 2$



(b) *Hamiltonian plot*  
 $H(\mathbf{p}) = \text{const.} = \mathbf{p} \cdot \mathbf{M}^{-1} \cdot \mathbf{p} / 2$



(c) *Overlapping plots*



# → *Introducing the Poincare' and Legendre contact transformations*

*Geometry of Legendre contact transformation*

*Example from thermodynamics*

*Legendre transform: special case of General Contact Transformation (lights, camera, **ACTION!**)*

# *Introducing the Poincare' and Legendre contact transformations*

*Given matrix relation:  $\mathbf{p}=\mathbf{M}\cdot\mathbf{v}$  or its inverse:  $\mathbf{v}=\mathbf{M}^{-1}\cdot\mathbf{p}$  you might be tempted to rewrite*

*Q-forms  $L(\mathbf{v}..)=(1/2)\mathbf{v}\cdot\mathbf{M}\cdot\mathbf{v}$  or  $H(\mathbf{p}..)=(1/2)\mathbf{p}\cdot\mathbf{M}^{-1}\cdot\mathbf{p}$  to be  $H=(1/2)\mathbf{p}\cdot\mathbf{v}$  or equivalently  $L=(1/2)\mathbf{v}\cdot\mathbf{p}$ .*

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*Numerically-CORRECT, but Differentially-WRONG!*

# *Introducing the Poincare' and Legendre contact transformations*

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*Numerically-CORRECT, but Differentially-WRONG!*

*Instead try:  $H(\mathbf{p}..)=\mathbf{p}\cdot\mathbf{v}-(1/2)\mathbf{v}\cdot\mathbf{p}=\mathbf{p}\cdot\mathbf{v}-L(\mathbf{v}..)$  or else:  $L(\mathbf{v}..)=\mathbf{p}\cdot\mathbf{v}-H(\mathbf{p}..)$*



# Introducing the Poincare' and Legendre contact transformations

Given matrix relation:  $\mathbf{p}=\mathbf{M}\cdot\mathbf{v}$  or its inverse:  $\mathbf{v}=\mathbf{M}^{-1}\cdot\mathbf{p}$  you might be tempted to rewrite

Q-forms  $L(\mathbf{v}..)=\frac{1}{2}\mathbf{v}\cdot\mathbf{M}\cdot\mathbf{v}$  or  $H(\mathbf{p}..)=\frac{1}{2}\mathbf{p}\cdot\mathbf{M}^{-1}\cdot\mathbf{p}$  to be  $H=\frac{1}{2}\mathbf{p}\cdot\mathbf{v}$  or equivalently  $L=\frac{1}{2}\mathbf{v}\cdot\mathbf{p}$ .

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Instead try:  $H(\mathbf{p}..)=\mathbf{p}\cdot\mathbf{v}-\frac{1}{2}\mathbf{v}\cdot\mathbf{p}=\mathbf{p}\cdot\mathbf{v}-L(\mathbf{v}..)$  or else:  $L(\mathbf{v}..)=\mathbf{p}\cdot\mathbf{v}-H(\mathbf{p}..)$

## Legendre contact transformation

$$L(\mathbf{v}) = \mathbf{p} \cdot \mathbf{v} - H(\mathbf{p}) \qquad H(\mathbf{p}) = \mathbf{p} \cdot \mathbf{v} - L(\mathbf{v})$$

Now explicit dependency (non)-relations give the right derivatives

$$\begin{aligned} \frac{\partial L(\mathbf{v})}{\partial \mathbf{p}} &= \frac{\partial}{\partial \mathbf{p}} \mathbf{p} \cdot \mathbf{v} - \frac{\partial H(\mathbf{p})}{\partial \mathbf{p}} & \frac{\partial H(\mathbf{p})}{\partial \mathbf{v}} &= \frac{\partial}{\partial \mathbf{v}} \mathbf{p} \cdot \mathbf{v} - \frac{\partial L(\mathbf{v})}{\partial \mathbf{v}} \\ 0 &= \mathbf{v} - \frac{\partial H(\mathbf{p})}{\partial \mathbf{p}} & 0 &= \mathbf{p} - \frac{\partial L(\mathbf{v})}{\partial \mathbf{v}} \end{aligned}$$

That is *Hamilton's 1<sup>st</sup> equation(s)* and *Lagrange's 1<sup>st</sup> equation(s)*

# *Introducing the Poincare' and Legendre contact transformations*

→ *Geometry of Legendre contact transformation*

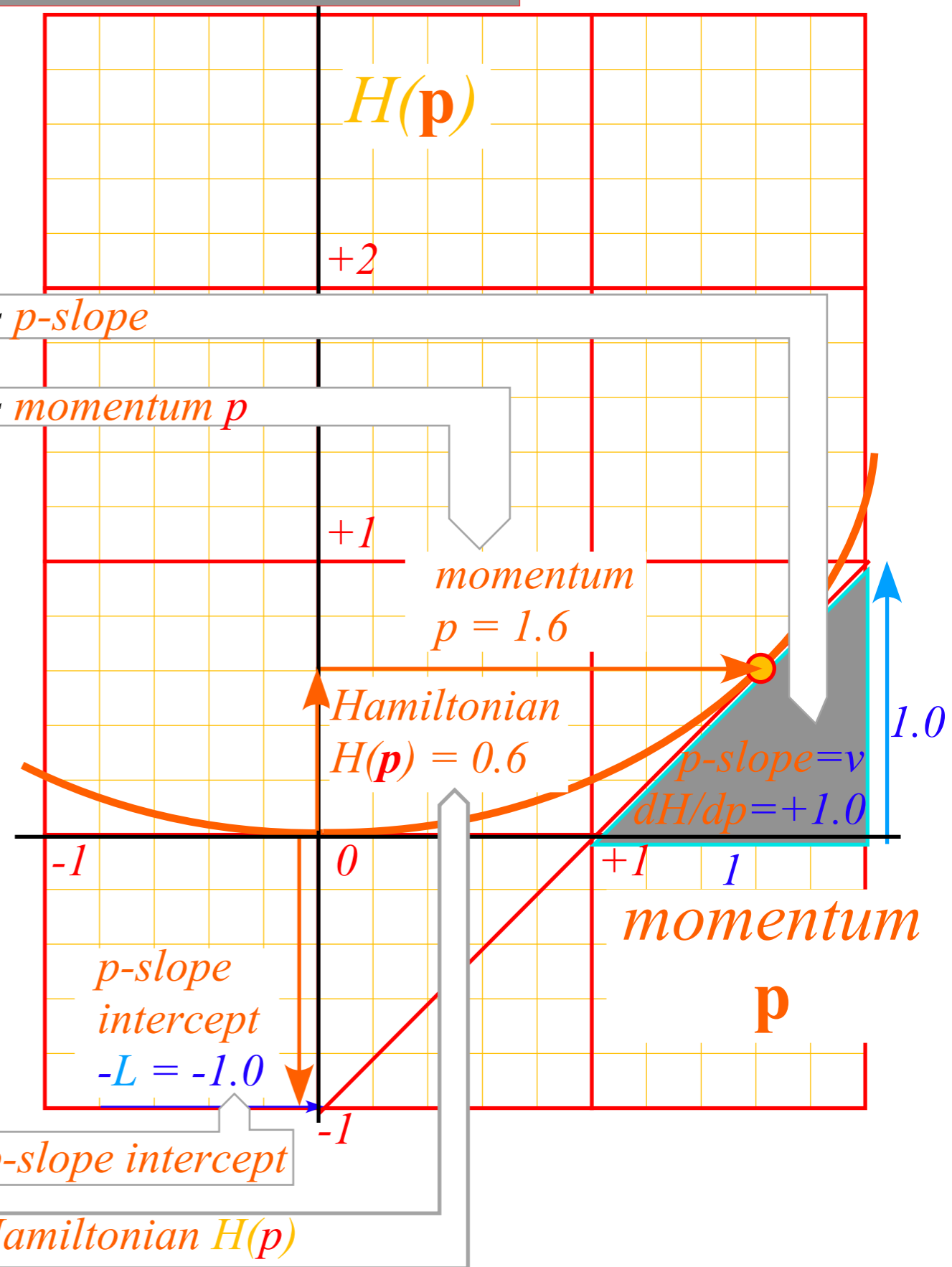
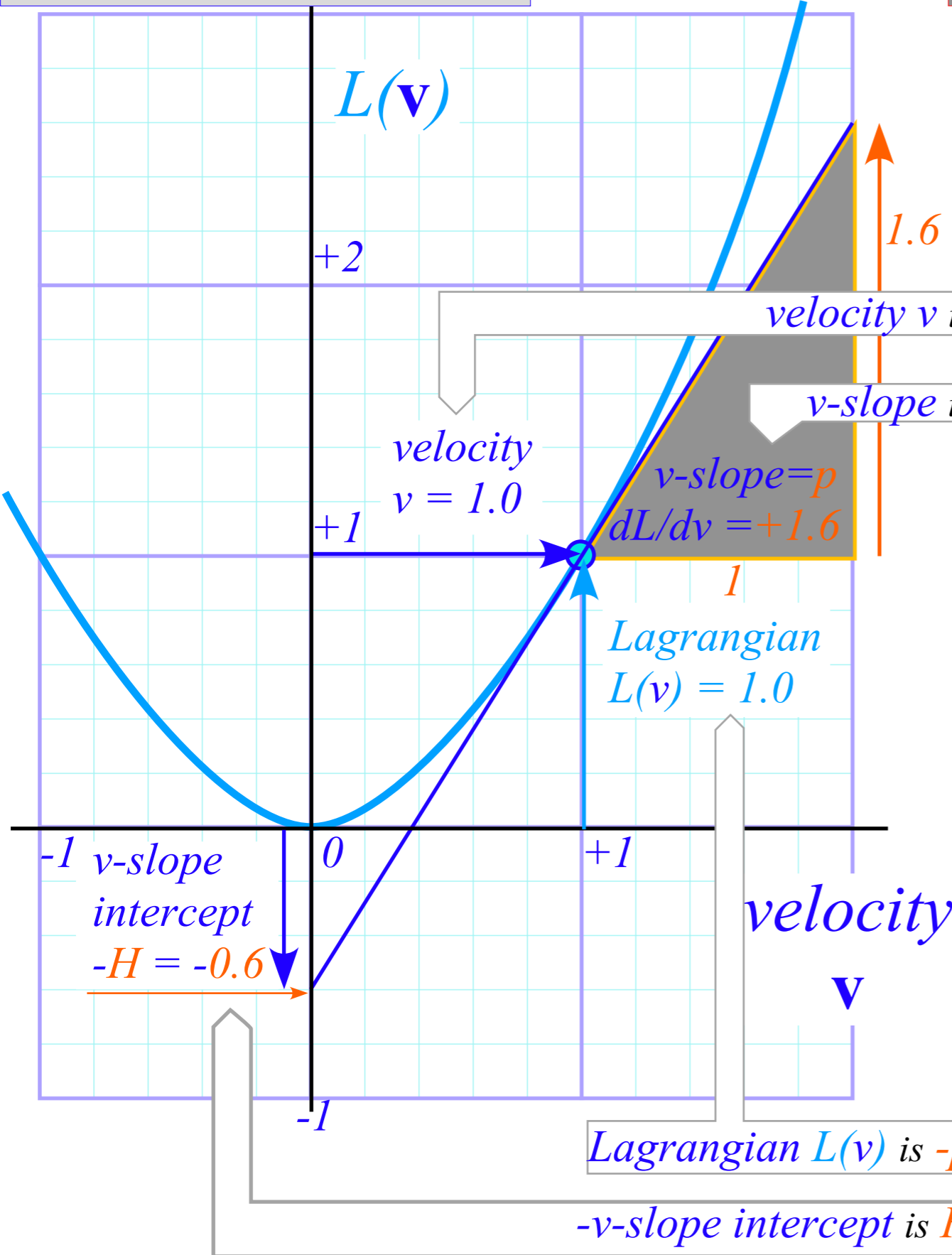
*Example from thermodynamics*

*Legendre transform: special case of General Contact Transformation (lights, camera, ACTION!)*

Unit 1  
Fig. 12.3

(a) *Lagrangian plot*  
 $L(\mathbf{v}) = \mathbf{v} \cdot \mathbf{p} - H(\mathbf{p})$

(b) *Hamiltonian plot*  
 $H(\mathbf{p}) = \mathbf{p} \cdot \mathbf{v} - L(\mathbf{v})$



Lagrangian  $L(v)$  is  $-p$ -slope intercept

$-v$ -slope intercept is Hamiltonian  $H(p)$

# How Legendre contact transformations work... (to make $\frac{\partial H}{\partial v} = 0$ or $\frac{\partial L}{\partial p} = 0$ )

Secant lines  $L(\mathbf{v}) = \mathbf{p} \cdot \mathbf{v} - H$  of fixed slope  $\mathbf{p} = \frac{\partial L}{\partial \mathbf{v}}$   
and decreasing intercept  $-H(v_{-2}) > -H(v_{-1}) > \dots$

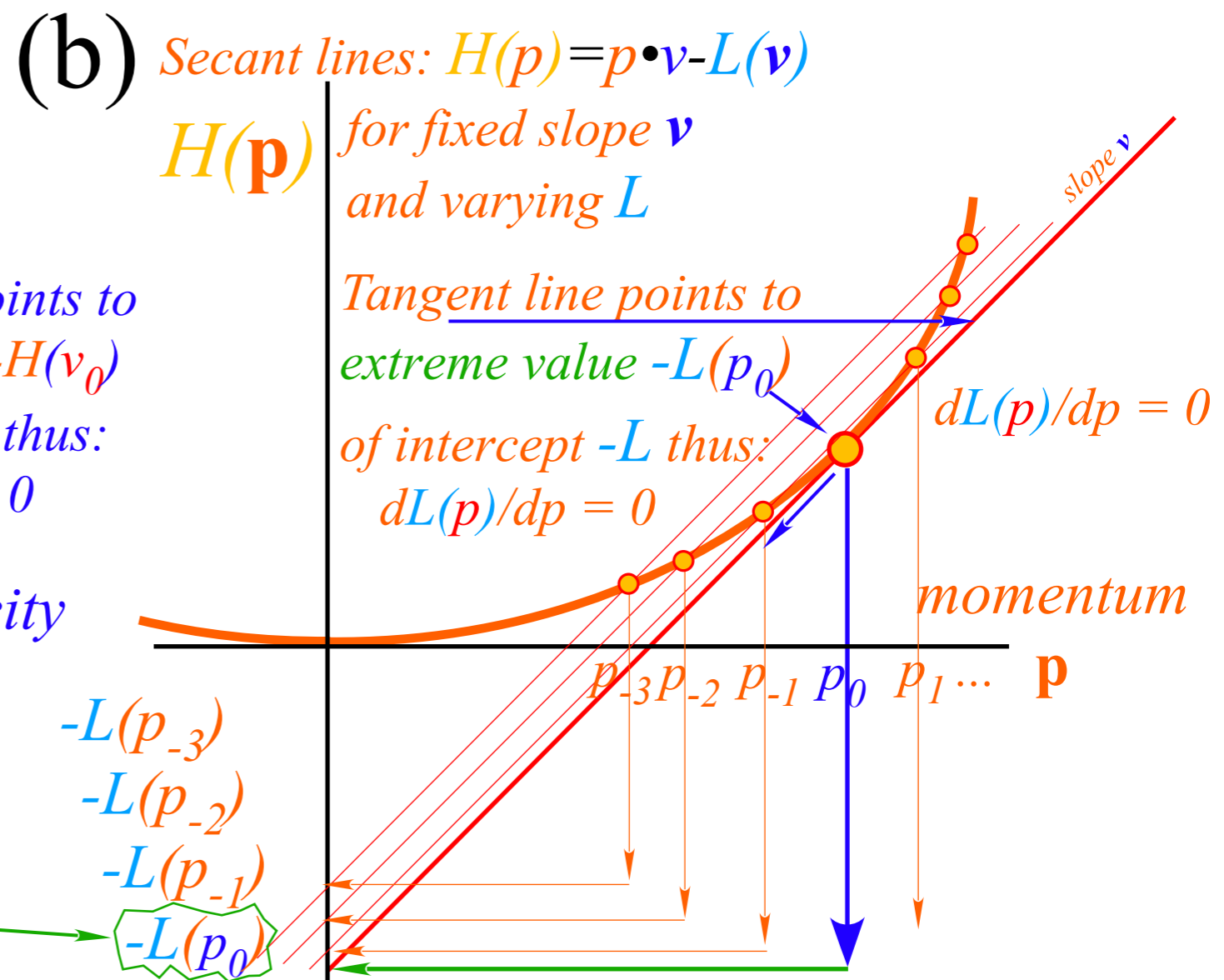
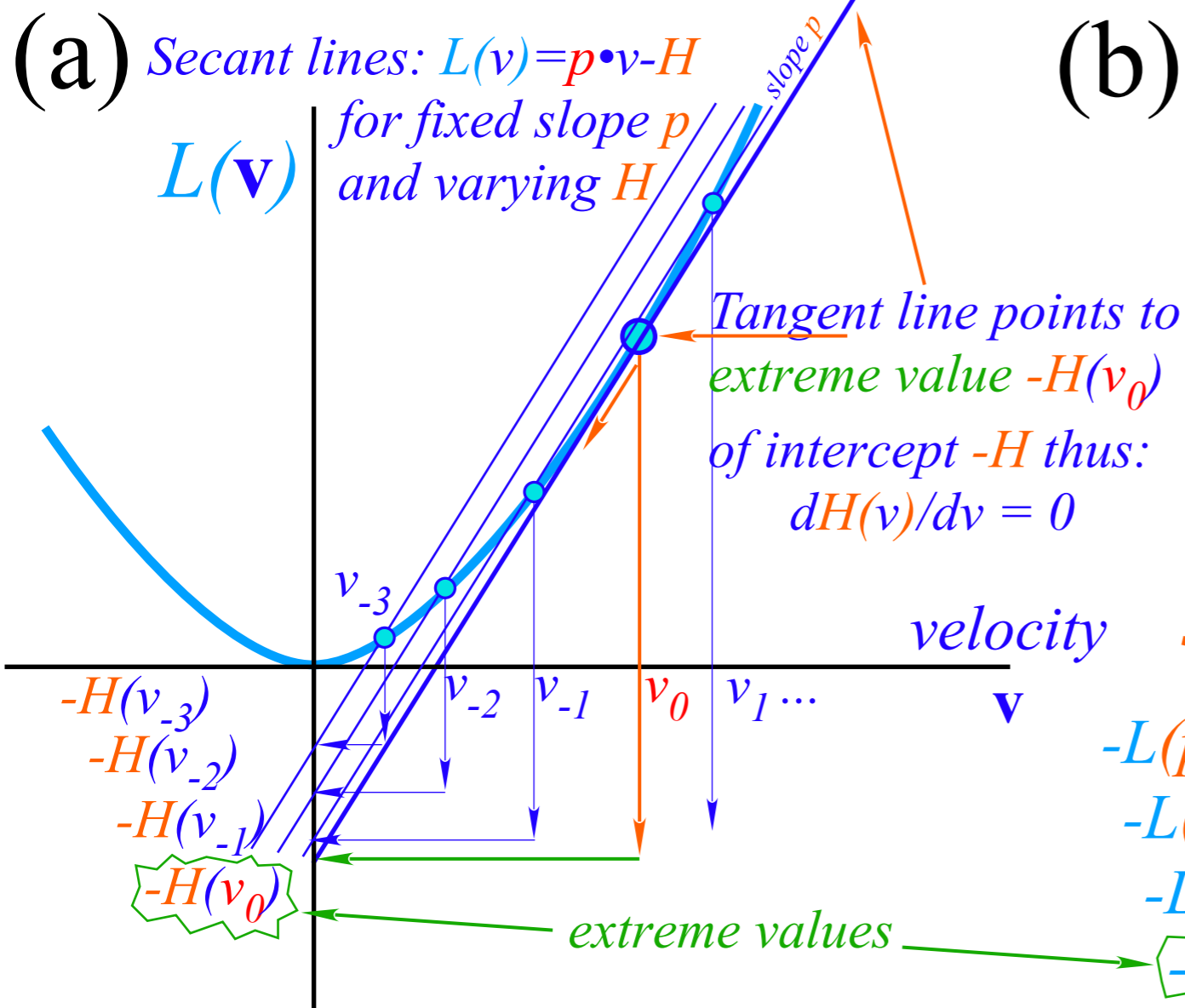
for increasing velocity  $v_{-2} > v_{-1} > \dots > v_0$

lead to unique tangent to  $L(\mathbf{v})$ -curve at the  
tangent contact point  $\mathbf{v} = \mathbf{v}_0$  that has max  $H(\mathbf{p}, \mathbf{v}_0)$

Thus  $\frac{\partial H}{\partial v} = 0$

(Similarly...)

Unit 1  
Fig. 12.4



$$\frac{\partial H}{\partial v} = 0 \text{ at each point } \mathbf{v} = \frac{\partial H}{\partial \mathbf{p}} \text{ of } L(\mathbf{v}) \text{ with slope } \mathbf{p} = \frac{\partial L}{\partial \mathbf{v}}$$

$$\frac{\partial L}{\partial p} = 0 \text{ at each point } \mathbf{p} = \frac{\partial L}{\partial \mathbf{v}} \text{ of } H(\mathbf{p}) \text{ with slope } \mathbf{v} = \frac{\partial H}{\partial \mathbf{p}}$$

# *Introducing the Poincare' and Legendre contact transformations*

*Geometry of Legendre contact transformation*

 *Example from thermodynamics*

*Legendre transform: special case of General Contact Transformation (lights, camera, ACTION!)*

## *Example of Legendre contact transformation in thermodynamics*

*Internal energy*  $U(S, V)$  is defined as a function of entropy  $S$  and volume  $V$ .

A new function *enthalpy*  $H(S, P)$  depends on entropy and *pressure*  $P$ .

It is a Legendre transform  $H(S, P) = P \cdot V + U$  of energy  $U(S, V)$  to new variable  $P = -\left(\frac{\partial U}{\partial V}\right)_S$  .

Except for  $\pm$  signs, it's our Hamiltonian  $H(\mathbf{p}) = \mathbf{p} \cdot \mathbf{v} - L(\mathbf{v})$  going from Lagrangian  $L(\mathbf{v})$

to use new variable momentum  $\mathbf{p} = \left(\frac{\partial L}{\partial \mathbf{v}}\right)_x$  .

# *Introducing the Poincare' and Legendre contact transformations*

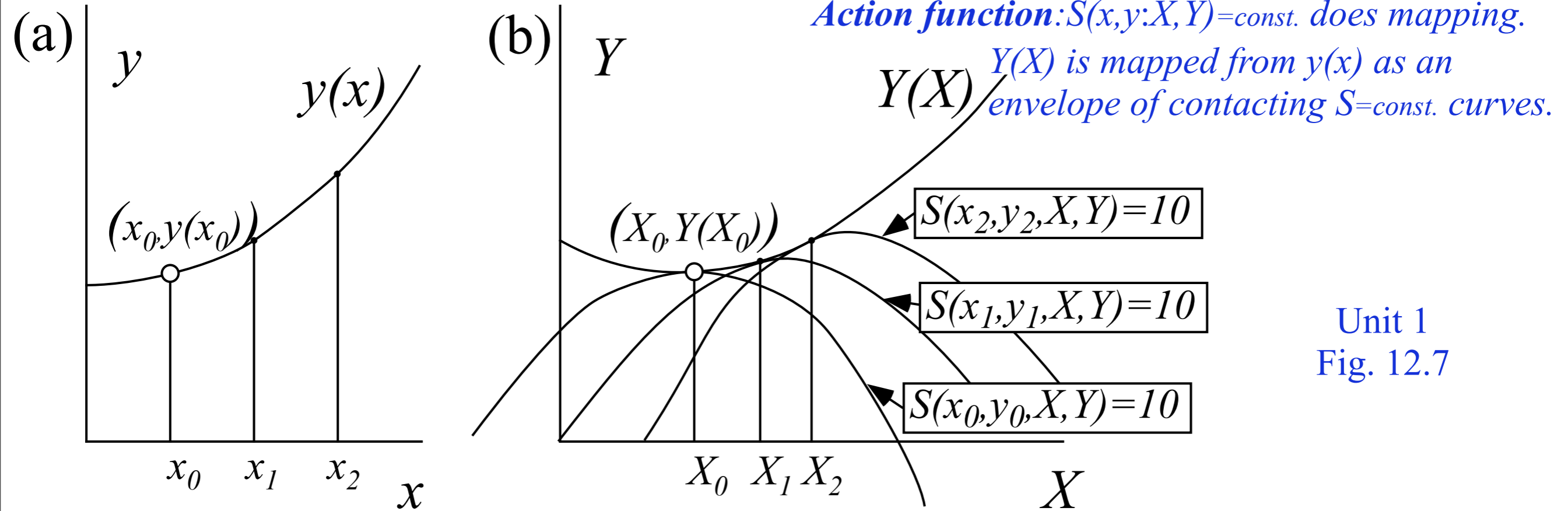
*Geometry of Legendre contact transformation*

*Example from thermodynamics*

 *Legendre transform: special case of General Contact Transformation (lights, camera, **ACTION!**)*

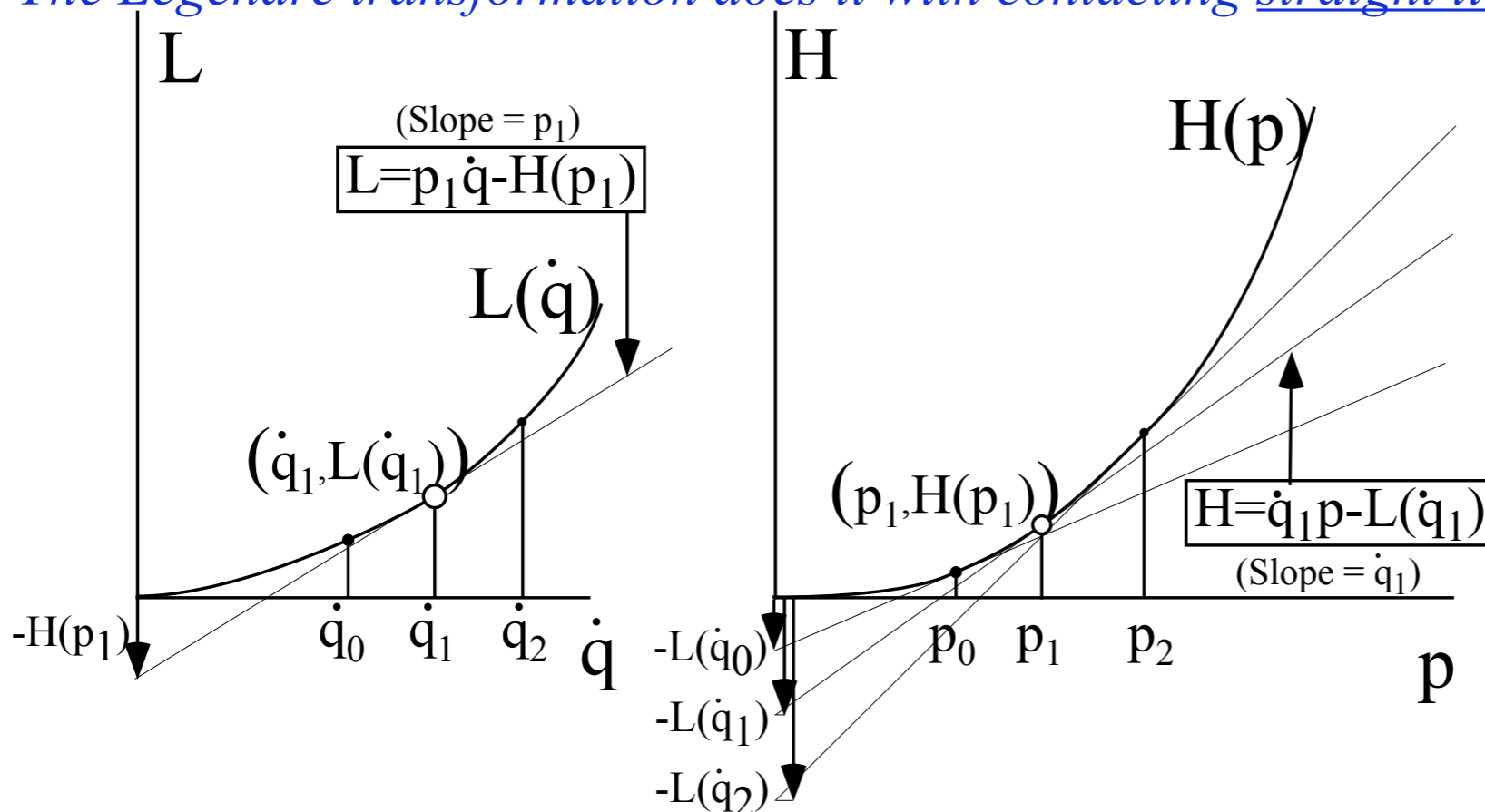
# Legendre transform: special case of General Contact Transformation

Active-Contact-Transformation Generator or  
**Action function:**  $S(x,y;X,Y)=const.$  does mapping.



Unit 1  
 Fig. 12.7

The Legendre transform does it with contacting straight line tangents.

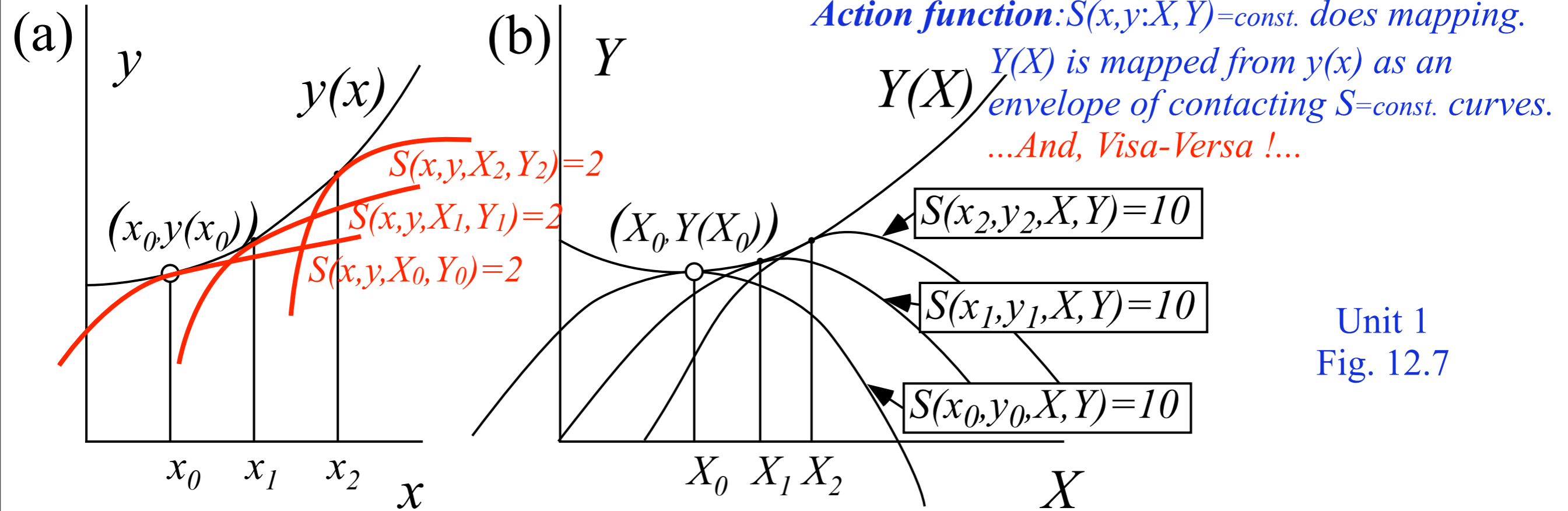


Unit 1  
 Fig. 12.9



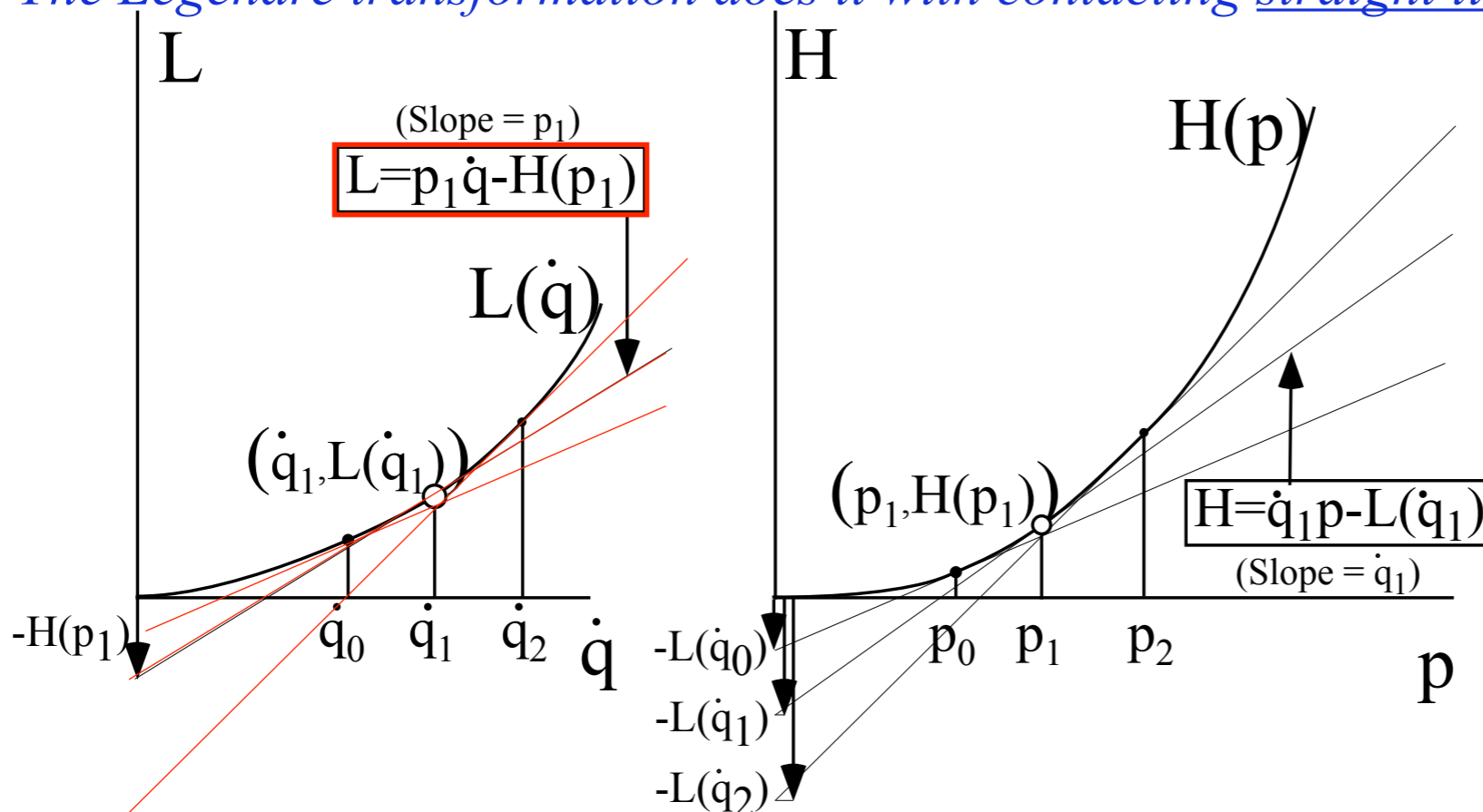
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Unit 1  
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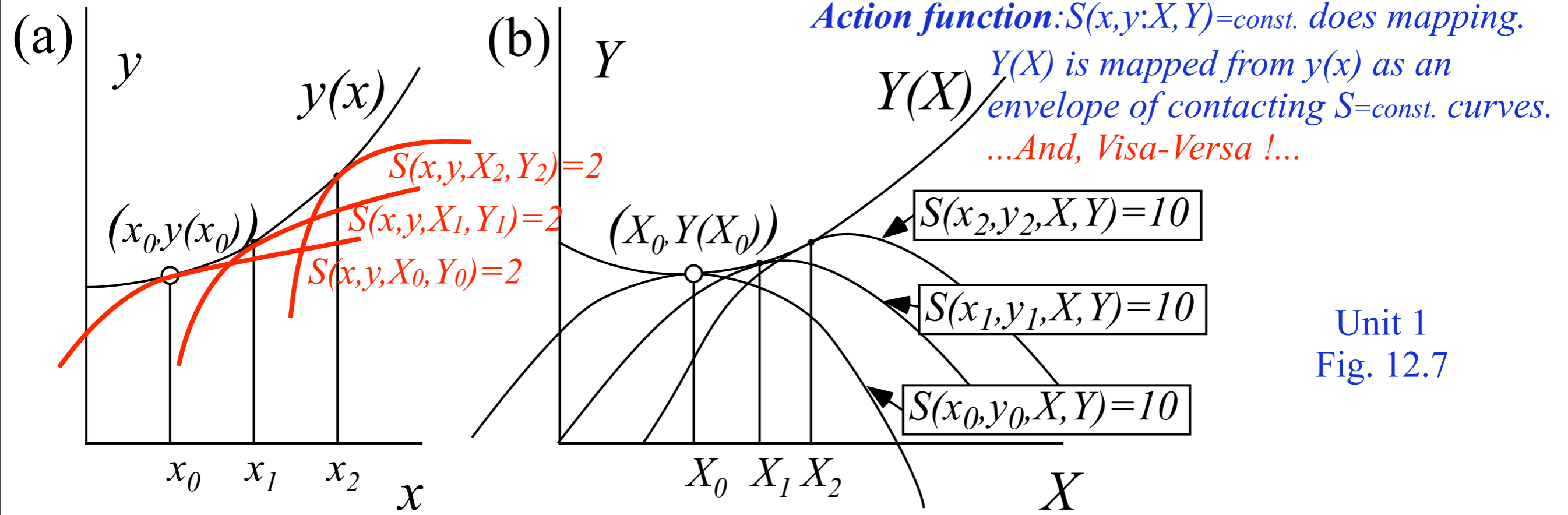
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Unit 1  
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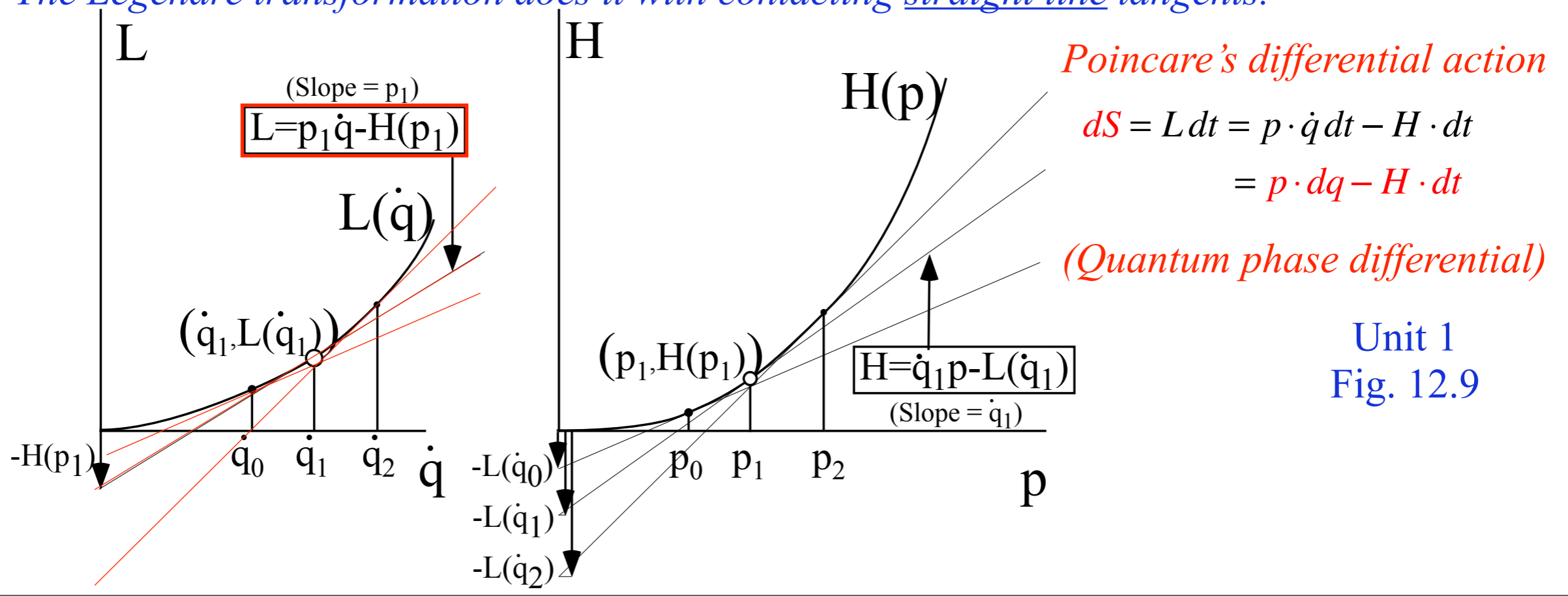
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Unit 1  
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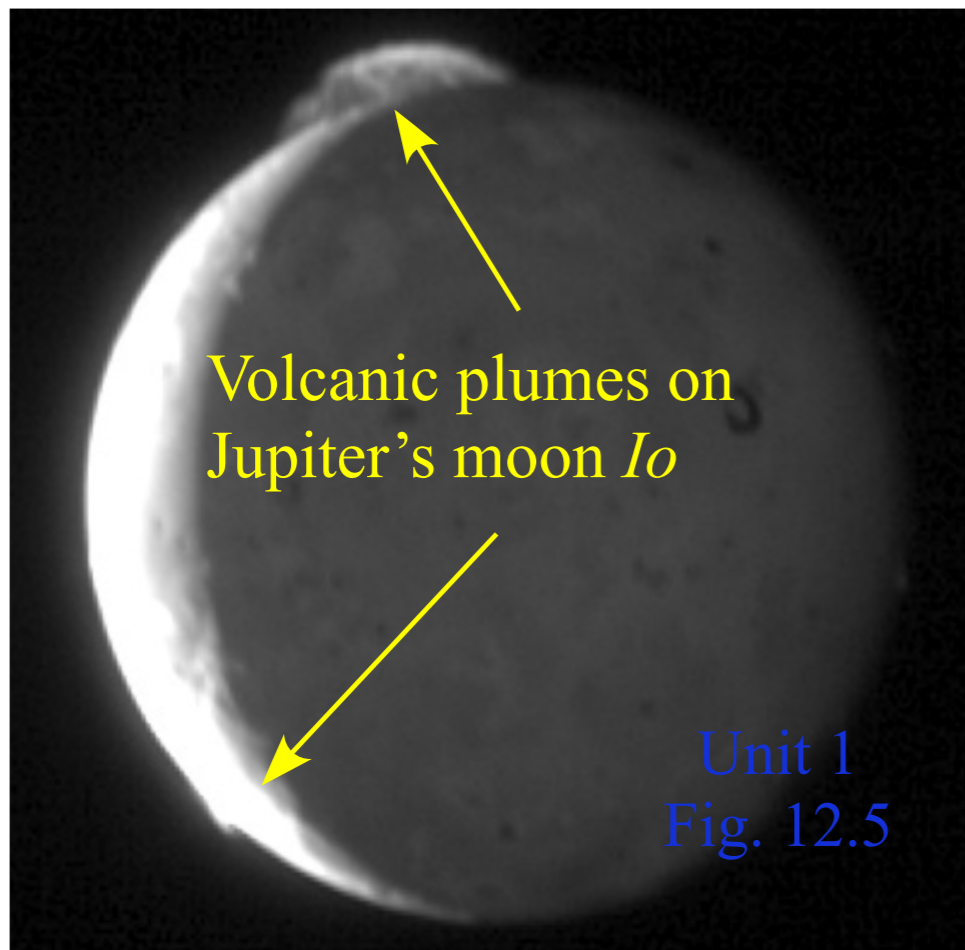
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Unit 1  
 Fig. 12.9

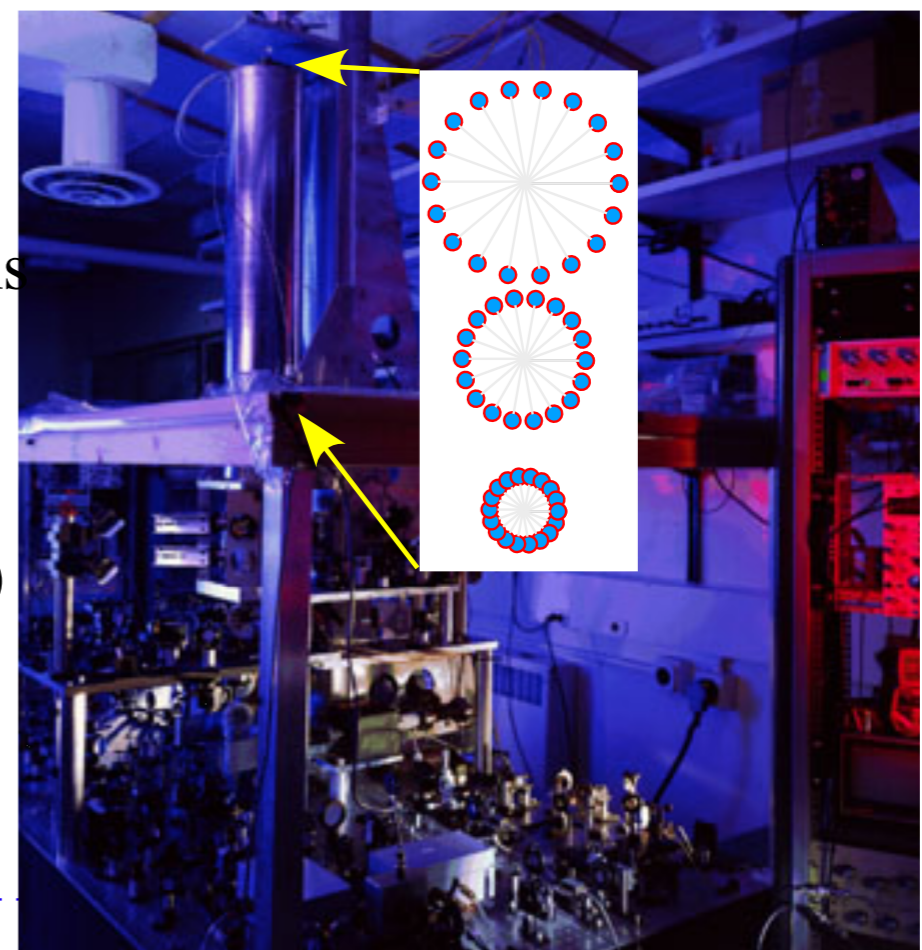
*A general contact transformation from sophomore physics*  
→ *Algebra-calculus development of “The Volcanoes of Io” and “The Atoms of NIST”*  
*Intuitive-geometric development of ” ” ” and ” ” ”*

(a)

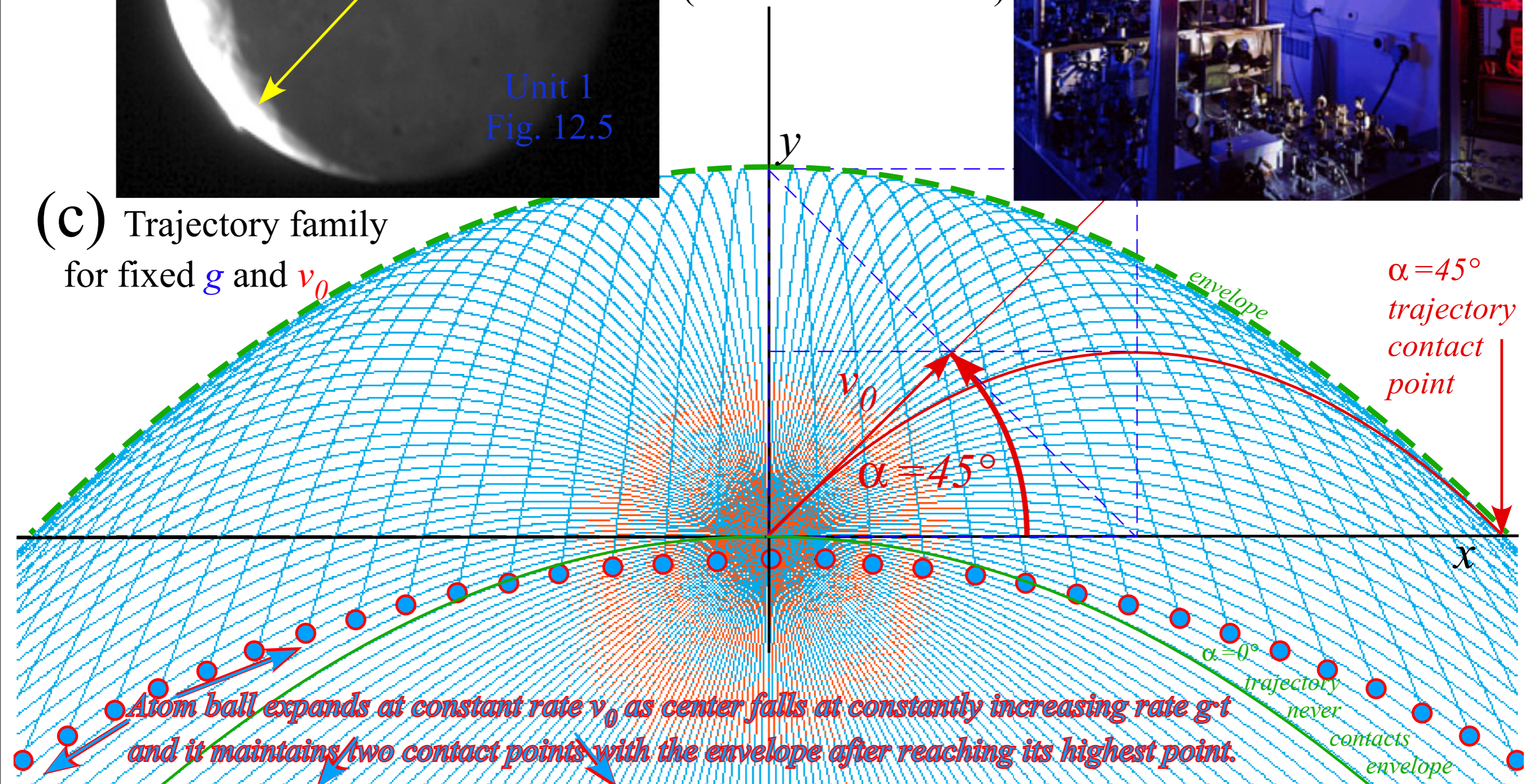


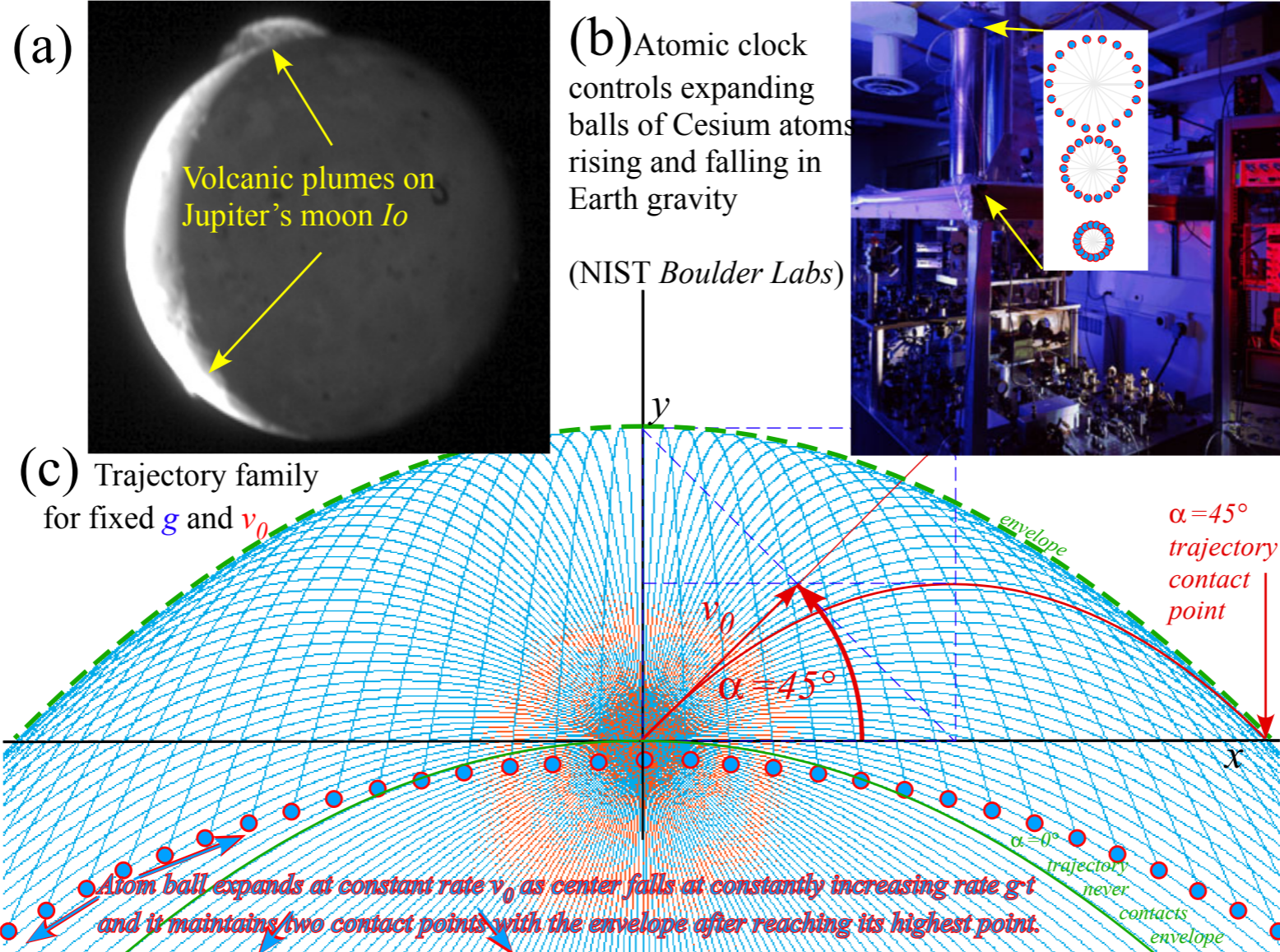
(b) Atomic clock controls expanding balls of Cesium atoms rising and falling in Earth gravity

(NIST Boulder Labs)



(c) Trajectory family for fixed  $g$  and  $v_0$





Unit 1  
Fig. 12.5

*UP-1 formulas for trajectories in constant gravity g*

$$\begin{aligned}
 x(t) &= (v_0 \cos \alpha)t & y(t) &= (v_0 \sin \alpha)t - \frac{1}{2}gt^2 \\
 \dot{x}(0) = v_x(0) &= v_0 \cos \alpha & \dot{y}(0) = v_y(0) &= v_0 \sin \alpha
 \end{aligned}$$

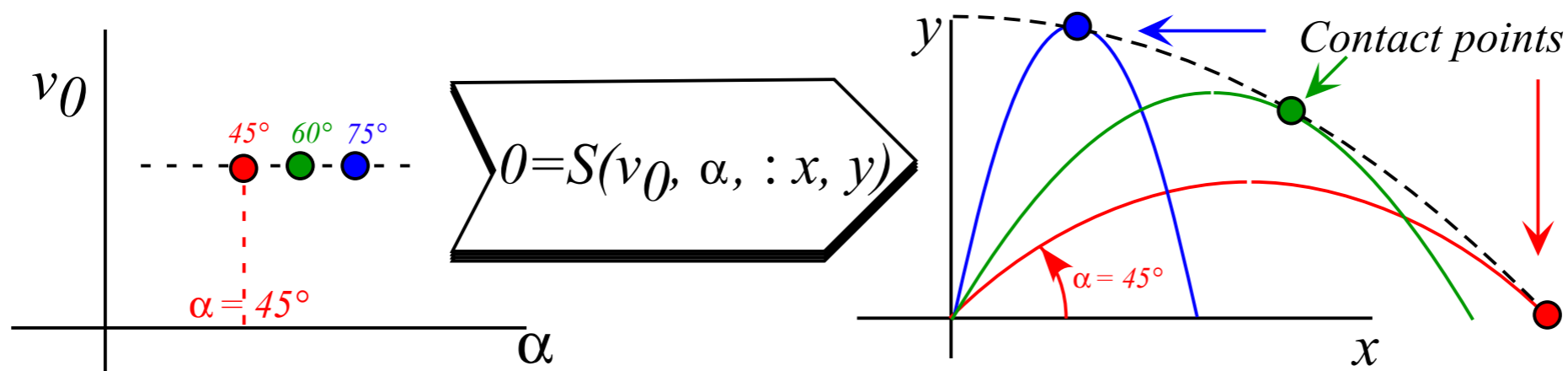
Substitute time  $t=x/(v_0 \cos \alpha)$  into  $y(t)$

$$y(x) = \frac{v_0 \sin \alpha}{v_0 \cos \alpha} x - \frac{gx^2}{2v_0^2 \cos^2 \alpha}$$

$$y(x) = x \tan \alpha - \frac{gx^2}{2v_0^2 \cos^2 \alpha}$$

Convert  $y(x)$  solution into Active Contact Transformation Generator  $S(v_0, \alpha: x, y)$

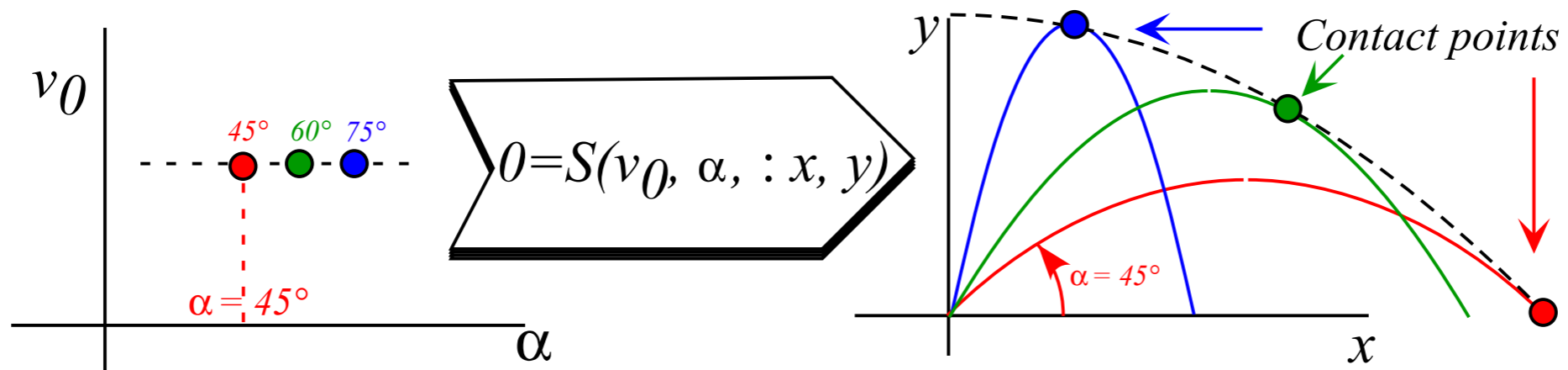
$$y(x) = x \tan \alpha - \frac{gx^2}{2v_0^2 \cos^2 \alpha} \quad \text{becomes:} \quad S(v_0, \alpha: x, y) = -y + x \tan \alpha - \frac{gx^2}{2v_0^2 \cos^2 \alpha} = 0$$



Unit 1  
Fig. 12.6

Convert  $y(x)$  solution into Active Contact Transformation Generator  $S(v_0, \alpha: x, y)$

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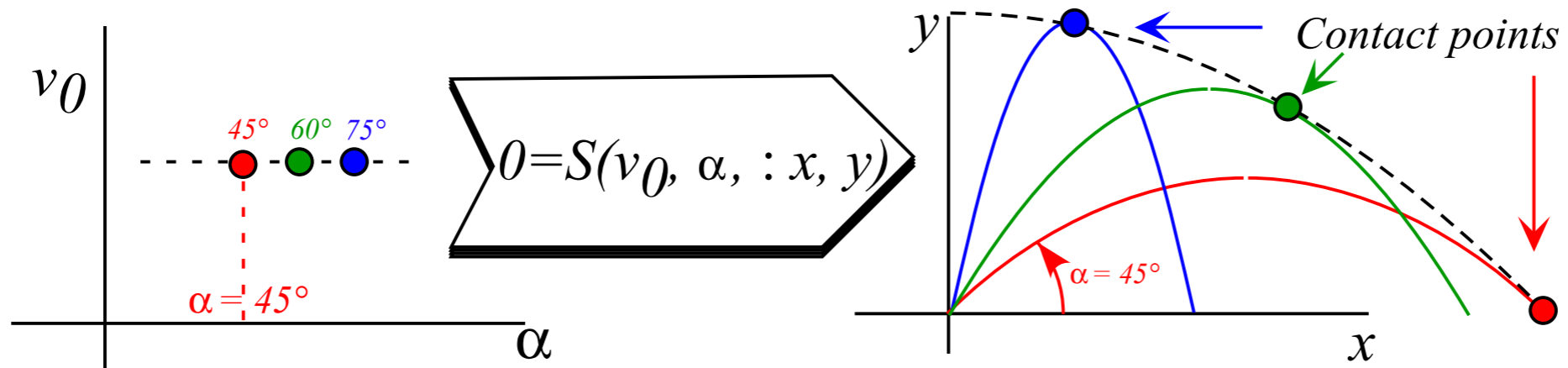
Unit 1  
Fig. 12.6

*Envelopes* of the  $v_0$ -trajectory region contain extremal *contact points* with each trajectory

where:  $\frac{\partial S(v_0, \alpha: x, y)}{\partial \alpha} = 0$

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Unit 1  
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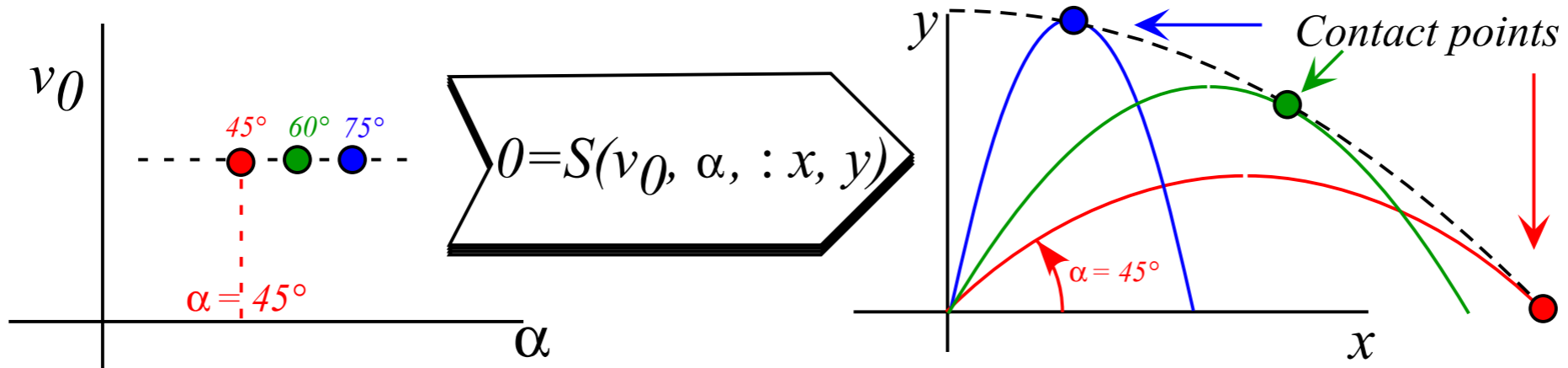
where:  $\frac{\partial S(v_0, \alpha: x, y)}{\partial \alpha} = 0$

$$x \frac{\partial \tan \alpha}{\partial \alpha} - \frac{gx^2}{2v_0^2} \frac{\partial \cos^{-2} \alpha}{\partial \alpha} = 0 = \frac{x}{\cos^2 \alpha} - \frac{gx^2}{2v_0^2} \frac{2 \sin \alpha}{\cos^3 \alpha}$$



Convert  $y(x)$  solution into Active Contact Transformation Generator  $S(v_0, \alpha: x, y)$

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Unit 1  
Fig. 12.6

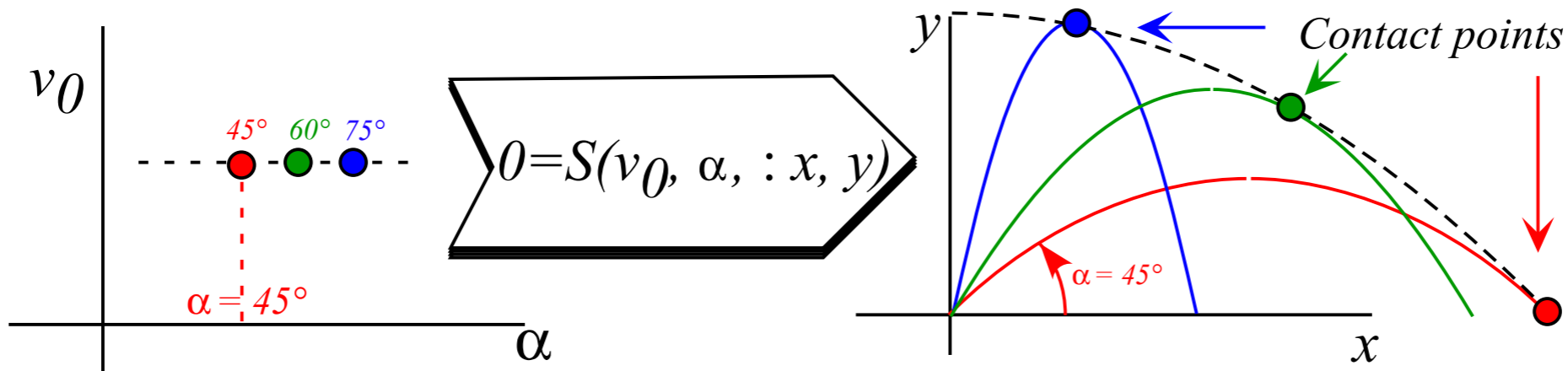
*Envelopes* of the  $v_0$ -trajectory region contain extremal *contact points* with each trajectory

where:  $\frac{\partial S(v_0, \alpha: x, y)}{\partial \alpha} = 0$

$$x \frac{\partial \tan \alpha}{\partial \alpha} - \frac{gx^2}{2v_0^2} \frac{\partial \cos^{-2} \alpha}{\partial \alpha} = 0 = \frac{x}{\cos^2 \alpha} - \frac{gx^2}{2v_0^2} \frac{2 \sin \alpha}{\cos^3 \alpha} \quad \text{gives:} \quad \tan \alpha = \frac{v_0^2}{gx} \quad \text{or:} \quad x = \frac{v_0^2}{g \tan \alpha}$$

Convert  $y(x)$  solution into Active Contact Transformation Generator  $S(v_0, \alpha: x, y)$

$$y(x) = x \tan \alpha - \frac{gx^2}{2v_0^2 \cos^2 \alpha} \quad \text{becomes:} \quad S(v_0, \alpha: x, y) = -y + x \tan \alpha - \frac{gx^2}{2v_0^2 \cos^2 \alpha} = 0$$



Unit 1  
Fig. 12.6

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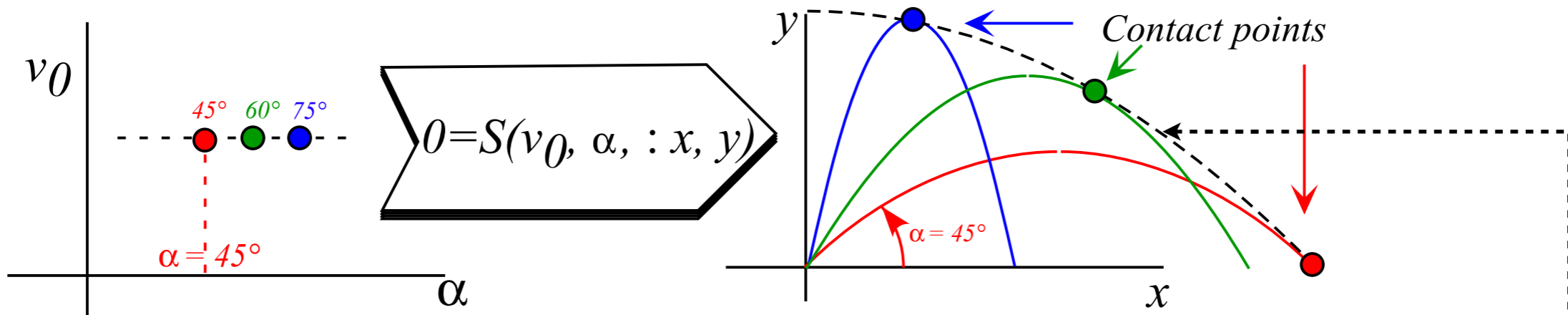
$$x \frac{\partial \tan \alpha}{\partial \alpha} - \frac{gx^2}{2v_0^2} \frac{\partial \cos^{-2} \alpha}{\partial \alpha} = 0 = \frac{x}{\cos^2 \alpha} - \frac{gx^2}{2v_0^2} \frac{2 \sin \alpha}{\cos^3 \alpha}$$

$$\tan \alpha = \frac{v_0^2}{gx} \quad \text{or:} \quad x = \frac{v_0^2}{g \tan \alpha}$$

$$y_{env}(x) = x \tan \alpha - \frac{gx^2}{2v_0^2} (1 + \tan^2 \alpha) \Rightarrow y_{env}(x) = x \frac{v_0^2}{gx} - \frac{gx^2}{2v_0^2} \left( 1 + \frac{v_0^4}{g^2 x^2} \right)$$

Convert  $y(x)$  solution into Active Contact Transformation Generator  $S(v_0, \alpha: x, y)$

$$y(x) = x \tan \alpha - \frac{gx^2}{2v_0^2 \cos^2 \alpha} \quad \text{becomes:} \quad S(v_0, \alpha: x, y) = -y + x \tan \alpha - \frac{gx^2}{2v_0^2 \cos^2 \alpha} = 0$$



*Envelopes* of the  $v_0$ -trajectory region contain extremal *contact points* with each trajectory

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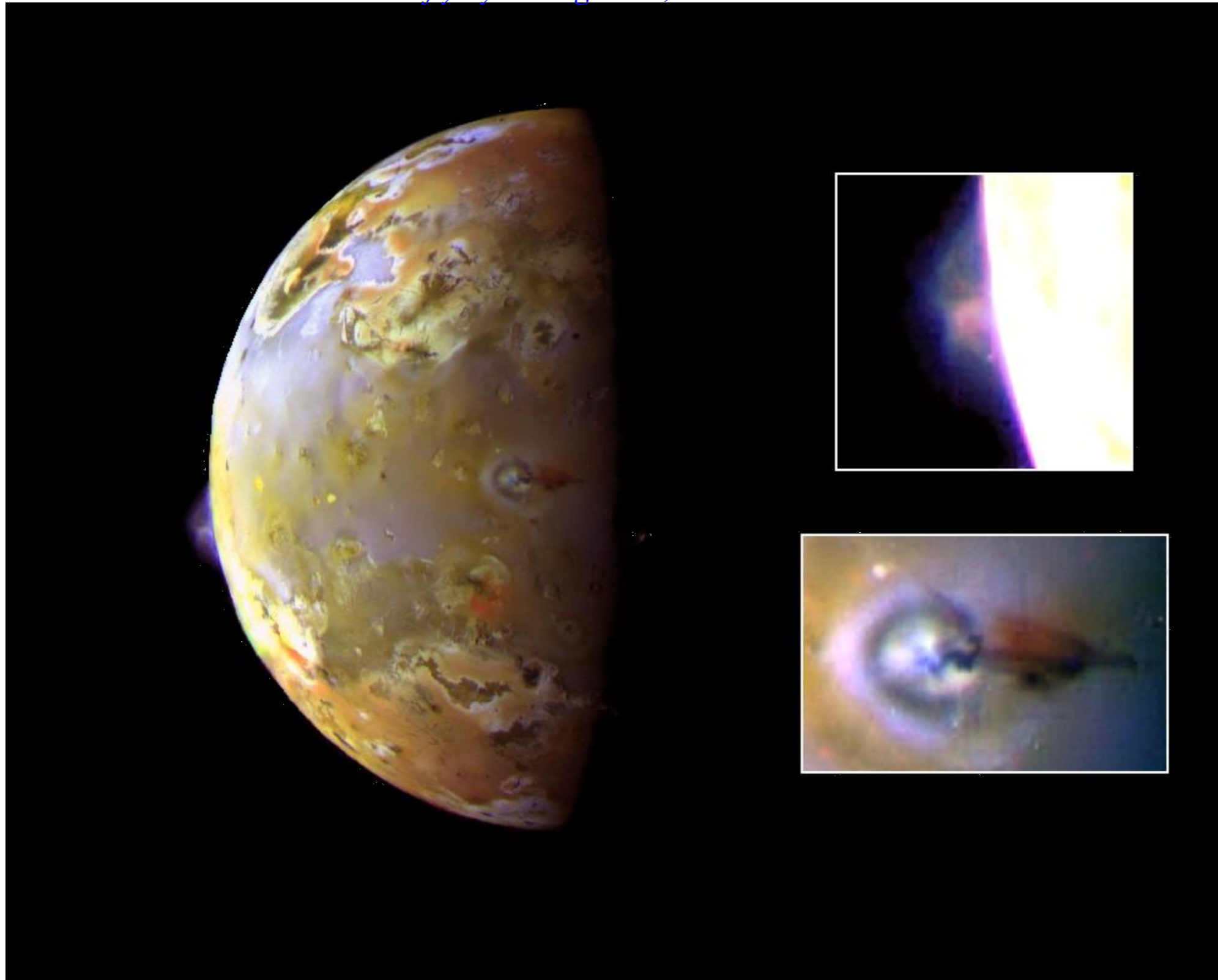
$$y_{env}(x) = \frac{v_0^2}{g} - \frac{gx^2}{2v_0^2} - \frac{gx^2}{2v_0^2} \frac{v_0^4}{g^2 x^2} = \frac{v_0^2}{2g} - \frac{gx^2}{2v_0^2}$$

*Envelope function*

# *The Plumes of Prometheus*

*NASA-Galileo Project*

*Io fly-by on August 18, 1997*



[http://antwrp.gsfc.nasa.gov/apod/image/9708/prometheus\\_gal\\_big.jpg](http://antwrp.gsfc.nasa.gov/apod/image/9708/prometheus_gal_big.jpg)

<http://antwrp.gsfc.nasa.gov/apod/ap970818.html>

[http://science.nasa.gov/science-news/science-at-nasa/1999/ast04oct99\\_1/](http://science.nasa.gov/science-news/science-at-nasa/1999/ast04oct99_1/)

# IO'S ALIEN VOLCANOES



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## IO'S ALIEN VOLCANOES

SCIENTISTS ARE EAGER FOR A CLOSER LOOK AT THE SOLAR SYSTEM'S STRANGEST AND MOST ACTIVE VOLCANOES WHEN GALILEO FLIES BY IO ON OCTOBER 11.

**October 4, 1999:** Thirty years ago, before the Voyager probes visited Jupiter, if you had described Io to a literary critic it would have been declared overwrought science fiction. Jupiter's strange moon is literally bursting with volcanoes. Dozens of active vents pepper the landscape which also includes gigantic frosty plains, towering mountains and volcanic rings the size of California. The volcanoes themselves are the hottest spots in the solar system with temperatures exceeding 1800 K (1527 C). The plumes which rise 300 km into space are so large they can be seen from Earth by the Hubble Space Telescope. Confounding common sense, these high-rising ejecta seem to be made up of, not blisteringly hot lava, but frozen sulfur dioxide. And to top it all off, Io bears a striking resemblance to a pepperoni pizza. Simply unbelievable.

**Right:** Digital Radiance simulation of Pillan Patera just before the Galileo flyby. [click for animation](#) → .



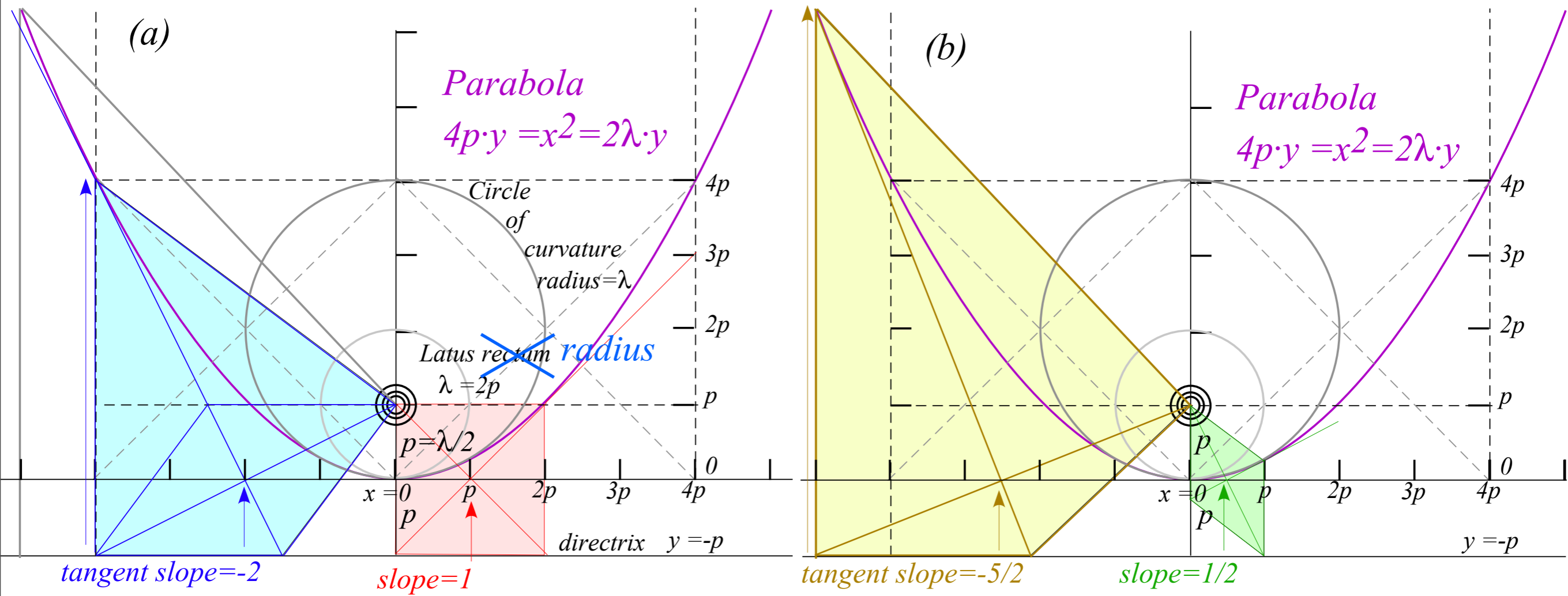
*Pretty bad sketch of plumes  
(Las Vegas model of planetary ejecta?)*

*Do these guys need a geometry lesson?*

*Go fly a kite?*

...conventional parabolic geometry...carried to extremes...


Recall Lecture 6 p.29



Unit 1  
Fig. 9.4

## *A general contact transformation from sophomore physics*

*Algebra-calculus development of “The Volcanoes of Io” and “The Atoms of NIST”*

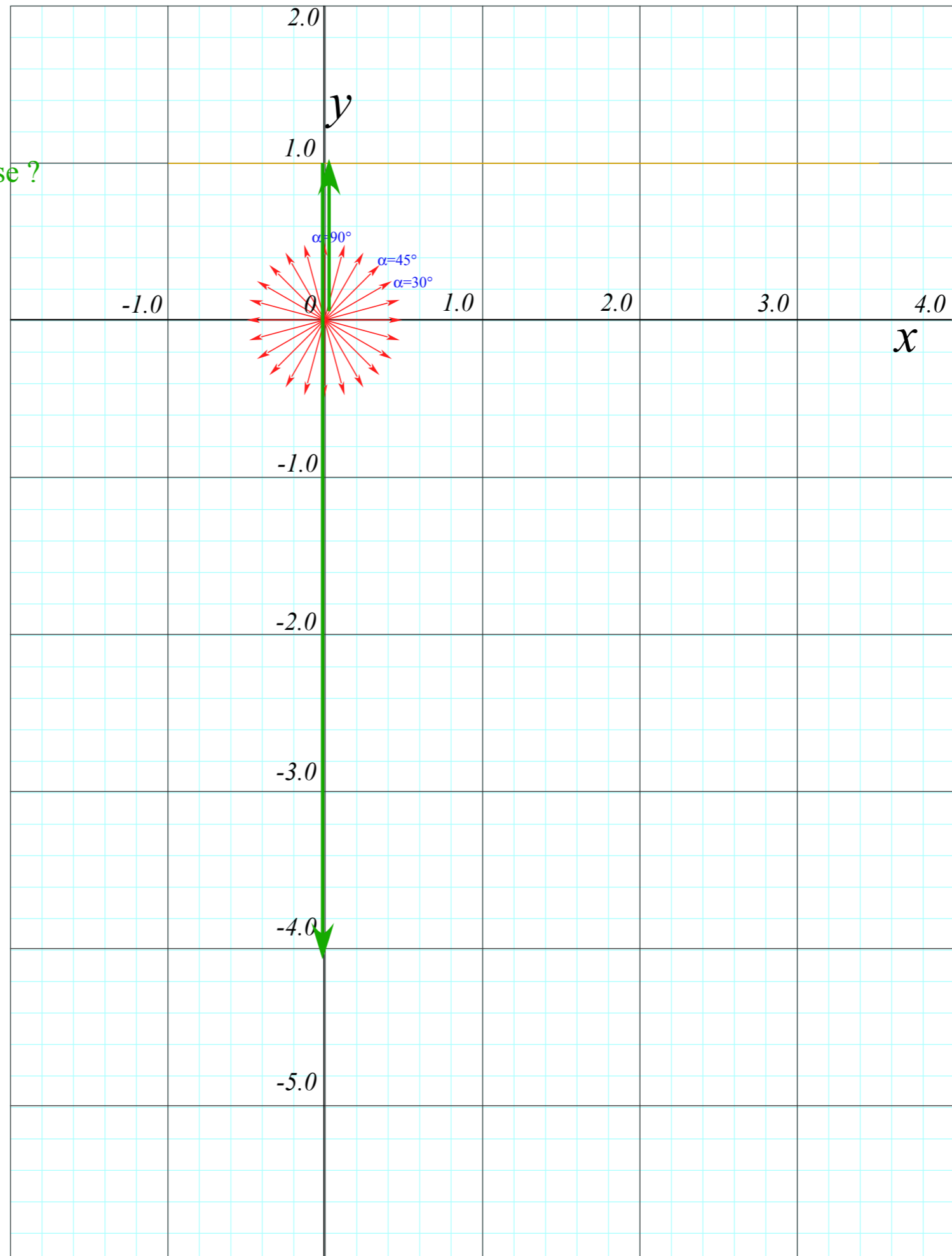
 *Intuitive-geometric development of ” ” ” and ” ” ”*

Say  $\alpha=90^\circ$  path rises to 1.0  
then drops. When at  $y=1.0$ ...

Q1. ...where is its focus?

Q2. ...where is the **blast wave**?

Q3. ...how high can  $\alpha=45^\circ$  path path rise ?







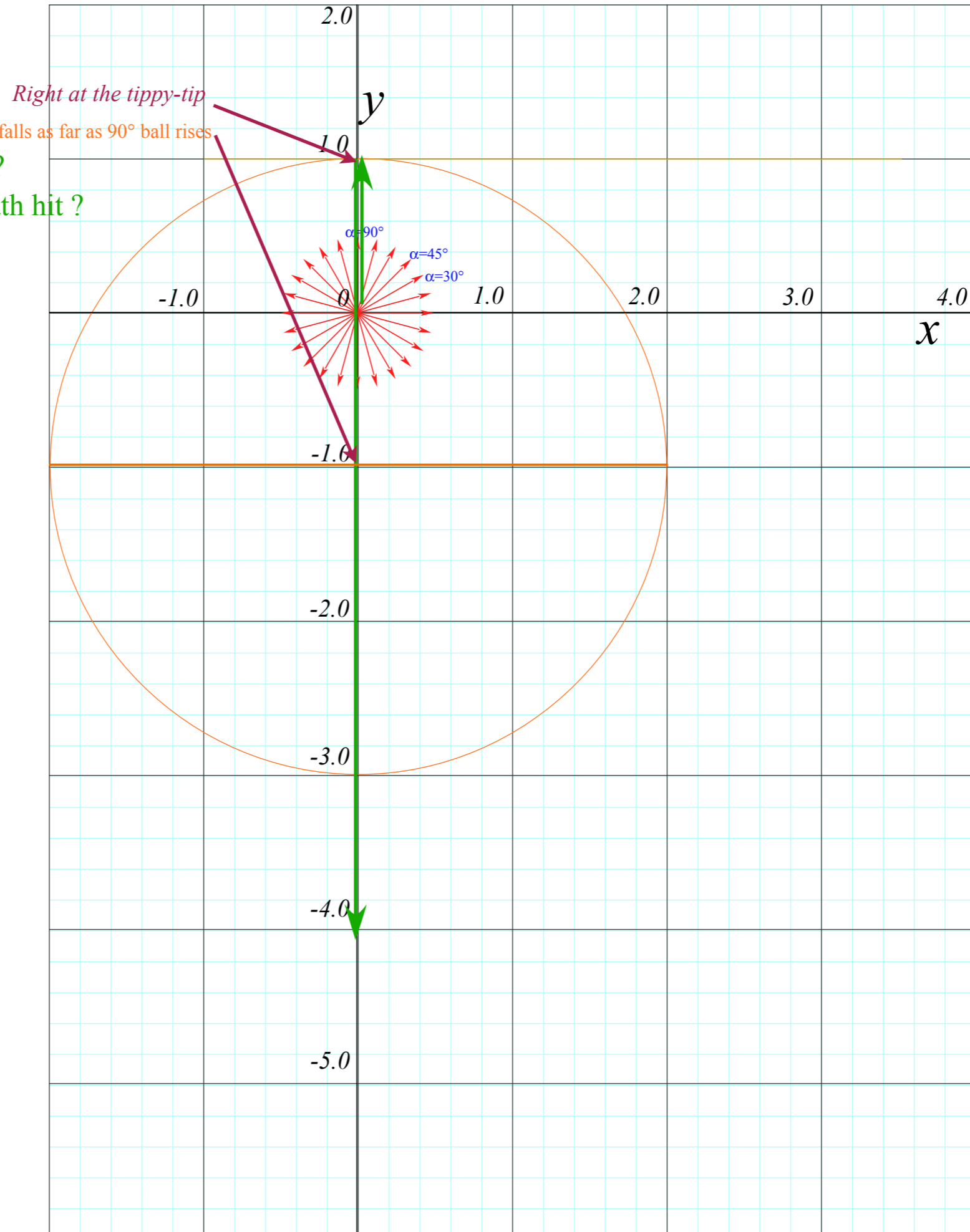
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Q2. ...where is the **blast wave**? center falls as far as  $90^\circ$  ball rises

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Q4. Where on  $x$ -axis does  $\alpha=45^\circ$  path hit ?



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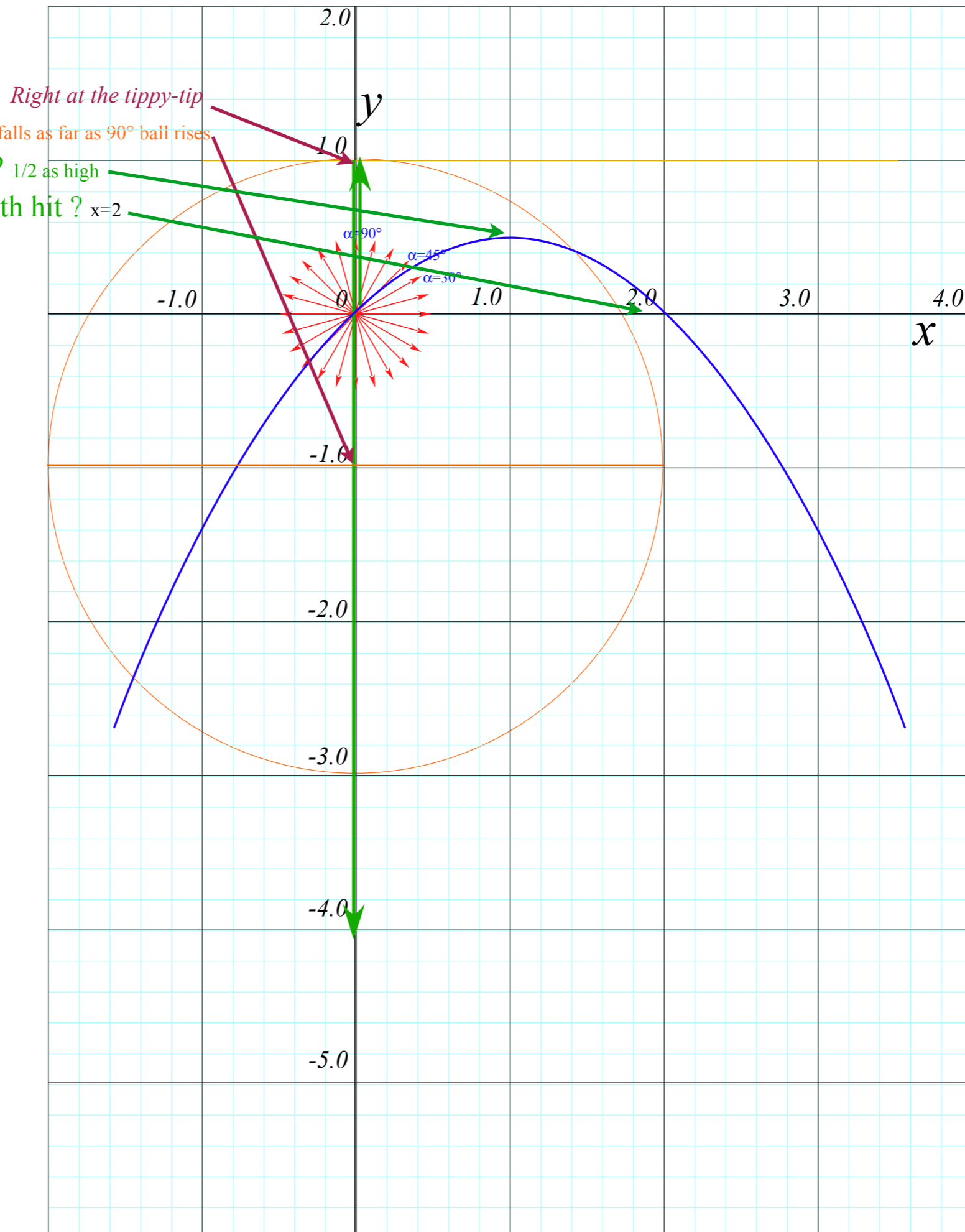
Q3. How high can  $\alpha=45^\circ$  path rise? 1/2 as high

Q4. Where on  $x$ -axis does  $\alpha=45^\circ$  path hit?  $x=2$

Q5. Where is **blast wave** then?

Q6. Where is  $\alpha=45^\circ$  path focus?

Q7. Guess for **all-path envelope**? and its focus? directrix?







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Q5. Where is **blast wave** then? centered on  $45^\circ$  normal

Q6 Where is  $\alpha=45^\circ$  path focus?  $x=1, y=0$

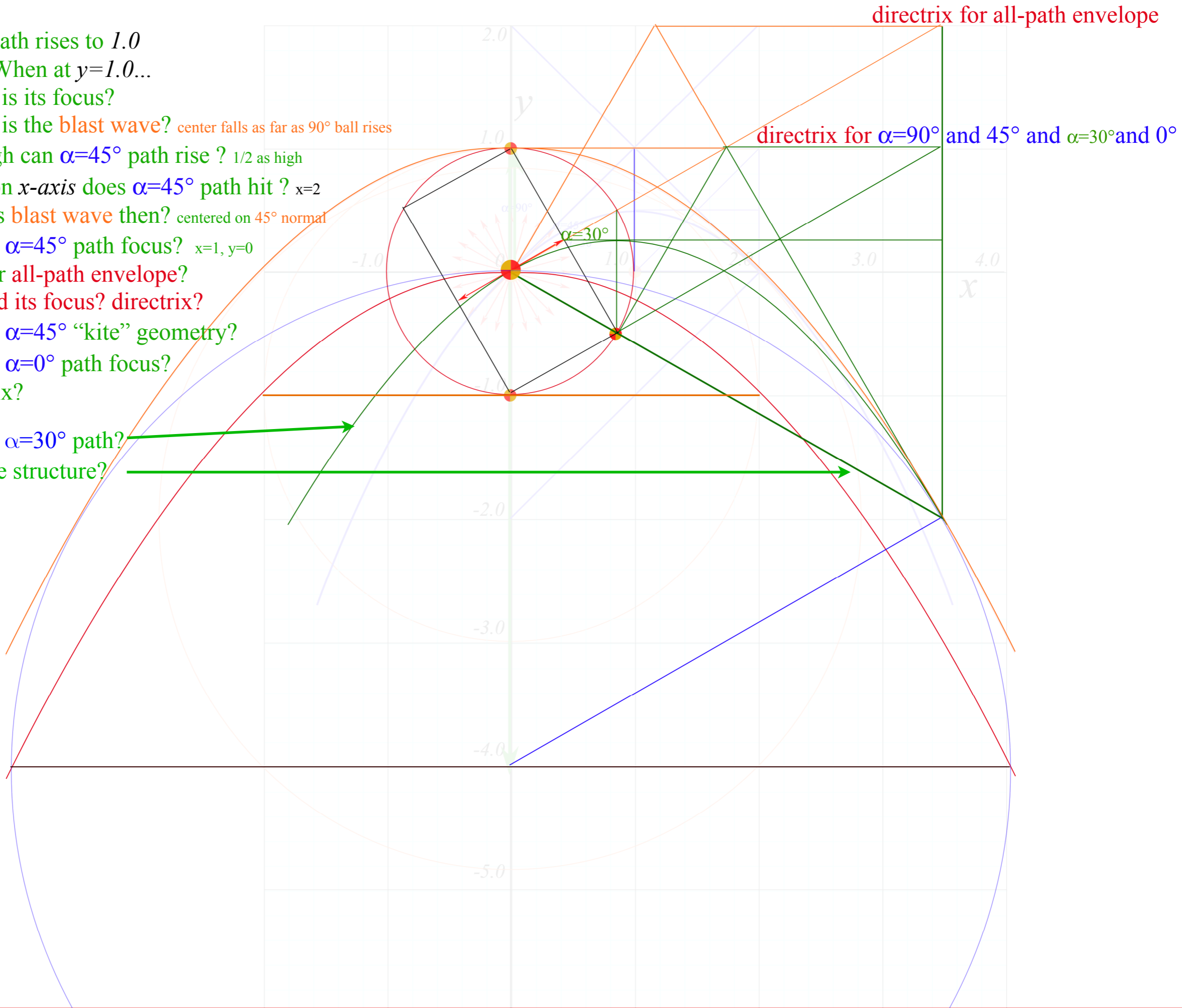
Q7 Guess for **all-path envelope**?  
and its focus? directrix?

Q7 Where is  $\alpha=45^\circ$  "kite" geometry?

Q8 Where is  $\alpha=0^\circ$  path focus?  
directrix?

Where is  $\alpha=30^\circ$  path?

...and kite structure?



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