

Lecture 7
Tue. 9.13.2016

Kepler Geometry of IHO (Isotropic Harmonic Oscillator) Elliptical Orbits

(Ch. 9 and Ch. 11 of Unit 1)

Constructing 2D IHO orbital phasor “clock” dynamics in uniform-body

Constructing 2D IHO orbits using Kepler anomaly plots

Mean-anomaly and eccentric-anomaly geometry

Calculus and vector geometry of IHO orbits

A confusing introduction to Coriolis-centrifugal force geometry (Derived better in Ch. 12)

Some Kepler’s “laws” for all central (isotropic) force $F(r)$ fields

Angular momentum invariance of IHO: $F(r)=-k\cdot r$ with $U(r)=k\cdot r^2/2$ (Derived here)

Angular momentum invariance of Coulomb: $F(r)=-GMm/r^2$ with $U(r)=-GMm\cdot/r$ (Derived in Unit 5)

Total energy $E=KE+PE$ invariance of IHO: $F(r)=-k\cdot r$ (Derived here)

Total energy $E=KE+PE$ invariance of Coulomb: $F(r)=-GMm/r^2$ (Derived in Unit 5)

Introduction to dual matrix operator contact geometry (based on IHO orbits)

Quadratic form ellipse $\mathbf{r}\cdot\mathbf{Q}\cdot\mathbf{r}=1$ vs. inverse form ellipse $\mathbf{p}\cdot\mathbf{Q}^{-1}\cdot\mathbf{p}=1$

Duality norm relations ($\mathbf{r}\cdot\mathbf{p}=1$)

Q -Ellipse tangents \mathbf{r}' normal to dual Q^{-1} -ellipse position \mathbf{p} ($\mathbf{r}'\cdot\mathbf{p}=0=\mathbf{r}\cdot\mathbf{p}'$)

Operator geometric sequences and eigenvectors

Alternative scaling of matrix operator geometry

Vector calculus of tensor operation

Q: Where is this headed? A: Lagrangian-Hamiltonian duality

[Link \$\Rightarrow\$ BoxIt simulation of IHO orbits](#)

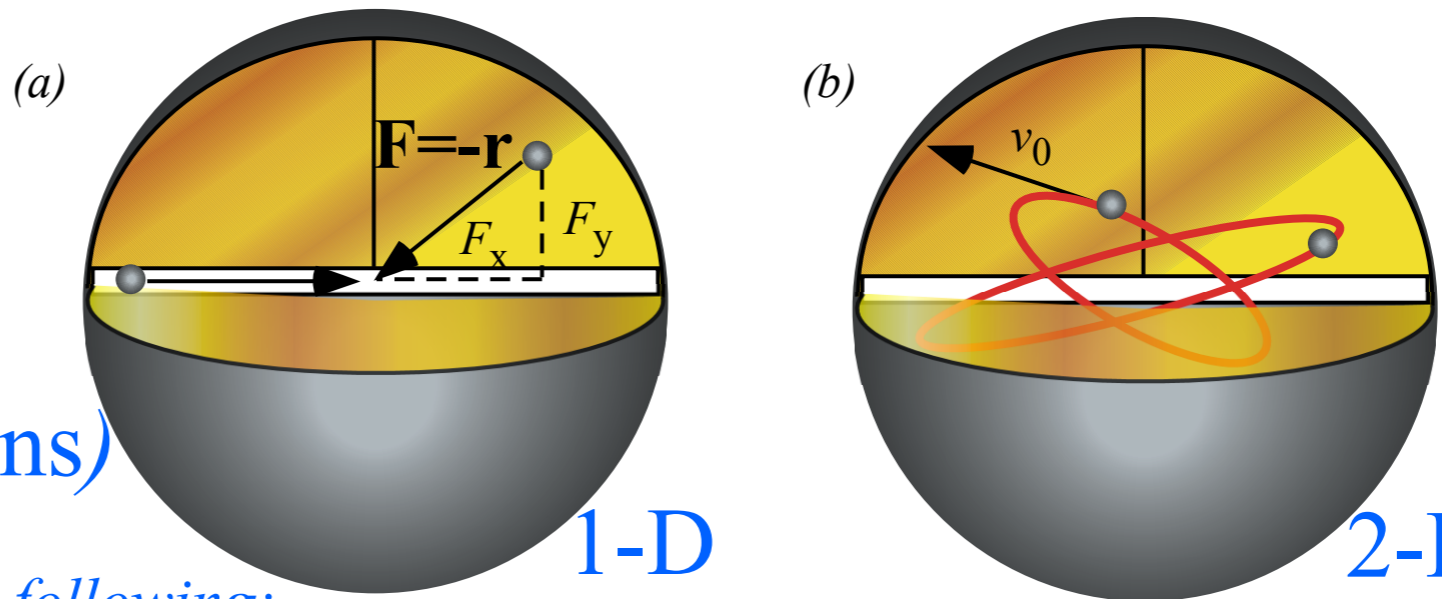
[Link \$\rightarrow\$ IHO orbital time rates of change](#)

[Link \$\rightarrow\$ IHO Exegesis Plot](#)

→ *Review of IHO orbital phasor “clock” dynamics in uniform-body with two “movie” examples*

Review of IHO orbital phase dynamics in uniform-body

Unit 1
Fig. 9.10



1-D

2-D or 3-D

(Paths are *always* 2-D ellipses if viewed right!)

I.H.O. Force law

$F = -x$ (1-Dimension)

$F = -r$ (2 or 3-Dimensions)

Each dimension x , y , or z obeys the following:

$$\text{Total } E = KE + PE = \frac{1}{2}mv^2 + U(x) = \frac{1}{2}mv^2 + \frac{1}{2}kx^2 = \text{const.}$$

Equations for x -motion

$[x(t)$ and $v_x=v(t)]$ are given first. They apply as well to dimensions

$[y(t)$ and $v_y=v(t)]$ and $[z(t)$ and $v_z=v(t)]$ in the ideal isotropic case.

$$1 = \frac{mv^2}{2E} + \frac{kx^2}{2E} = \left(\frac{v}{\sqrt{2E/m}} \right)^2 + \left(\frac{x}{\sqrt{2E/k}} \right)^2$$

$$1 = \frac{mv^2}{2E} + \frac{kx^2}{2E} = (\cos\theta)^2 + (\sin\theta)^2$$

Another example of the old “scale-a-circle” trick...

Let : **(1)** $v = \sqrt{2E/m} \cos\theta$, and : **(2)** $x = \sqrt{2E/k} \sin\theta$ angular velocity: $\omega = \frac{d\theta}{dt}$ def. **(3)**

$$\sqrt{\frac{2E}{m}} \cos\theta = v = \frac{dx}{dt} = \frac{d\theta}{dt} \frac{dx}{d\theta} = \omega \frac{dx}{d\theta} = \omega \sqrt{\frac{2E}{k}} \cos\theta$$

by (1) by def. (3) by (2)

$$\omega = \frac{d\theta}{dt} = \sqrt{\frac{k}{m}}$$

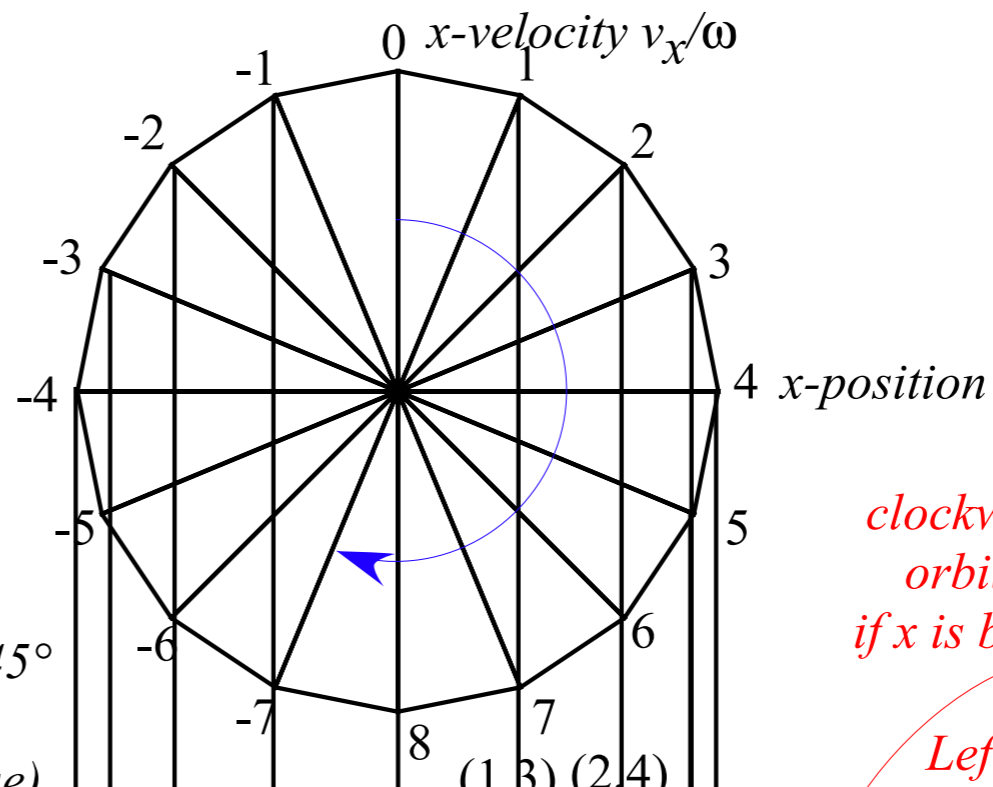
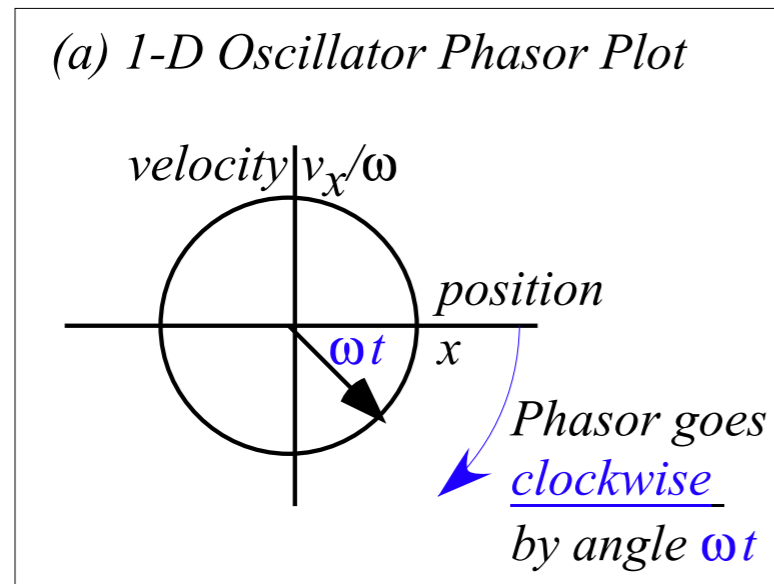
by def. (3) divide this by (1)

$$\theta = \int \omega \cdot dt = \omega \cdot t + \alpha$$

by integration given constant ω .

Review of IHO orbital phase dynamics in uniform-body

Unit 1
Fig. 9.10



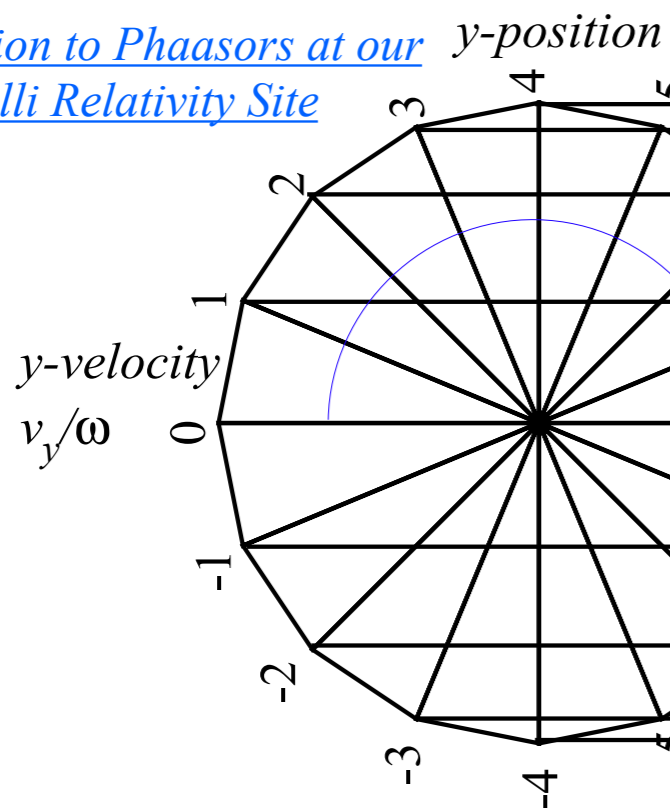
clockwise orbit if x is behind y

Left-handed

(b) 2-D Oscillator Phasor Plot

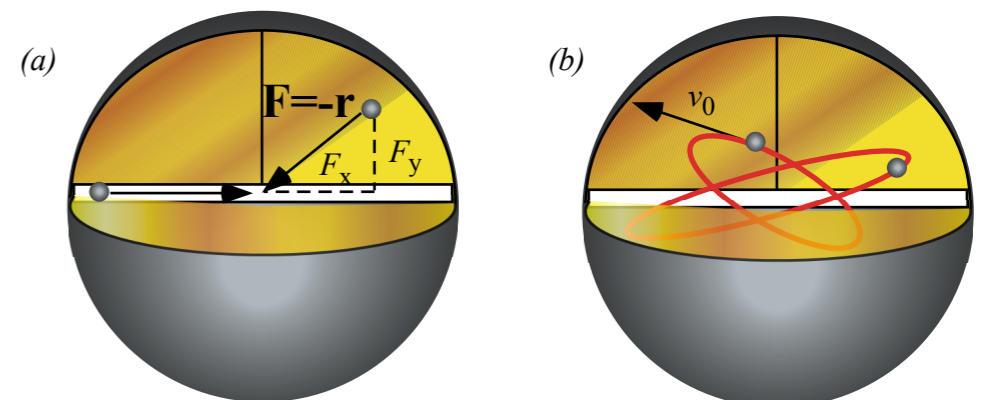
(x -Phase 45° behind the y -Phase)

[Introduction to Phasors at our Pirelli Relativity Site](#)



counter-clockwise if y is behind x

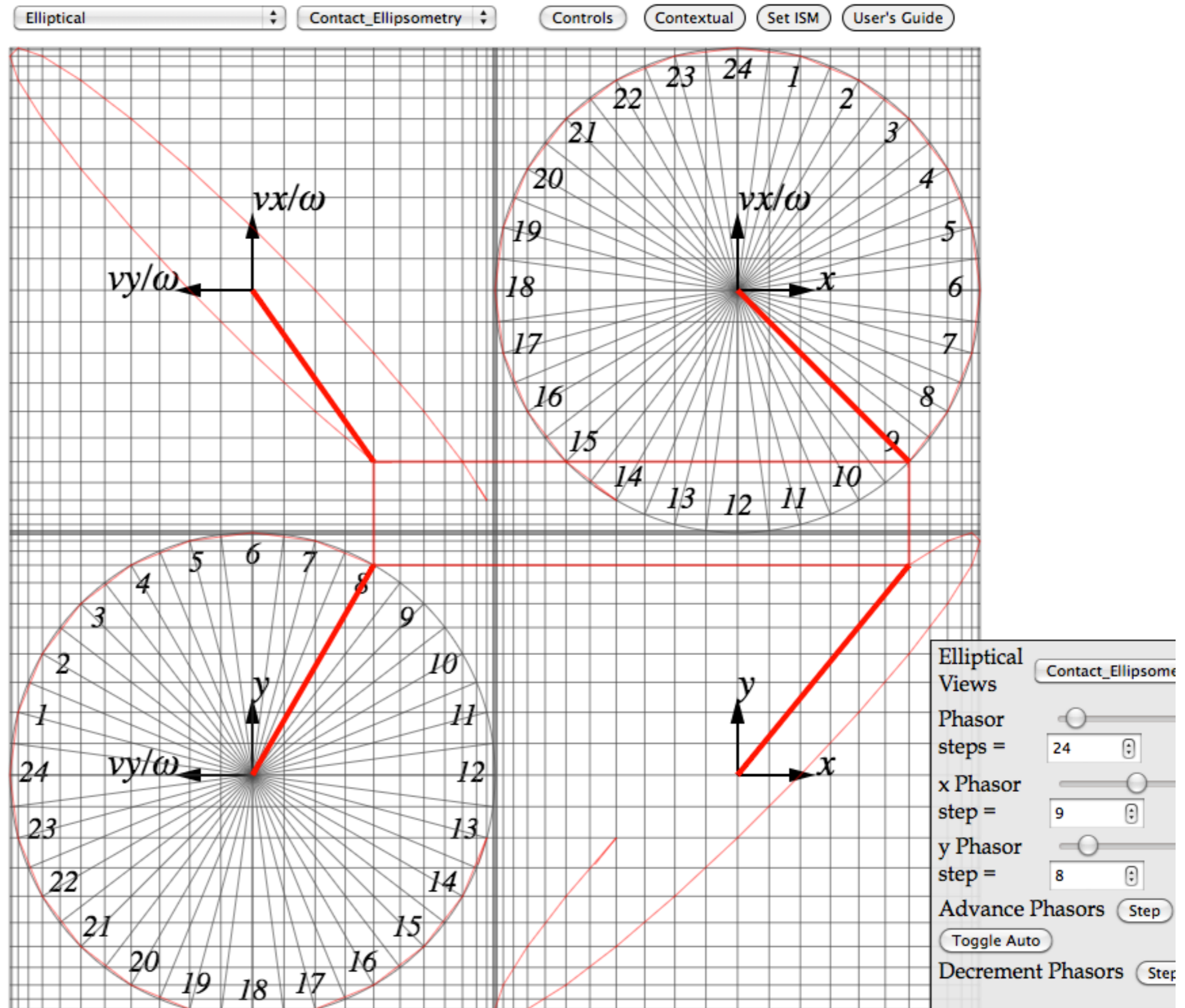
Right-handed



[BoxIt web simulation - With \$y\$ -Phasor is on other side of \$xy\$ plot](#)

[RelaWavity web simulation - Contact ellipsometry](#)

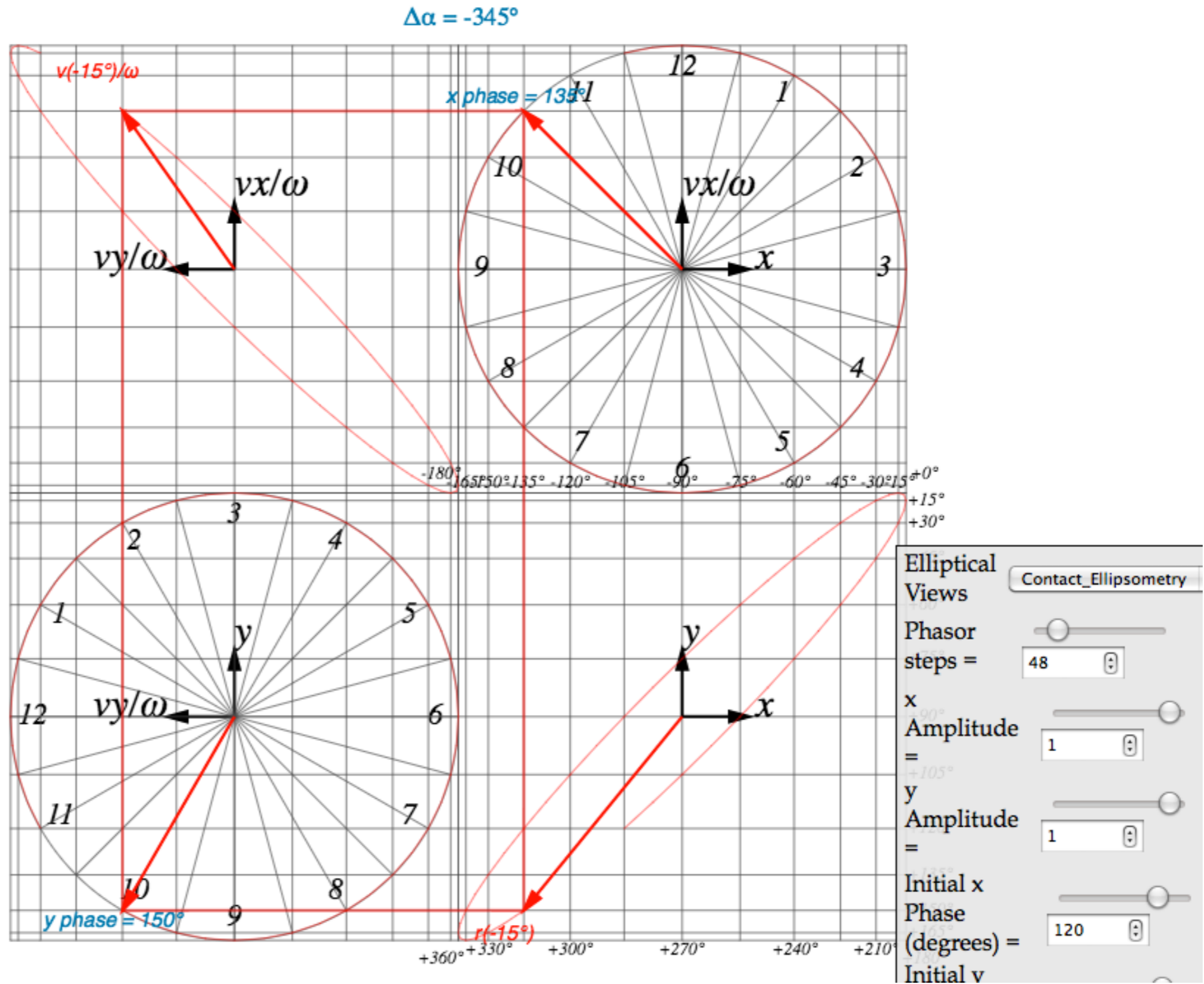
*RelaWavity
ellipsometry
web-app*



[RelaWavity Web Simulation](#)
[Ellipsometry](#)

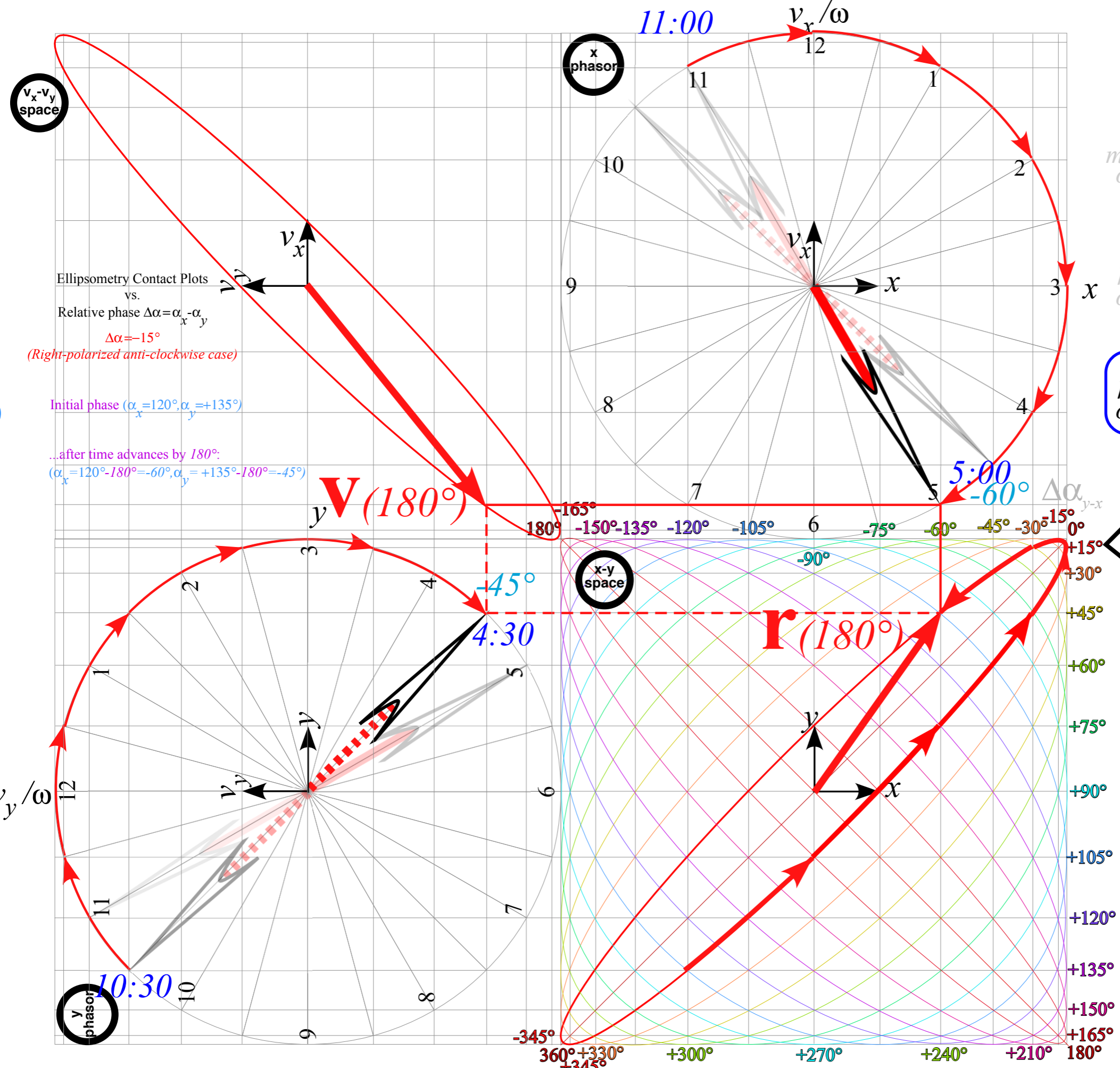
Geometry of Kepler anomalies for vectors $[\mathbf{r}(\phi), \mathbf{v}(\phi)]$ in coordinate (x, y) space rendered by animation web-apps BoxIt and RelaWavity described below after p.70.

RelaWavity
ellipsometry
web-app



RelaWavity Web Simulation
Ellipsometry

Geometry of Kepler anomalies for vectors $[\mathbf{r}(\phi), \mathbf{v}(\phi)]$ in coordinate (x,y) space rendered by animation web-apps BoxIt and RelaWavity described below after p.7 and p.17.



phase lag:
 $\Delta\alpha = \alpha_x - \alpha_y = 15^\circ$
 1 minute orbit (2.5 seconds for second hand)
 or
 1 hour orbit (2.5 minutes for minute hand)
 or
 12 hour orbit (1/2 hour for hour hand)

$\Delta\alpha = \alpha_x - \alpha_y = 15^\circ$

*Constructing 2D IHO orbits using **Kepler anomaly plots***

 *Mean-anomaly and eccentric-anomaly geometry*

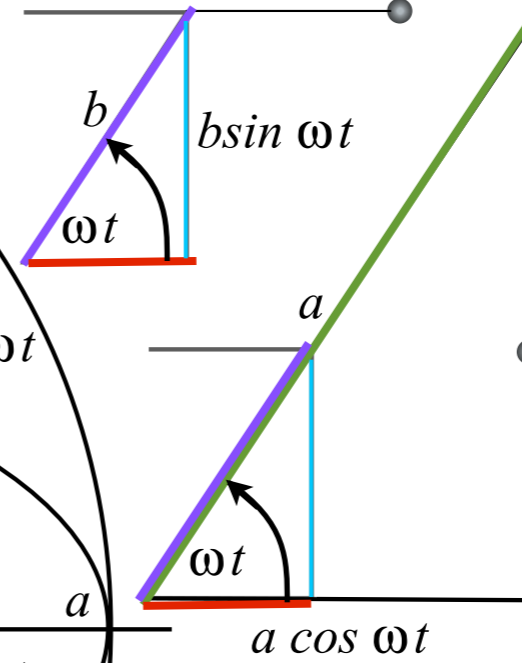
Calculus and vector geometry of IHO orbits

A confusing introduction to Coriolis-centrifugal force geometry (Derived better in Ch. 12)

Linear Harmonic
Force-Field
Orbits

Kepler's
Mean Anomaly Line
(slope angle $\theta = \omega t$)

Kepler's
Eccentric Anomaly Line
(slope is polar angle $\phi = \text{atan}[y/x]$)

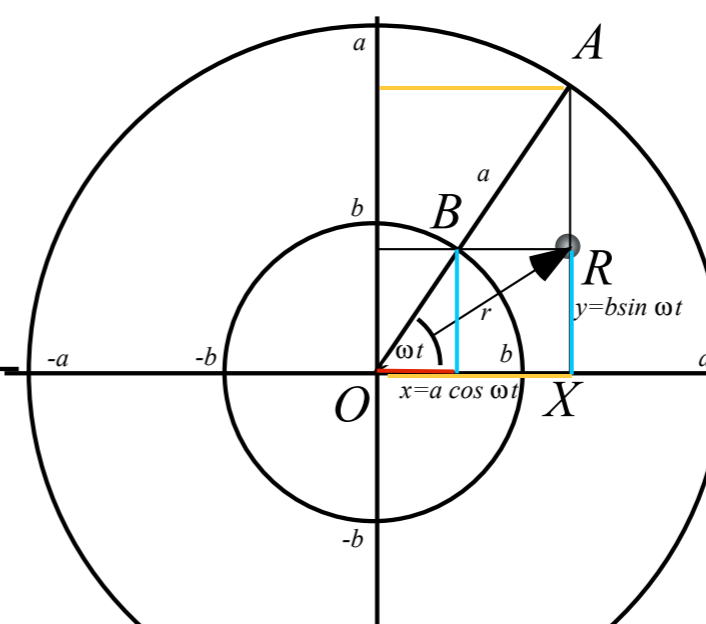
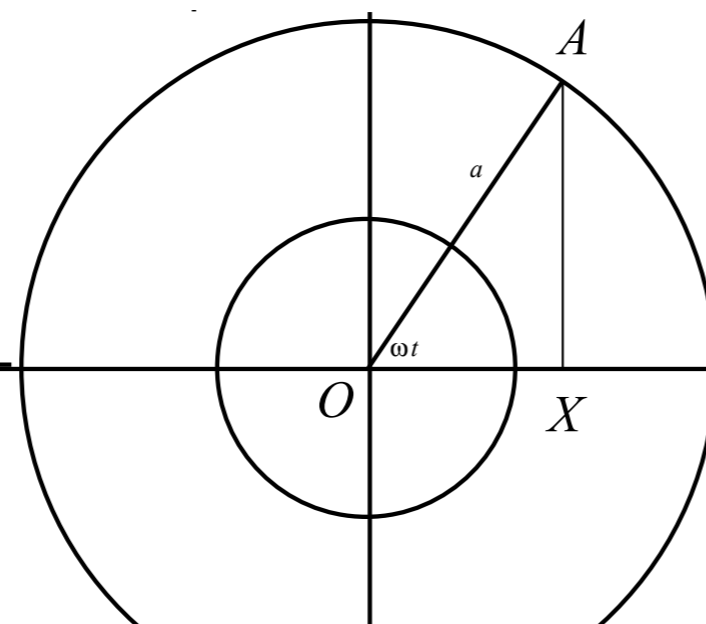
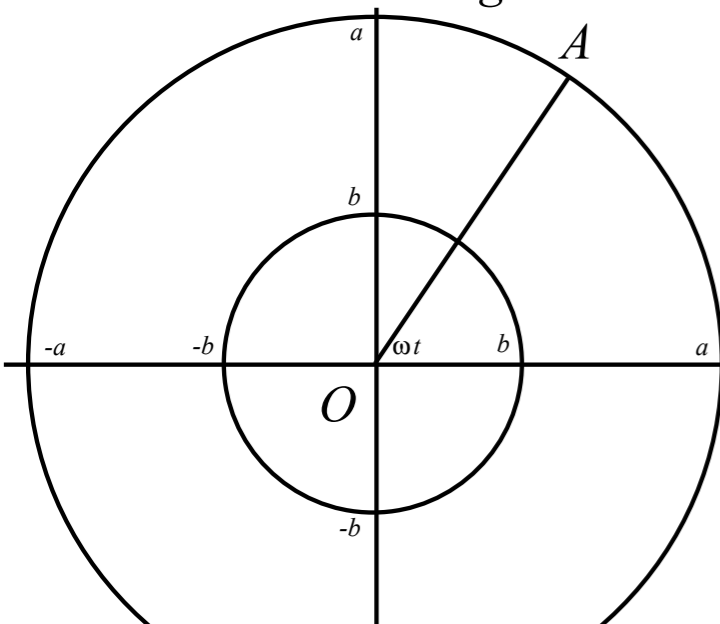


Unit 1
Fig. 11.1
(top 2/3's)

Step 1. Draw concentric circles of radius a and b and a radius OA at angle ωt

Step 2. Draw vertical line AX from a -circle at ωt to x -axis

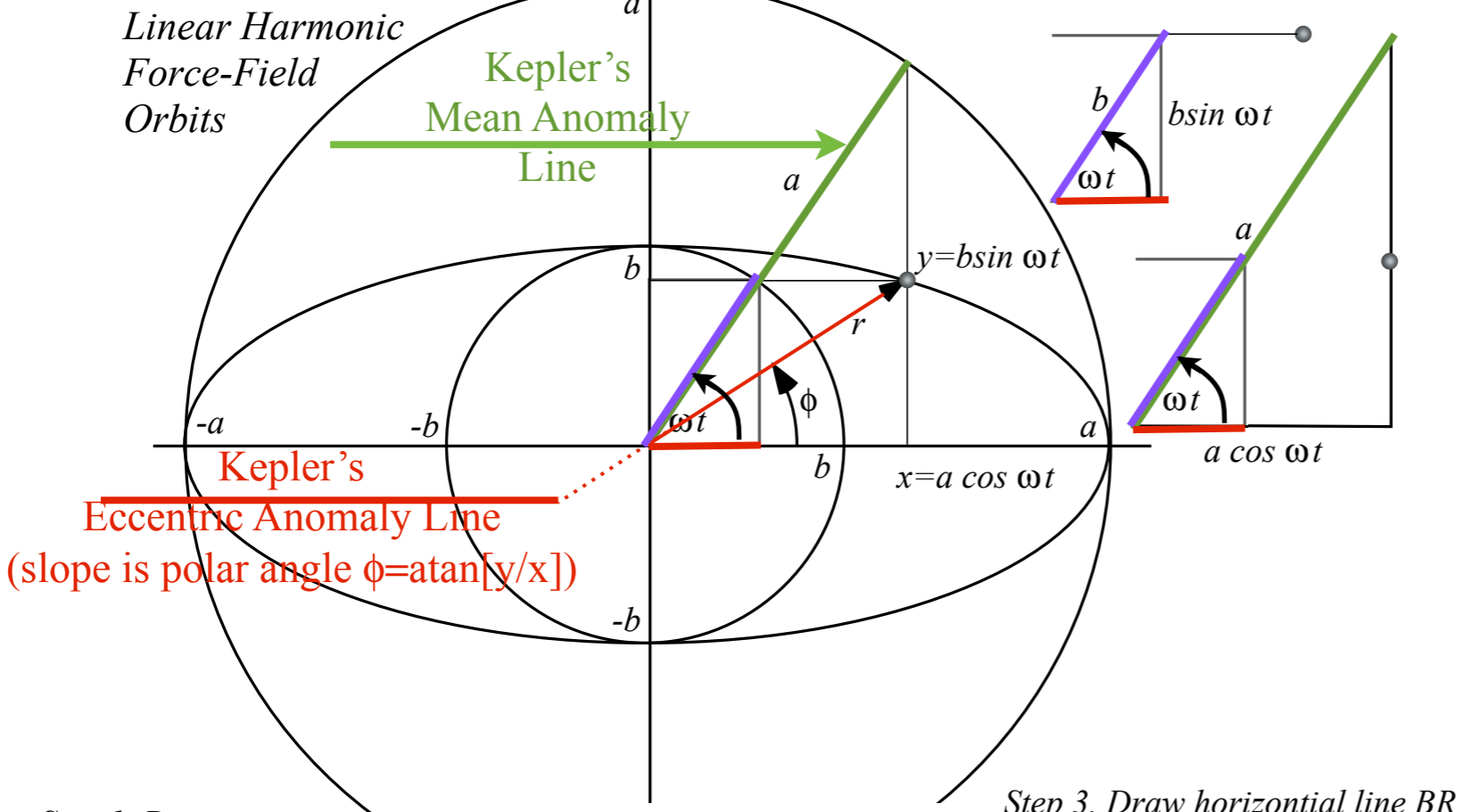
Step 3. Draw horizontal line BR from b -circle at ωt to line AX . Intersection is orbit point R .



Linear Harmonic
Force-Field
Orbits

Kepler's
Mean Anomaly
Line

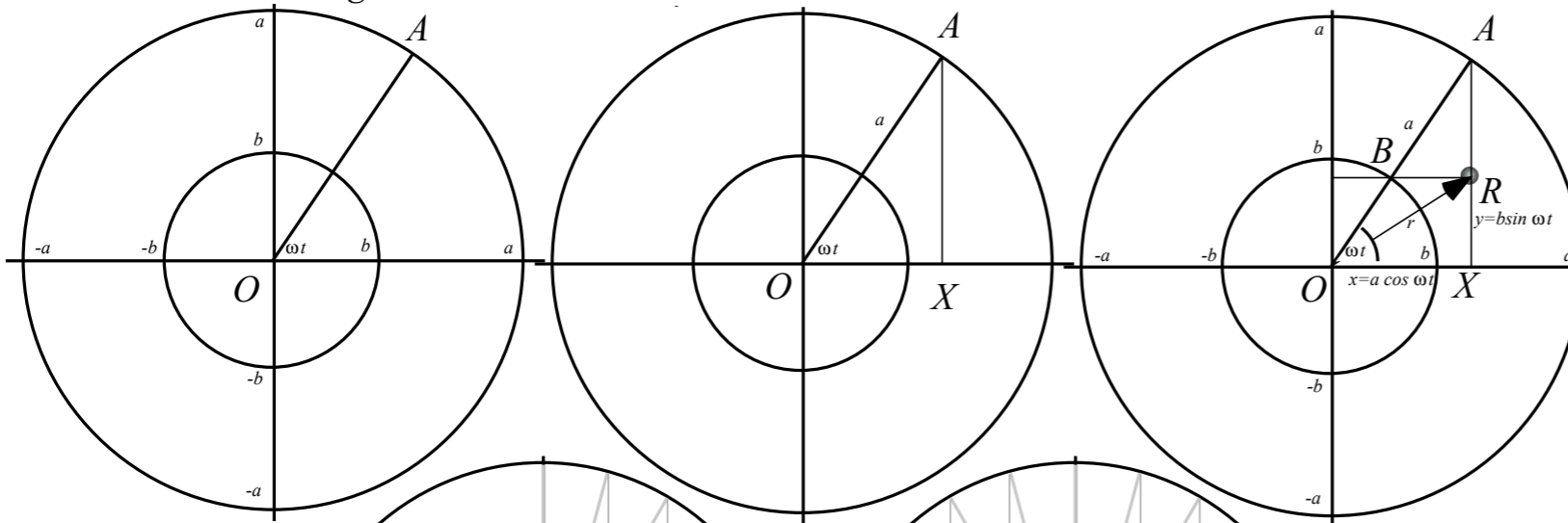
Kepler's
Eccentric Anomaly Line
(slope is polar angle $\phi = \text{atan}[y/x]$)



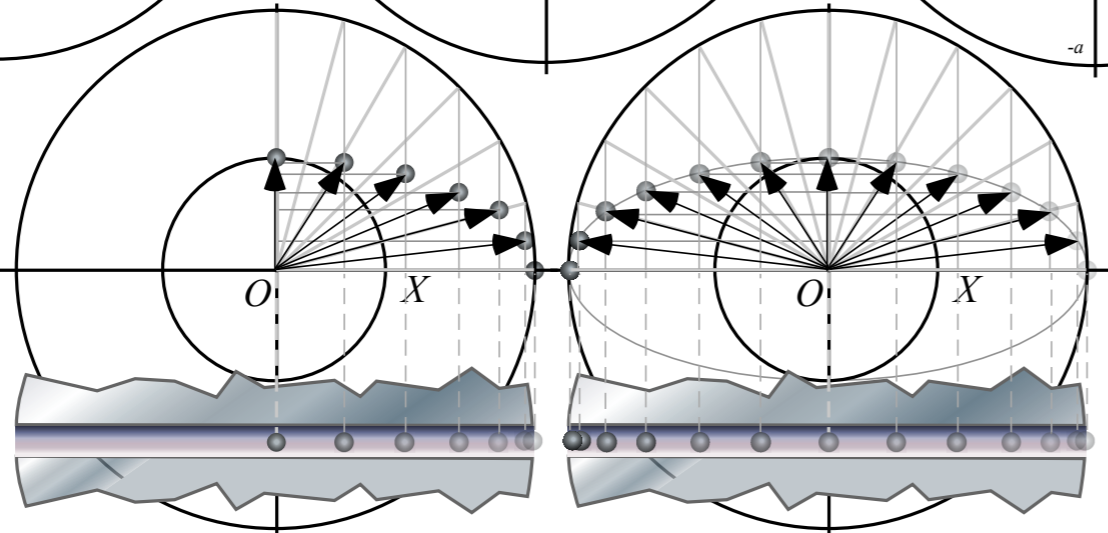
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Step 3. Draw horizontal line BR from b -circle at ωt to line AX . Intersection is orbit point R .



Step 4-N
Repeat
as often
as needed



Unit 1
Fig. 11.1

*Constructing 2D IHO orbits using **Kepler anomaly plots***

Mean-anomaly and eccentric-anomaly geometry

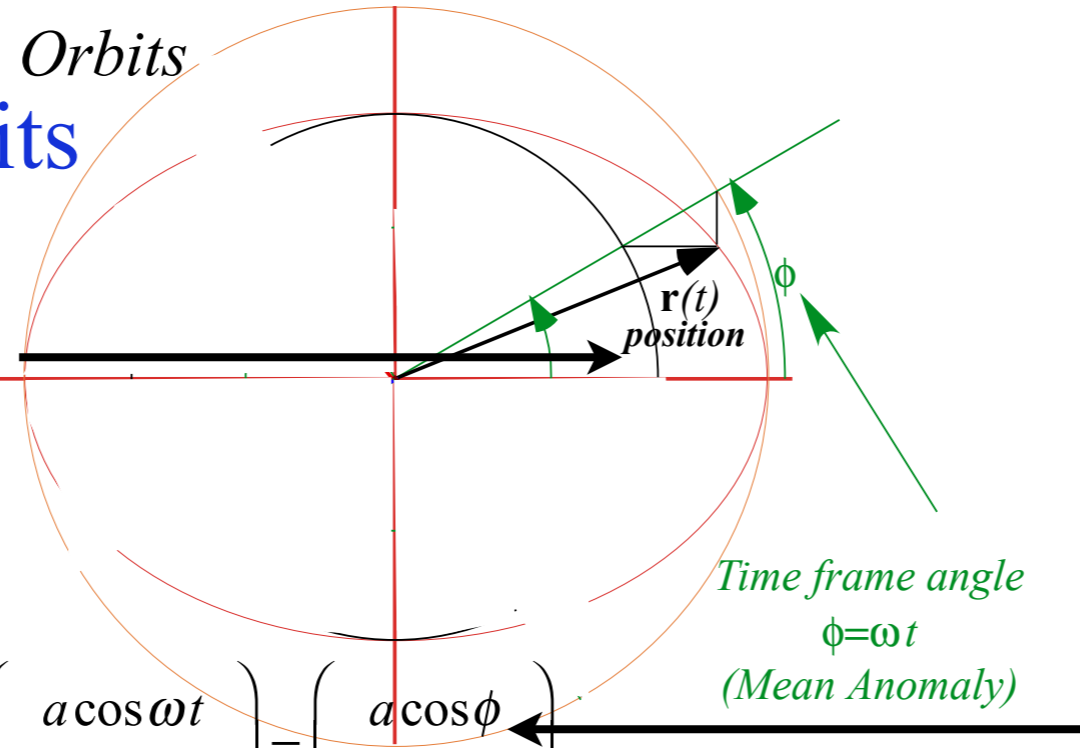
 *Calculus and vector geometry of IHO orbits*

A confusing introduction to Coriolis-centrifugal force geometry (Derived better in Ch. 12)

Calculus of IHO orbits

(a) Orbits

mean-anomaly ϕ of position vector \mathbf{r}



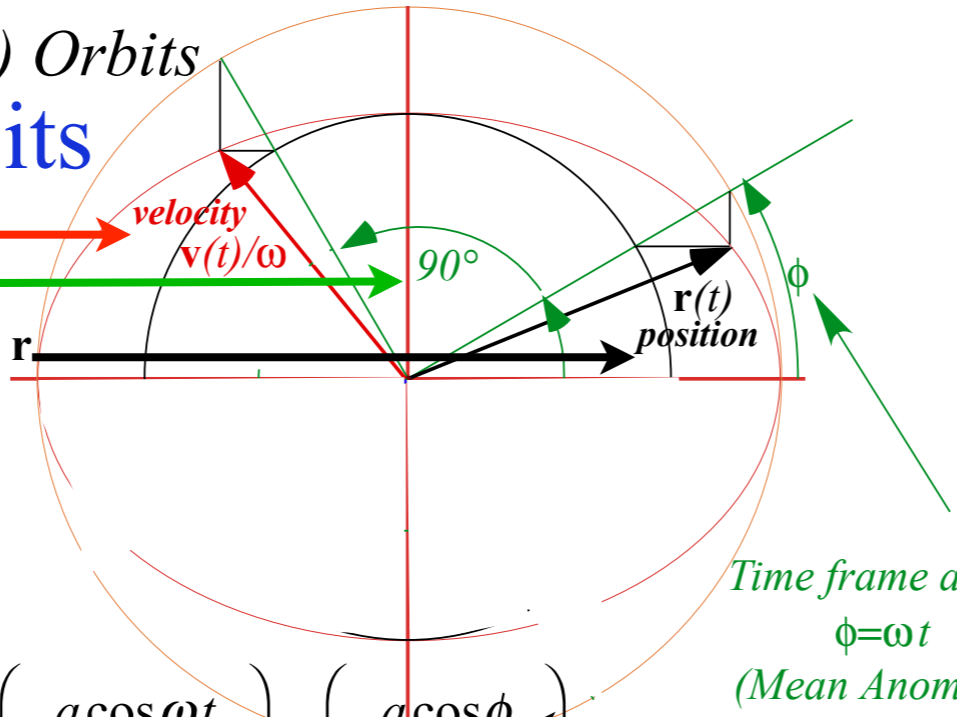
radius vector : $\mathbf{r} = \begin{pmatrix} r_x \\ r_y \end{pmatrix} = \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} a \cos \omega t \\ b \sin \omega t \end{pmatrix} = \begin{pmatrix} a \cos \phi \\ b \sin \phi \end{pmatrix}$

Unit 1
Fig. 11.5

Calculus of IHO orbits

(a) Orbits

To make velocity vector \mathbf{v} just rotate by $\pi/2$ or 90° the mean-anomaly ϕ of position vector \mathbf{r}



Time frame angle
 $\phi = \omega t$
 (Mean Anomaly)

$$\text{radius vector : } \mathbf{r} = \begin{pmatrix} r_x \\ r_y \end{pmatrix} = \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} a \cos \omega t \\ b \sin \omega t \end{pmatrix} = \begin{pmatrix} a \cos \phi \\ b \sin \phi \end{pmatrix}$$

mean-anomaly ϕ of position vector \mathbf{r} rotated by $\pi/2$ or 90° is *m.a.* of vector \mathbf{v}

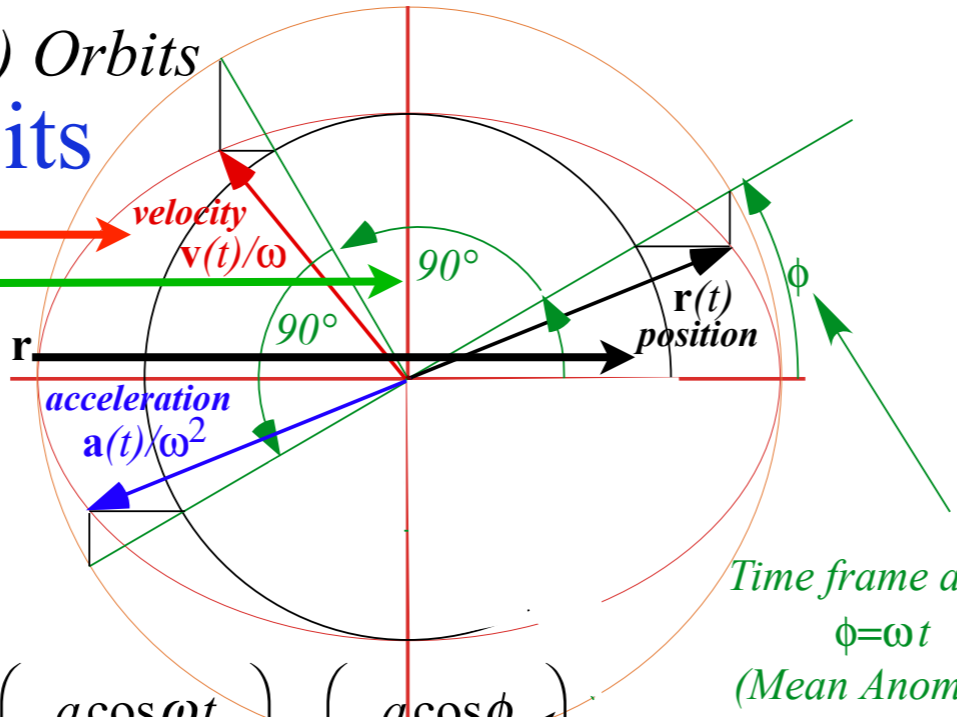
$$\text{velocity vector : } \mathbf{v} = \begin{pmatrix} v_x \\ v_y \end{pmatrix} = \begin{pmatrix} -a\omega \sin \omega t \\ b\omega \cos \omega t \end{pmatrix} = \frac{d\mathbf{r}}{dt} = \dot{\mathbf{r}} = \begin{pmatrix} a \cos \left(\phi + \frac{\pi}{2} \right) \\ b \sin \left(\phi + \frac{\pi}{2} \right) \end{pmatrix} \text{ (for } \omega = 1 \text{)}$$

Unit 1
 Fig. 11.5

Calculus of IHO orbits

(a) Orbits

To make velocity vector \mathbf{v} just rotate by $\pi/2$ or 90° the mean-anomaly ϕ of position vector \mathbf{r}



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m.a. $\phi + \pi/2$ of vector \mathbf{v} rotated by another $\pi/2$ is m.a. of vector \mathbf{a}

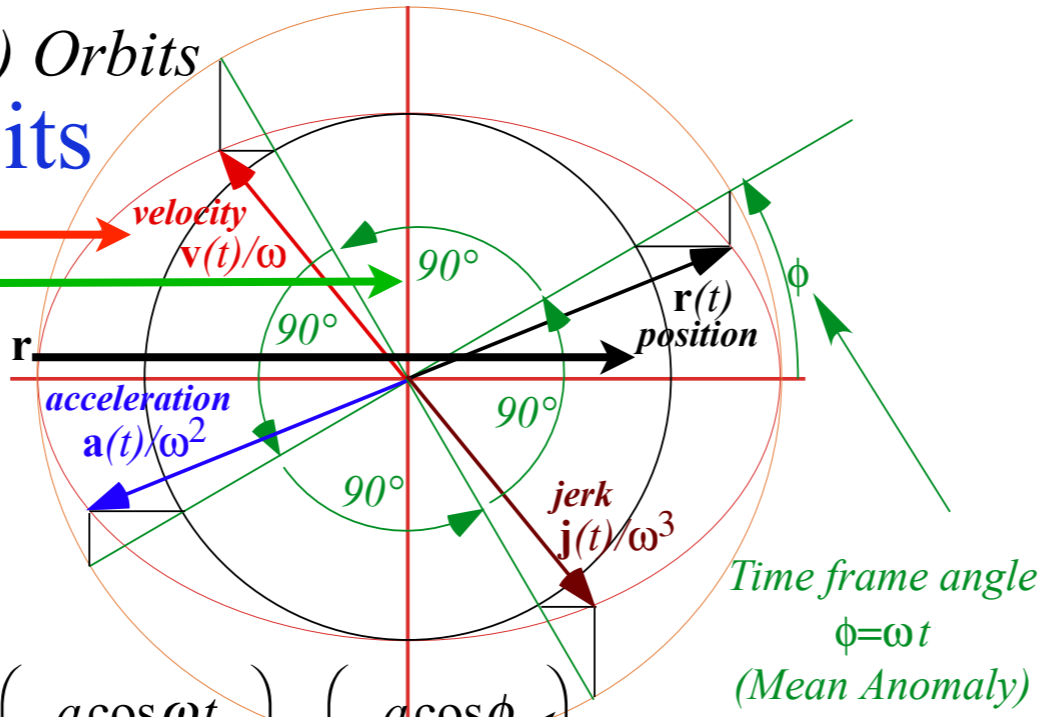
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Unit 1
Fig. 11.5

Calculus of IHO orbits

(a) Orbits

To make velocity vector \mathbf{v} just rotate by $\pi/2$ or 90° the mean-anomaly ϕ of position vector \mathbf{r}



$$\text{radius vector : } \mathbf{r} = \begin{pmatrix} r_x \\ r_y \end{pmatrix} = \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} a \cos \omega t \\ b \sin \omega t \end{pmatrix} = \begin{pmatrix} a \cos \phi \\ b \sin \phi \end{pmatrix}$$

mean-anomaly ϕ of position vector \mathbf{r} rotated by $\pi/2$ or 90° is *m.a.* of vector \mathbf{v}

Unit 1
Fig. 11.5

$$\text{velocity vector : } \mathbf{v} = \begin{pmatrix} v_x \\ v_y \end{pmatrix} = \begin{pmatrix} -a\omega \sin \omega t \\ b\omega \cos \omega t \end{pmatrix} = \frac{d\mathbf{r}}{dt} = \dot{\mathbf{r}} = \begin{pmatrix} a \cos \left(\phi + \frac{\pi}{2} \right) \\ b \sin \left(\phi + \frac{\pi}{2} \right) \end{pmatrix} \quad (\text{for } \omega = 1)$$

m.a. $\phi + \pi/2$ of vector \mathbf{v} rotated by another $\pi/2$ is *m.a.* of vector \mathbf{a}

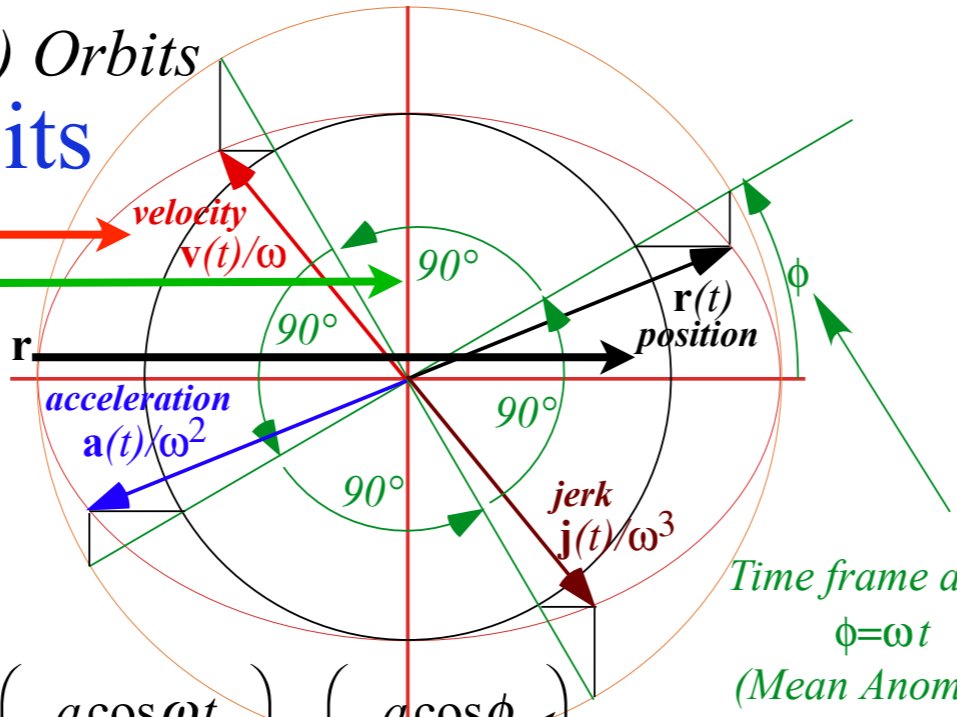
$$\text{acceleration or force vector : } \frac{\mathbf{F}}{m} = \mathbf{a} = \begin{pmatrix} a_x \\ a_y \end{pmatrix} = \begin{pmatrix} -a\omega^2 \cos \omega t \\ -b\omega^2 \sin \omega t \end{pmatrix} = \frac{d\mathbf{v}}{dt} = \dot{\mathbf{v}} = \ddot{\mathbf{r}} = \frac{d^2\mathbf{r}}{dt^2} = \begin{pmatrix} a \cos \left(\phi + \frac{2\pi}{2} \right) \\ b \sin \left(\phi + \frac{2\pi}{2} \right) \end{pmatrix}$$

$$\text{jerk or change of acceleration : } \mathbf{j} = \begin{pmatrix} j_x \\ j_y \end{pmatrix} = \begin{pmatrix} +a\omega^3 \sin \omega t \\ -b\omega^3 \cos \omega t \end{pmatrix} = \frac{d\mathbf{a}}{dt} = \dot{\mathbf{a}} = \ddot{\mathbf{v}} = \ddot{\mathbf{r}} = \frac{d^3\mathbf{r}}{dt^3} = \begin{pmatrix} a \cos \left(\phi + \frac{3\pi}{2} \right) \\ b \sin \left(\phi + \frac{3\pi}{2} \right) \end{pmatrix} \quad \dots \text{and so forth...}$$

Calculus of IHO orbits

(a) Orbits

To make velocity vector \mathbf{v} just rotate by $\pi/2$ or 90° the mean-anomaly ϕ of position vector \mathbf{r}



Time frame angle $\phi = \omega t$
(Mean Anomaly)

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mean-anomaly ϕ of position vector \mathbf{r} rotated by $\pi/2$ or 90° is *m.a.* of vector \mathbf{v}

Unit 1
Fig. 11.5

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m.a. $\phi + \pi/2$ of vector \mathbf{v} rotated by another $\pi/2$ is *m.a.* of vector \mathbf{a}

$$\text{acceleration or force vector : } \frac{\mathbf{F}}{m} = \mathbf{a} = \begin{pmatrix} a_x \\ a_y \end{pmatrix} = \begin{pmatrix} -a\omega^2 \cos \omega t \\ -b\omega^2 \sin \omega t \end{pmatrix} = \frac{d\mathbf{v}}{dt} = \dot{\mathbf{v}} = \ddot{\mathbf{r}} = \frac{d^2\mathbf{r}}{dt^2} = \begin{pmatrix} a \cos \left(\phi + \frac{2\pi}{2} \right) \\ b \sin \left(\phi + \frac{2\pi}{2} \right) \end{pmatrix}$$

...and so forth...

$$\text{jerk or change of acceleration : } \mathbf{j} = \begin{pmatrix} j_x \\ j_y \end{pmatrix} = \begin{pmatrix} +a\omega^3 \sin \omega t \\ -b\omega^3 \cos \omega t \end{pmatrix} = \frac{d\mathbf{a}}{dt} = \dot{\mathbf{a}} = \ddot{\mathbf{v}} = \ddot{\ddot{\mathbf{r}}} = \frac{d^3\mathbf{r}}{dt^3} = \begin{pmatrix} a \cos \left(\phi + \frac{3\pi}{2} \right) \\ b \sin \left(\phi + \frac{3\pi}{2} \right) \end{pmatrix}$$

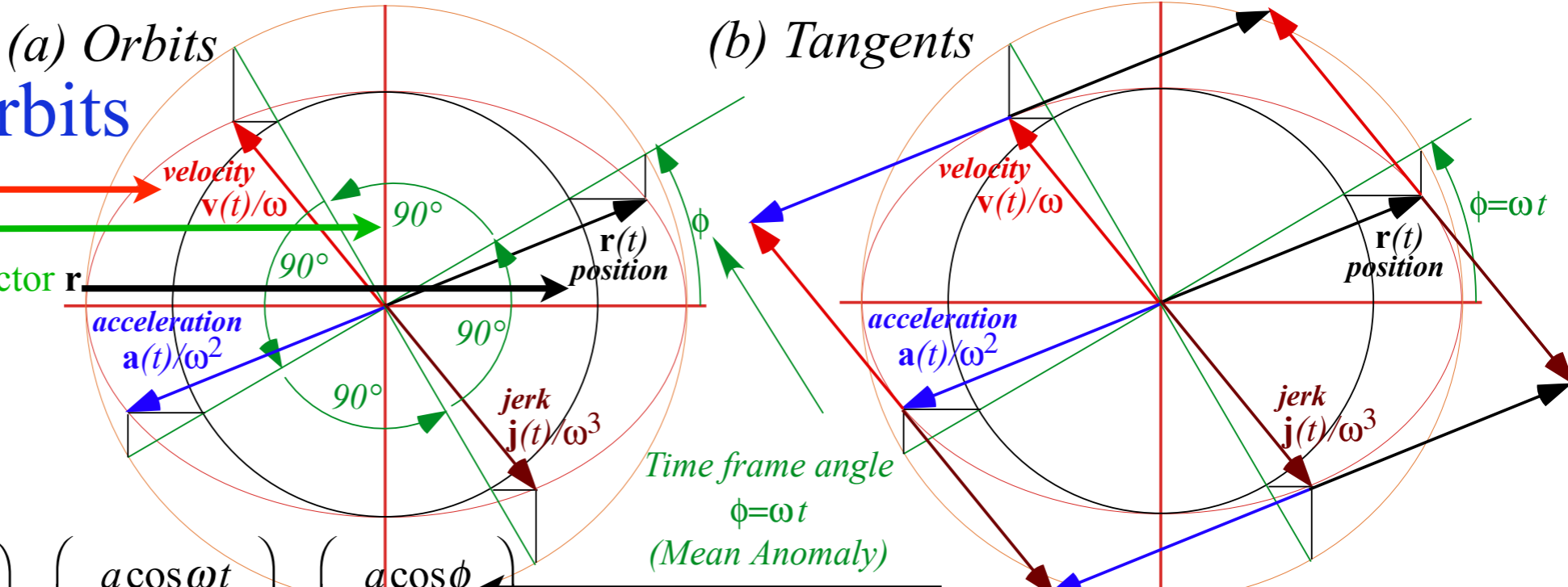
...and so on...
...But, now it repeats after 4 *t*-derivatives

$$\text{inauguration or change of jerk : } \mathbf{i} = \begin{pmatrix} i_x \\ i_y \end{pmatrix} = \begin{pmatrix} +a\omega^4 \cos \omega t \\ +b\omega^4 \sin \omega t \end{pmatrix} = \frac{d\mathbf{j}}{dt} = \dot{\mathbf{j}} = \ddot{\mathbf{a}} = \ddot{\ddot{\mathbf{v}}} = \ddot{\ddot{\ddot{\mathbf{r}}}} = \frac{d^4\mathbf{r}}{dt^4} = \begin{pmatrix} a \cos \left(\phi + \frac{4\pi}{2} \right) \\ b \sin \left(\phi + \frac{4\pi}{2} \right) \end{pmatrix}$$

Calculus of IHO orbits

To make velocity vector \mathbf{v} just rotate by $\pi/2$ or 90° the mean-anomaly ϕ of position vector \mathbf{r}

[Link \$\Rightarrow\$ BoxIt simulation of IHO orbits](#)



$$\text{radius vector : } \mathbf{r} = \begin{pmatrix} r_x \\ r_y \end{pmatrix} = \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} a \cos \omega t \\ b \sin \omega t \end{pmatrix} = \begin{pmatrix} a \cos \phi \\ b \sin \phi \end{pmatrix}$$

mean-anomaly ϕ of position vector \mathbf{r} rotated by $\pi/2$ or 90° is *m.a.* of vector \mathbf{v}

Unit 1
Fig. 11.5

[Link \$\rightarrow\$ IHO Exegesis Plot](#)

[Link \$\rightarrow\$ IHO orbital time rates of change](#)

$$\text{velocity vector : } \mathbf{v} = \begin{pmatrix} v_x \\ v_y \end{pmatrix} = \begin{pmatrix} -a\omega \sin \omega t \\ b\omega \cos \omega t \end{pmatrix} = \frac{d\mathbf{r}}{dt} = \dot{\mathbf{r}} = \begin{pmatrix} a \cos \left(\phi + \frac{\pi}{2} \right) \\ b \sin \left(\phi + \frac{\pi}{2} \right) \end{pmatrix} \quad (\text{for } \omega = 1)$$

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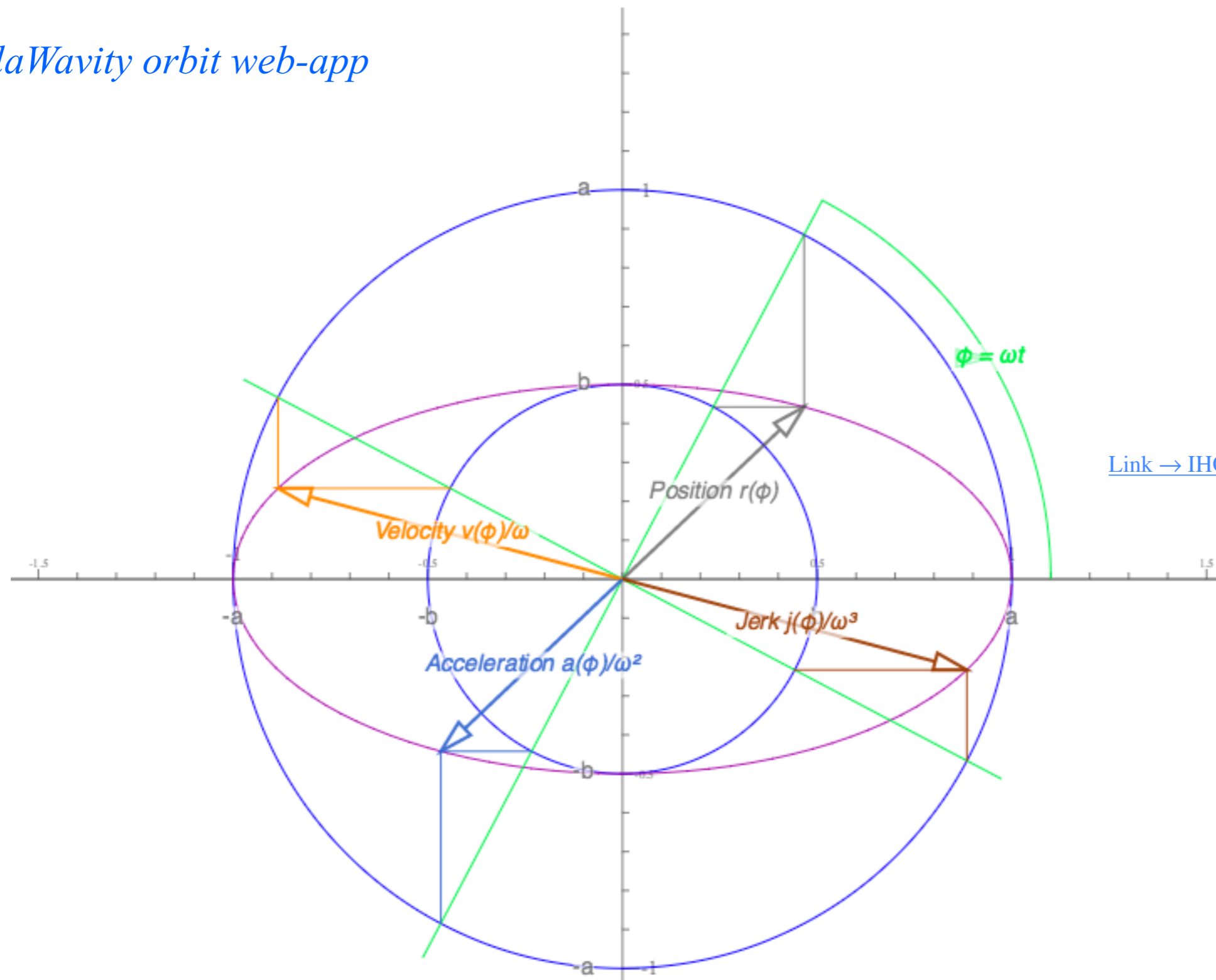
...and so forth...

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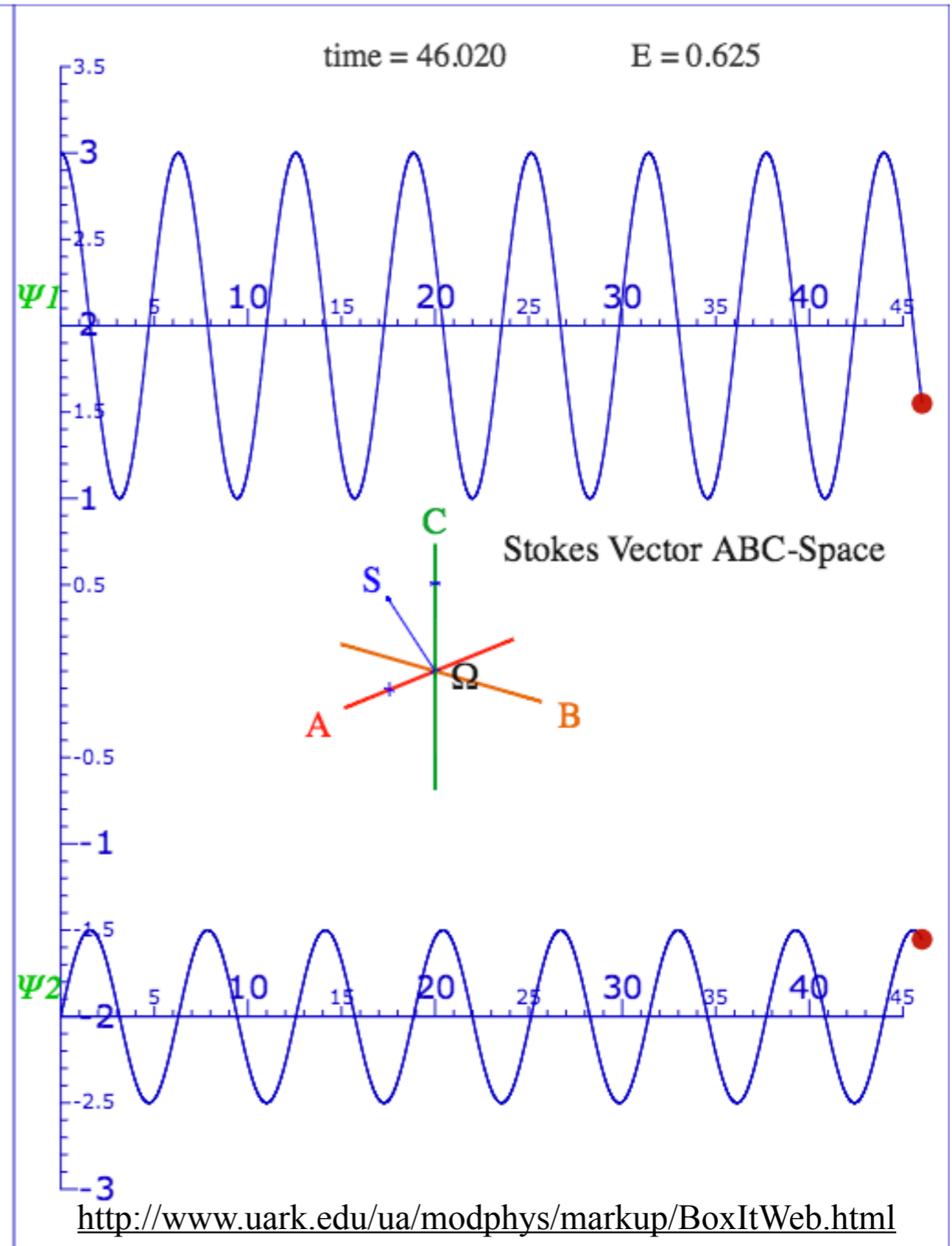
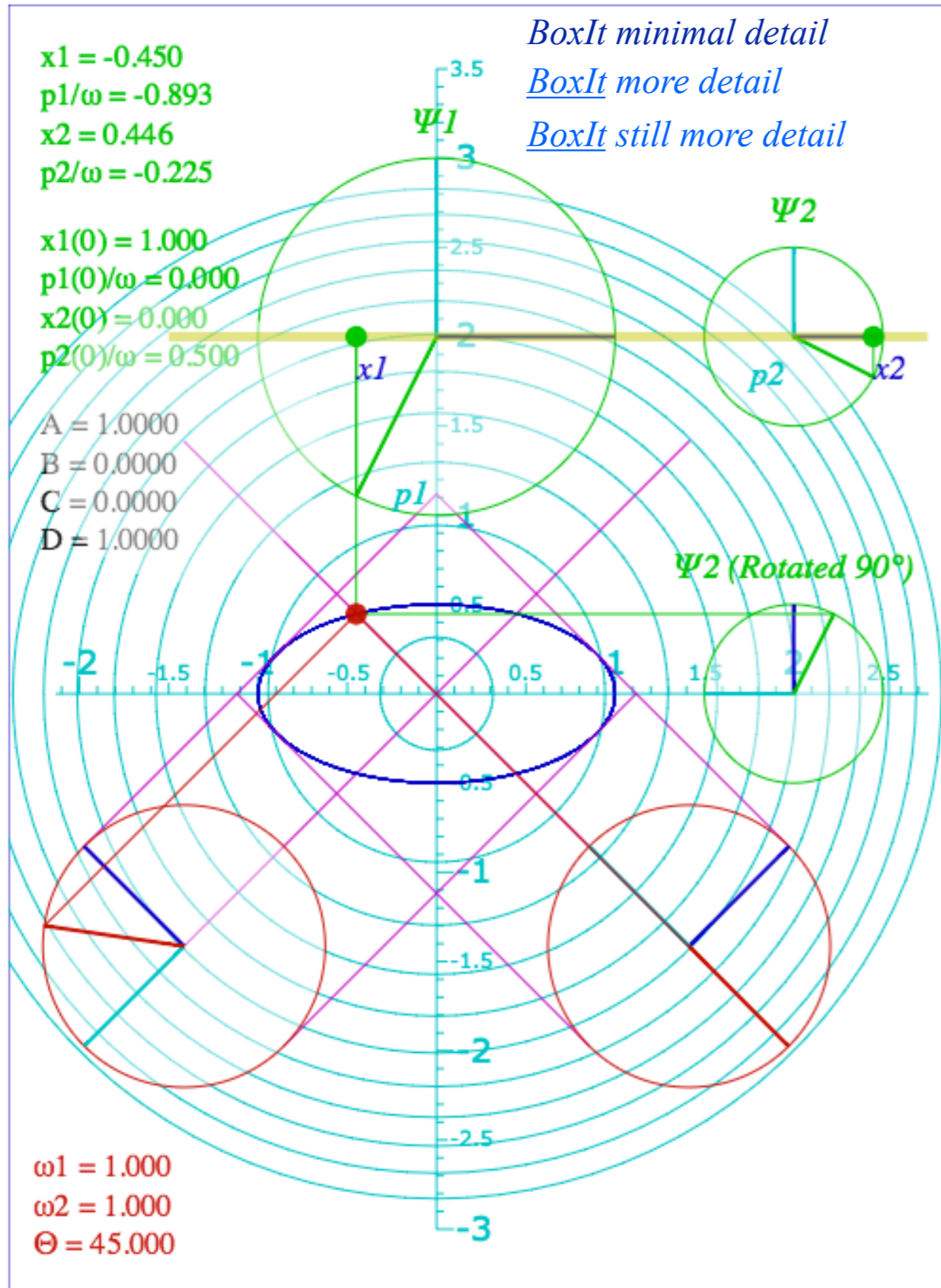
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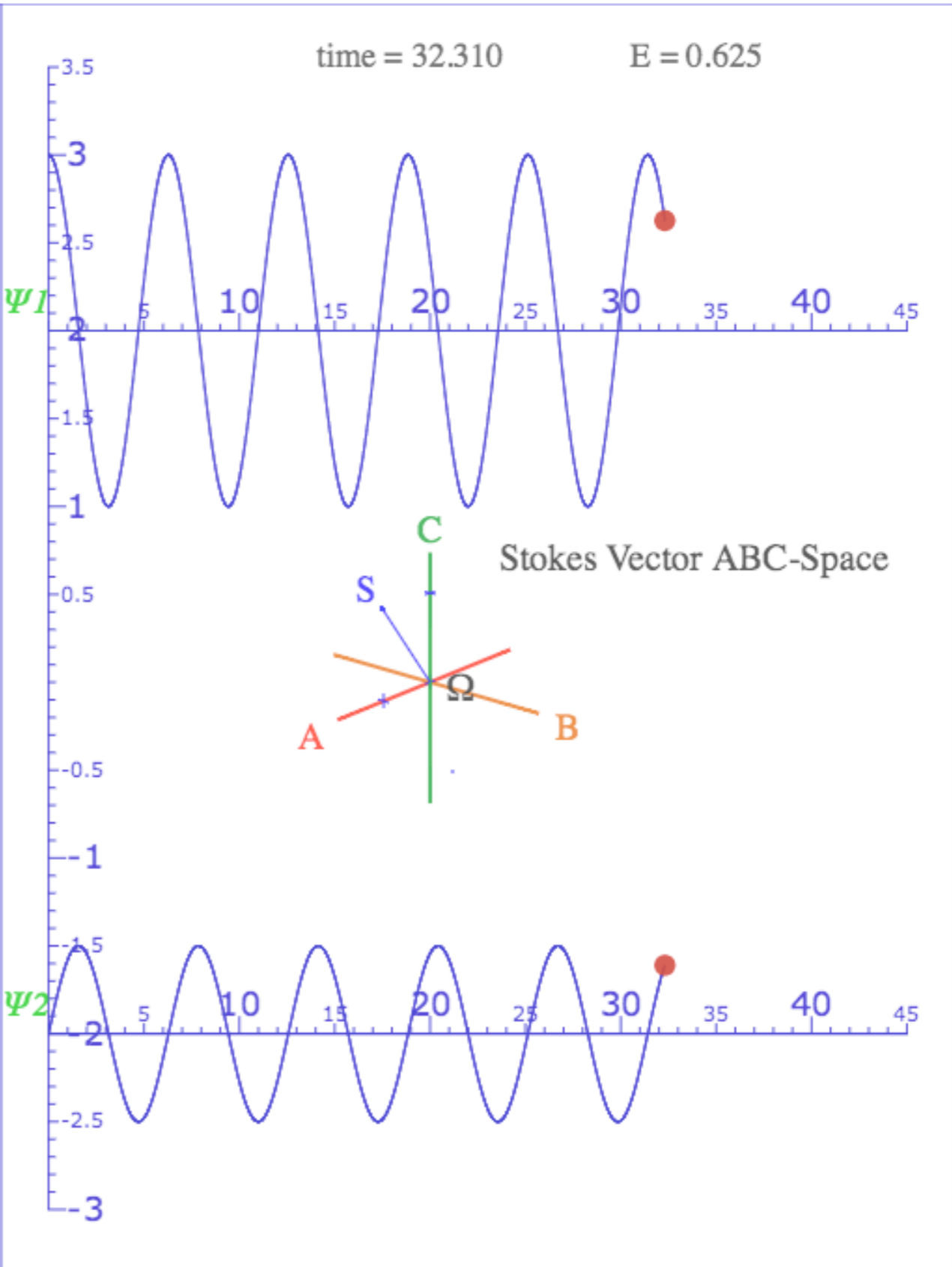
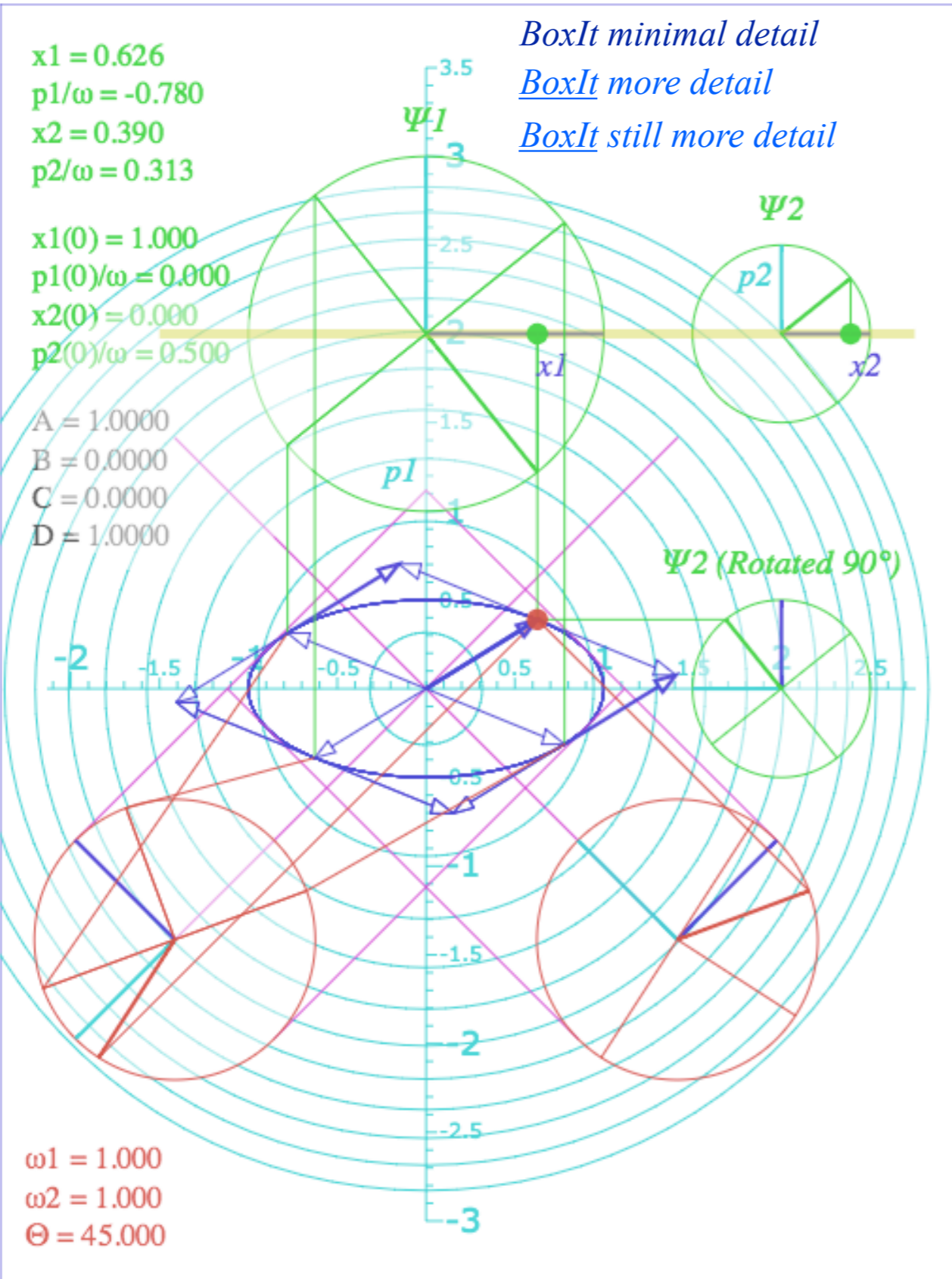
RelaWavity orbit web-app



Geometry of Kepler anomalies for vectors $[\mathbf{r}(\phi), \mathbf{v}(\phi), \mathbf{a}(\phi), \mathbf{j}(\phi),]$ in coordinate (x,y) space rendered by animation web-apps BoxIt and RelaWavity.



Geometry of Kepler anomalies for vectors $[\mathbf{r}(\phi)]$ in coordinate (x,y) space and 2-particle (x_1,x_2) space rendered by animation web-apps BoxIt.

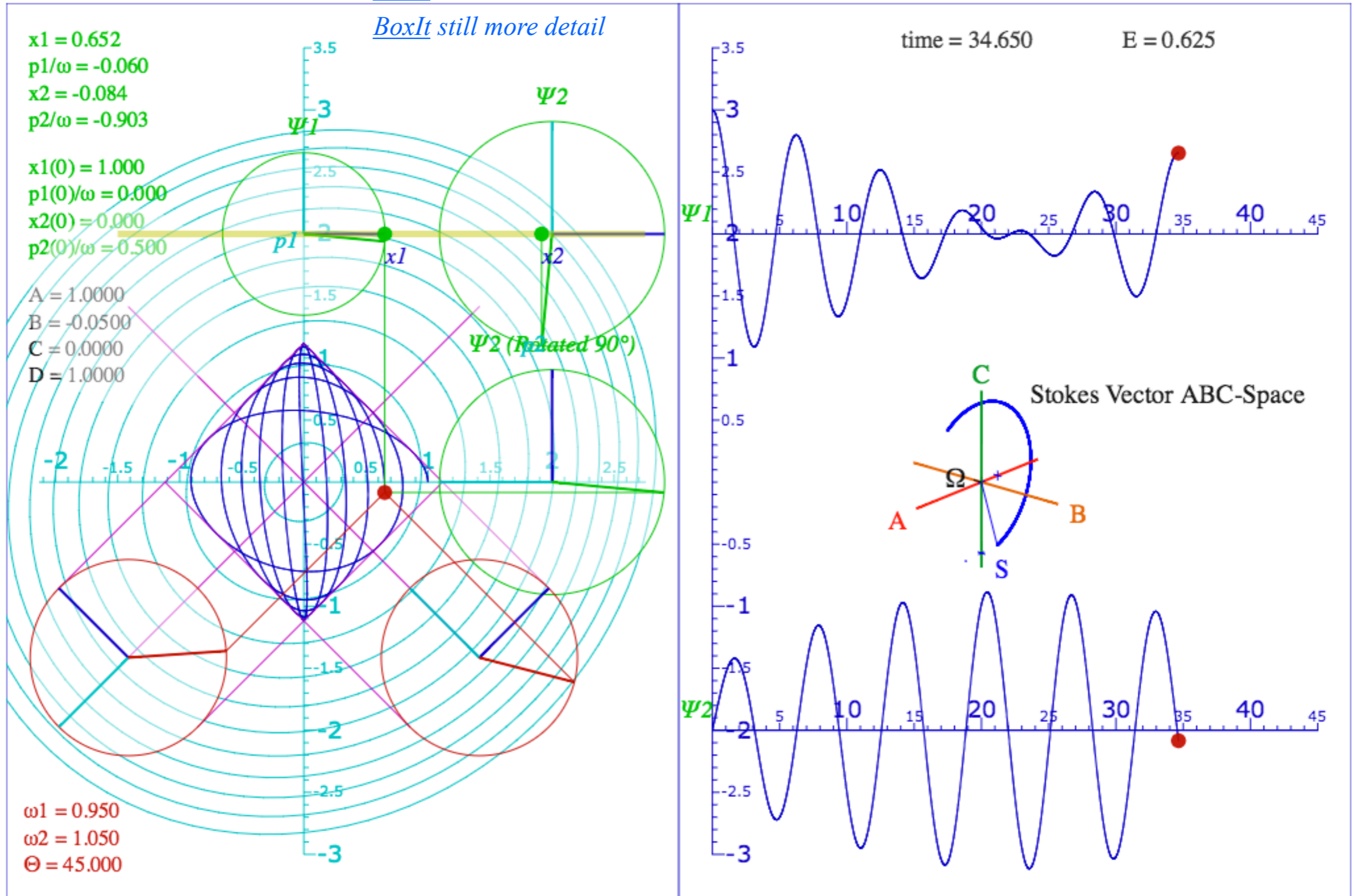


Geometry of Kepler anomalies for vectors $[\mathbf{r}(\phi), \mathbf{v}(\phi), \mathbf{a}(\phi), \mathbf{j}(\phi),]$ in coordinate (x,y) space and 2-particle (x_1,x_2) space rendered by animation web-apps BoxIt.

[BoxIt minimal detail](#)

[BoxIt more detail](#)

[BoxIt still more detail](#)



*Geometry of vectors $[\mathbf{r}(\phi), \mathbf{p}(\phi)]$ and quantum spin S -space
and 2-particle (x_1, x_2) space rendered by animation web-apps BoxIt.*

[BoxIt Web Simulation - B-Type Motion](#)

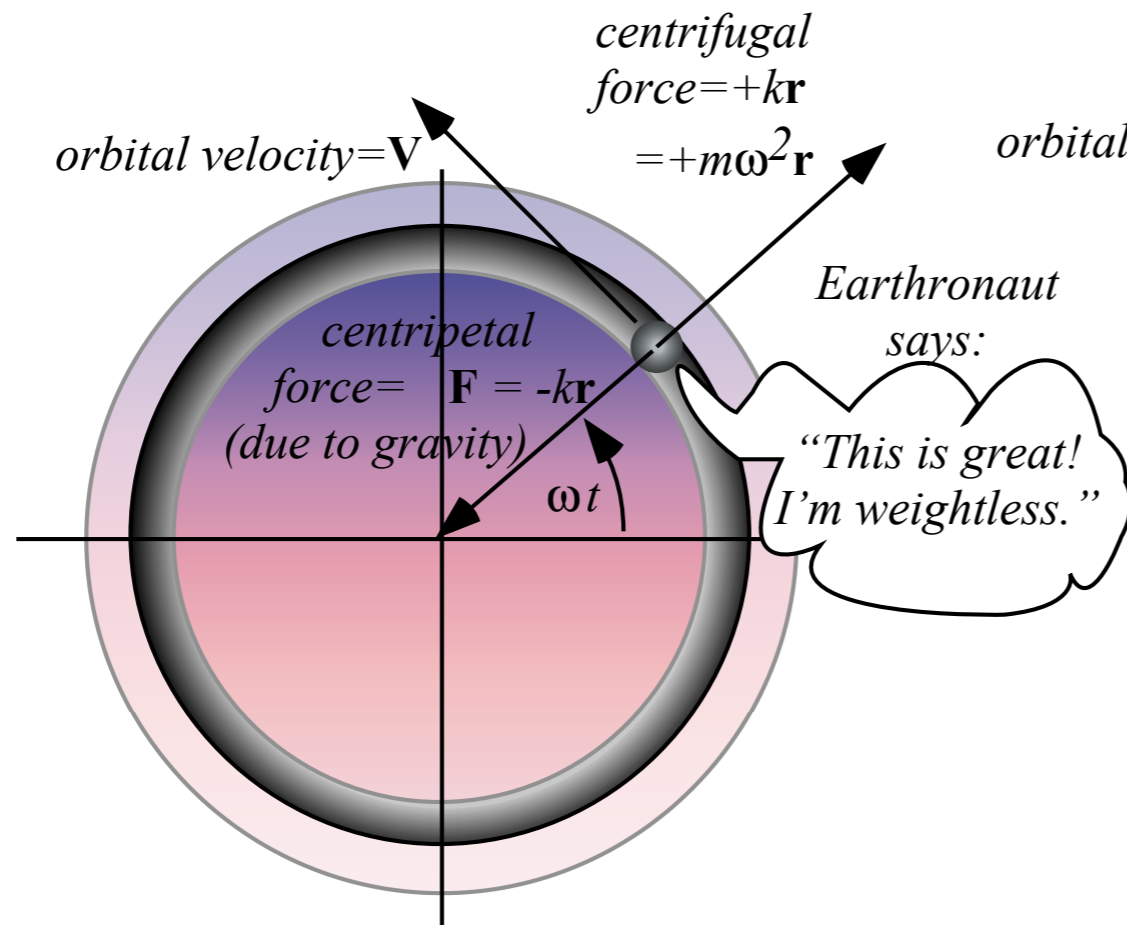
*Constructing 2D IHO orbits using **Kepler anomaly plots***

Mean-anomaly and eccentric-anomaly geometry

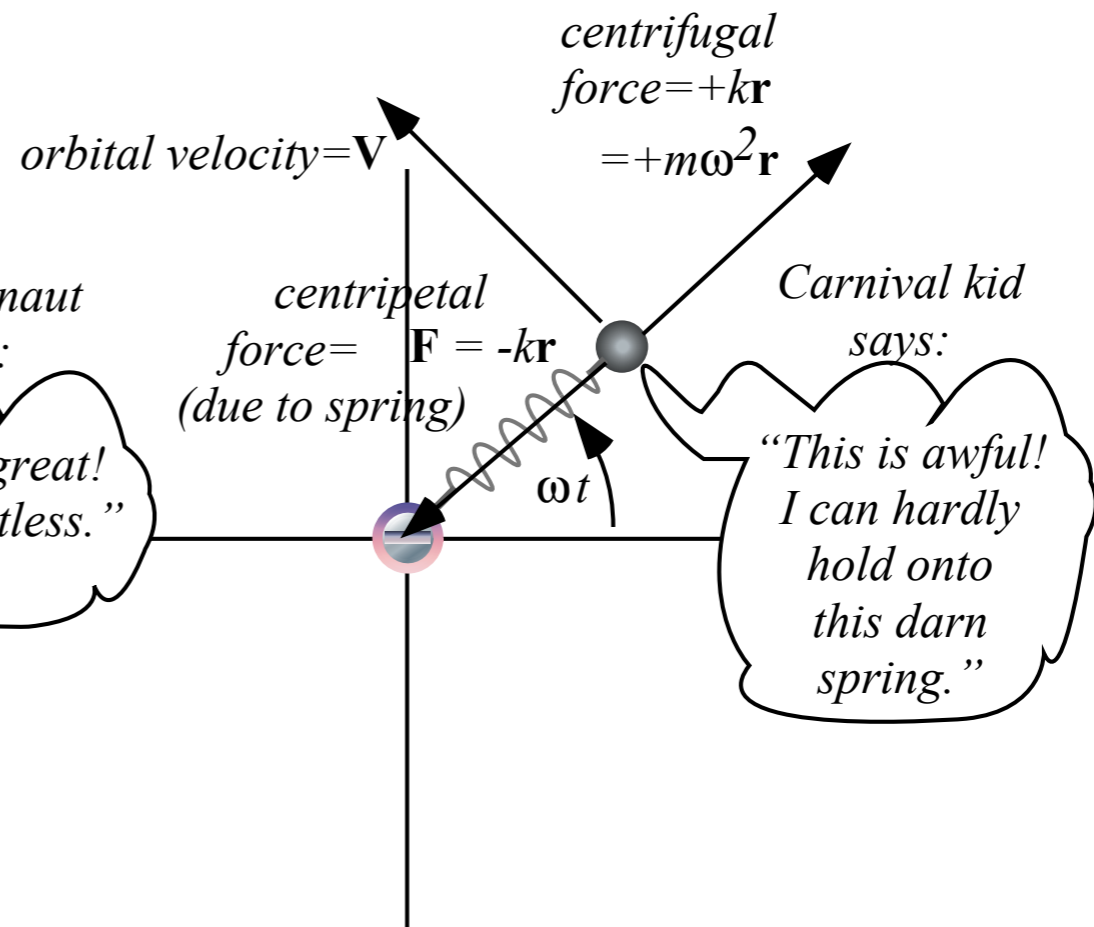
Calculus and vector geometry of IHO orbits

 *A confusing introduction to Coriolis-centrifugal force geometry* (Derived better in Ch. 12)

(a) "Earthronaut" orbiting tunnel inside Earth

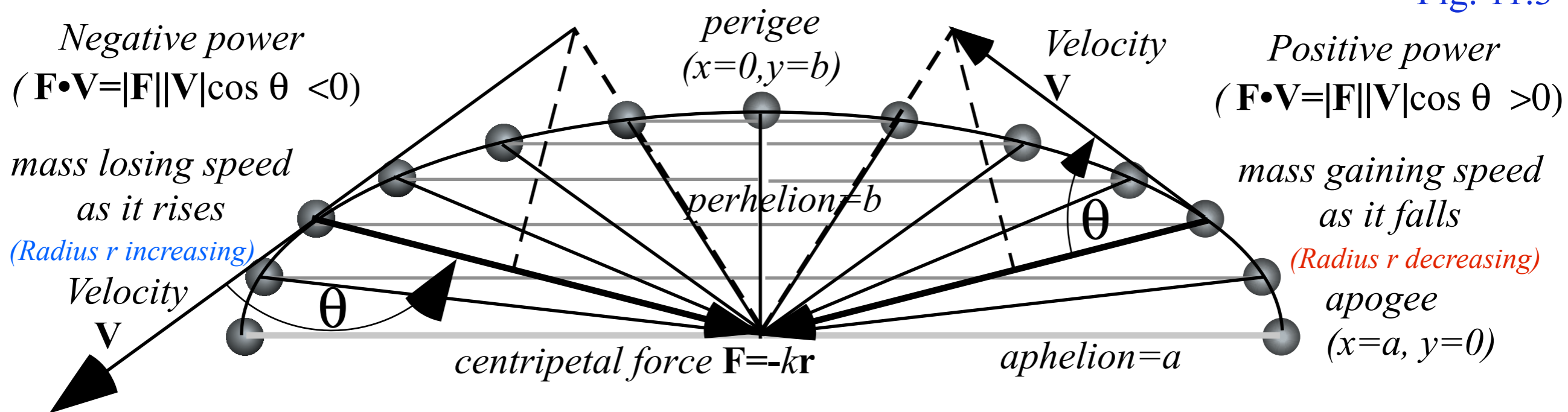


(b) "Carnival kid" orbiting in space attached to a spring



Unit 1
 Fig. 11.2

Unit 1
 Fig. 11.3

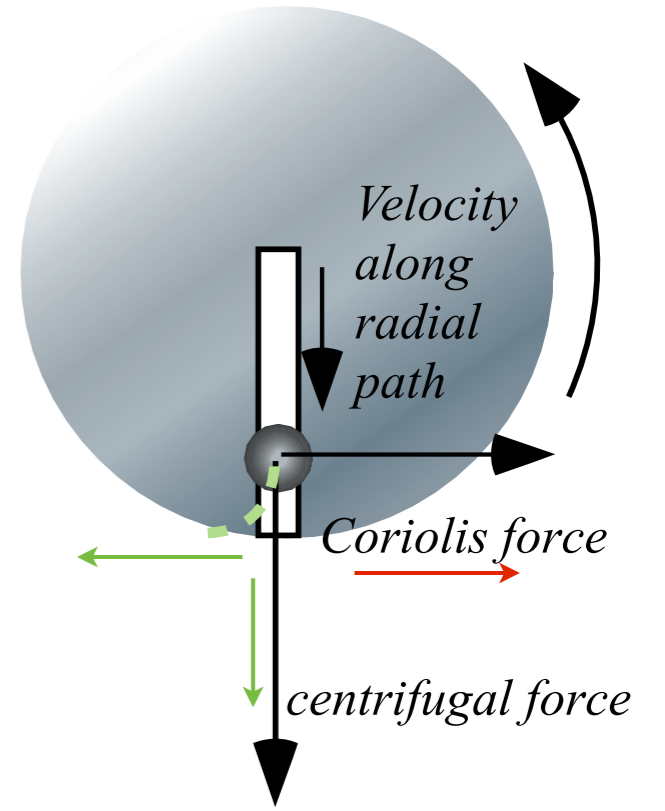
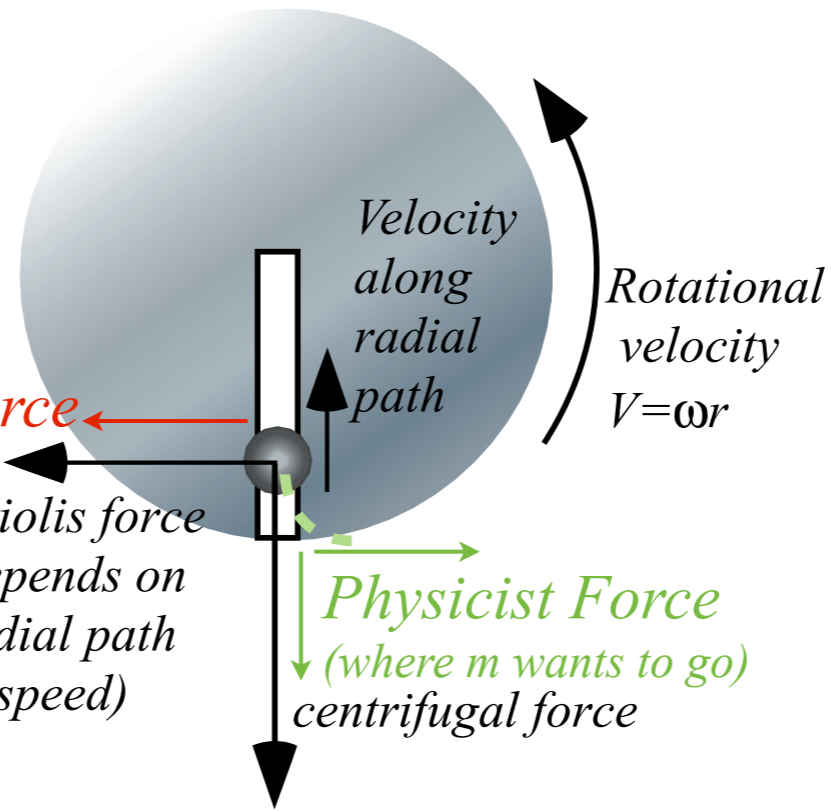


(a) Centrifugal and Coriolis Forces on Merry-Go-Round

Mathematician Force
(to hold m back)
Constraint force
keeps m in radial slot

Coriolis force
(depends on
radial path
speed)

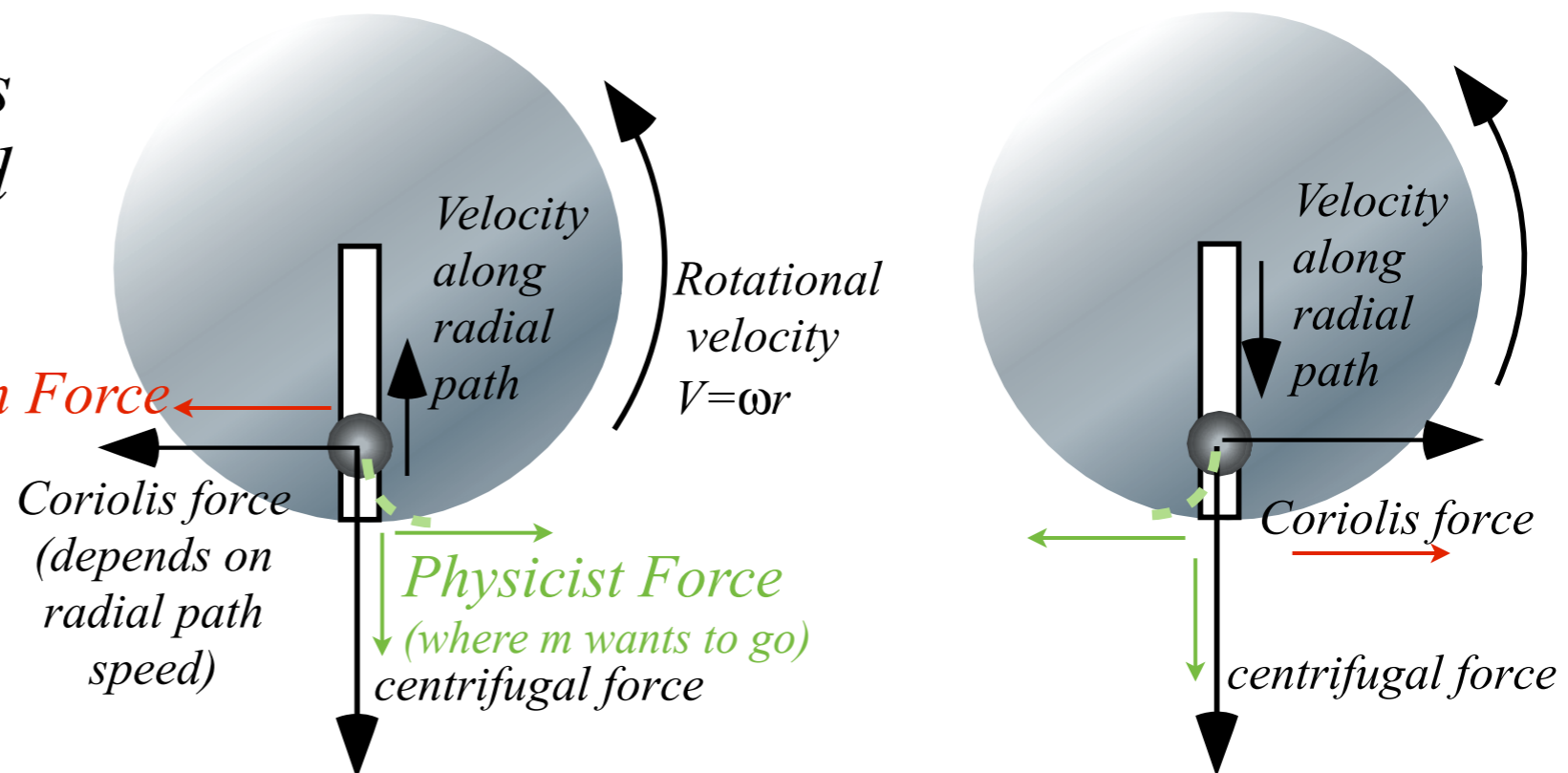
Physicist Force
(where m wants to go)
centrifugal force



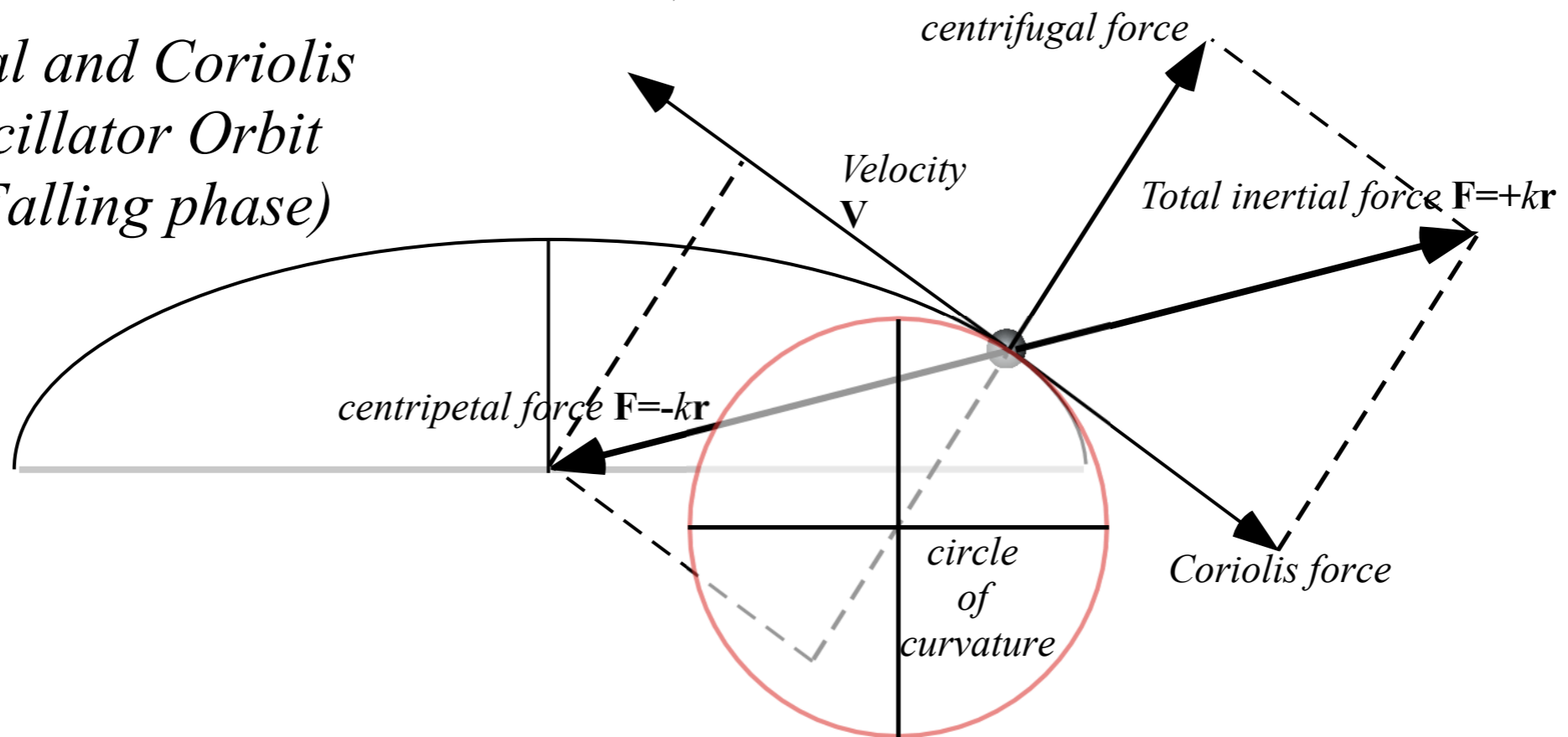
(a) Centrifugal and Coriolis Forces on Merry-Go-Round

Mathematician Force
(to hold m back)

Constraint force
keeps m in radial slot



(b) Centrifugal and Coriolis Forces on Oscillator Orbit (Falling phase)

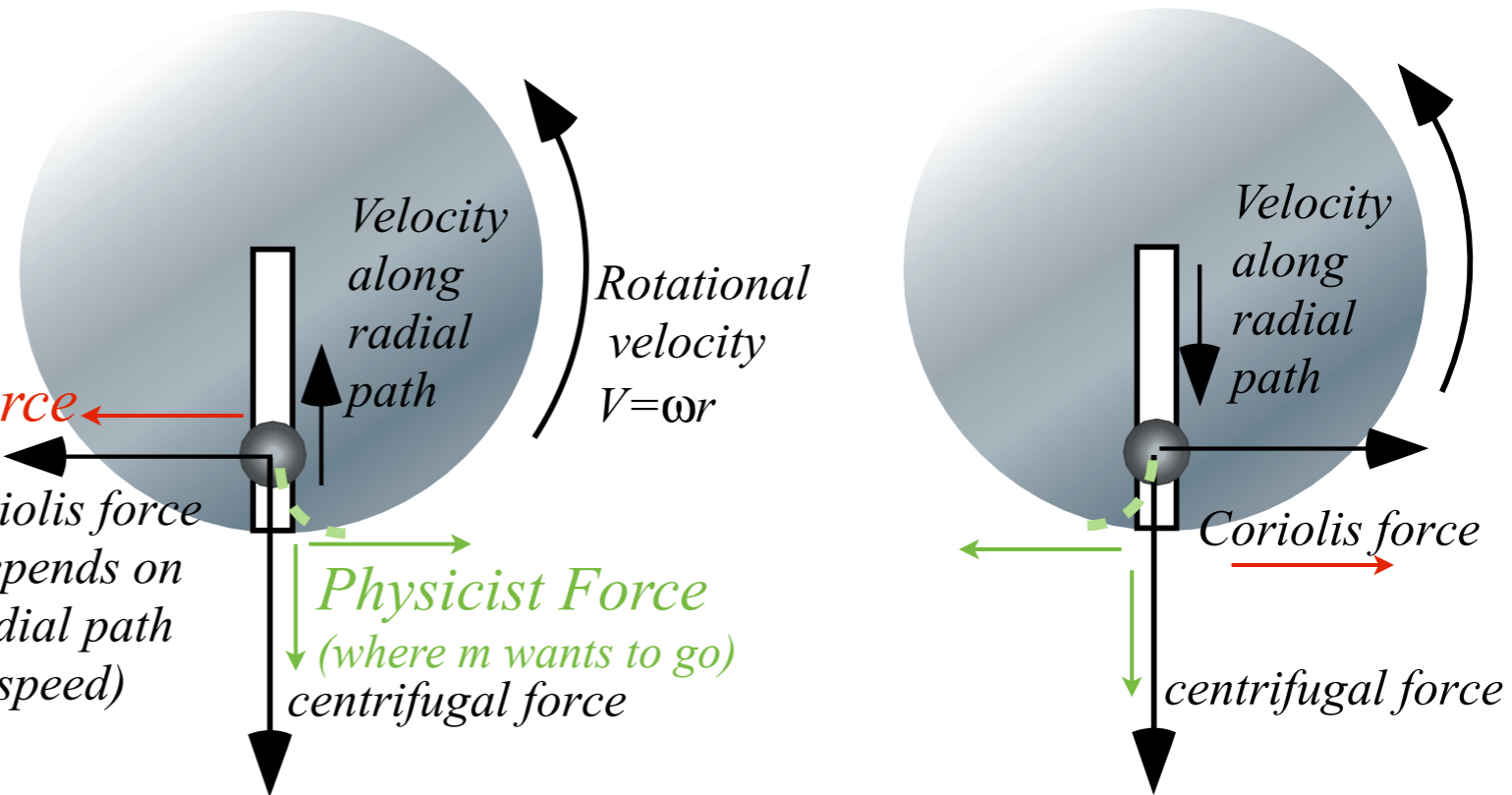


(a) Centrifugal and Coriolis Forces on Merry-Go-Round

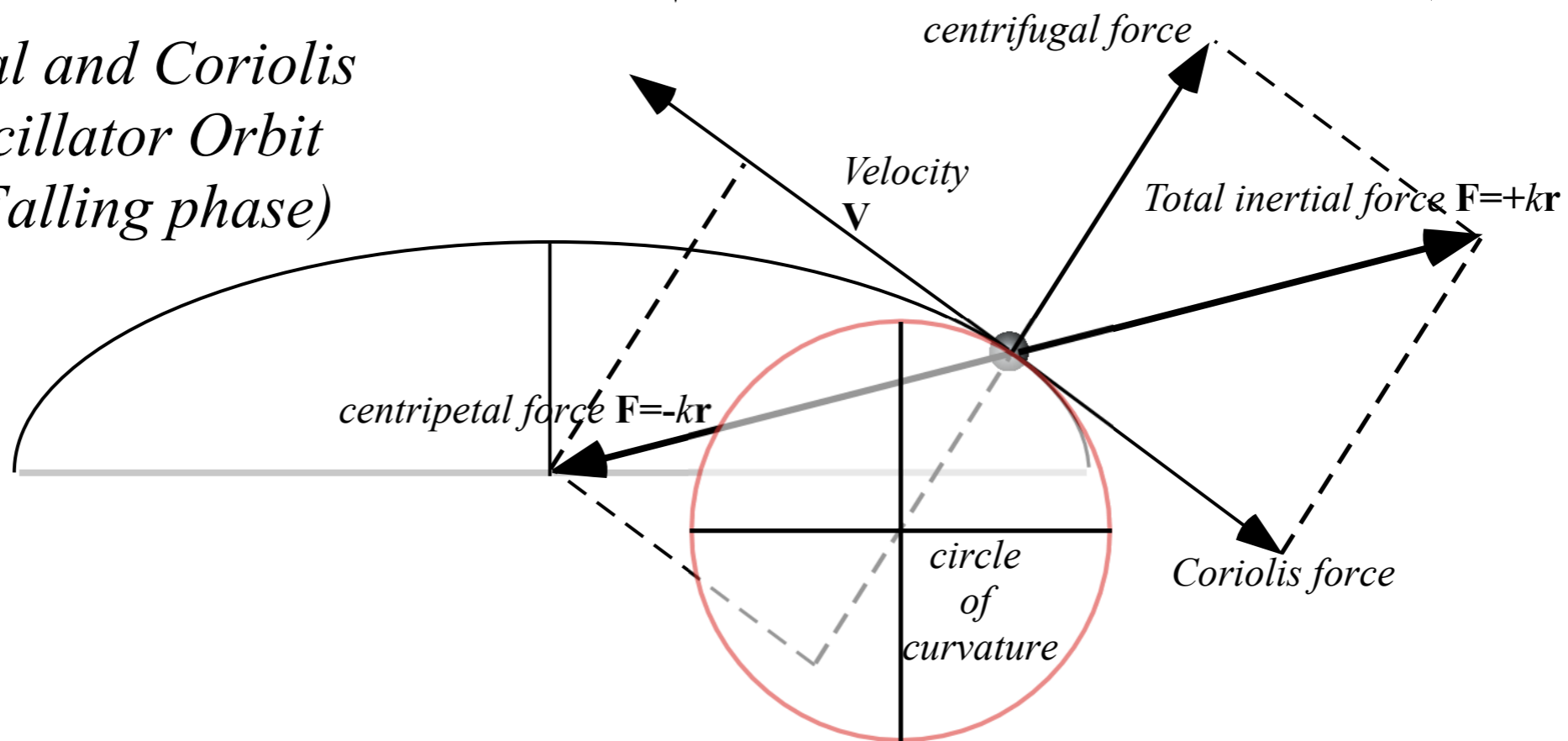
Mathematician Force
(to hold m back)
Constraint force
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Coriolis force
(depends on radial path speed)

Physicist Force
(where m wants to go)
centrifugal force



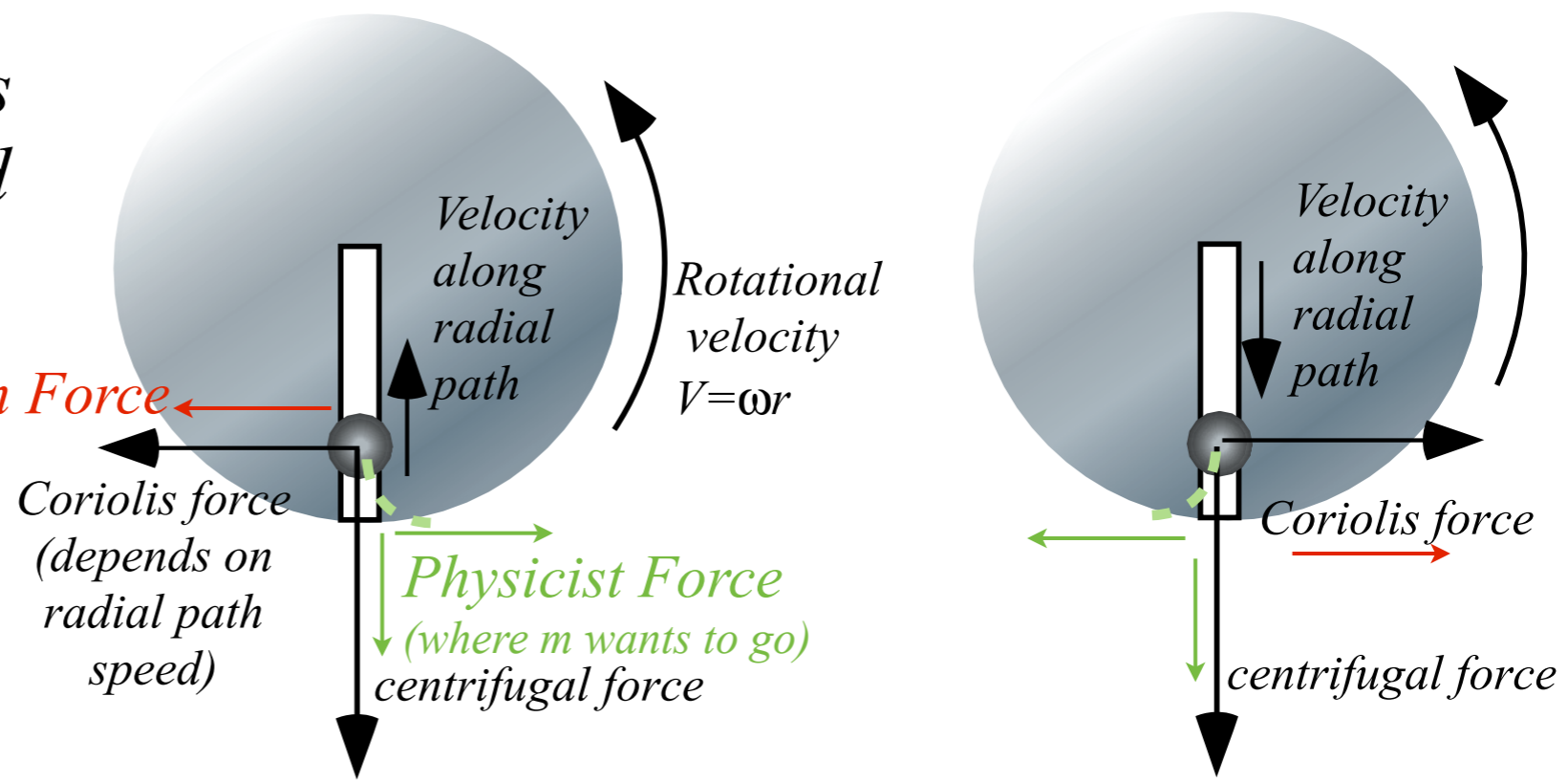
(b) Centrifugal and Coriolis Forces on Oscillator Orbit (Falling phase)



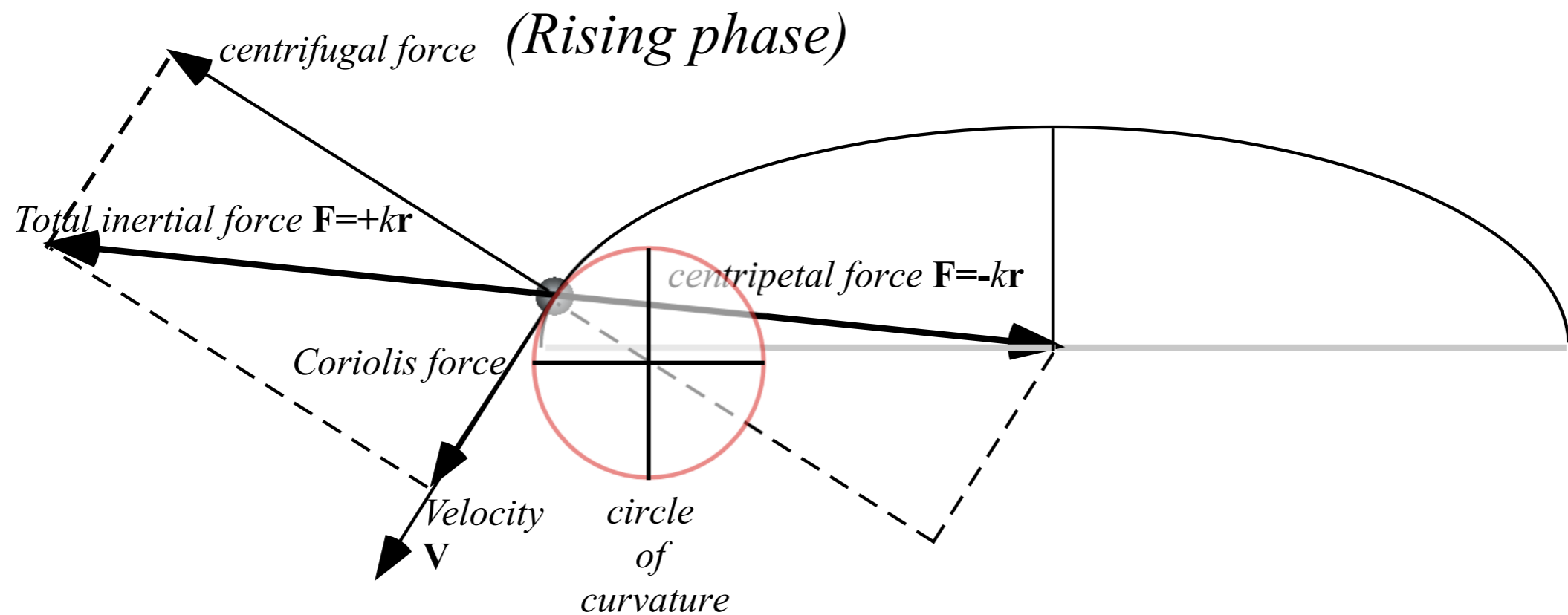
(a) Centrifugal and Coriolis Forces on Merry-Go-Round

Mathematician Force
(to hold m back)

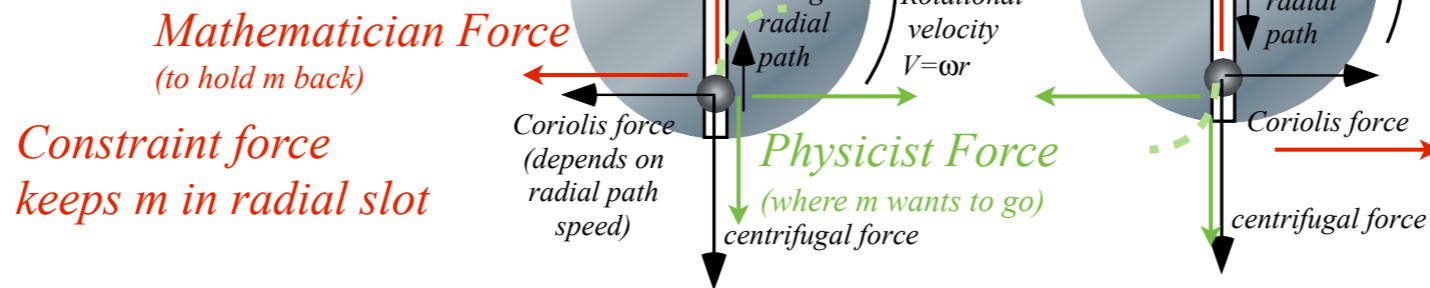
Constraint force
keeps m in radial slot



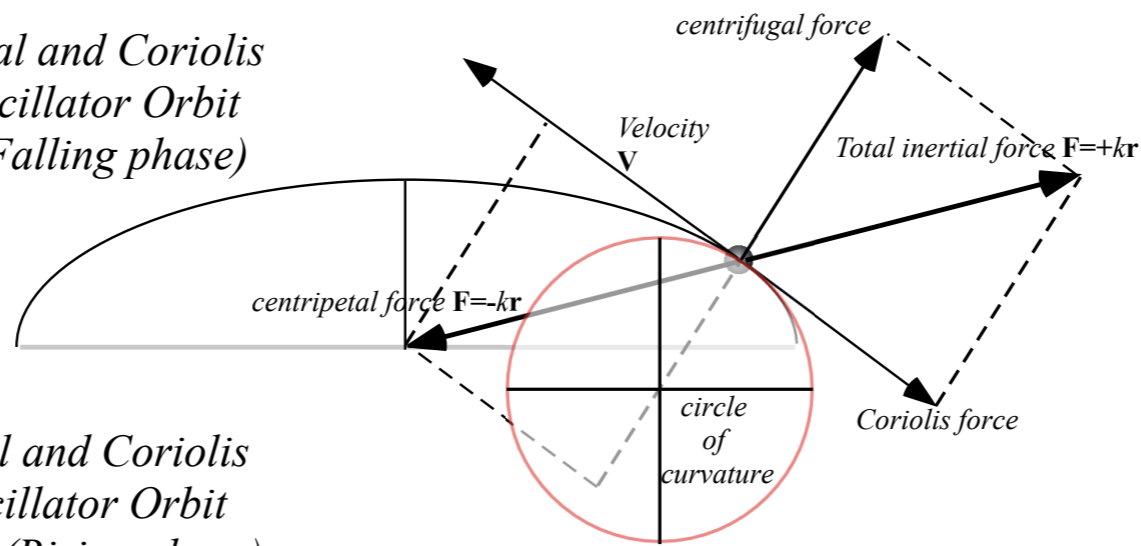
(c) Centrifugal and Coriolis Forces on Oscillator Orbit



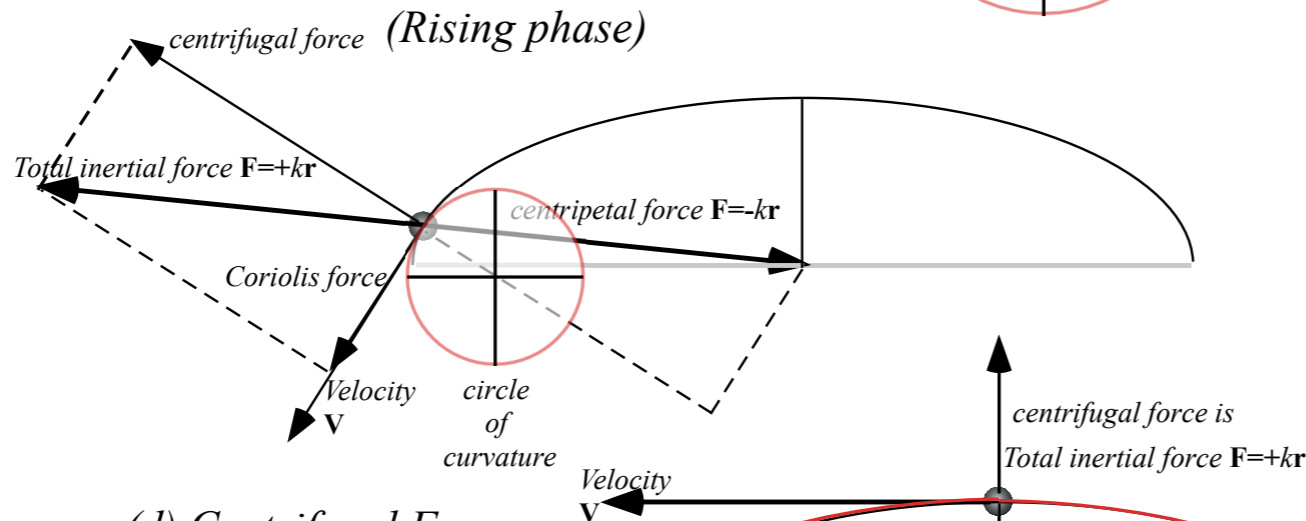
(a) Centrifugal and Coriolis Forces on Merry-Go-Round



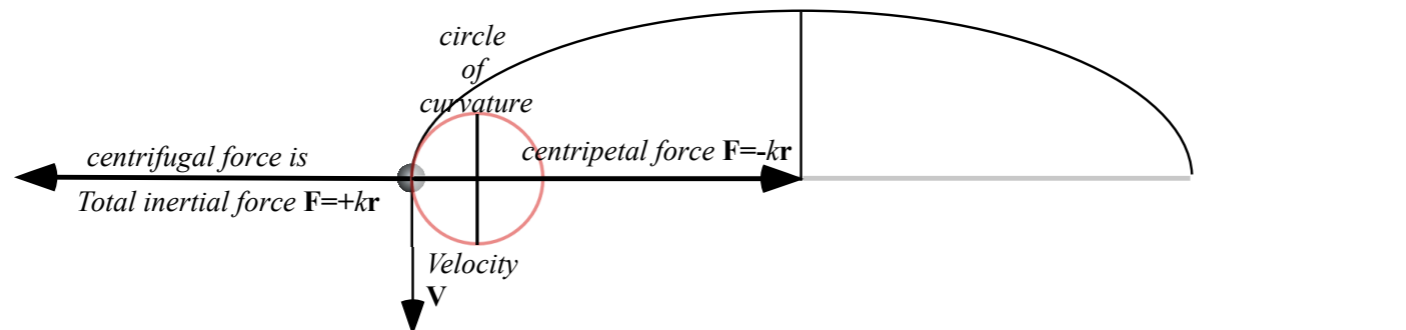
(b) Centrifugal and Coriolis Forces on Oscillator Orbit (Falling phase)



(c) Centrifugal and Coriolis Forces on Oscillator Orbit (Rising phase)




(d) Centrifugal Force on Oscillator Orbit (apogee and perigee)



Unit 1
Fig. 11.4
a-d

*Quite confusing?
Discussion of Coriolis
forces will be done more elegantly
and made more physically intuitive
in Ch. 12 of Unit 1 and in Unit 6.*

Some Kepler's "laws" for all central (isotropic) force $F(r)$ fields

-  *Angular momentum invariance of IHO: $F(r) = -k \cdot r$ with $U(r) = k \cdot r^2 / 2$ (Derived here)*
- Angular momentum invariance of Coulomb: $F(r) = -GMm/r^2$ with $U(r) = -GMm \cdot /r$ (Derived in Unit 5)*
- Total energy $E = KE + PE$ invariance of IHO: $F(r) = -k \cdot r$ (Derived here)*
- Total energy $E = KE + PE$ invariance of Coulomb: $F(r) = -GMm/r^2$ (Derived in Unit 5)*

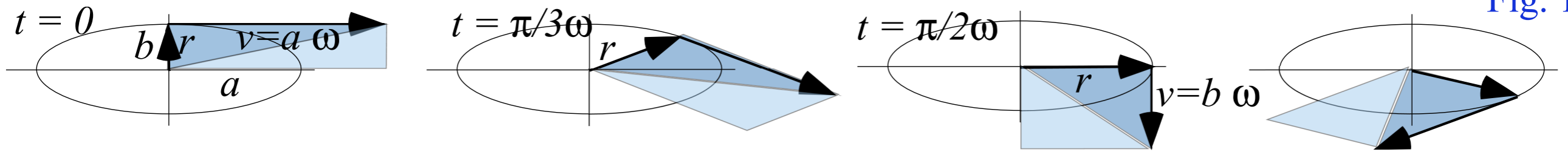
Some Kepler's "laws" for central (isotropic) force $F(r)$

...and certainly apply to the IHO: $F(r) = -k \cdot r$ with $U(r) = k \cdot r^2/2$

(Recall from Lecture 6: $k = Gm \frac{4\pi}{3} \rho_{\oplus}$)

Unit 1

Fig. 11.8



1. Area of triangle $\triangle_r^v = \mathbf{r} \times \mathbf{v}/2$ is constant

$$\mathbf{r} \times \mathbf{v} = r_x v_y - r_y v_x = a \cos \omega t \cdot (b \omega \cos \omega t) - b \sin \omega t \cdot (-a \omega \sin \omega t) = ab \cdot \omega (\cos^2 \omega t + \sin^2 \omega t) \quad \checkmark \text{ for IHO}$$

$$\begin{pmatrix} r_x \\ r_y \end{pmatrix} = \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} a \cos \omega t \\ b \sin \omega t \end{pmatrix} \quad \begin{pmatrix} v_x \\ v_y \end{pmatrix} = \begin{pmatrix} -a \omega \sin \omega t \\ b \omega \cos \omega t \end{pmatrix}$$

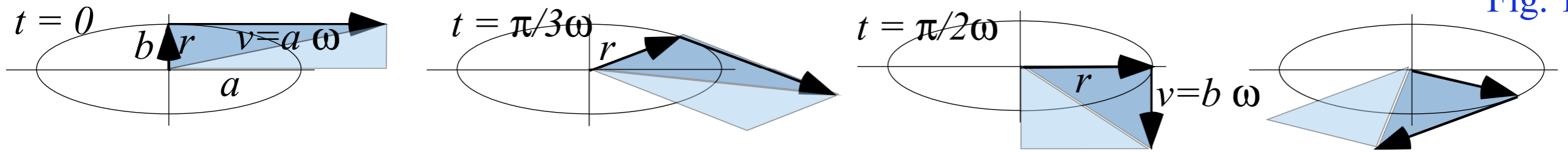
Some Kepler's "laws" that apply to any central (isotropic) force $F(r)$

...and certainly apply to the IHO: $F(r) = -k \cdot r$ with $U(r) = k \cdot r^2 / 2$

(Recall from Lecture 6: $k = Gm \frac{4\pi}{3} \rho_{\oplus}$)

Unit 1

Fig. 11.8



1. Area of triangle $\Delta_r^v = \mathbf{r} \times \mathbf{v} / 2$ is constant

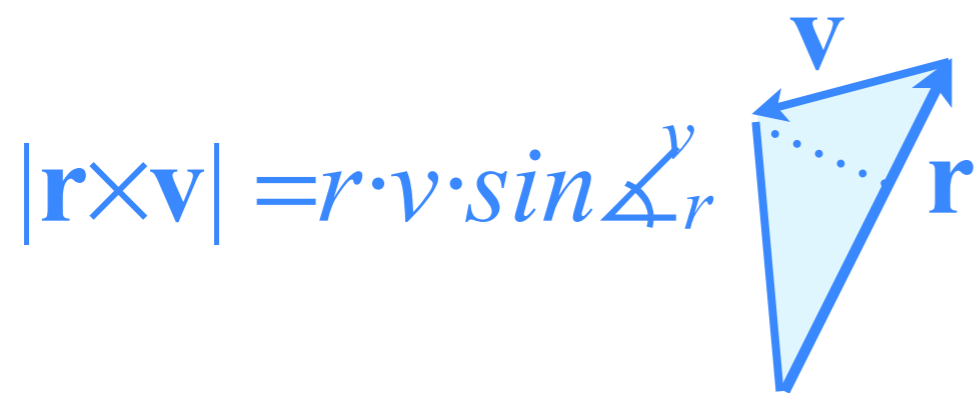
$$\mathbf{r} \times \mathbf{v} = r_x v_y - r_y v_x = a \cos \omega t \cdot (b \omega \cos \omega t) - a \sin \omega t \cdot (-b \omega \sin \omega t) = ab \cdot \omega$$

✓ for IHO

2. Angular momentum $\mathbf{L} = m \mathbf{r} \times \mathbf{v}$ is conserved

$$L = m |\mathbf{r} \times \mathbf{v}| = m (r_x v_y - r_y v_x) = m \cdot ab \cdot \omega$$

✓ for IHO



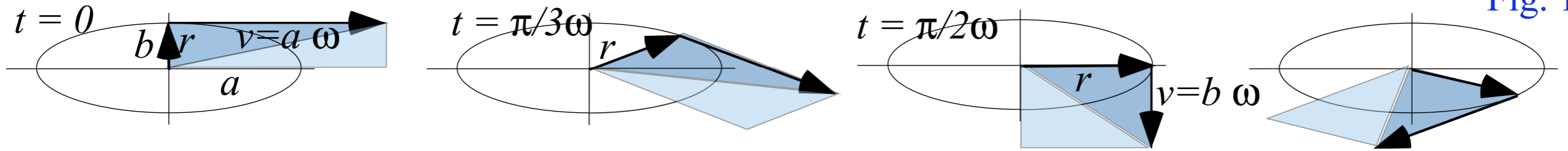
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Unit 1

Fig. 11.8



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✓ for IHO

3. Equal area is swept by radius vector in each equal time interval T

$$A_T = \int_0^T \frac{\mathbf{r} \times d\mathbf{r}}{2} = \int_0^T \frac{\mathbf{r} \times \frac{d\mathbf{r}}{dt}}{2} dt = \int_0^T \frac{\mathbf{r} \times \mathbf{v}}{2} dt = \frac{L}{2m} \int_0^T dt = \frac{L}{2m} T$$

✓ for IHO

$$|\mathbf{r} \times d\mathbf{r}| = r \cdot dr \cdot \sin \angle_r$$

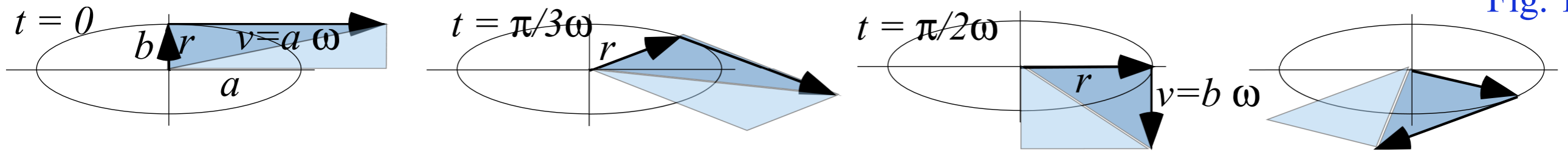
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Unit 1

Fig. 11.8



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$$L = m \mathbf{r} \times \mathbf{v} = m (r_x v_y - r_y v_x) = m \cdot ab \cdot \omega = m \cdot ab \cdot \frac{2\pi}{\tau}$$

✓ for IHO

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✓ for IHO

In one period: $\tau = \frac{1}{\nu} = \frac{2\pi}{\omega} = \frac{2mA_\tau}{L}$ the area is: $A_\tau = \frac{L\tau}{2m}$ ($= ab \cdot \pi$ for ellipse orbit)

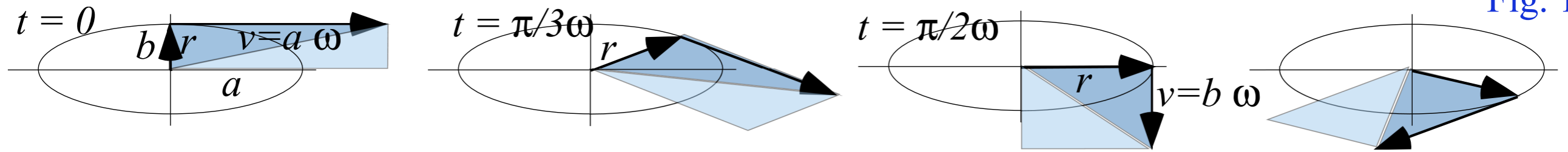
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(Recall from Lecture 6: $k = Gm \frac{4\pi}{3} \rho_{\oplus}$)

Unit 1

Fig. 11.8



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In one period: $\tau = \frac{1}{\nu} = \frac{2\pi}{\omega} = \frac{2mA_\tau}{L}$ the area is: $A_\tau = \frac{L\tau}{2m}$ ($= ab \cdot \pi$ for ellipse orbit)

(Recall from Lecture 6: $\omega = \sqrt{k/m} = \sqrt{G\rho_{\oplus} 4\pi/3}$)

(G IHO formulas from Lect. 6 p.70-79)

Some Kepler's "laws" for all central (isotropic) force $F(r)$ fields

Angular momentum invariance of IHO: $F(r) = -k \cdot r$ with $U(r) = k \cdot r^2 / 2$ (Derived here)

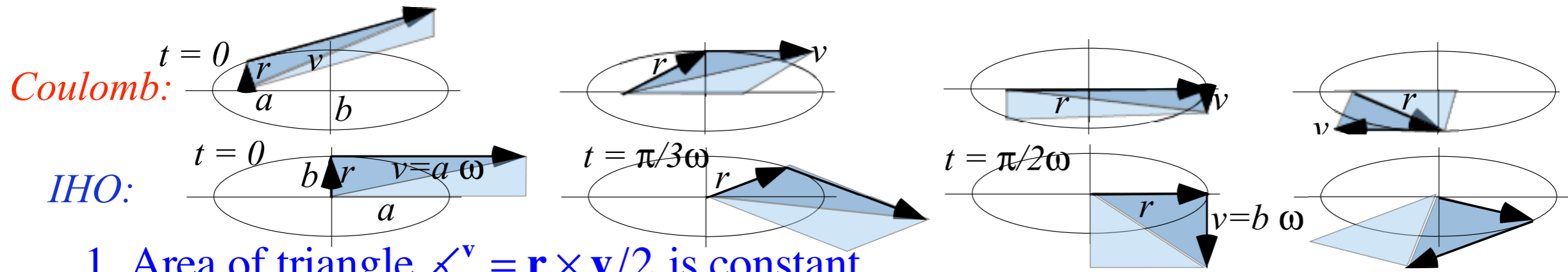
 *Angular momentum invariance of Coulomb: $F(r) = -GMm/r^2$ with $U(r) = -GMm \cdot /r$ (Derived in Unit 5)*

Total energy $E = KE + PE$ invariance of IHO: $F(r) = -k \cdot r$ (Derived here)

Total energy $E = KE + PE$ invariance of Coulomb: $F(r) = -GMm/r^2$ (Derived in Unit 5)

Some Kepler's "laws" that apply to any central (isotropic) force $F(r)$

Apply to IHO: $F(r) = -k \cdot r$ with $U(r) = k \cdot r^2 / 2$ and Coulomb: $F(r) = -GMm/r^2$ with $U(r) = -GMm \cdot / r$



1. Area of triangle $\triangle_r^v = \mathbf{r} \times \mathbf{v} / 2$ is constant

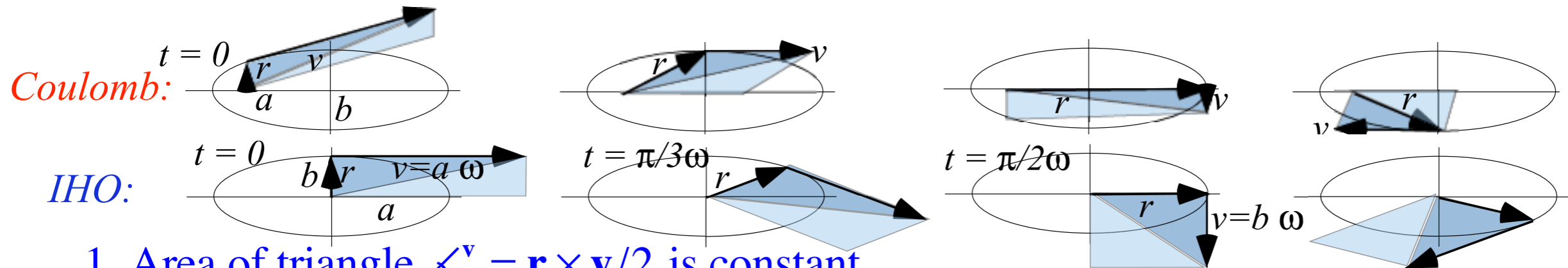
$$\mathbf{r} \times \mathbf{v} = r_x v_y - r_y v_x = \begin{cases} ab \cdot \sqrt{G\rho_{\oplus} 4\pi / 3} & \text{for IHO} \\ a^{-1/2} b \sqrt{GM_{\oplus}} & \text{for Coul.} \end{cases}$$

✓ for IHO
✓ for Coul.

(Derived in Unit 5)

Some Kepler's "laws" that apply to any central (isotropic) force $F(r)$

Apply to IHO: $F(r) = -k \cdot r$ with $U(r) = k \cdot r^2 / 2$ and Coulomb: $F(r) = -GMm/r^2$ with $U(r) = -GMm \cdot / r$



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$$\mathbf{r} \times \mathbf{v} = r_x v_y - r_y v_x = \begin{cases} ab \cdot \sqrt{G\rho_{\oplus} 4\pi / 3} & \text{for IHO} \\ a^{-1/2} b \sqrt{GM_{\oplus}} & \text{for Coul.} \end{cases}$$

✓ for IHO
(Derived in Unit 5) ✓ for Coul.

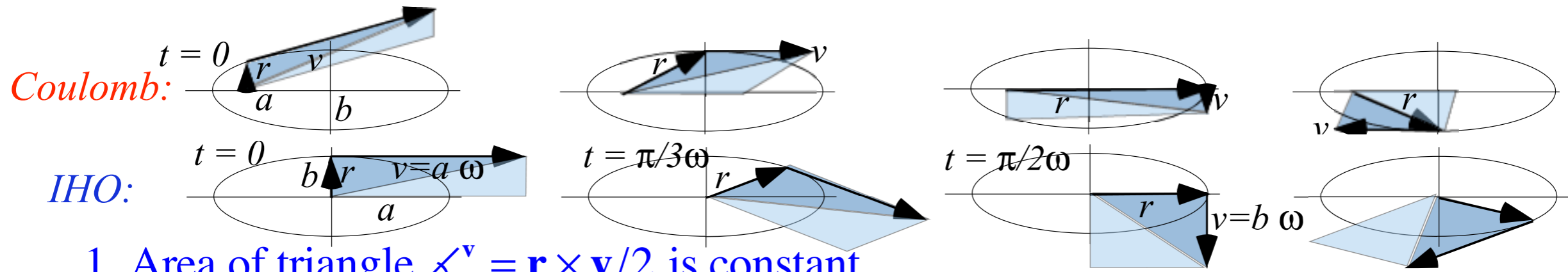
2. Angular momentum $L = m \mathbf{r} \times \mathbf{v}$ is conserved

$$L = m \mathbf{r} \times \mathbf{v} = m (r_x v_y - r_y v_x) = \begin{cases} m \cdot ab \cdot \sqrt{G\rho_{\oplus} 4\pi / 3} & \text{for IHO} \\ m \cdot a^{-1/2} b \sqrt{GM_{\oplus}} & \text{for Coul. (... in Unit 5)} \end{cases}$$

✓ for IHO
✓ for Coul.

Some Kepler's "laws" that apply to any central (isotropic) force $F(r)$

Apply to IHO: $F(r) = -k \cdot r$ with $U(r) = k \cdot r^2 / 2$ and Coulomb: $F(r) = -GMm/r^2$ with $U(r) = -GMm \cdot /r$



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✓ for IHO
(Derived in Unit 5) ✓ for Coul.

2. Angular momentum $L = m \mathbf{r} \times \mathbf{v}$ is conserved

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✓ for IHO
✓ for Coul.

3. Equal area is swept by radius vector in each equal time interval T

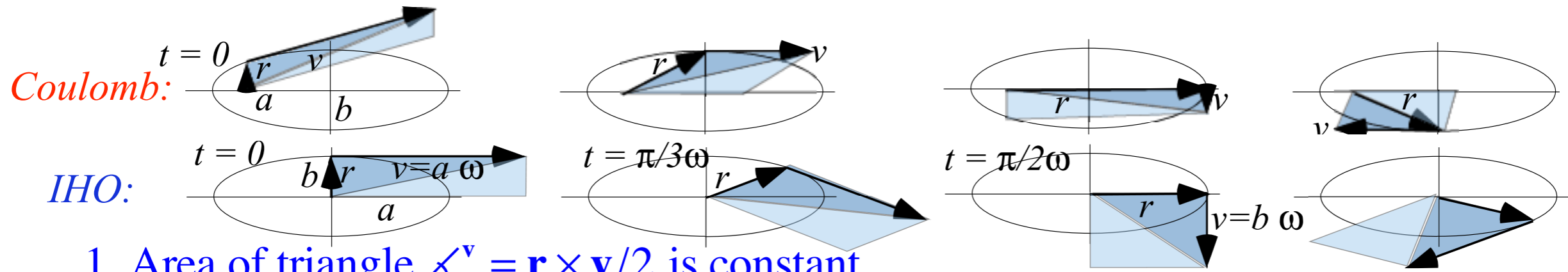
In one period:

$$\tau = \frac{1}{\nu} = \frac{2\pi}{\omega} = \frac{2mA_{\tau}}{L} = \frac{2m \cdot ab \cdot \pi}{L} = \begin{cases} \frac{2m \cdot ab \cdot \pi}{m \cdot ab \cdot \sqrt{G\rho_{\oplus} 4\pi / 3}} \\ \frac{2m \cdot ab \cdot \pi}{m \cdot a^{-1/2} b \sqrt{GM_{\oplus}}} \end{cases}$$

(G IHO formulas from Lect. 6 p.70-79)

Some Kepler's "laws" that apply to any central (isotropic) force $F(r)$

Apply to IHO: $F(r) = -k \cdot r$ with $U(r) = k \cdot r^2 / 2$ and Coulomb: $F(r) = -GMm/r^2$ with $U(r) = -GMm \cdot /r$



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$$\mathbf{r} \times \mathbf{v} = r_x v_y - r_y v_x = \begin{cases} ab \cdot \sqrt{G\rho_{\oplus} 4\pi / 3} & \text{for IHO} \\ a^{-1/2} b \sqrt{GM_{\oplus}} & \text{for Coul.} \end{cases}$$

✓ for IHO
(Derived in Unit 5) ✓ for Coul.

2. Angular momentum $L = m \mathbf{r} \times \mathbf{v}$ is conserved

$$L = m \mathbf{r} \times \mathbf{v} = m (r_x v_y - r_y v_x) = \begin{cases} m \cdot ab \cdot \sqrt{G\rho_{\oplus} 4\pi / 3} & \text{for IHO} \\ m \cdot a^{-1/2} b \sqrt{GM_{\oplus}} & \text{for Coul. (... in Unit 5)} \end{cases}$$

✓ for IHO
✓ for Coul.

3. Equal area is swept by radius vector in each equal time interval T

In one period:

$$\tau = \frac{1}{\nu} = \frac{2\pi}{\omega} = \frac{2mA_{\tau}}{L} = \frac{2m \cdot ab \cdot \pi}{L}$$

Applies to any central $F(r)$

$$= \begin{cases} \frac{2m \cdot ab \cdot \pi}{m \cdot ab \cdot \sqrt{G\rho_{\oplus} 4\pi / 3}} = \frac{2\pi}{\sqrt{G\rho_{\oplus} 4\pi / 3}} & \text{for IHO} \\ \frac{2m \cdot ab \cdot \pi}{m \cdot a^{-1/2} b \sqrt{GM_{\oplus}}} = \frac{2\pi}{a^{-3/2} \sqrt{GM_{\oplus}}} & \text{for Coul.} \end{cases}$$

(not a function of a or b)
that is ω_{IHO}
(not a function of b)
that is ω_{Coul}

Some Kepler's "laws" for all central (isotropic) force $F(r)$ fields

Angular momentum invariance of IHO: $F(r) = -k \cdot r$ with $U(r) = k \cdot r^2 / 2$

(Derived here)

Angular momentum invariance of Coulomb: $F(r) = -GMm/r^2$ with $U(r) = -GMm \cdot /r$ (Derived in Unit 5)



Total energy $E = KE + PE$ invariance of IHO: $F(r) = -k \cdot r$

(Derived here)

Total energy $E = KE + PE$ invariance of Coulomb: $F(r) = -GMm/r^2$

(Derived in Unit 5)

Kepler laws involve \mathbf{L} -momentum conservation in isotropic force $F(r)$

Now consider orbital energy conservation of the IHO: $F(r)=-k \cdot r$ with $U(r)=k \cdot r^2/2$

Total energy= $KE + PE$ is constant

$$\begin{aligned} KE + PE &= \frac{1}{2} \mathbf{v} \cdot \mathbf{M} \cdot \mathbf{v} + \frac{1}{2} \mathbf{r} \cdot \mathbf{K} \cdot \mathbf{r} \\ &= \frac{1}{2} \begin{pmatrix} v_x & v_y \end{pmatrix} \cdot \begin{pmatrix} m & 0 \\ 0 & m \end{pmatrix} \cdot \begin{pmatrix} v_x \\ v_y \end{pmatrix} + \begin{pmatrix} r_x & r_y \end{pmatrix} \cdot \begin{pmatrix} k & 0 \\ 0 & k \end{pmatrix} \cdot \begin{pmatrix} r_x \\ r_y \end{pmatrix} \\ &= \frac{1}{2} m v_x^2 + \frac{1}{2} m v_y^2 + \frac{1}{2} k r_x^2 + \frac{1}{2} k r_y^2 \\ &= \frac{1}{2} m (-a\omega \sin \omega t)^2 + \frac{1}{2} m (b\omega \cos \omega t)^2 + \frac{1}{2} k (a \cos \omega t)^2 + \frac{1}{2} k (b \sin \omega t)^2 \end{aligned}$$

$$\begin{pmatrix} v_x \\ v_y \end{pmatrix} = \begin{pmatrix} -a\omega \sin \omega t \\ b\omega \cos \omega t \end{pmatrix}$$

$$\begin{pmatrix} r_x \\ r_y \end{pmatrix} = \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} a \cos \omega t \\ b \sin \omega t \end{pmatrix}$$

Kepler laws involve Δ -momentum conservation in isotropic force $F(r)$

Now consider orbital energy conservation of the IHO: $F(r) = -k \cdot r$ with $U(r) = k \cdot r^2 / 2$

Total IHO energy = KE + PE is constant

$$\begin{aligned} KE + PE &= \frac{1}{2} \mathbf{v} \cdot \mathbf{M} \cdot \mathbf{v} + \frac{1}{2} \mathbf{r} \cdot \mathbf{K} \cdot \mathbf{r} \\ &= \frac{1}{2} \begin{pmatrix} v_x & v_y \end{pmatrix} \cdot \begin{pmatrix} m & 0 \\ 0 & m \end{pmatrix} \cdot \begin{pmatrix} v_x \\ v_y \end{pmatrix} + \begin{pmatrix} r_x & r_y \end{pmatrix} \cdot \begin{pmatrix} k & 0 \\ 0 & k \end{pmatrix} \cdot \begin{pmatrix} r_x \\ r_y \end{pmatrix} \\ &= \frac{1}{2} m v_x^2 + \frac{1}{2} m v_y^2 + \frac{1}{2} k r_x^2 + \frac{1}{2} k r_y^2 \\ &= \frac{1}{2} m (-a\omega \sin \omega t)^2 + \frac{1}{2} m (b\omega \cos \omega t)^2 + \frac{1}{2} k (a \cos \omega t)^2 + \frac{1}{2} k (b \sin \omega t)^2 \\ &= \frac{1}{2} m a^2 \omega^2 (\sin^2 \omega t) + \frac{1}{2} m b^2 \omega^2 (\cos^2 \omega t) + \frac{1}{2} k a^2 (\cos^2 \omega t) + \frac{1}{2} k b^2 (\sin^2 \omega t) \\ &= \frac{1}{2} m \omega^2 (a^2 + b^2) \quad \text{Given : } k = m\omega^2 \end{aligned}$$

Kepler laws involve Δ -momentum conservation in isotropic force $F(r)$

Now consider orbital energy conservation of the IHO: $F(r)=-k \cdot r$ with $U(r)=k \cdot r^2/2$

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$$\begin{aligned}
 KE + PE &= \frac{1}{2} \mathbf{v} \cdot \mathbf{M} \cdot \mathbf{v} + \frac{1}{2} \mathbf{r} \cdot \mathbf{K} \cdot \mathbf{r} \\
 &= \frac{1}{2} \begin{pmatrix} v_x & v_y \end{pmatrix} \cdot \begin{pmatrix} m & 0 \\ 0 & m \end{pmatrix} \cdot \begin{pmatrix} v_x \\ v_y \end{pmatrix} + \begin{pmatrix} r_x & r_y \end{pmatrix} \cdot \begin{pmatrix} k & 0 \\ 0 & k \end{pmatrix} \cdot \begin{pmatrix} r_x \\ r_y \end{pmatrix} \\
 &= \frac{1}{2} m v_x^2 + \frac{1}{2} m v_y^2 + \frac{1}{2} k r_x^2 + \frac{1}{2} k r_y^2 \\
 &= \frac{1}{2} m (-a\omega \sin \omega t)^2 + \frac{1}{2} m (b\omega \cos \omega t)^2 + \frac{1}{2} k (a \cos \omega t)^2 + \frac{1}{2} k (b \sin \omega t)^2 \\
 &= \frac{1}{2} m a^2 \omega^2 (\sin^2 \omega t) + \frac{1}{2} m b^2 \omega^2 (\cos^2 \omega t) + \frac{1}{2} k a^2 (\cos^2 \omega t) + \frac{1}{2} k b^2 (\sin^2 \omega t) \\
 &= \frac{1}{2} m \omega^2 (a^2 + b^2) \quad \text{Given : } k = m\omega^2
 \end{aligned}$$

$$E = KE + PE = \frac{1}{2} m \omega^2 (a^2 + b^2) = \frac{1}{2} k (a^2 + b^2) \quad \text{since: } \omega = \sqrt{\frac{k}{m}} = \sqrt{G \rho_{\oplus} 4\pi / 3} \quad \text{or: } m\omega^2 = k$$

Some Kepler's "laws" for all central (isotropic) force $F(r)$ fields

Angular momentum invariance of IHO: $F(r) = -k \cdot r$ with $U(r) = k \cdot r^2 / 2$ (Derived here)

Angular momentum invariance of Coulomb: $F(r) = -GMm/r^2$ with $U(r) = -GMm \cdot /r$ (Derived in Unit 5)

Total energy $E = KE + PE$ invariance of IHO: $F(r) = -k \cdot r$ (Derived here)

 *Total energy $E = KE + PE$ invariance of Coulomb: $F(r) = -GMm/r^2$ (Derived in Unit 5)*

Kepler laws involve ∇ -momentum conservation in isotropic force $F(r)$

Now consider orbital energy conservation of the IHO: $F(r)=-k \cdot r$ with $U(r)=k \cdot r^2/2$

Total IHO energy= $KE + PE$ is constant

$$\begin{aligned}
 KE + PE &= \frac{1}{2} \mathbf{v} \cdot \mathbf{M} \cdot \mathbf{v} + \frac{1}{2} \mathbf{r} \cdot \mathbf{K} \cdot \mathbf{r} \\
 &= \frac{1}{2} \begin{pmatrix} v_x & v_y \end{pmatrix} \cdot \begin{pmatrix} m & 0 \\ 0 & m \end{pmatrix} \cdot \begin{pmatrix} v_x \\ v_y \end{pmatrix} + \begin{pmatrix} r_x & r_y \end{pmatrix} \cdot \begin{pmatrix} k & 0 \\ 0 & k \end{pmatrix} \cdot \begin{pmatrix} r_x \\ r_y \end{pmatrix} \\
 &= \frac{1}{2} m v_x^2 + \frac{1}{2} m v_y^2 + \frac{1}{2} k r_x^2 + \frac{1}{2} k r_y^2 \\
 &= \frac{1}{2} m (-a\omega \sin \omega t)^2 + \frac{1}{2} m (b\omega \cos \omega t)^2 + \frac{1}{2} k (a \cos \omega t)^2 + \frac{1}{2} k (b \sin \omega t)^2 \\
 &= \frac{1}{2} m a^2 \omega^2 (\sin^2 \omega t) + \frac{1}{2} m b^2 \omega^2 (\cos^2 \omega t) + \frac{1}{2} k a^2 (\cos^2 \omega t) + \frac{1}{2} k b^2 (\sin^2 \omega t) \\
 &= \frac{1}{2} m \omega^2 (a^2 + b^2) \quad \text{Given : } k = m\omega^2
 \end{aligned}$$

$$E = KE + PE = \frac{1}{2} m \omega^2 (a^2 + b^2) = \frac{1}{2} k (a^2 + b^2) \quad \text{since: } \omega = \sqrt{\frac{k}{m}} = \sqrt{G \rho_{\oplus} 4\pi / 3} \quad \text{or: } m\omega^2 = k$$

We'll see that the Coul. orbits are simpler:

(like the period...not a function of b)

Kepler laws involve ∇ -momentum conservation in isotropic force $F(r)$

Now consider orbital energy conservation of the IHO: $F(r)=-k \cdot r$ with $U(r)=k \cdot r^2/2$

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 KE + PE &= \frac{1}{2} \mathbf{v} \cdot \mathbf{M} \cdot \mathbf{v} + \frac{1}{2} \mathbf{r} \cdot \mathbf{K} \cdot \mathbf{r} \\
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 &= \frac{1}{2} m v_x^2 + \frac{1}{2} m v_y^2 + \frac{1}{2} k r_x^2 + \frac{1}{2} k r_y^2 \\
 &= \frac{1}{2} m (-a\omega \sin \omega t)^2 + \frac{1}{2} m (b\omega \cos \omega t)^2 + \frac{1}{2} k (a \cos \omega t)^2 + \frac{1}{2} k (b \sin \omega t)^2 \\
 &= \frac{1}{2} m a^2 \omega^2 (\sin^2 \omega t) + \frac{1}{2} m b^2 \omega^2 (\cos^2 \omega t) + \frac{1}{2} k a^2 (\cos^2 \omega t) + \frac{1}{2} k b^2 (\sin^2 \omega t) \\
 &= \frac{1}{2} m \omega^2 (a^2 + b^2) \quad \text{Given : } k = m\omega^2
 \end{aligned}$$

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We'll see that the Coul. orbits are simpler:

(like the period...not a function of b)

$$E = KE + PE = \frac{1}{2} m v_x^2 + \frac{1}{2} m v_y^2 - \frac{k}{r} = \frac{1}{2} m v_x^2 + \frac{1}{2} m v_y^2 - \frac{GM_{\oplus} m}{r} = -\frac{GM_{\oplus} m}{a}$$

- Introduction to dual matrix operator contact geometry (based on IHO orbits)
- Quadratic form ellipse $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = 1$ vs. inverse form ellipse $\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p} = 1$
 - Duality norm relations ($\mathbf{r} \cdot \mathbf{p} = 1$)
 - \mathbf{Q} -Ellipse tangents \mathbf{r}' normal to dual \mathbf{Q}^{-1} -ellipse position \mathbf{p} ($\mathbf{r}' \cdot \mathbf{p} = 0 = \mathbf{r} \cdot \mathbf{p}'$)
 - Operator geometric sequences and eigenvectors
 - Alternative scaling of matrix operator geometry
 - Vector calculus of tensor operation

Quadratic forms and tangent contact geometry of their ellipses

A matrix Q that generates an ellipse by $\mathbf{r} \bullet Q \bullet \mathbf{r} = 1$ is called positive-definite (if $\mathbf{r} \bullet Q \bullet \mathbf{r}$ always > 0)

$$\begin{pmatrix} x & y \end{pmatrix} \bullet \overbrace{\begin{pmatrix} \frac{1}{a^2} & 0 \\ 0 & \frac{1}{b^2} \end{pmatrix}}^{\mathbf{r} \bullet Q \bullet \mathbf{r}} \bullet \begin{pmatrix} x \\ y \end{pmatrix} = 1 = \overbrace{\begin{pmatrix} x & y \end{pmatrix}}^{\mathbf{r}} \bullet \overbrace{\begin{pmatrix} \frac{x}{a^2} \\ \frac{y}{b^2} \end{pmatrix}}^{Q \bullet \mathbf{r}} = \frac{x^2}{a^2} + \frac{y^2}{b^2}$$

A inverse matrix Q^{-1} generates an ellipse by $\mathbf{p} \bullet Q^{-1} \bullet \mathbf{p} = 1$ called inverse or dual ellipse:

$$\begin{pmatrix} p_x & p_y \end{pmatrix} \bullet \overbrace{\begin{pmatrix} a^2 & 0 \\ 0 & b^2 \end{pmatrix}}^{\mathbf{p} \bullet Q^{-1} \bullet \mathbf{p}} \bullet \begin{pmatrix} p_x \\ p_y \end{pmatrix} = 1 = \overbrace{\begin{pmatrix} p_x & p_y \end{pmatrix}}^{\mathbf{p}} \bullet \overbrace{\begin{pmatrix} a^2 p_x \\ b^2 p_y \end{pmatrix}}^{Q^{-1} \bullet \mathbf{p}} = a^2 p_x^2 + b^2 p_y^2$$

Quadratic forms and tangent contact geometry of their ellipses

A matrix Q that generates an ellipse by $\mathbf{r} \bullet Q \bullet \mathbf{r} = 1$ is called positive-definite (if $\mathbf{r} \bullet Q \bullet \mathbf{r}$ always > 0)

$$\left(x \quad y \right) \bullet \overbrace{\begin{pmatrix} \frac{1}{a^2} & 0 \\ 0 & \frac{1}{b^2} \end{pmatrix}}^{\mathbf{r} \bullet Q \bullet \mathbf{r}} \bullet \begin{pmatrix} x \\ y \end{pmatrix} = 1 = \overbrace{\begin{pmatrix} x & y \end{pmatrix}}^{\mathbf{r}} \bullet \overbrace{\begin{pmatrix} \frac{x}{a^2} \\ \frac{y}{b^2} \end{pmatrix}}^{Q \bullet \mathbf{r} = \mathbf{p}} = \frac{x^2}{a^2} + \frac{y^2}{b^2}$$

Defined mapping between ellipses

A inverse matrix Q^{-1} generates an ellipse by $\mathbf{p} \bullet Q^{-1} \bullet \mathbf{p} = 1$ called inverse or dual ellipse:

$$\left(p_x \quad p_y \right) \bullet \overbrace{\begin{pmatrix} a^2 & 0 \\ 0 & b^2 \end{pmatrix}}^{\mathbf{p} \bullet Q^{-1} \bullet \mathbf{p}} \bullet \begin{pmatrix} p_x \\ p_y \end{pmatrix} = 1 = \overbrace{\begin{pmatrix} p_x & p_y \end{pmatrix}}^{\mathbf{p}} \bullet \overbrace{\begin{pmatrix} a^2 p_x \\ b^2 p_y \end{pmatrix}}^{Q^{-1} \bullet \mathbf{p} = \mathbf{r}} = a^2 p_x^2 + b^2 p_y^2$$

Introduction to dual matrix operator contact geometry (based on IHO orbits)

→ *Quadratic form ellipse $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = 1$ vs. inverse form ellipse $\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p} = 1$*

Duality norm relations ($\mathbf{r} \cdot \mathbf{p} = 1$)

\mathbf{Q} -Ellipse tangents \mathbf{r}' normal to dual \mathbf{Q}^{-1} -ellipse position \mathbf{p} ($\mathbf{r}' \cdot \mathbf{p} = 0 = \mathbf{r} \cdot \mathbf{p}'$)

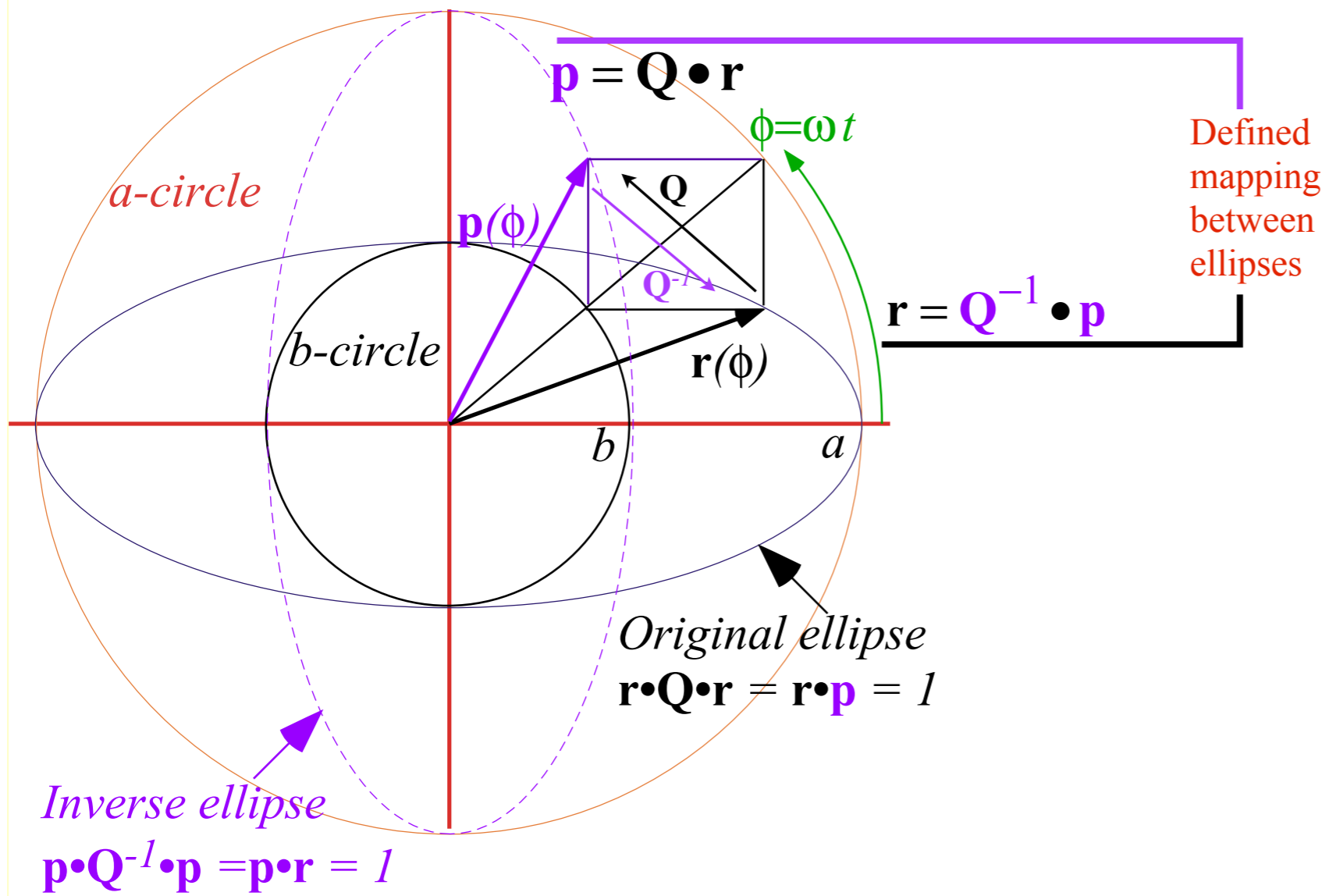
Operator geometric sequences and eigenvectors

Alternative scaling of matrix operator geometry

Vector calculus of tensor operation

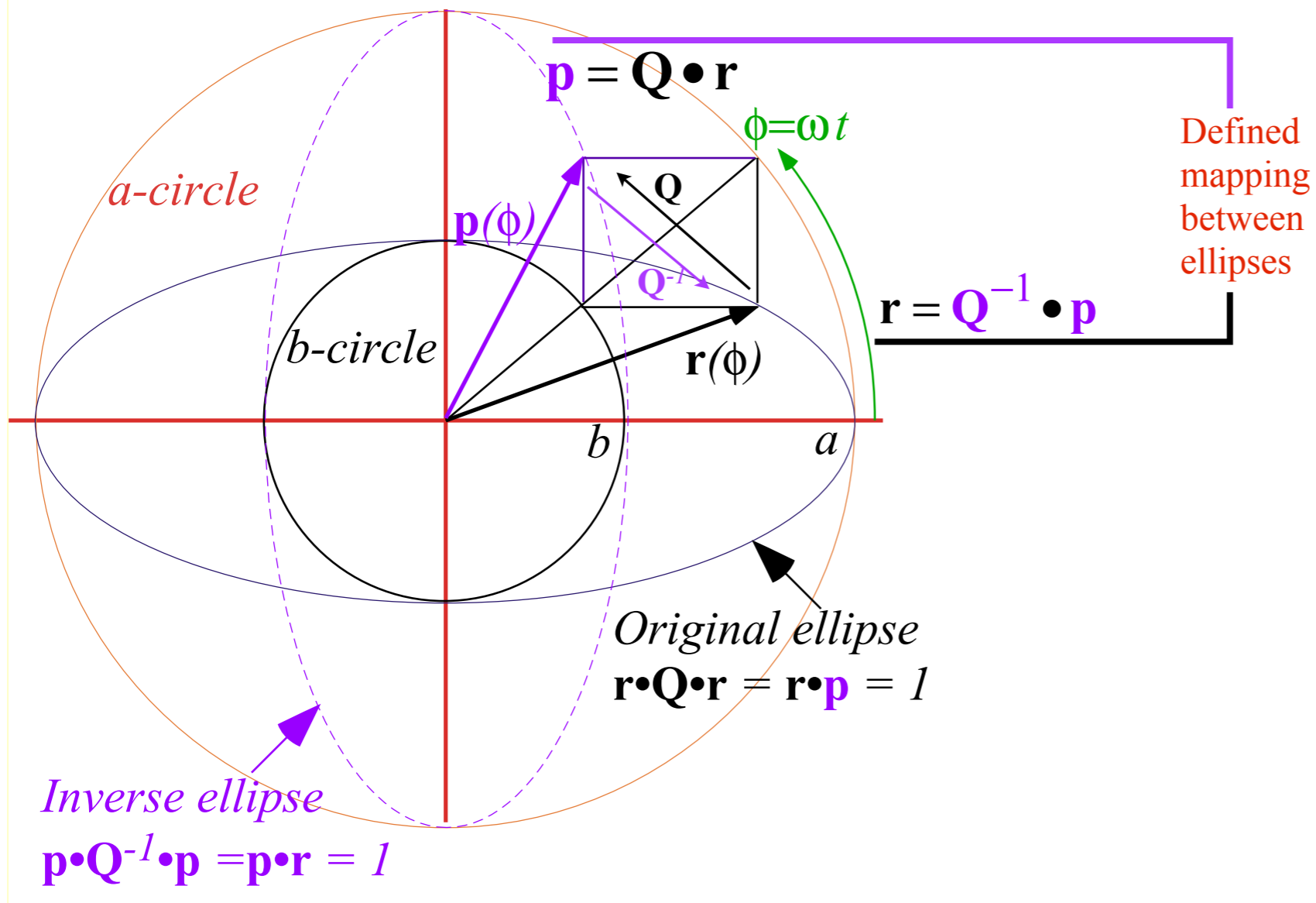
(a) Quadratic form ellipse and
Inverse quadratic form ellipse

based on
Unit 1
Fig. 11.6



(a) Quadratic form ellipse and
Inverse quadratic form ellipse

based on
Unit 1
Fig. 11.6



Here plot of \mathbf{p} -ellipse is re-scaled by scalefactor $S = a \cdot b$

\mathbf{p} -ellipse x -radius = $1/a$ plotted at: $S(1/a) = b$ ($=1$ for $a=2, b=1$)

\mathbf{p} -ellipse y -radius = $1/b$ plotted at: $S(1/b) = a$ ($=2$ for $a=2, b=1$)

Introduction to dual matrix operator geometry (based on IHO orbits)

Quadratic form ellipse $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = 1$ vs. inverse form ellipse $\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p} = 1$



Duality norm relations ($\mathbf{r} \cdot \mathbf{p} = 1$)

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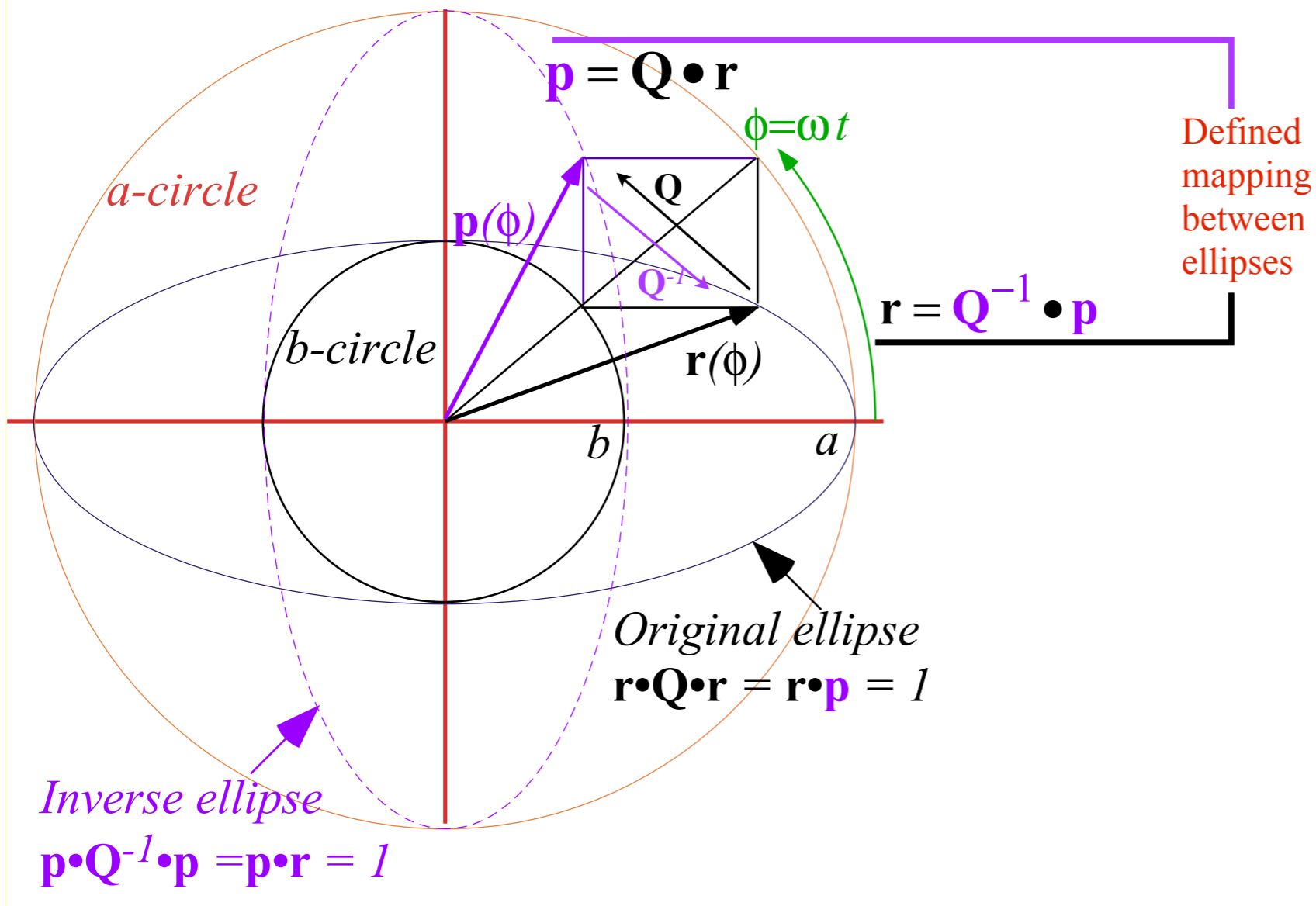
Operator geometric sequences and eigenvectors

Alternative scaling of matrix operator geometry

Vector calculus of tensor operation

(a) Quadratic form ellipse and
Inverse quadratic form ellipse

based on
Unit 1
Fig. 11.6



Quadratic form $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = 1$ has mutual duality relations with inverse form $\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p} = 1 = \mathbf{p} \cdot \mathbf{r}$

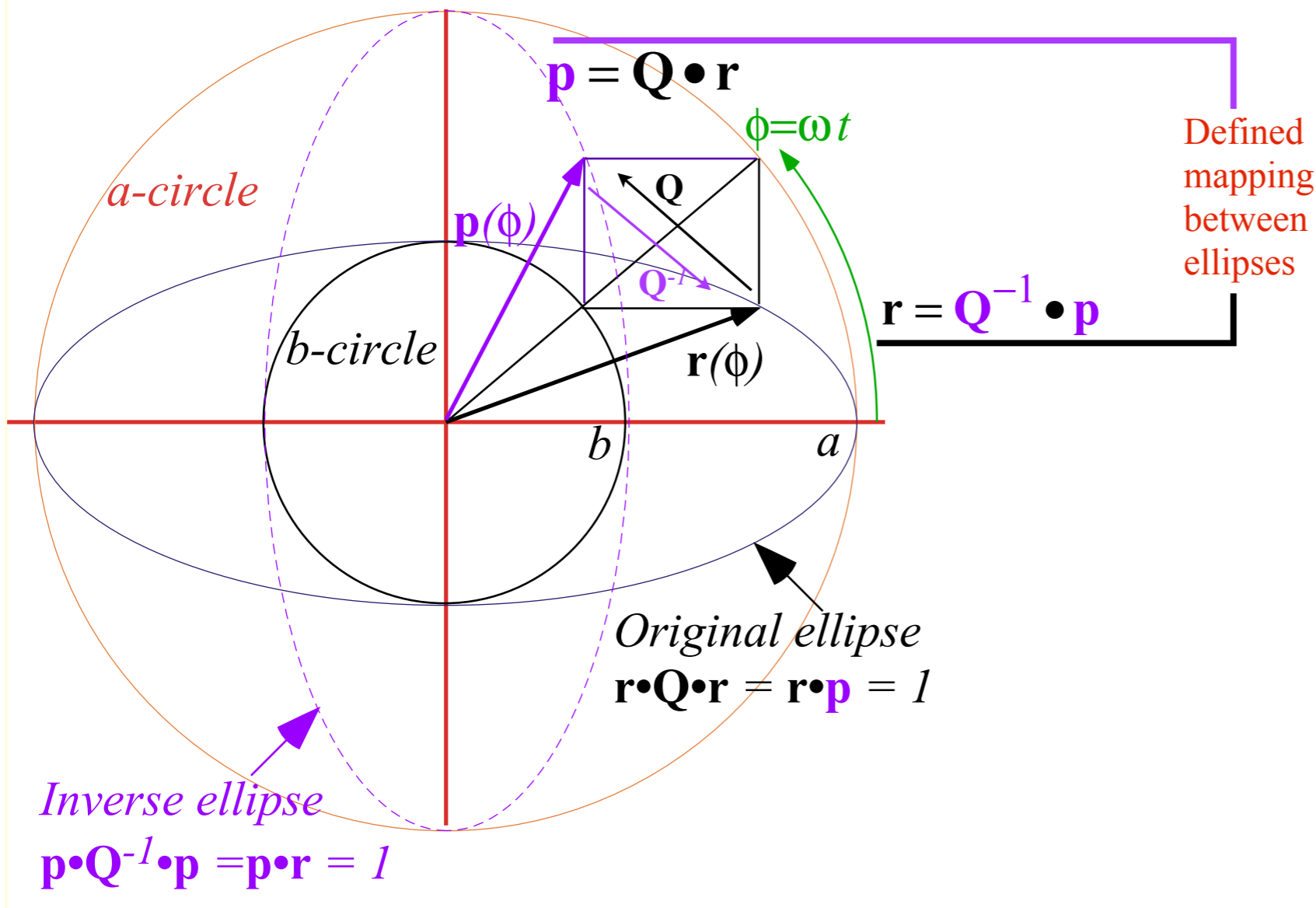
Here plot of \mathbf{p} -ellipse is re-scaled by scalefactor $S = a \cdot b$

\mathbf{p} -ellipse x -radius = $1/a$ plotted at: $S(1/a) = b$ ($=1$ for $a=2, b=1$)

\mathbf{p} -ellipse y -radius = $1/b$ plotted at: $S(1/b) = a$ ($=2$ for $a=2, b=1$)

(a) Quadratic form ellipse and Inverse quadratic form ellipse

based on
Unit 1
Fig. 11.6



Quadratic form $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = 1$ has mutual duality relations with inverse form $\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p} = 1 = \mathbf{p} \cdot \mathbf{r}$

$$\mathbf{p} = \mathbf{Q} \cdot \mathbf{r} = \begin{pmatrix} \overbrace{1/a^2}^{\mathbf{Q}} & 0 \\ 0 & \overbrace{1/b^2}^{\mathbf{Q}} \end{pmatrix} \cdot \begin{pmatrix} \overbrace{x}^{\mathbf{r}} \\ \overbrace{y}^{\mathbf{r}} \end{pmatrix} = \begin{pmatrix} \overbrace{x/a^2}^{\mathbf{p}} \\ \overbrace{y/b^2}^{\mathbf{p}} \end{pmatrix} = \begin{pmatrix} (1/a)\cos\phi \\ (1/b)\sin\phi \end{pmatrix} \text{ where: } \begin{matrix} x = r_x = a \cos\phi = a \cos\omega t \\ y = r_y = b \sin\phi = b \sin\omega t \end{matrix} \text{ so: } \boxed{\mathbf{p} \cdot \mathbf{r} = 1}$$

Here plot of \mathbf{p} -ellipse is re-scaled by scalefactor $S = a \cdot b$

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\mathbf{p} -ellipse y -radius $= 1/b$ plotted at: $S(1/b) = a$ ($= 2$ for $a = 2, b = 1$)

- [Link \$\Rightarrow\$ BoxIt simulation of IHO orbits](#)
- [Link \$\rightarrow\$ IHO orbital time rates of change](#)
- [Link \$\rightarrow\$ IHO Exegesis Plot](#)

Introduction to dual matrix operator geometry (based on IHO orbits)

Quadratic form ellipse $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = 1$ vs. inverse form ellipse $\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p} = 1$

Duality norm relations ($\mathbf{r} \cdot \mathbf{p} = 1$)

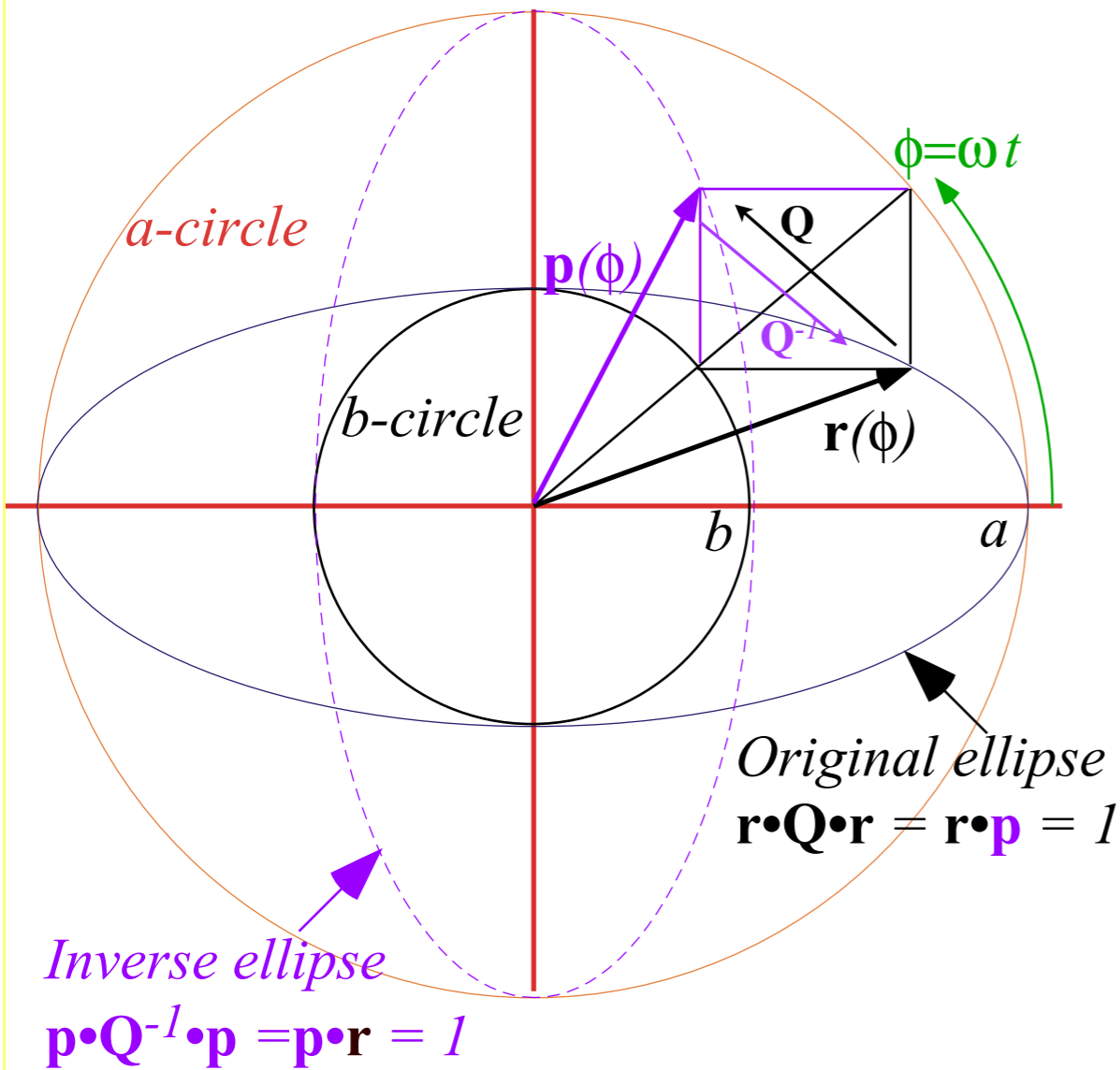
 *Q -Ellipse tangents \mathbf{r}' normal to dual Q^{-1} -ellipse position \mathbf{p} ($\mathbf{r}' \cdot \mathbf{p} = 0 = \mathbf{r} \cdot \mathbf{p}'$)*

Operator geometric sequences and eigenvectors

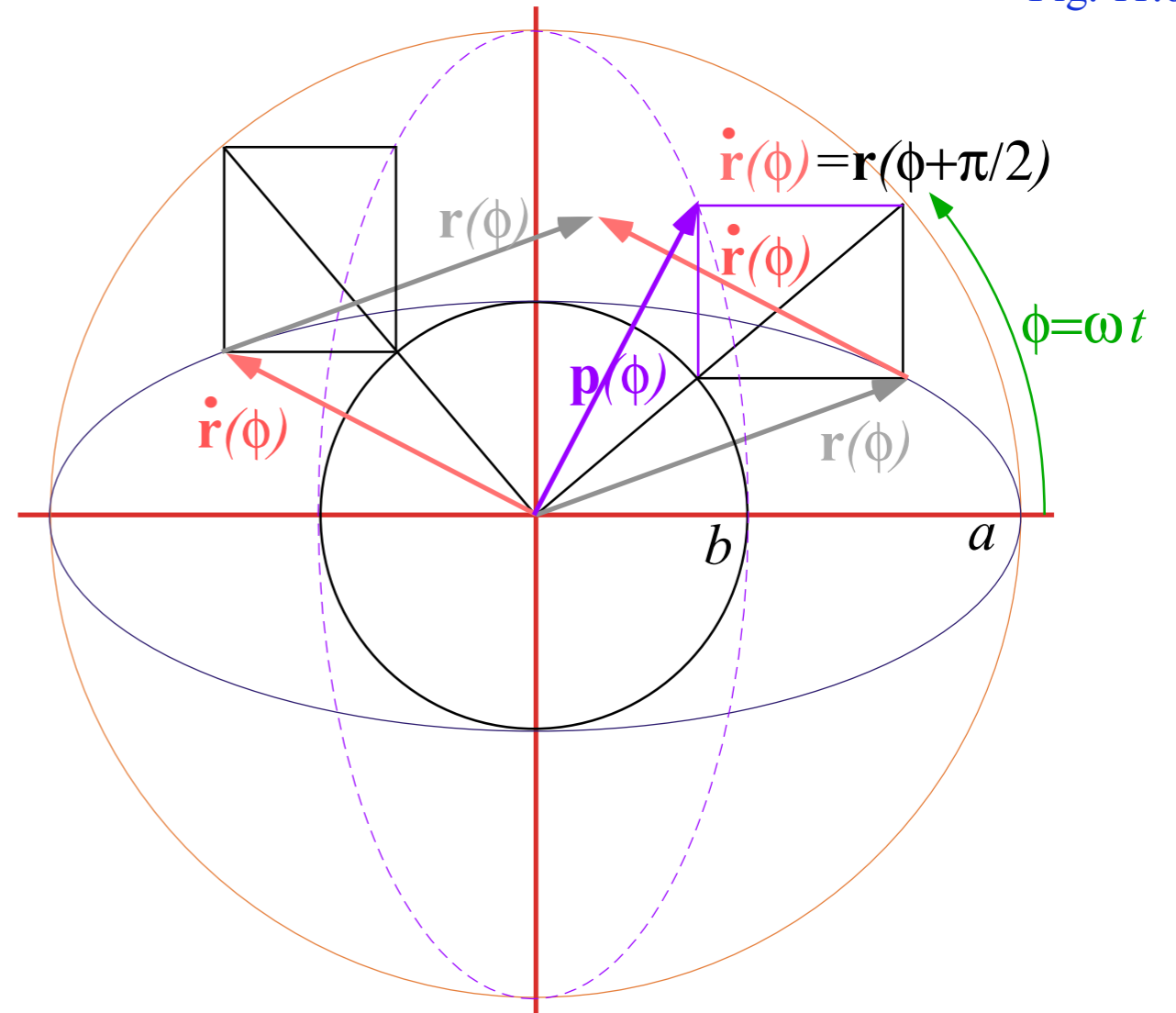
Alternative scaling of matrix operator geometry

Vector calculus of tensor operation

(a) Quadratic form ellipse and Inverse quadratic form ellipse



(b) Ellipse tangents



Quadratic form $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = 1$ has mutual duality relations with inverse form $\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p} = 1 = \mathbf{p} \cdot \mathbf{r}$

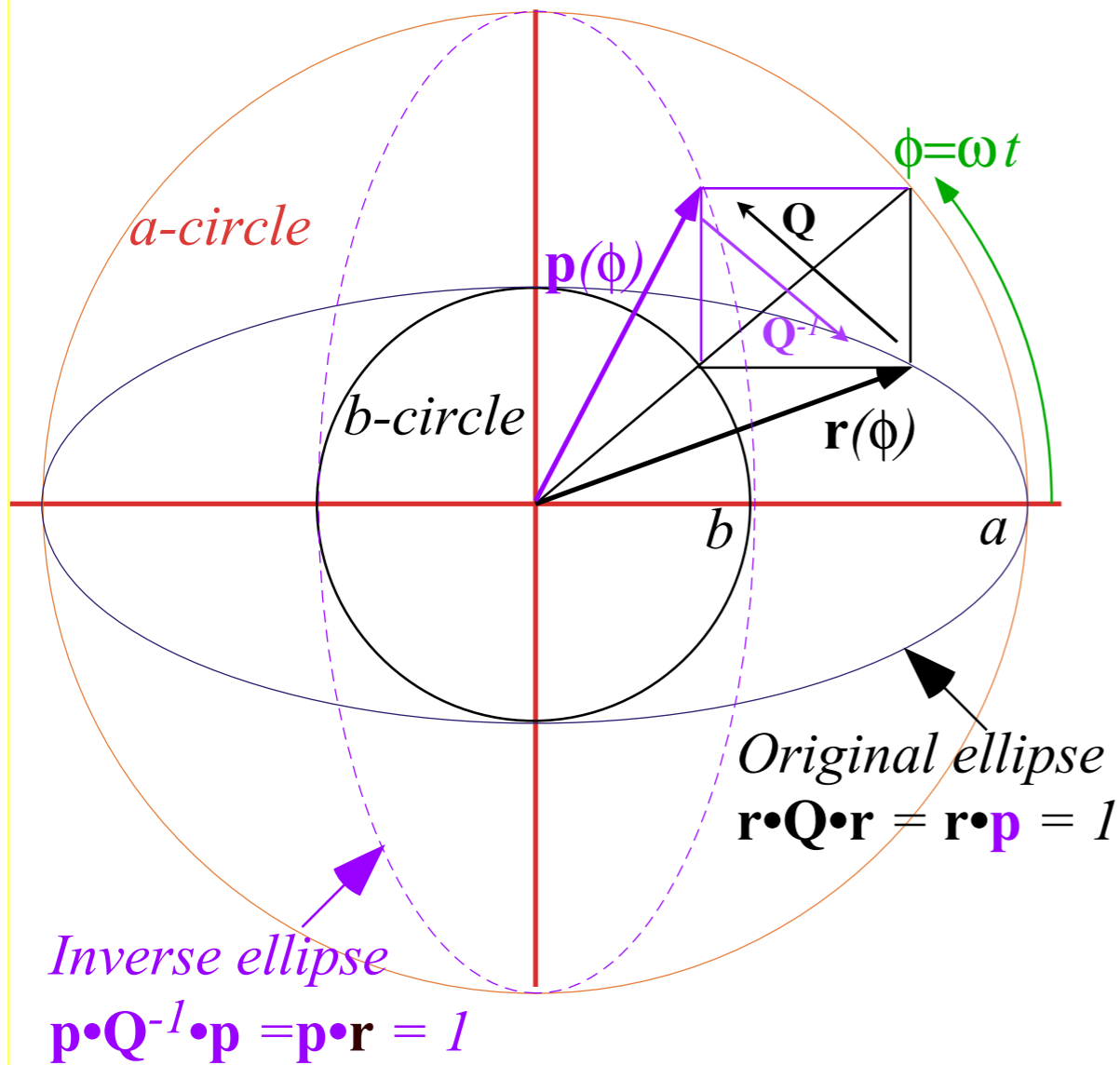
$$\mathbf{p} = \mathbf{Q} \cdot \mathbf{r} = \begin{pmatrix} 1/a^2 & 0 \\ 0 & 1/b^2 \end{pmatrix} \cdot \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a^2 \\ y/b^2 \end{pmatrix} = \begin{pmatrix} (1/a)\cos\phi \\ (1/b)\sin\phi \end{pmatrix} \text{ where: } \begin{matrix} x = r_x = a \cos\phi = a \cos\omega t \\ y = r_y = b \sin\phi = b \sin\omega t \end{matrix} \text{ so: } \boxed{\mathbf{p} \cdot \mathbf{r} = 1}$$

Here plot of \mathbf{p} -ellipse is re-scaled by scalefactor $S = a \cdot b$

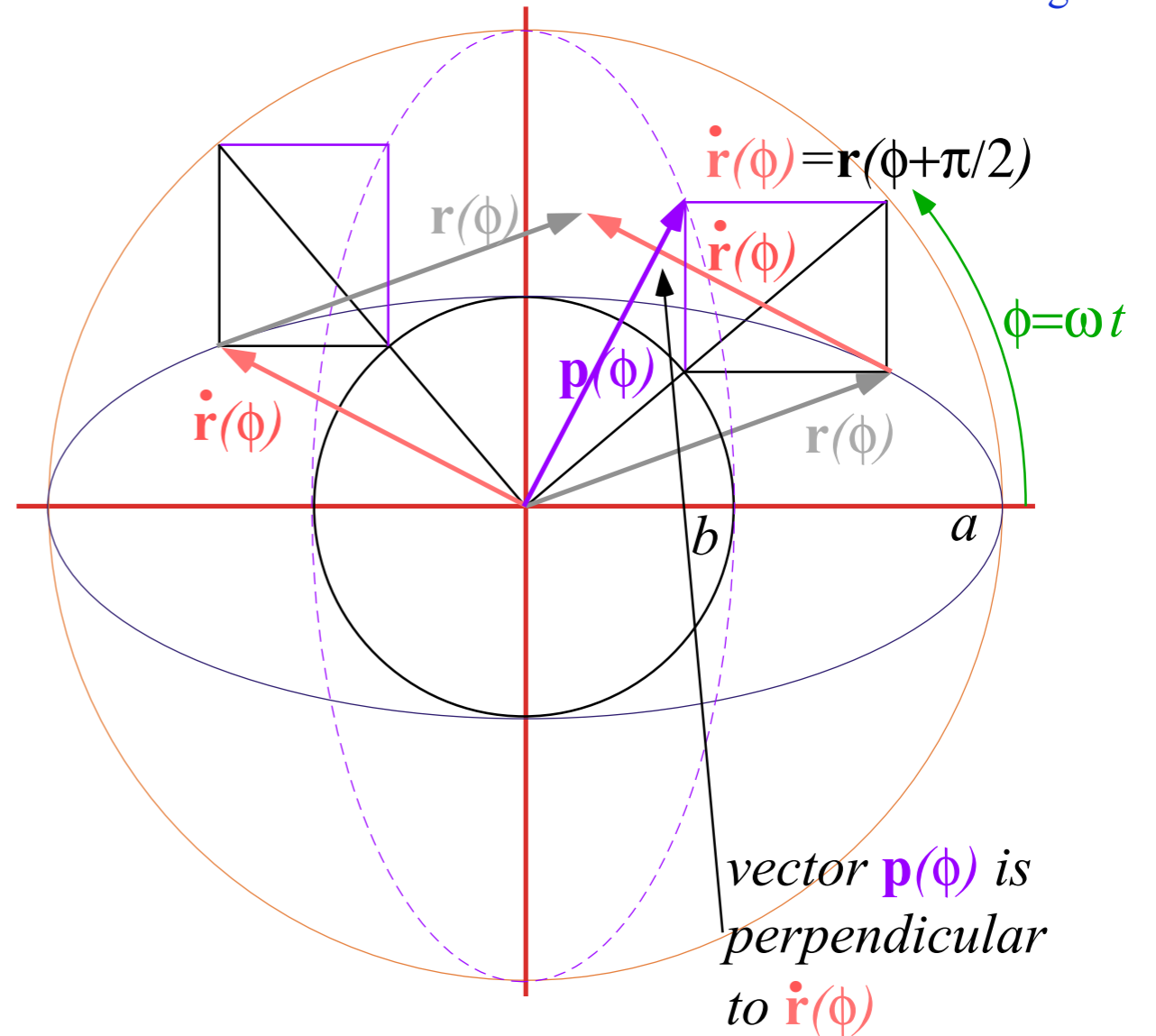
\mathbf{p} -ellipse x -radius $= 1/a$ plotted at: $S(1/a) = b$ ($= 1$ for $a=2, b=1$)

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(a) Quadratic form ellipse and Inverse quadratic form ellipse



(b) Ellipse tangents



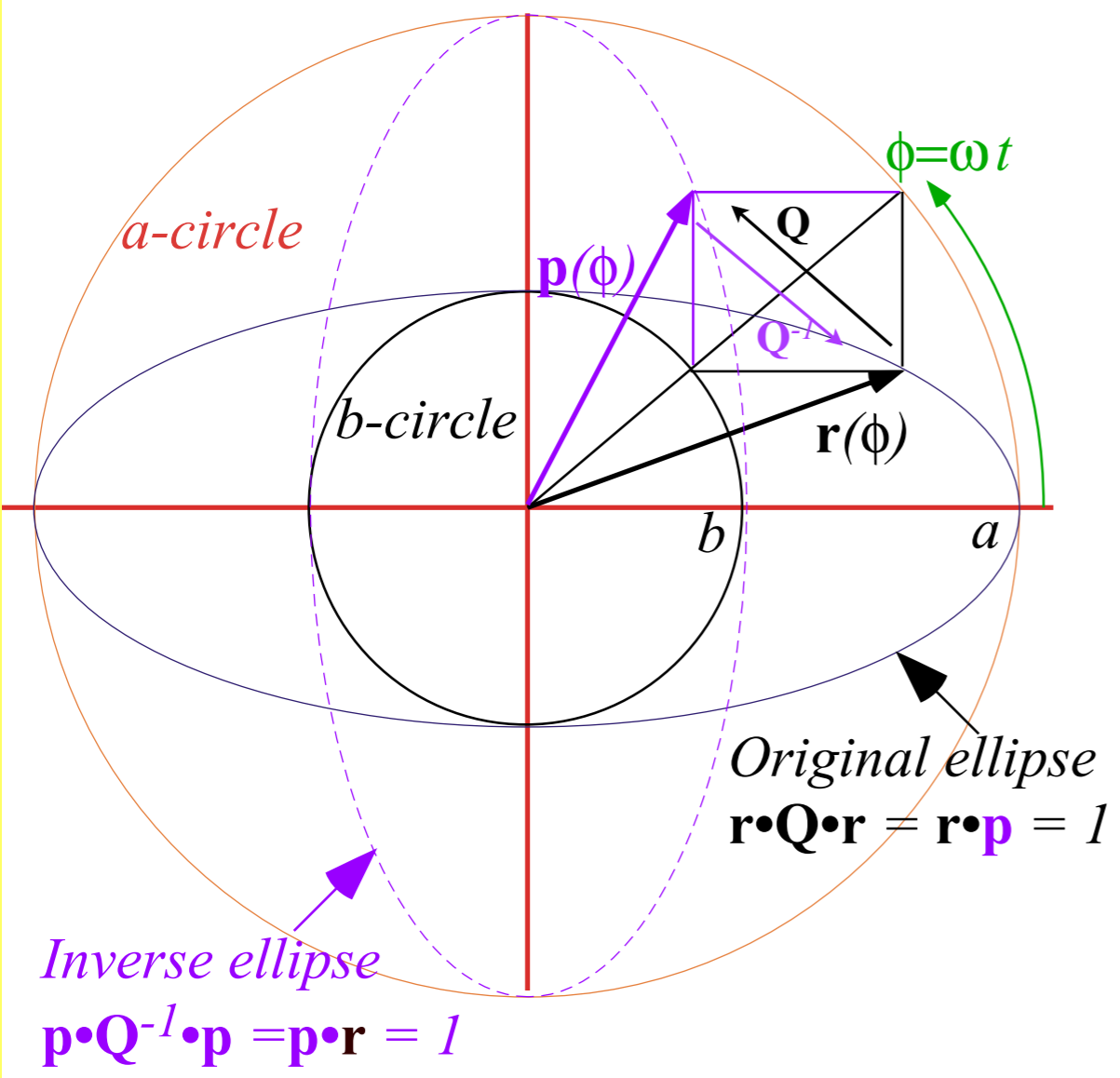
Quadratic form $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = 1$ has mutual duality relations with inverse form $\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p} = 1 = \mathbf{p} \cdot \mathbf{r}$

$$\mathbf{p} = \mathbf{Q} \cdot \mathbf{r} = \begin{pmatrix} 1/a^2 & 0 \\ 0 & 1/b^2 \end{pmatrix} \cdot \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a^2 \\ y/b^2 \end{pmatrix} = \begin{pmatrix} (1/a)\cos\phi \\ (1/b)\sin\phi \end{pmatrix} \quad \text{where:} \quad \begin{matrix} x = r_x = a \cos\phi = a \cos\omega t \\ y = r_y = b \sin\phi = b \sin\omega t \end{matrix} \quad \text{so: } \boxed{\mathbf{p} \cdot \mathbf{r} = 1}$$

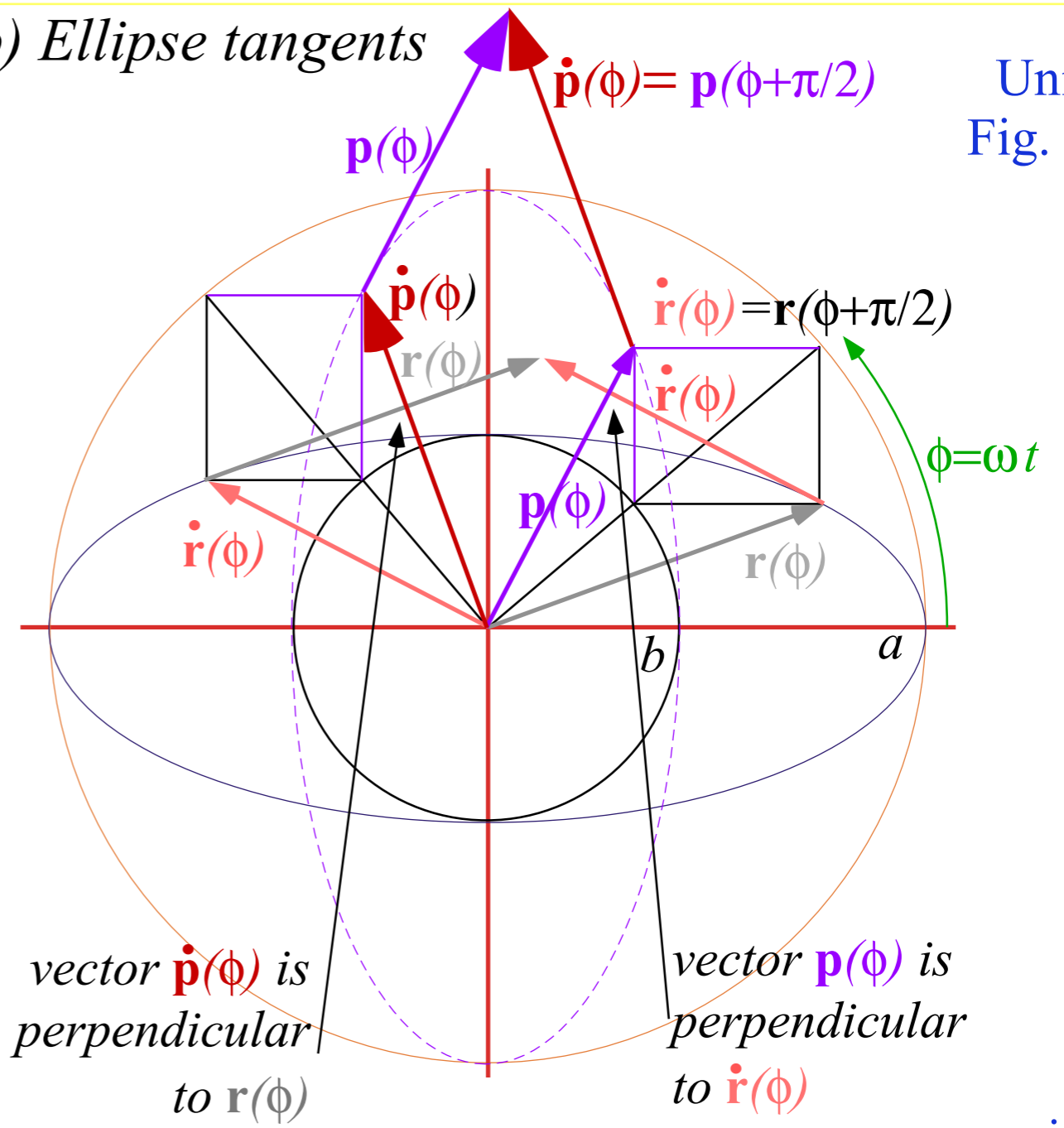
\mathbf{p} is perpendicular to velocity $\mathbf{v} = \dot{\mathbf{r}}$, a mutual orthogonality

$$\boxed{\dot{\mathbf{r}} \cdot \mathbf{p} = 0} = \begin{pmatrix} \dot{r}_x & \dot{r}_y \end{pmatrix} \cdot \begin{pmatrix} p_x \\ p_y \end{pmatrix} = \begin{pmatrix} -a \sin\phi & b \cos\phi \end{pmatrix} \cdot \begin{pmatrix} (1/a)\cos\phi \\ (1/b)\sin\phi \end{pmatrix} \quad \text{where:} \quad \begin{matrix} \dot{r}_x = -a \sin\phi \\ \dot{r}_y = b \cos\phi \end{matrix} \quad \text{and:} \quad \begin{matrix} p_x = (1/a)\cos\phi \\ p_y = (1/b)\sin\phi \end{matrix}$$

(a) Quadratic form ellipse and Inverse quadratic form ellipse



(b) Ellipse tangents



Quadratic form $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = 1$ has mutual duality relations with inverse form $\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p} = 1$

unit mutual projection

$$\mathbf{p} = \mathbf{Q} \cdot \mathbf{r} = \begin{pmatrix} 1/a^2 & 0 \\ 0 & 1/b^2 \end{pmatrix} \cdot \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a^2 \\ y/b^2 \end{pmatrix} = \begin{pmatrix} (1/a)\cos\phi \\ (1/b)\sin\phi \end{pmatrix} \text{ where: } \begin{matrix} x = r_x = a \cos\phi = a \cos\omega t \\ y = r_y = b \sin\phi = b \sin\omega t \end{matrix} \text{ so: } \boxed{\mathbf{p} \cdot \mathbf{r} = 1}$$

\mathbf{p} is perpendicular to velocity $\mathbf{v} = \dot{\mathbf{r}}$, a mutual orthogonality. So is \mathbf{r} perpendicular to $\dot{\mathbf{p}}$: $\boxed{\dot{\mathbf{p}} \cdot \mathbf{r} = 0}$

$$\boxed{\dot{\mathbf{r}} \cdot \mathbf{p} = 0} = \begin{pmatrix} \dot{r}_x & \dot{r}_y \end{pmatrix} \cdot \begin{pmatrix} p_x \\ p_y \end{pmatrix} = \begin{pmatrix} -a \sin\phi & b \cos\phi \end{pmatrix} \cdot \begin{pmatrix} (1/a)\cos\phi \\ (1/b)\sin\phi \end{pmatrix} \text{ where: } \begin{matrix} \dot{r}_x = -a \sin\phi \\ \dot{r}_y = b \cos\phi \end{matrix} \text{ and: } \begin{matrix} p_x = (1/a)\cos\phi \\ p_y = (1/b)\sin\phi \end{matrix}$$

Elliptical

Exegesis

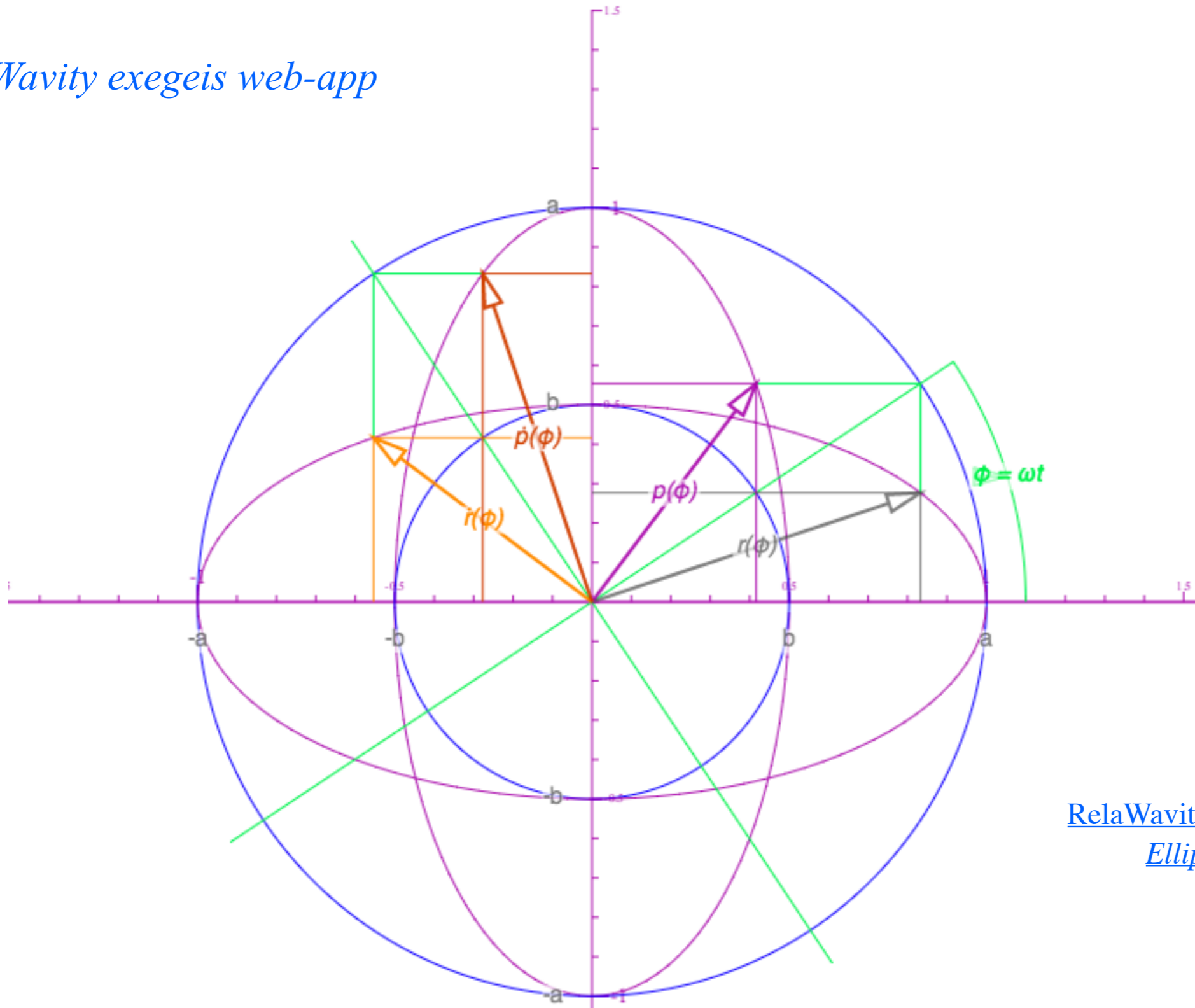
Controls

Contextual

Set ISM

User's Guide

RelaWavity exegeis web-app



RelaWavity Web Simulation
Ellipse/Exegesis

Geometry of dual ellipse Kepler anomalies for vectors $[\mathbf{r}(\phi), \mathbf{p}(\phi)]$ and $\frac{d}{dt}[\mathbf{r}(\phi), \mathbf{p}(\phi),]$ in coordinate (x,y) space rendered by animation web-app in RelaWavity and described in Lect. 12-advanced.

Introduction to dual matrix operator geometry (based on IHO orbits)

Quadratic form ellipse $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = 1$ vs. inverse form ellipse $\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p} = 1$

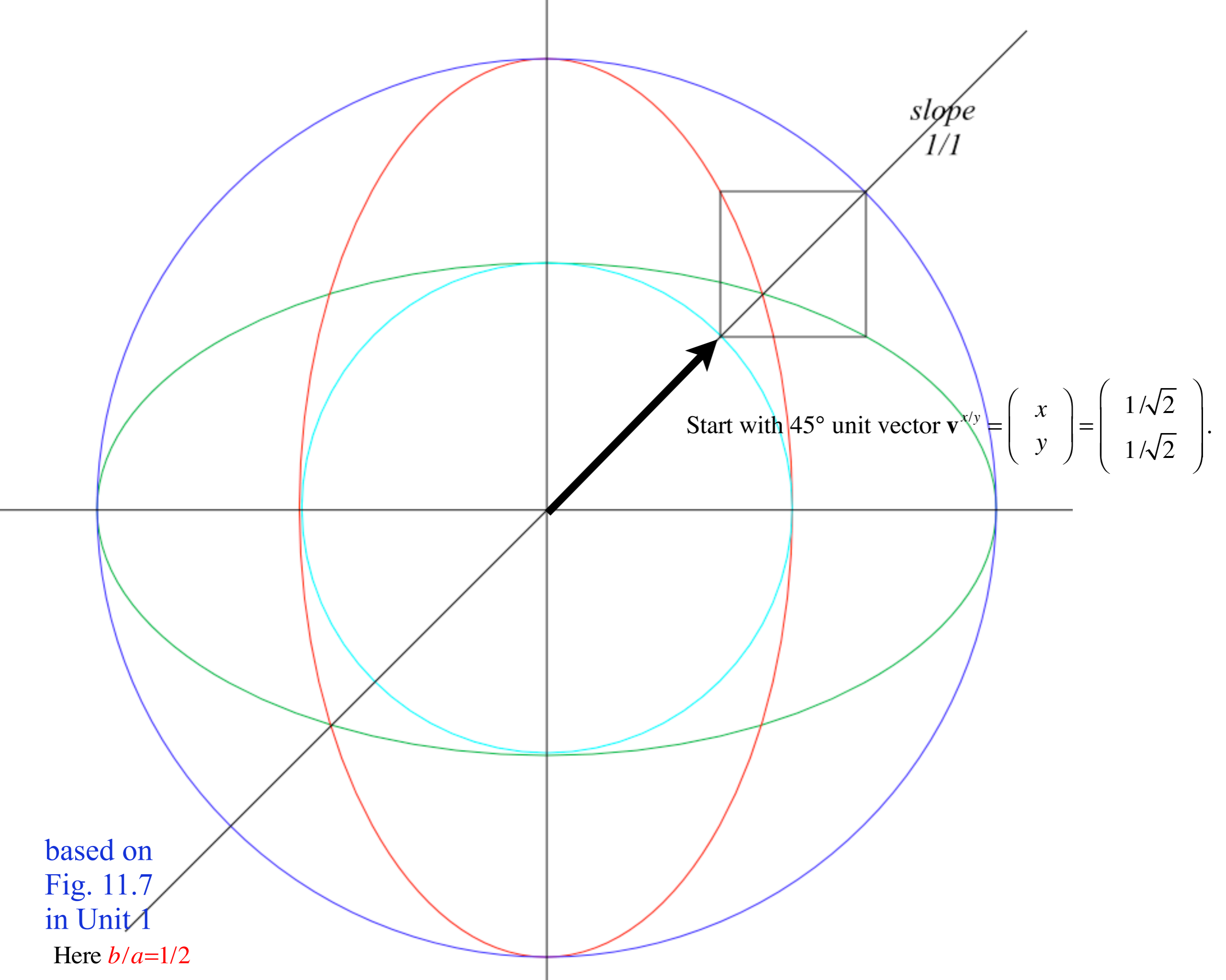
Duality norm relations ($\mathbf{r} \cdot \mathbf{p} = 1$)

\mathbf{Q} -Ellipse tangents \mathbf{r}' normal to dual \mathbf{Q}^{-1} -ellipse position \mathbf{p} ($\mathbf{r}' \cdot \mathbf{p} = 0 = \mathbf{r} \cdot \mathbf{p}'$)

 *Operator geometric sequences and eigenvectors*

Alternative scaling of matrix operator geometry

Vector calculus of tensor operation



Diagonal \mathbf{R} -matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $a/b=2$.

$$\mathbf{R} \cdot \mathbf{v}^{x/y} = \begin{pmatrix} 1/a & 0 \\ 0 & 1/b \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a \\ y/b \end{pmatrix}$$

(Slope increases if $a > b$.)

Action of "sqrt-" matrix $R = \sqrt{Q}$

slope a/b

slope $1/1$

slope b/a

Action of "sqrt⁻¹-" matrix $R^{-1} = \sqrt{Q^{-1}}$

Diagonal \mathbf{R}^{-1} -matrix acts on vector $\mathbf{v}^{x/y}$.

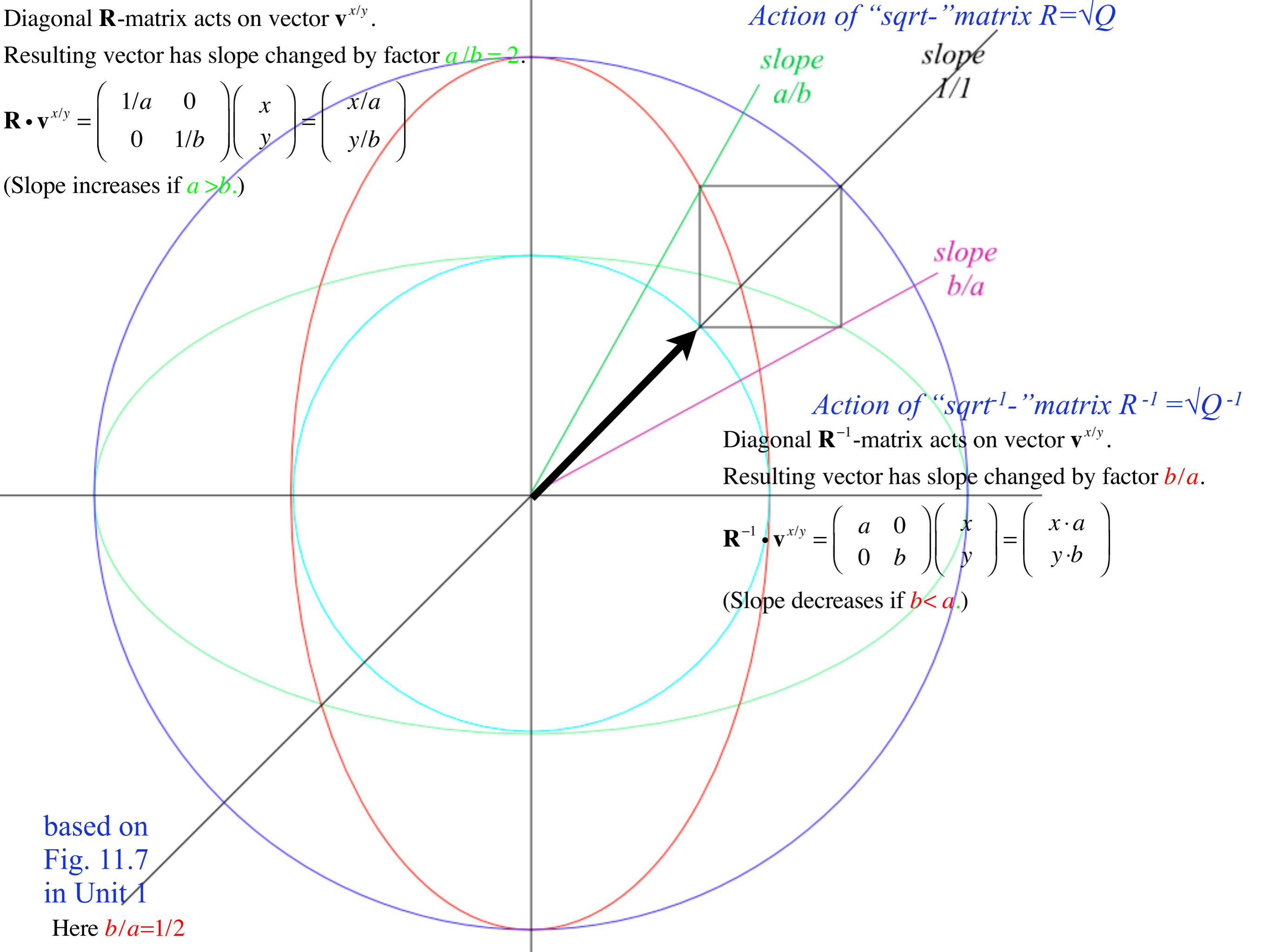
Resulting vector has slope changed by factor b/a .

$$\mathbf{R}^{-1} \cdot \mathbf{v}^{x/y} = \begin{pmatrix} a & 0 \\ 0 & b \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x \cdot a \\ y \cdot b \end{pmatrix}$$

(Slope decreases if $b < a$.)

based on
Fig. 11.7
in Unit 1

Here $b/a=1/2$



Diagonal \mathbf{R} -matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $a/b=2$.

$$\mathbf{R} \cdot \mathbf{v}^{x/y} = \begin{pmatrix} 1/a & 0 \\ 0 & 1/b \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a \\ y/b \end{pmatrix}$$

(It increases if $a > b$.)

Diagonal ($\mathbf{R}^2 = \mathbf{Q}$)-matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $a^2/b^2 = 4$.

$$\mathbf{Q} \cdot \mathbf{v}^{x/y} = \begin{pmatrix} 1/a^2 & 0 \\ 0 & 1/b^2 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a^2 \\ y/b^2 \end{pmatrix}$$

(It increases if $a > b$.)

Action of "sqrt-" matrix $R = \sqrt{Q}$

slope a^2/b^2

slope a/b

slope $1/1$

slope b/a

slope b^2/a^2

Action of "sqrt⁻¹-" matrix $R^{-1} = \sqrt{Q^{-1}}$

Diagonal \mathbf{R}^{-1} -matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $b/a=1/2$.

$$\mathbf{R}^{-1} \cdot \mathbf{v}^{x/y} = \begin{pmatrix} a & 0 \\ 0 & b \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x \cdot a \\ y \cdot b \end{pmatrix}$$

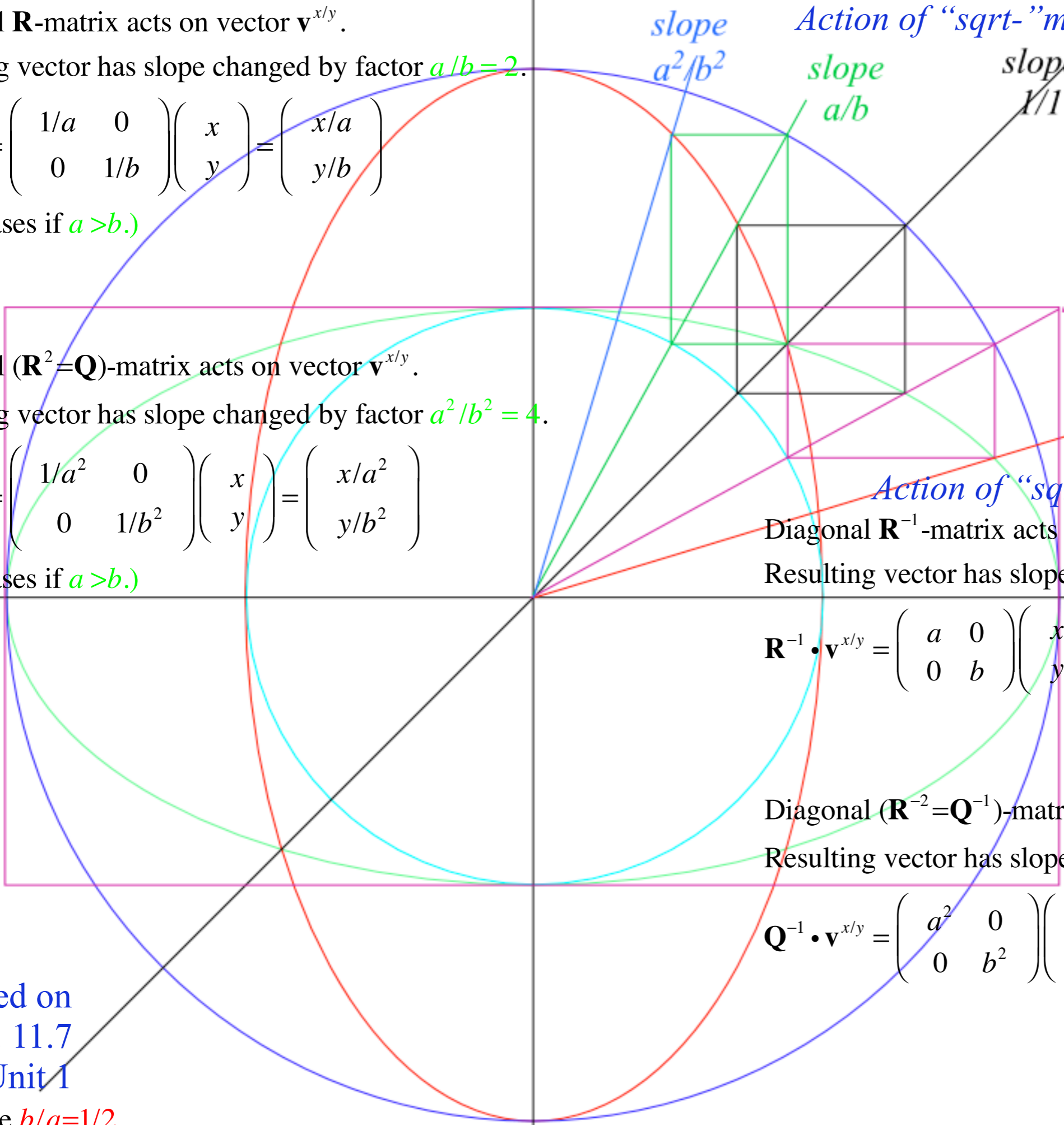
Diagonal ($\mathbf{R}^{-2} = \mathbf{Q}^{-1}$)-matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $b^2/a^2=1/4$.

$$\mathbf{Q}^{-1} \cdot \mathbf{v}^{x/y} = \begin{pmatrix} a^2 & 0 \\ 0 & b^2 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x \cdot a^2 \\ y \cdot b^2 \end{pmatrix}$$

based on
Fig. 11.7
in Unit 1

Here $b/a=1/2$



Diagonal \mathbf{R} -matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $a/b = 2$.

$$\mathbf{R} \cdot \mathbf{v}^{x/y} = \begin{pmatrix} 1/a & 0 \\ 0 & 1/b \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a \\ y/b \end{pmatrix}$$

(It increases if $a > b$.)

Diagonal ($\mathbf{R}^2 = \mathbf{Q}$)-matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $a^2/b^2 = 4$.

$$\mathbf{Q} \cdot \mathbf{v}^{x/y} = \begin{pmatrix} 1/a^2 & 0 \\ 0 & 1/b^2 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a^2 \\ y/b^2 \end{pmatrix}$$

(It increases if $a > b$.)

Either process can go on forever...

Diagonal ($\mathbf{R}^{2n} = \mathbf{Q}^n$)-matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $a^{2n}/b^{2n} = 4^n$.

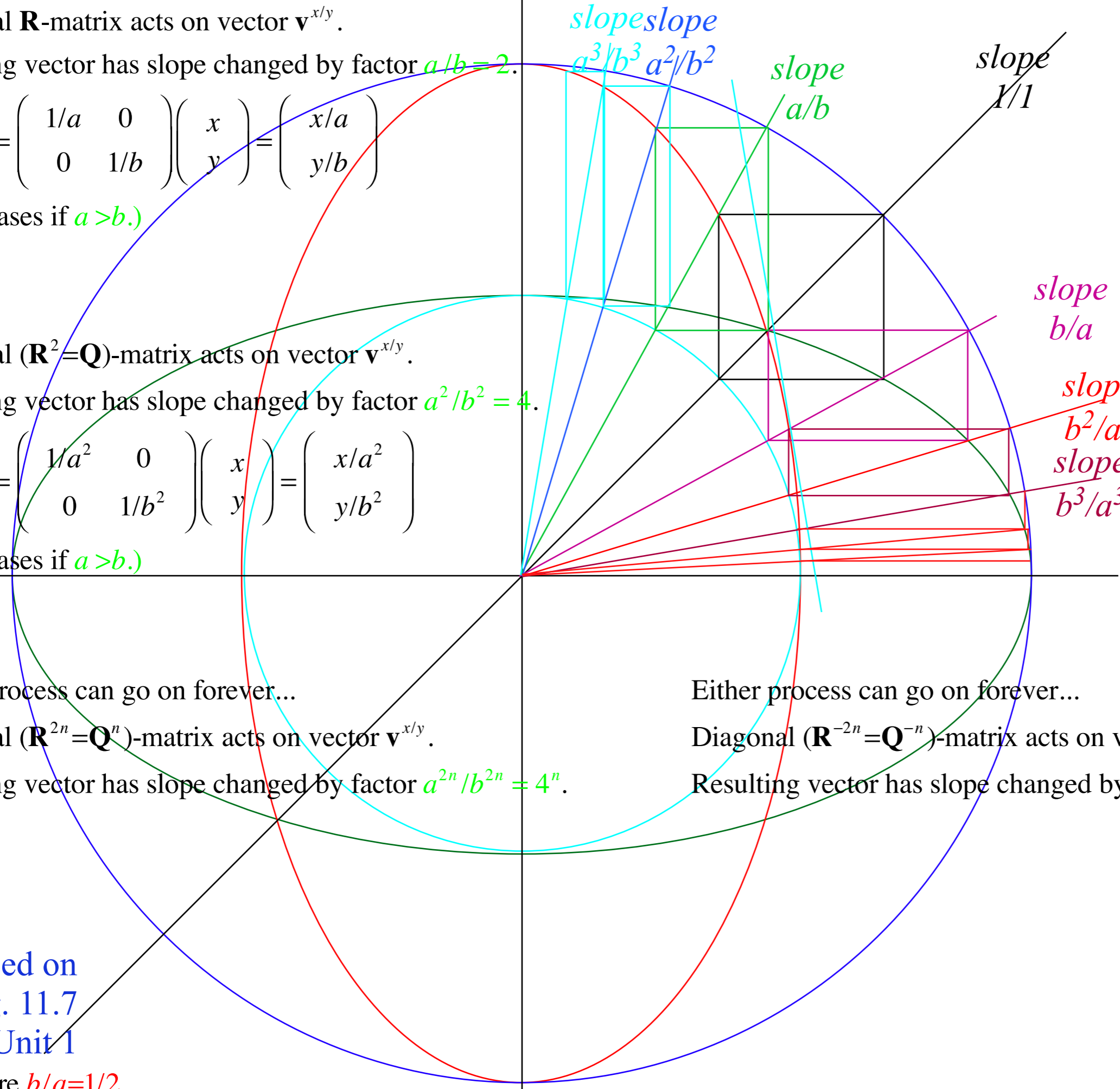
Either process can go on forever...

Diagonal ($\mathbf{R}^{-2n} = \mathbf{Q}^{-n}$)-matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $b^{2n}/a^{2n} = 4^{-n}$.

based on
Fig. 11.7
in Unit 1

Here $b/a = 1/2$



slopeslope
 a^3/b^3 a^2/b^2

slope
 $/a/b$

slope
 $1/1$

slope
 b/a

slope
 b^2/a^2

slope
 b^3/a^3

Diagonal \mathbf{R} -matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $a/b = 2$.

$$\mathbf{R} \cdot \mathbf{v}^{x/y} = \begin{pmatrix} 1/a & 0 \\ 0 & 1/b \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a \\ y/b \end{pmatrix}$$

(It increases if $a > b$.)

EIGENVECTOR

Diagonal ($\mathbf{R}^2 = \mathbf{Q}$)-matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $a^2/b^2 = 4$.

$$\mathbf{Q} \cdot \mathbf{v}^{x/y} = \begin{pmatrix} 1/a^2 & 0 \\ 0 & 1/b^2 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a^2 \\ y/b^2 \end{pmatrix}$$

(It increases if $a > b$.)

EIGENVECTOR

Either process can go on forever...

Diagonal ($\mathbf{R}^{2^n} = \mathbf{Q}^n$)-matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $a^{2^n}/b^{2^n} = 4^n$.

...Finally, the result approaches **EIGENVECTOR** $|y\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$

of ∞ -slope which is "immune" to \mathbf{R} , \mathbf{Q} or \mathbf{Q}^n :

$$\mathbf{R}|y\rangle = (1/b)|y\rangle \quad \mathbf{Q}^n|y\rangle = (1/b^2)^n|y\rangle$$

Here $b/a = 1/2$

slopeslope

a^3/b^3 a^2/b^2

slope
 $/a/b$

slope
 $1/1$

slope
 b/a

slope
 b^2/a^2

slope
 b^3/a^3

Either process can go on forever...

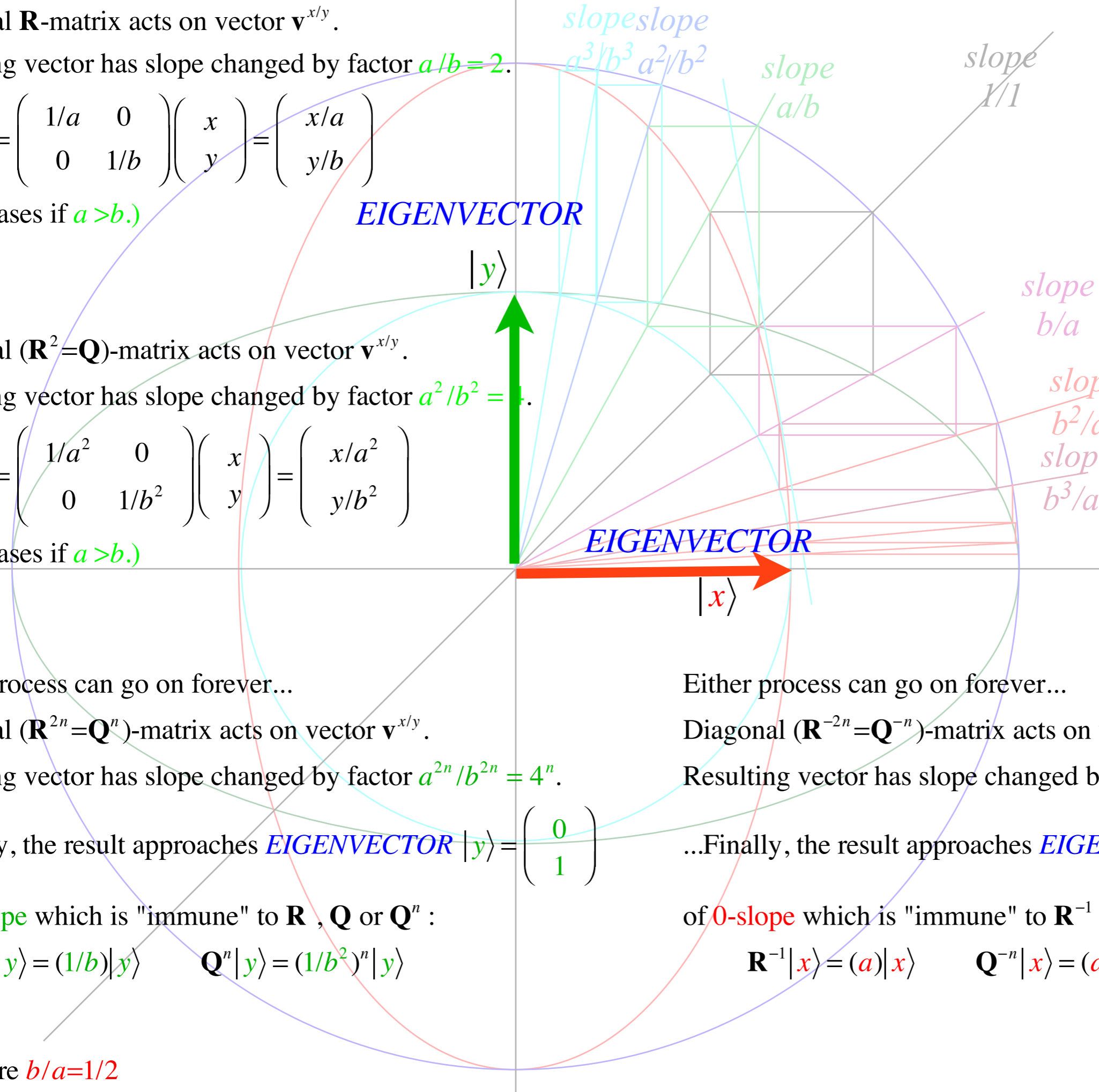
Diagonal ($\mathbf{R}^{-2^n} = \mathbf{Q}^{-n}$)-matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $b^{2^n}/a^{2^n} = 4^{-n}$.

...Finally, the result approaches **EIGENVECTOR** $|x\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$

of 0-slope which is "immune" to \mathbf{R}^{-1} , \mathbf{Q}^{-1} or \mathbf{Q}^{-n} :

$$\mathbf{R}^{-1}|x\rangle = (a)|x\rangle \quad \mathbf{Q}^{-n}|x\rangle = (a^2)^n|x\rangle$$



Diagonal \mathbf{R} -matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $a/b=2$.

$$\mathbf{R} \cdot \mathbf{v}^{x/y} = \begin{pmatrix} 1/a & 0 \\ 0 & 1/b \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a \\ y/b \end{pmatrix}$$

(It increases if $a > b$.)

EIGENVECTOR

Diagonal ($\mathbf{R}^2=\mathbf{Q}$)-matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $a^2/b^2=4$.

$$\mathbf{Q} \cdot \mathbf{v}^{x/y} = \begin{pmatrix} 1/a^2 & 0 \\ 0 & 1/b^2 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a^2 \\ y/b^2 \end{pmatrix}$$

(It increases if $a > b$.)

EIGENVECTOR

Either process can go on forever...

Diagonal ($\mathbf{R}^{2^n}=\mathbf{Q}^n$)-matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $a^{2^n}/b^{2^n}=4^n$.

...Finally, the result approaches **EIGENVECTOR** $|y\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$

of ∞ -slope which is "immune" to \mathbf{R} , \mathbf{Q} or \mathbf{Q}^n :

$$\mathbf{R}|y\rangle = (1/b)|y\rangle \quad \mathbf{Q}^n|y\rangle = (1/b^2)^n|y\rangle$$

Eigenvalues

Eigensolution Relations

Either process can go on forever...

Diagonal ($\mathbf{R}^{-2^n}=\mathbf{Q}^{-n}$)-matrix acts on vector $\mathbf{v}^{x/y}$.

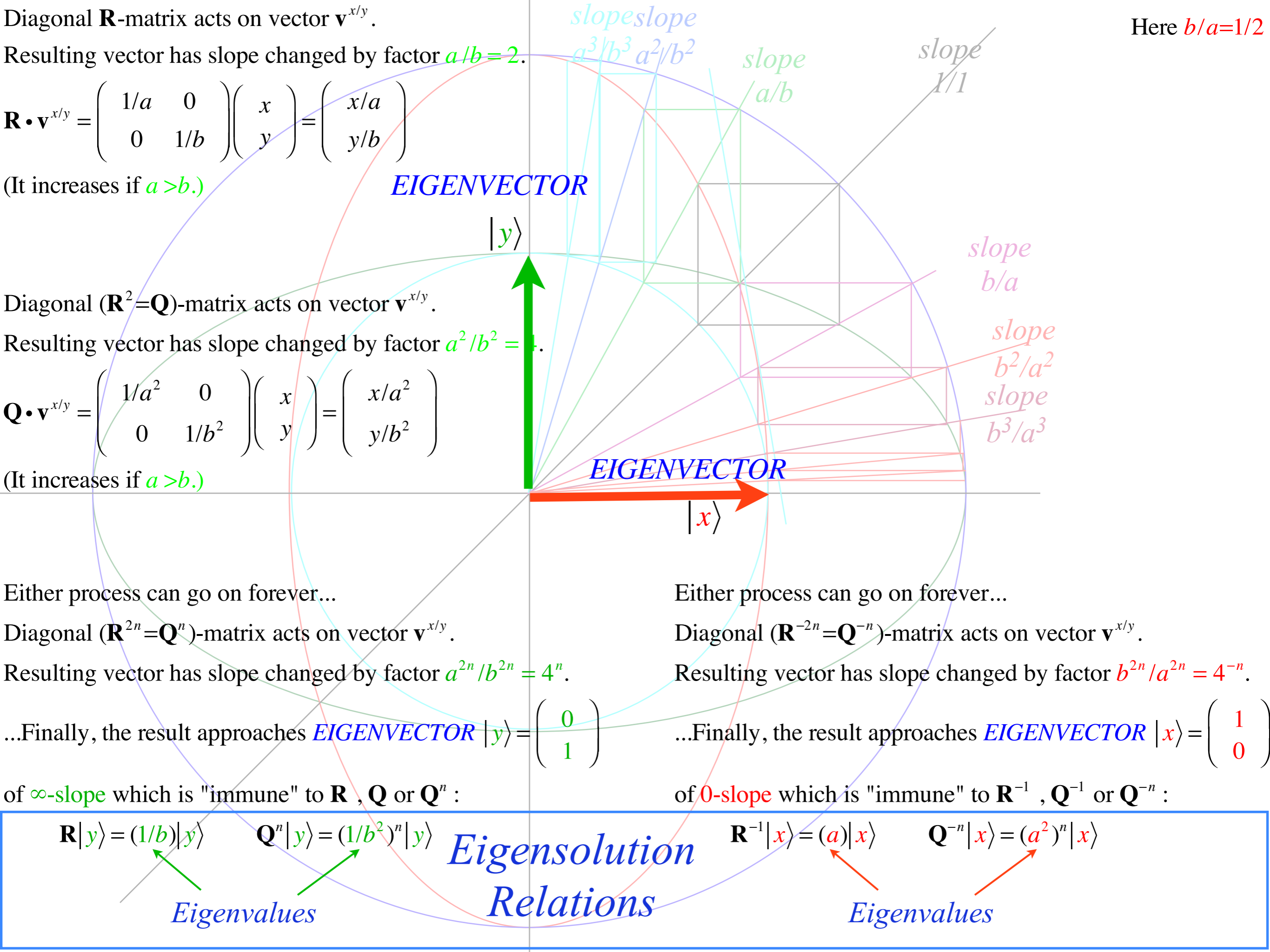
Resulting vector has slope changed by factor $b^{2^n}/a^{2^n}=4^{-n}$.

...Finally, the result approaches **EIGENVECTOR** $|x\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$

of 0-slope which is "immune" to \mathbf{R}^{-1} , \mathbf{Q}^{-1} or \mathbf{Q}^{-n} :

$$\mathbf{R}^{-1}|x\rangle = (a)|x\rangle \quad \mathbf{Q}^{-n}|x\rangle = (a^2)^n|x\rangle$$

Eigenvalues



Introduction to dual matrix operator geometry (based on IHO orbits)

Quadratic form ellipse $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = 1$ vs. inverse form ellipse $\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p} = 1$

Duality norm relations ($\mathbf{r} \cdot \mathbf{p} = 1$)

\mathbf{Q} -Ellipse tangents \mathbf{r}' normal to dual \mathbf{Q}^{-1} -ellipse position \mathbf{p} ($\mathbf{r}' \cdot \mathbf{p} = 0 = \mathbf{r} \cdot \mathbf{p}'$)

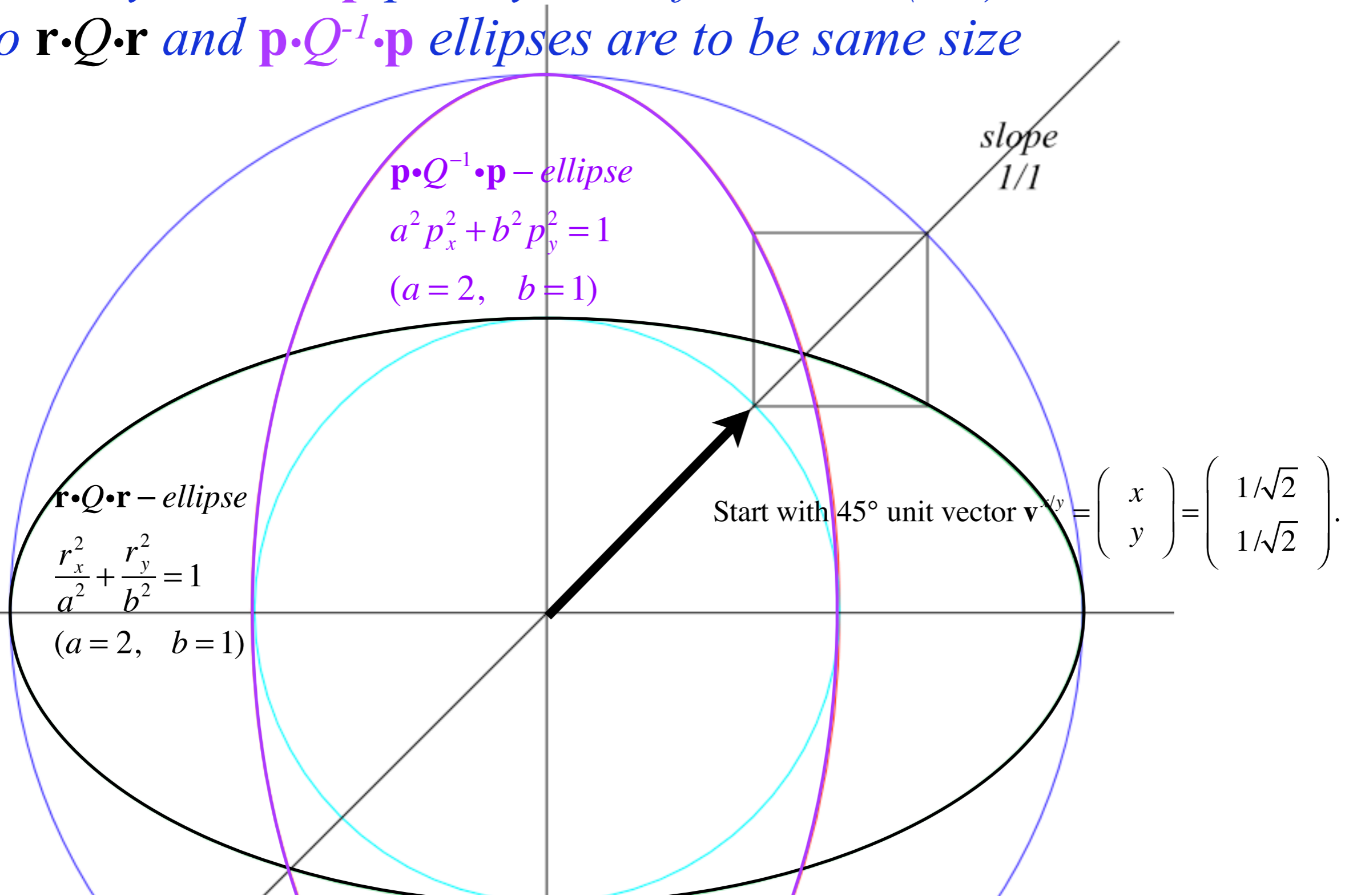
Operator geometric sequences and eigenvectors



Alternative scaling of matrix operator geometry

Vector calculus of tensor operation

You may rescale **p**-plot by scale factor $S=(a \cdot b)$
 so $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r}$ and $\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p}$ ellipses are to be same size



Here plot of **p**-ellipse is re-scaled by scalefactor $S=a \cdot b$

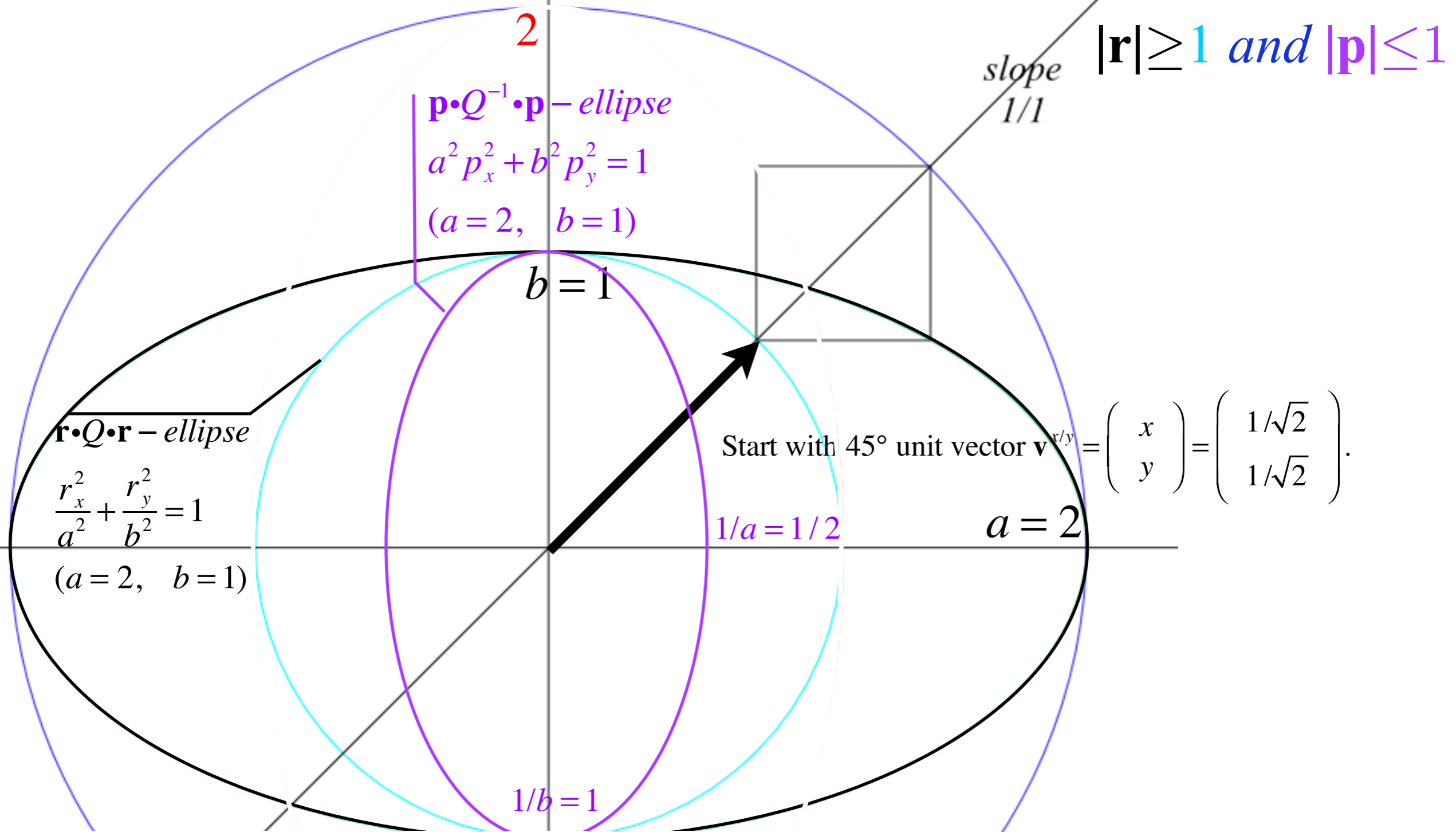
p-ellipse x -radius= $1/a$ plotted at: $S(1/a)=b$ ($=1$ for $a=2, b=1$)

p-ellipse y -radius= $1/b$ plotted at: $S(1/b)=a$ ($=2$ for $a=2, b=1$)

..or else rescale **p**-plot by scale factor $S=b$

Here $b/a=1/2$

to separate $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r}$ and $\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p}$ ellipses into different regions



Here plot of **p**-ellipse is re-scaled by scalefactor $S=b$

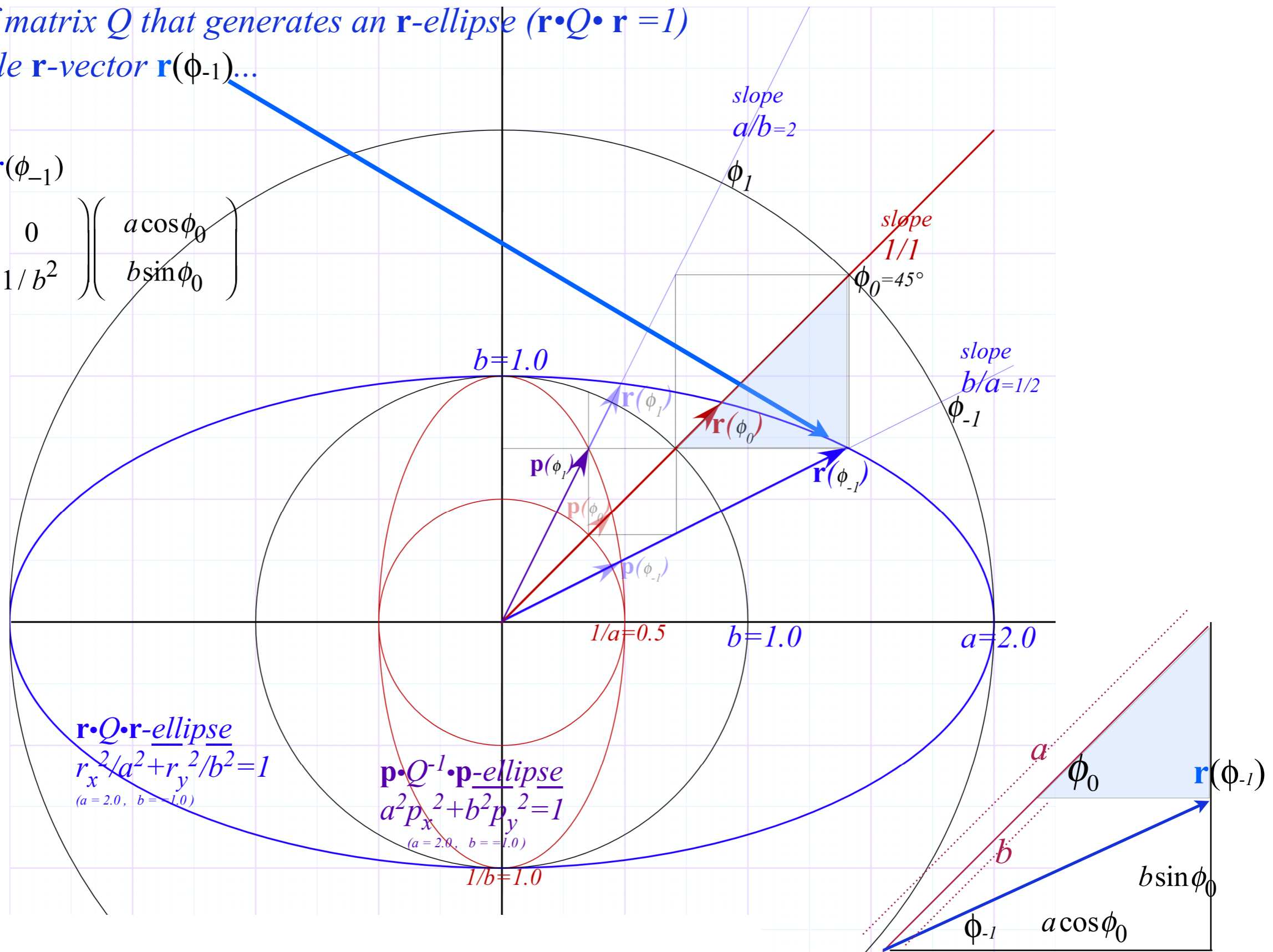
p-ellipse x -radius= $1/a$ plotted at: $S(1/a)=b/a$ ($=1/2$ for $a=2, b=1$)

p-ellipse y -radius= $1/b$ plotted at: $S(1/b)=1$

Action of matrix Q that generates an \mathbf{r} -ellipse ($\mathbf{r} \cdot Q \cdot \mathbf{r} = 1$) on a single \mathbf{r} -vector $\mathbf{r}(\phi_{-1})$...

$$\mathbf{p}(\phi_1) = \mathbf{Q} \cdot \mathbf{r}(\phi_{-1})$$

$$= \begin{pmatrix} 1/a^2 & 0 \\ 0 & 1/b^2 \end{pmatrix} \begin{pmatrix} a \cos \phi_0 \\ b \sin \phi_0 \end{pmatrix}$$



Variation of Fig. 11.7 in Unit 1

$\mathbf{r} \cdot Q \cdot \mathbf{r}$ -ellipse
 $r_x^2/a^2 + r_y^2/b^2 = 1$
 (a = 2.0, b = 1.0)

$\mathbf{p} \cdot Q^{-1} \cdot \mathbf{p}$ -ellipse
 $a^2 p_x^2 + b^2 p_y^2 = 1$
 (a = 2.0, b = 1.0)

Here plot of \mathbf{p} -ellipse is re-scaled by scalefactor $S=b$

\mathbf{p} -ellipse x-radius= $1/a$ plotted at: $S(1/a)=b/a$ (=1/2 for $a=2, b=1$)

\mathbf{p} -ellipse y-radius= $1/b$ plotted at: $S(1/b)=1$

Action of matrix Q that generates an \mathbf{r} -ellipse ($\mathbf{r} \cdot Q \cdot \mathbf{r} = 1$)

on a single \mathbf{r} -vector $\mathbf{r}(\phi_{-1})$... is to rotate it to a new vector \mathbf{p} on the \mathbf{p} -ellipse ($\mathbf{p} \cdot Q^{-1} \cdot \mathbf{p} = 1$), that is, $Q \cdot \mathbf{r}(\phi_{-1}) = \mathbf{p}(\phi_{+1})$

$$\mathbf{p}(\phi_1) = Q \cdot \mathbf{r}(\phi_{-1})$$

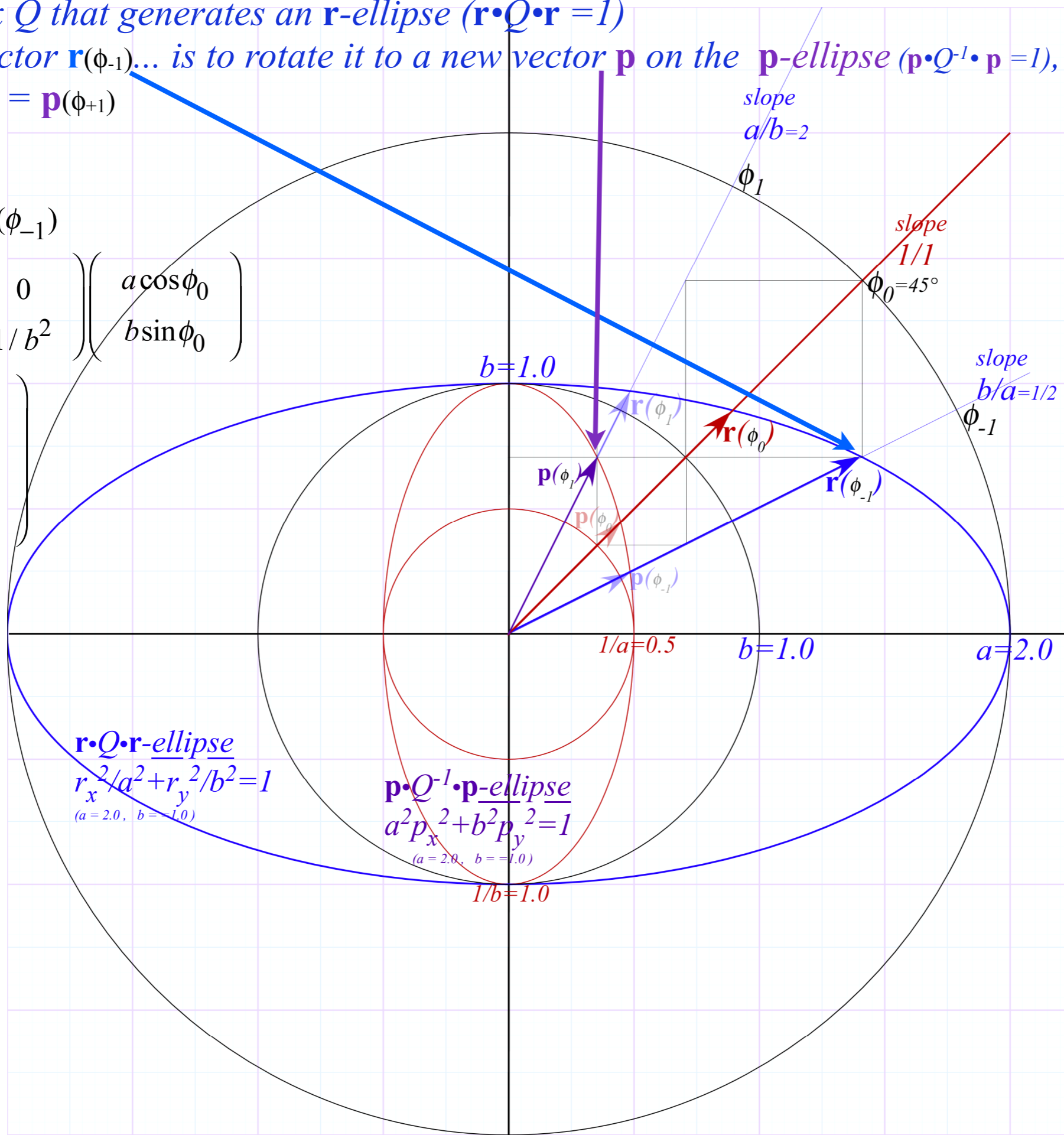
$$= \begin{pmatrix} 1/a^2 & 0 \\ 0 & 1/b^2 \end{pmatrix} \begin{pmatrix} a \cos \phi_0 \\ b \sin \phi_0 \end{pmatrix}$$

$$= \begin{pmatrix} \frac{1}{a} \cos \phi_0 \\ \frac{1}{b} \sin \phi_0 \end{pmatrix}$$

$$= \begin{pmatrix} \frac{1}{2} & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & \frac{1}{2} \end{pmatrix}$$

$\mathbf{r} \cdot Q \cdot \mathbf{r}$ -ellipse
 $r_x^2/a^2 + r_y^2/b^2 = 1$
 ($a = 2.0, b = 1.0$)

$\mathbf{p} \cdot Q^{-1} \cdot \mathbf{p}$ -ellipse
 $a^2 p_x^2 + b^2 p_y^2 = 1$
 ($a = 2.0, b = 1.0$)



Variation of
 Fig. 11.7
 in Unit 1

Action of matrix Q that generates an \mathbf{r} -ellipse ($\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = 1$) on a single \mathbf{r} -vector $\mathbf{r}(\phi_{-1}) \dots$ is to rotate it to a new vector \mathbf{p} on the \mathbf{p} -ellipse ($\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p} = 1$), that is, $\mathbf{Q} \cdot \mathbf{r}(\phi_{-1}) = \mathbf{p}(\phi_{+1})$

$$\mathbf{p}(\phi_1) = \mathbf{Q} \cdot \mathbf{r}(\phi_{-1})$$

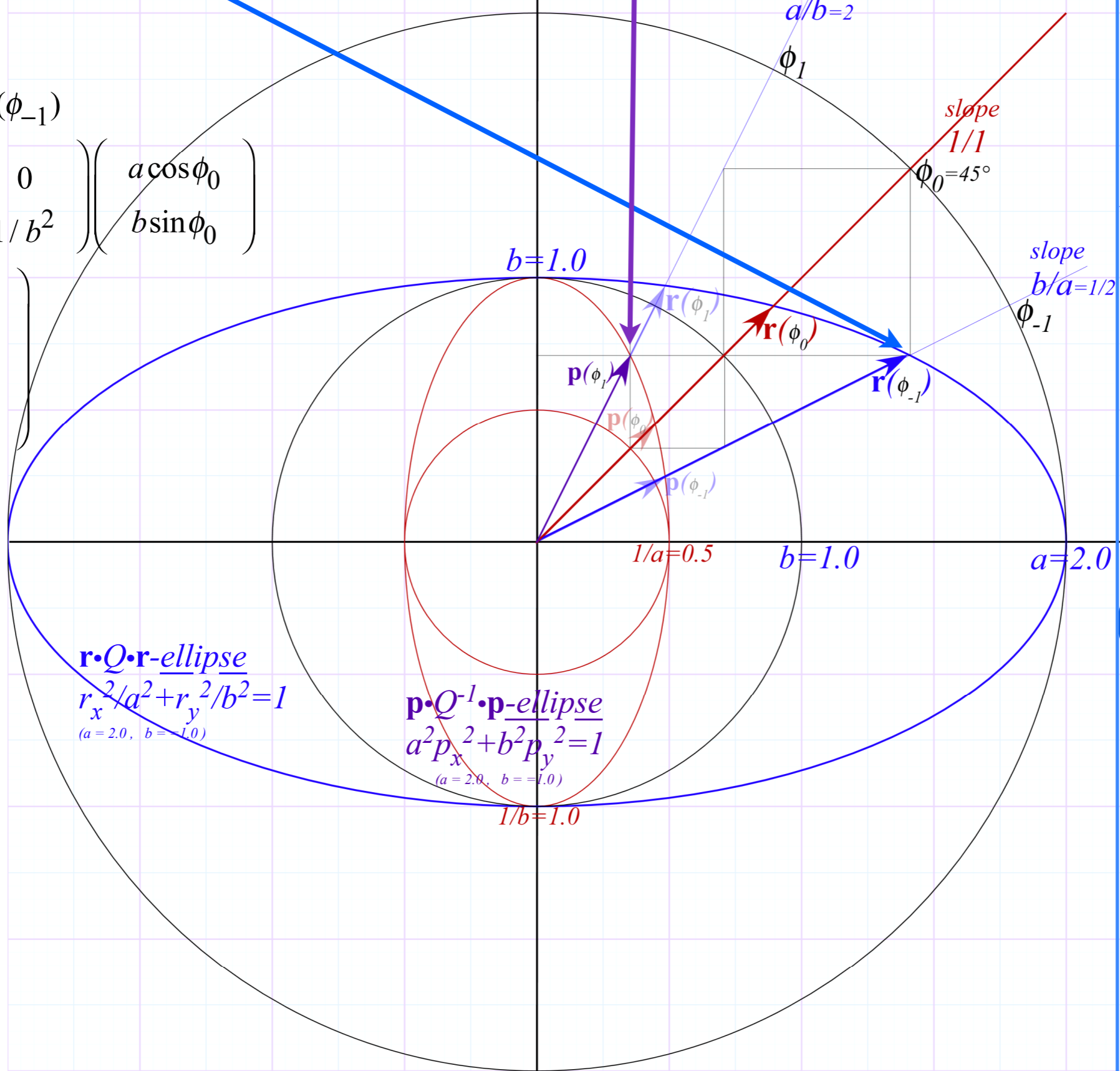
$$= \begin{pmatrix} 1/a^2 & 0 \\ 0 & 1/b^2 \end{pmatrix} \begin{pmatrix} a \cos \phi_0 \\ b \sin \phi_0 \end{pmatrix}$$

$$= \begin{pmatrix} \frac{1}{a} \cos \phi_0 \\ \frac{1}{b} \sin \phi_0 \end{pmatrix}$$

$$= \begin{pmatrix} \frac{1}{2} & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & \frac{1}{2} \end{pmatrix}$$

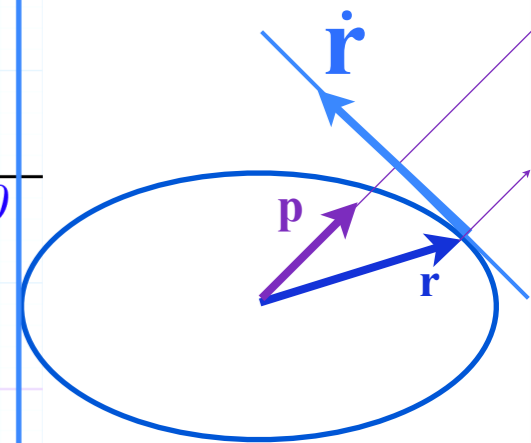
$\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r}$ -ellipse
 $r_x^2/a^2 + r_y^2/b^2 = 1$
 ($a = 2.0, b = 1.0$)

$\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p}$ -ellipse
 $a^2 p_x^2 + b^2 p_y^2 = 1$
 ($a = 2.0, b = 1.0$)



Key points of matrix geometry:

Matrix Q maps any vector \mathbf{r} to a new vector \mathbf{p} normal to the tangent $\dot{\mathbf{r}}$ to its $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r}$ -ellipse.



Variation of Fig. 11.7 in Unit 1

Action of matrix Q that generates an \mathbf{r} -ellipse ($\mathbf{r} \cdot Q \cdot \mathbf{r} = 1$) on a single \mathbf{r} -vector $\mathbf{r}(\phi_{-1})$... is to rotate it to a new vector \mathbf{p} on the \mathbf{p} -ellipse ($\mathbf{p} \cdot Q^{-1} \cdot \mathbf{p} = 1$), that is, $Q \cdot \mathbf{r}(\phi_{-1}) = \mathbf{p}(\phi_{+1})$

$$\mathbf{p}(\phi_1) = Q \cdot \mathbf{r}(\phi_{-1})$$

$$= \begin{pmatrix} 1/a^2 & 0 \\ 0 & 1/b^2 \end{pmatrix} \begin{pmatrix} a \cos \phi_0 \\ b \sin \phi_0 \end{pmatrix}$$

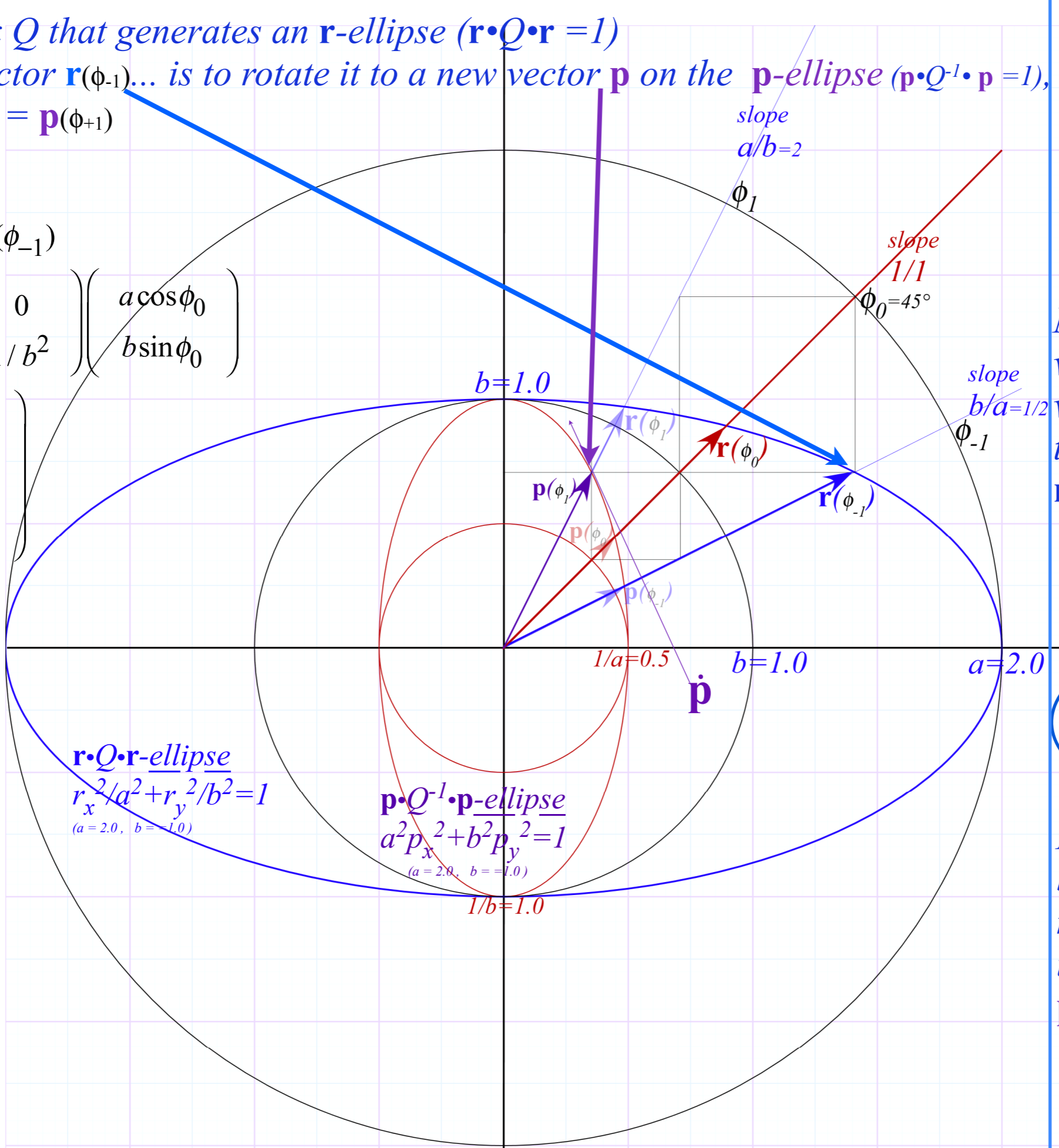
$$= \begin{pmatrix} \frac{1}{a} \cos \phi_0 \\ \frac{1}{b} \sin \phi_0 \end{pmatrix}$$

$$= \begin{pmatrix} \frac{1}{2} & \frac{1}{\sqrt{2}} \\ \frac{1}{1} & \frac{1}{\sqrt{2}} \end{pmatrix}$$

$\mathbf{r} \cdot Q \cdot \mathbf{r}$ -ellipse
 $r_x^2/a^2 + r_y^2/b^2 = 1$
 (a = 2.0, b = 1.0)

$\mathbf{p} \cdot Q^{-1} \cdot \mathbf{p}$ -ellipse
 $a^2 p_x^2 + b^2 p_y^2 = 1$
 (a = 2.0, b = 1.0)

Variation of Fig. 11.7 in Unit 1



Key points of matrix geometry:

Matrix Q maps any vector \mathbf{r} to a new vector \mathbf{p} normal to the tangent $\dot{\mathbf{r}}$ to its $\mathbf{r} \cdot Q \cdot \mathbf{r}$ -ellipse.

Matrix Q^{-1} maps \mathbf{p} back to \mathbf{r} that is normal to the tangent $\dot{\mathbf{p}}$ to its $\mathbf{p} \cdot Q^{-1} \cdot \mathbf{p}$ -ellipse.

Introduction to dual matrix operator geometry (based on IHO orbits)

Quadratic form ellipse $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = 1$ vs. inverse form ellipse $\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p} = 1$

Duality norm relations ($\mathbf{r} \cdot \mathbf{p} = 1$)

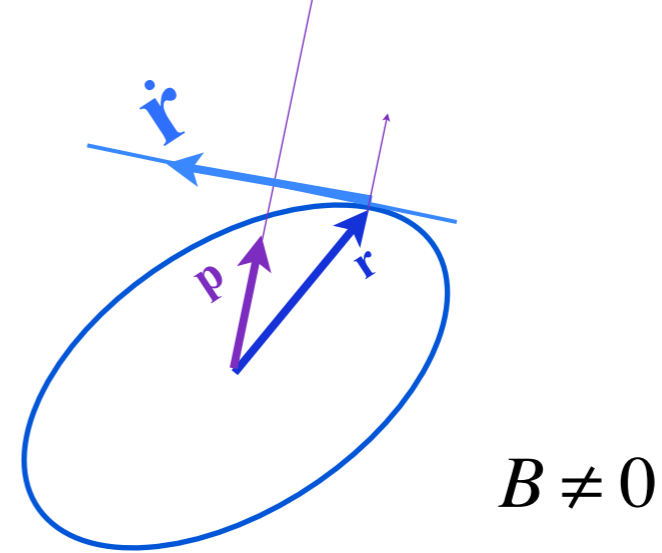
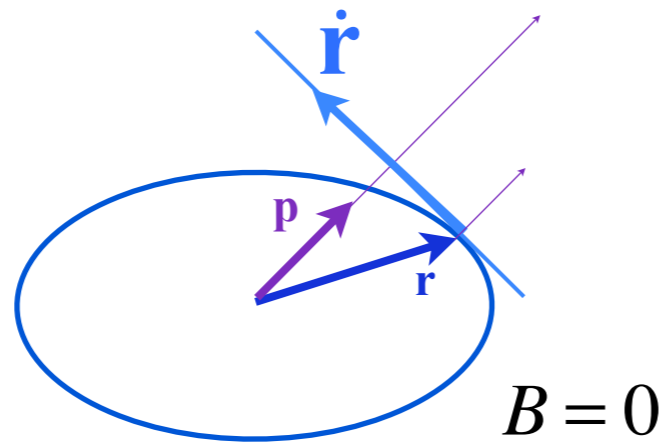
\mathbf{Q} -Ellipse tangents \mathbf{r}' normal to dual \mathbf{Q}^{-1} -ellipse position \mathbf{p} ($\mathbf{r}' \cdot \mathbf{p} = 0 = \mathbf{r} \cdot \mathbf{p}'$)

Operator geometric sequences and eigenvectors

Alternative scaling of matrix operator geometry



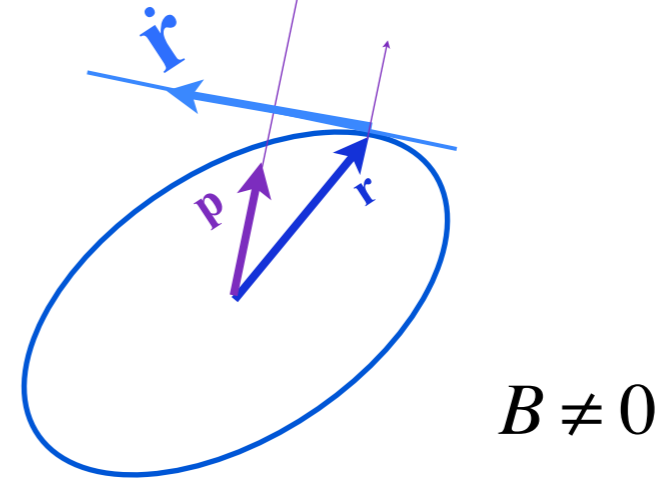
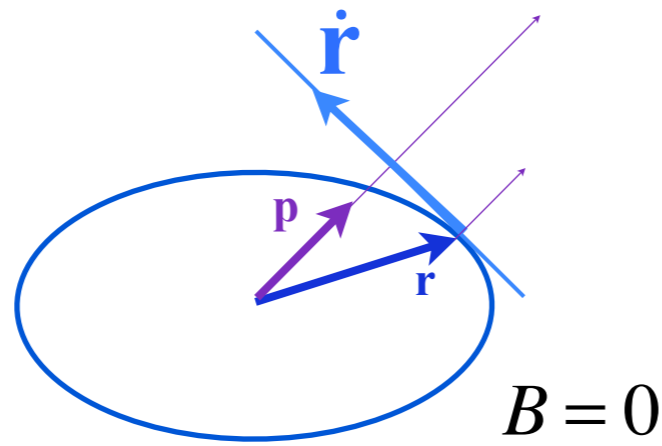
Vector calculus of tensor operation



Derive matrix “normal-to-ellipse” geometry by vector calculus:

Let matrix $Q = \begin{pmatrix} A & B \\ B & D \end{pmatrix}$

define the ellipse $1 = \mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = \begin{pmatrix} x & y \end{pmatrix} \cdot \begin{pmatrix} A & B \\ B & D \end{pmatrix} \cdot \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x & y \end{pmatrix} \cdot \begin{pmatrix} A \cdot x + B \cdot y \\ B \cdot x + D \cdot y \end{pmatrix} = A \cdot x^2 + 2B \cdot xy + D \cdot y^2 = 1$



Derive matrix “normal-to-ellipse” geometry by vector calculus:

Let matrix $Q = \begin{pmatrix} A & B \\ B & D \end{pmatrix}$

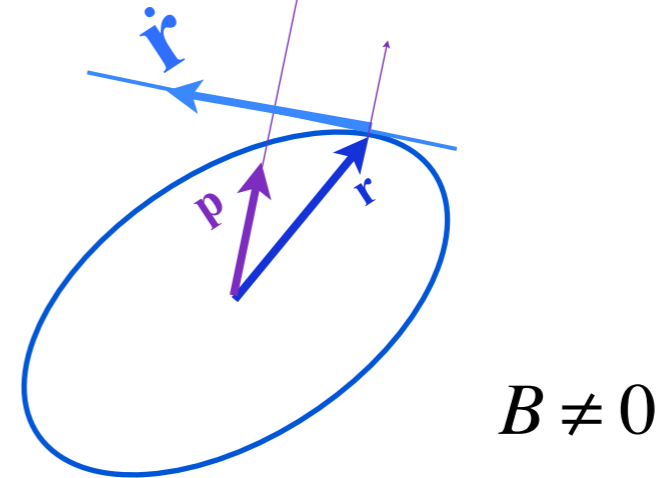
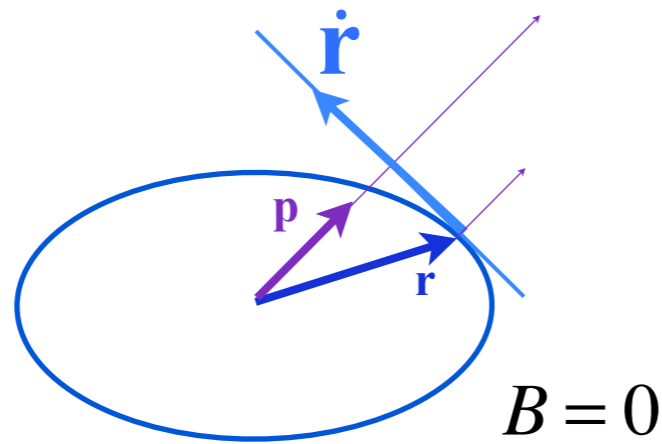
define the ellipse $1 = \mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = \begin{pmatrix} x & y \end{pmatrix} \cdot \begin{pmatrix} A & B \\ B & D \end{pmatrix} \cdot \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x & y \end{pmatrix} \cdot \begin{pmatrix} A \cdot x + B \cdot y \\ B \cdot x + D \cdot y \end{pmatrix} = A \cdot x^2 + 2B \cdot xy + D \cdot y^2 = 1$

Compare operation by Q on vector \mathbf{r} with vector derivative or gradient of $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r}$

$$\begin{pmatrix} A & B \\ B & D \end{pmatrix} \cdot \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} A \cdot x + B \cdot y \\ B \cdot x + D \cdot y \end{pmatrix}$$

$$\frac{\partial}{\partial \mathbf{r}} (\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r}) = \nabla (\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r})$$

$$\begin{pmatrix} \frac{\partial}{\partial x} \\ \frac{\partial}{\partial y} \end{pmatrix} (A \cdot x^2 + 2B \cdot xy + D \cdot y^2) = \begin{pmatrix} 2A \cdot x + 2B \cdot y \\ 2B \cdot x + 2D \cdot y \end{pmatrix}$$



Derive matrix “normal-to-ellipse” geometry by vector calculus:

Let matrix $Q = \begin{pmatrix} A & B \\ B & D \end{pmatrix}$

define the ellipse $1 = \mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = \begin{pmatrix} x & y \end{pmatrix} \cdot \begin{pmatrix} A & B \\ B & D \end{pmatrix} \cdot \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x & y \end{pmatrix} \cdot \begin{pmatrix} A \cdot x + B \cdot y \\ B \cdot x + D \cdot y \end{pmatrix} = A \cdot x^2 + 2B \cdot xy + D \cdot y^2 = 1$

Compare operation by Q on vector \mathbf{r} with vector derivative or gradient of $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r}$

$$\begin{pmatrix} A & B \\ B & D \end{pmatrix} \cdot \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} A \cdot x + B \cdot y \\ B \cdot x + D \cdot y \end{pmatrix} \quad \frac{\partial}{\partial \mathbf{r}} (\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r}) = \nabla (\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r})$$

$$\begin{pmatrix} \frac{\partial}{\partial x} \\ \frac{\partial}{\partial y} \end{pmatrix} (A \cdot x^2 + 2B \cdot xy + D \cdot y^2) = \begin{pmatrix} 2A \cdot x + 2B \cdot y \\ 2B \cdot x + 2D \cdot y \end{pmatrix}$$

Very simple result:

$$\frac{\partial}{\partial \mathbf{r}} \left(\frac{\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r}}{2} \right) = \nabla \left(\frac{\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r}}{2} \right) = \mathbf{Q} \cdot \mathbf{r}$$

Introduction to dual matrix operator geometry (based on IHO orbits)

Quadratic form ellipse $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = 1$ vs. inverse form ellipse $\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p} = 1$

Duality norm relations ($\mathbf{r} \cdot \mathbf{p} = 1$)

Q-Ellipse tangents \mathbf{r}' normal to dual \mathbf{Q}^{-1} -ellipse position \mathbf{p} ($\mathbf{r}' \cdot \mathbf{p} = 0 = \mathbf{r} \cdot \mathbf{p}'$)

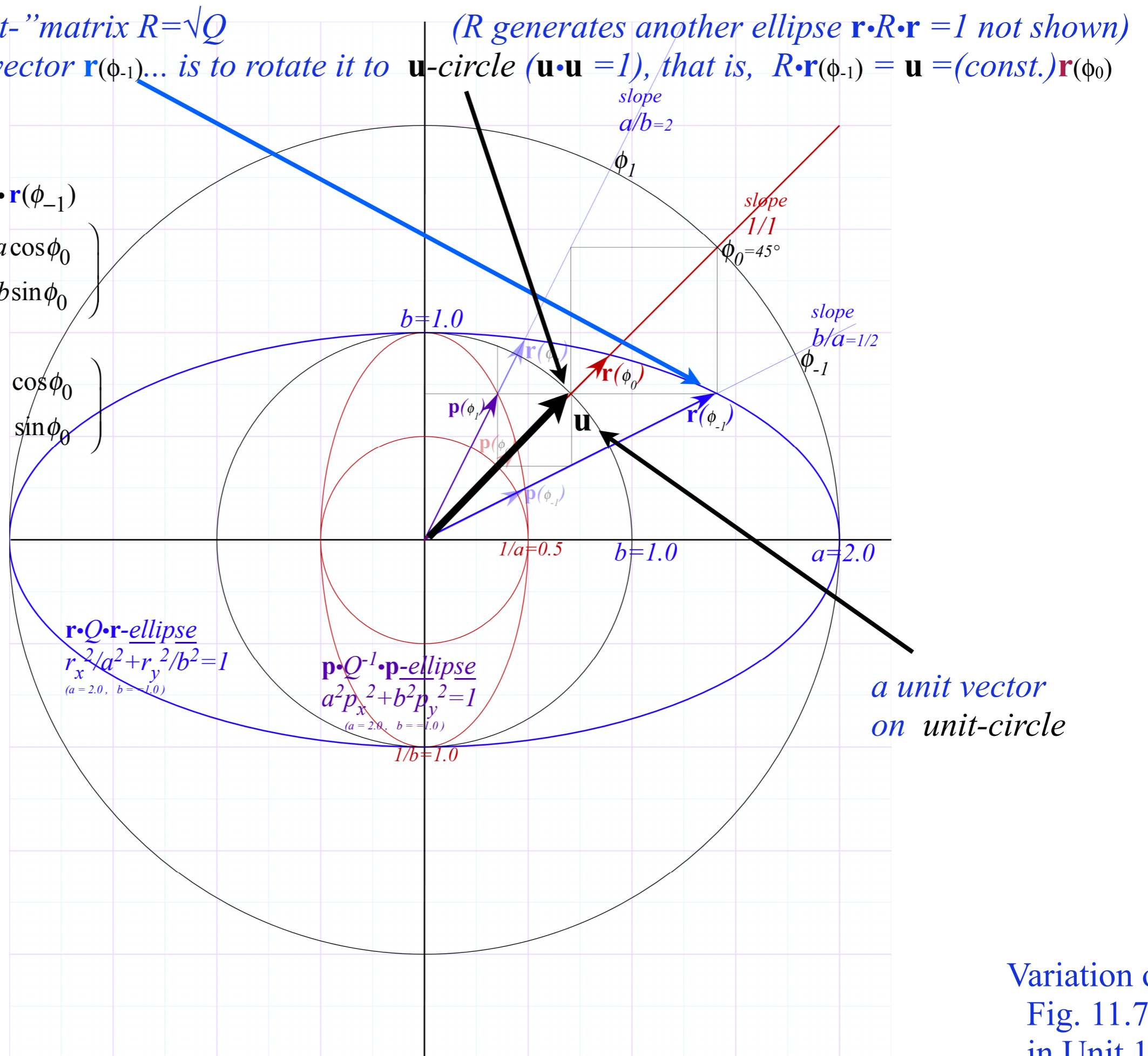
(Still more) Operator geometric sequences and eigenvectors

Alternative scaling of matrix operator geometry

Vector calculus of tensor operation

Action of "sqrt-" matrix $R=\sqrt{Q}$ (R generates another ellipse $\mathbf{r}\cdot R\cdot\mathbf{r} = 1$ not shown) on a single \mathbf{r} -vector $\mathbf{r}(\phi_{-1})$... is to rotate it to \mathbf{u} -circle ($\mathbf{u}\cdot\mathbf{u} = 1$), that is, $R\cdot\mathbf{r}(\phi_{-1}) = \mathbf{u} = (\text{const.})\mathbf{r}(\phi_0)$

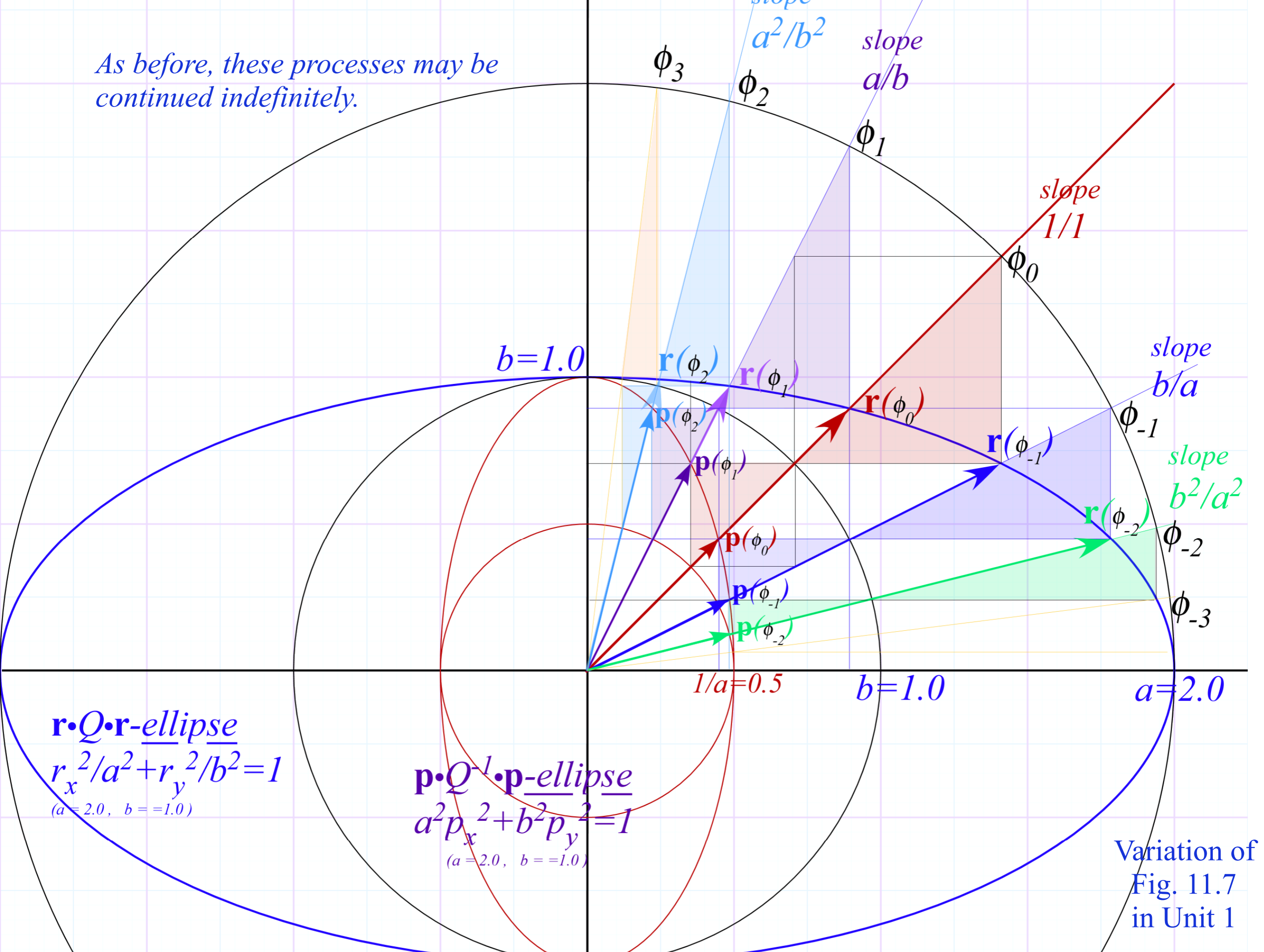
$$\begin{aligned} \mathbf{u} &= \sqrt{\mathbf{Q}} \cdot \mathbf{r}(\phi_{-1}) = \mathbf{R} \cdot \mathbf{r}(\phi_{-1}) \\ &= \begin{pmatrix} 1/a & 0 \\ 0 & 1/b \end{pmatrix} \begin{pmatrix} a \cos \phi_0 \\ b \sin \phi_0 \end{pmatrix} \\ &= \begin{pmatrix} \frac{1}{a} a \cos \phi_0 \\ \frac{1}{b} b \sin \phi_0 \end{pmatrix} = \begin{pmatrix} \cos \phi_0 \\ \sin \phi_0 \end{pmatrix} \\ &= \begin{pmatrix} \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} \end{pmatrix} \end{aligned}$$



a unit vector on unit-circle

Variation of Fig. 11.7 in Unit 1

As before, these processes may be continued indefinitely.

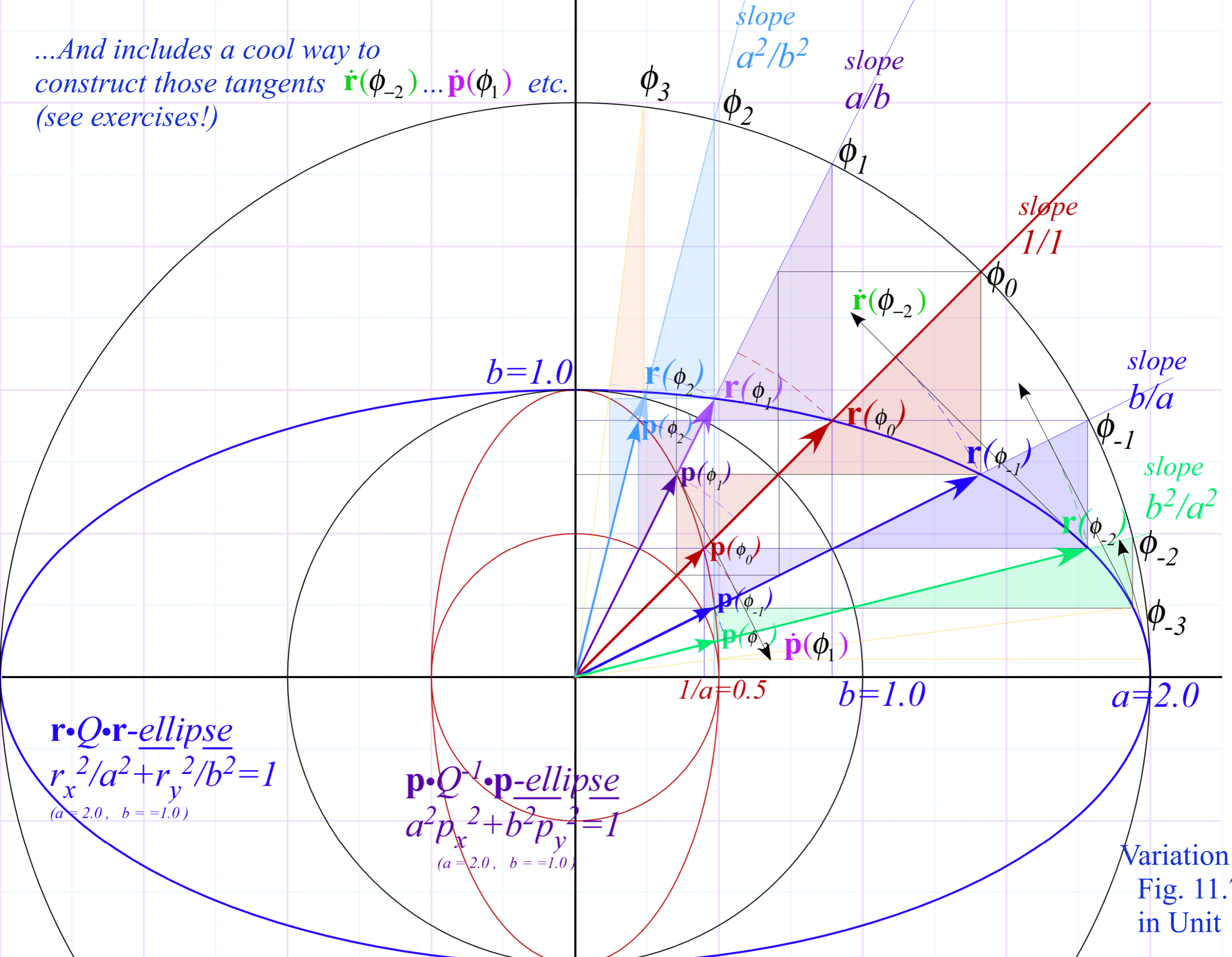


$r \cdot Q \cdot r$ -ellipse
 $r_x^2/a^2 + r_y^2/b^2 = 1$
 ($a=2.0, b=1.0$)

$p \cdot Q^{-1} \cdot p$ -ellipse
 $a^2 p_x^2 + b^2 p_y^2 = 1$
 ($a=2.0, b=1.0$)

Variation of Fig. 11.7 in Unit 1

...And includes a cool way to construct those tangents $\mathbf{r}(\phi_{-2}) \dots \mathbf{p}(\phi_1)$ etc. (see exercises!)



$\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r}$ -ellipse

$$r_x^2/a^2 + r_y^2/b^2 = 1$$

($a=2.0, b=1.0$)

$\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p}$ -ellipse

$$a^2 p_x^2 + b^2 p_y^2 = 1$$

($a=2.0, b=1.0$)

Variation of Fig. 11.7 in Unit 1

*Q: Where is this headed?
Preview of Lecture 8*

A: Lagrangian-Hamiltonian duality

The R and Q matrix transformations are like the mechanics rescaling matrices $\sqrt{\mathbf{M}}$ and \mathbf{M} :

Like $Q=R^2$:

$$\mathbf{M} = \begin{pmatrix} m_1 & 0 \\ 0 & m_2 \end{pmatrix} = \mathbf{R}^2$$

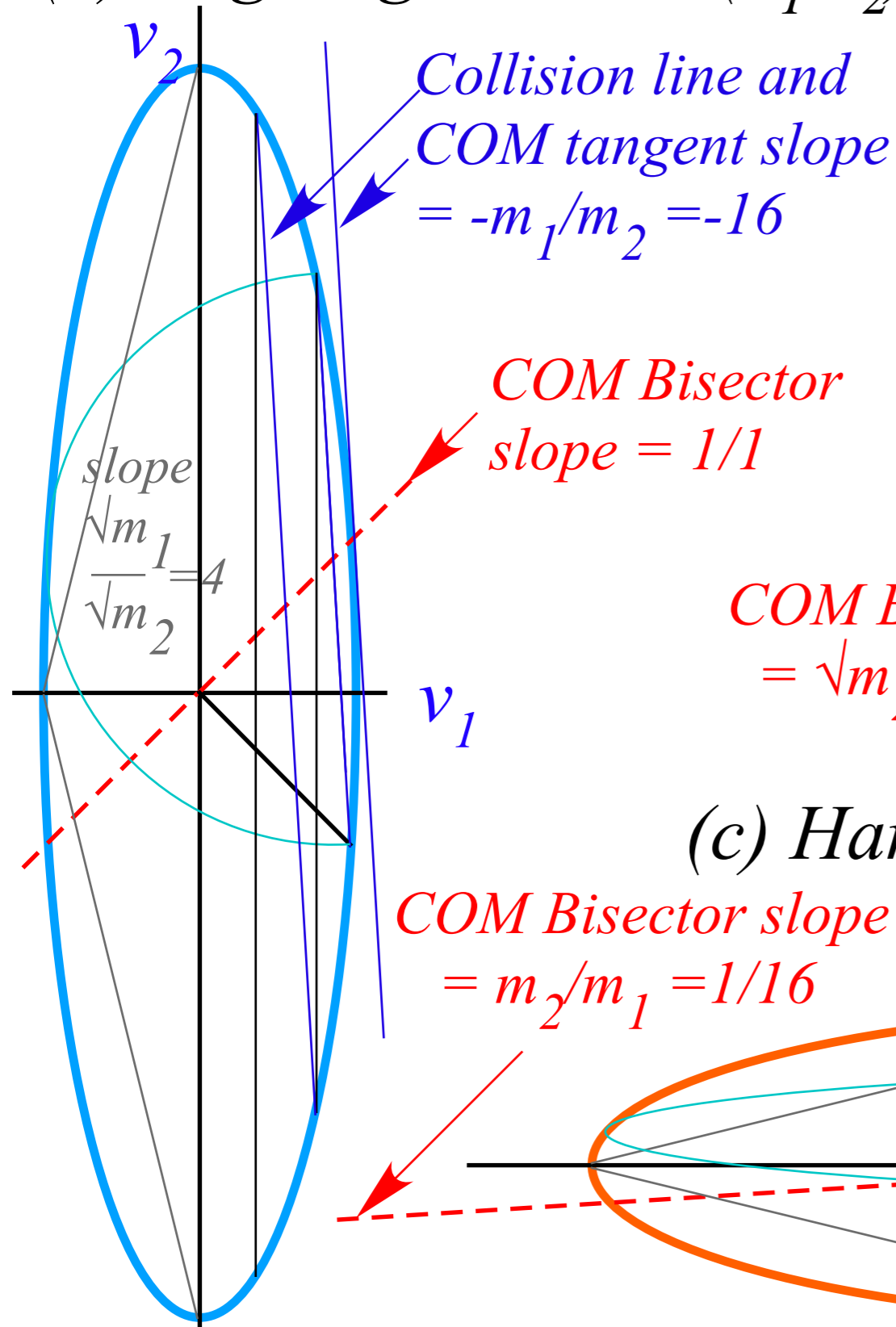
Like $\sqrt{Q}=R$:

$$\sqrt{\mathbf{M}} = \begin{pmatrix} \sqrt{m_1} & 0 \\ 0 & \sqrt{m_2} \end{pmatrix} = \mathbf{R}$$

Like $Q^{-1}=R^{-2}$:

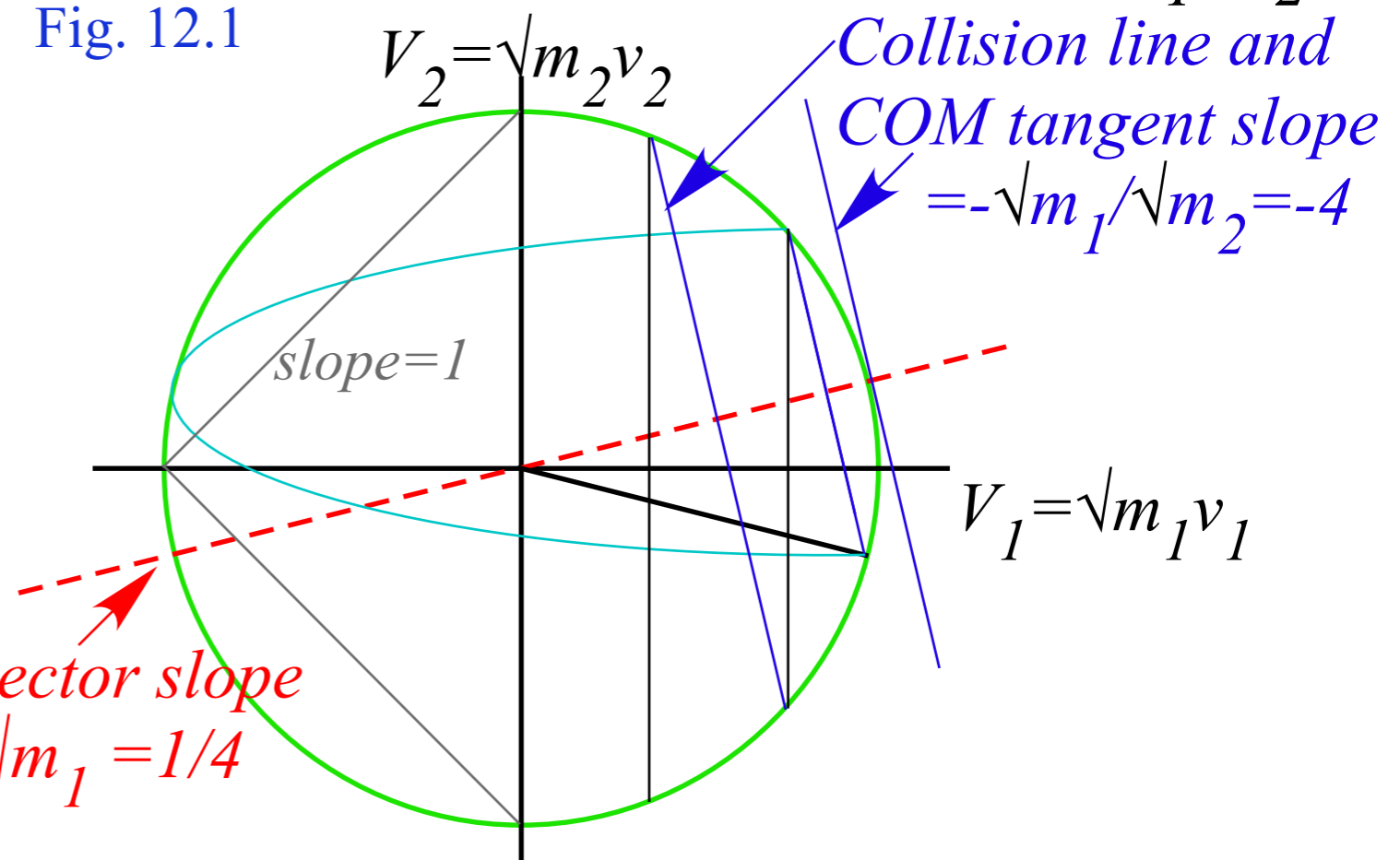
$$\mathbf{M}^{-1} = \begin{pmatrix} 1/m_1 & 0 \\ 0 & 1/m_2 \end{pmatrix} = \mathbf{R}^{-2}$$

(a) Lagrangian $L = L(v_1, v_2)$

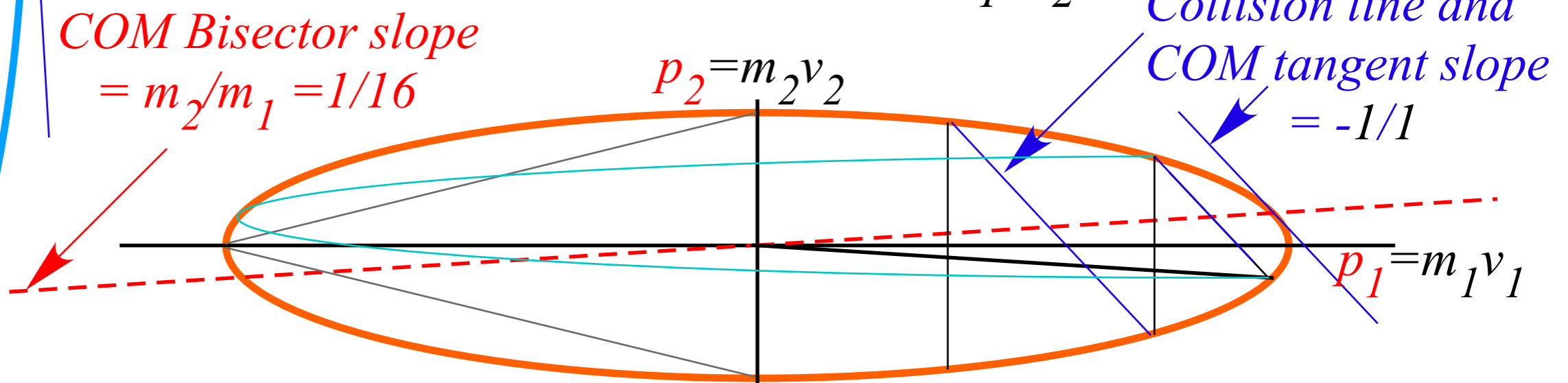


Unit 1
Fig. 12.1

(b) Estrangian $E = E(V_1, V_2)$

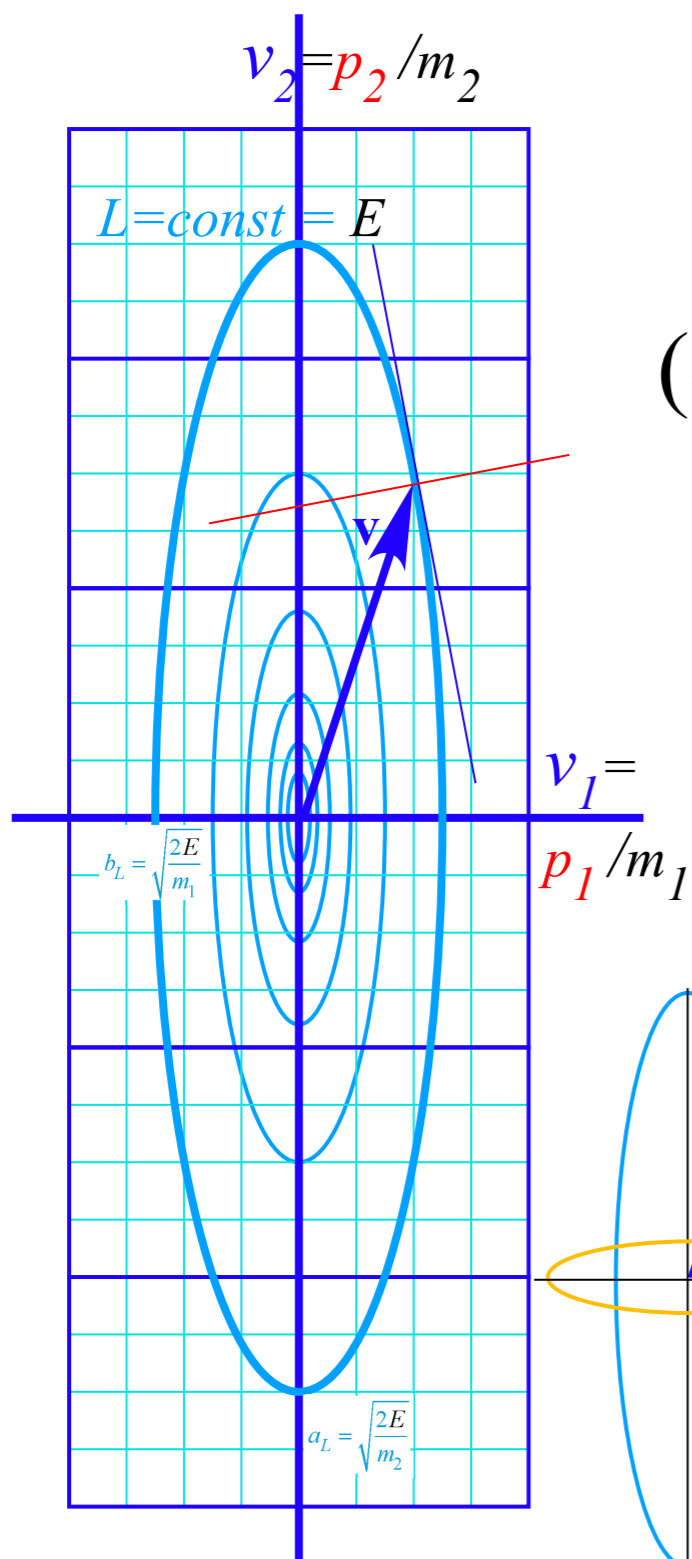


(c) Hamiltonian $H = H(p_1, p_2)$

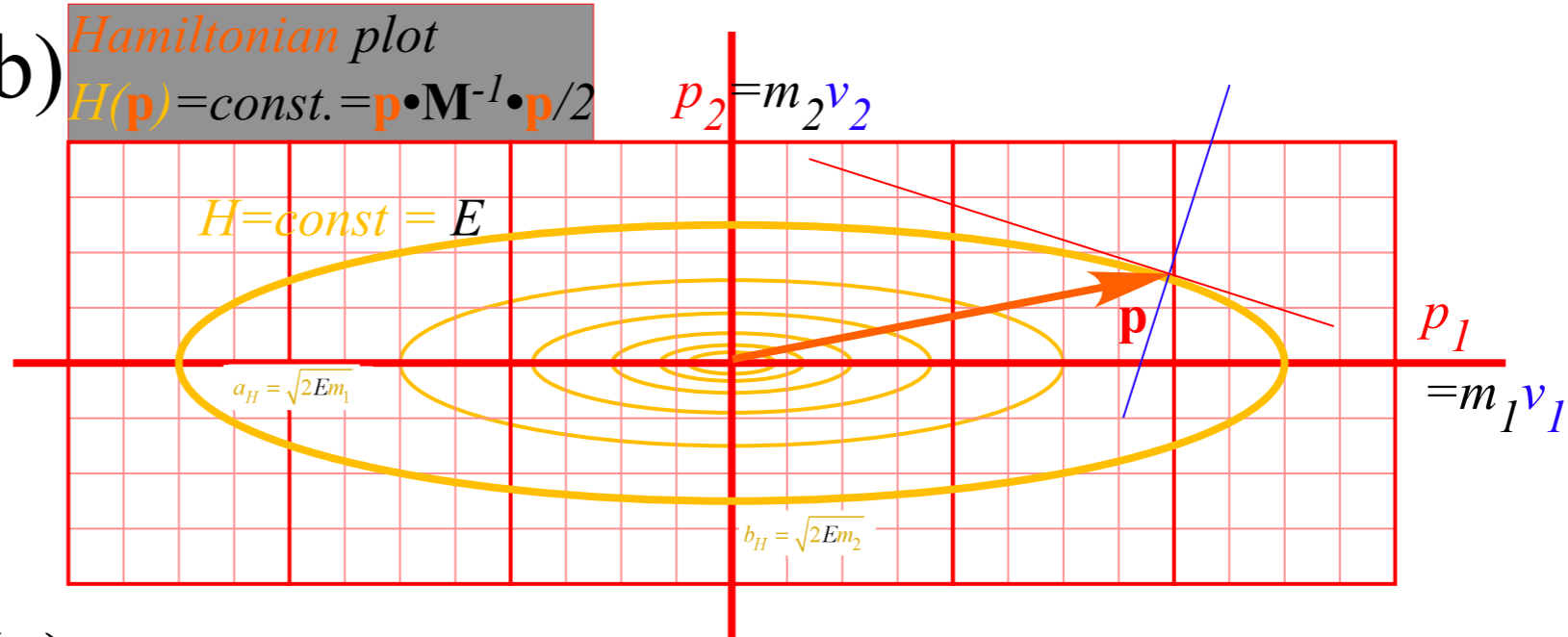


Unit 1
Fig. 12.2

(a) *Lagrangian plot*
 $L(\mathbf{v}) = \text{const.} = \mathbf{v} \cdot \mathbf{M} \cdot \mathbf{v} / 2$



(b) *Hamiltonian plot*
 $H(\mathbf{p}) = \text{const.} = \mathbf{p} \cdot \mathbf{M}^{-1} \cdot \mathbf{p} / 2$



(c) *Overlapping plots*

