

Group Theory in Quantum Mechanics

Lecture 2 (1.17.13)

Quantum amplitudes, analyzers, and axioms

(Quantum Theory for Computer Age - Ch. 1 of Unit 1)

(Principles of Symmetry, Dynamics, and Spectroscopy - Sec. 1-2 of Ch. 1)

Review: “Abstraction” of bra and ket vectors from a Transformation Matrix

Introducing scalar and matrix products

Planck's energy and N-quanta (Cavity/Beam wave mode)

Did Max Planck Goof? What's 1-photon worth?

Feynman amplitude axiom 1

What comes out of a beam sorter channel or branch-b?

Sample calculations

Feynman amplitude axioms 2-3

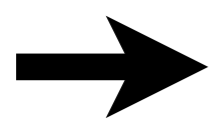
Beam analyzers: Sorter-unsorters

The “Do-Nothing” analyzer

Feynman amplitude axiom 4

Some “Do-Something” analyzers

Sorter-counter, Filter, 1/2-wave plate, 1/4-wave plate



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“Abstraction” of bra and ket vectors from a Transformation Matrix

Ket or column vectors

Bra or row vectors

Given
Transformation
Matrix $T_{m,n'}$:

$$\begin{pmatrix} \langle x|x' \rangle & \langle x|y' \rangle \\ \langle y|x' \rangle & \langle y|y' \rangle \end{pmatrix} = \begin{pmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{pmatrix}$$

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$$|x'\rangle = \begin{pmatrix} \langle x|x' \rangle \\ \langle y|x' \rangle \end{pmatrix} = \begin{pmatrix} \cos\theta \\ \sin\theta \end{pmatrix}$$

*Abstracting ket $|n'\rangle$ state vectors
from
Transformation Matrix*

$$T_{m,n'} = \langle m | n' \rangle$$

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Ket or column vectors

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$$\langle x| = \begin{pmatrix} \langle x|x' \rangle & \langle x|y' \rangle \end{pmatrix} = \begin{pmatrix} \cos\theta & -\sin\theta \end{pmatrix}$$

$$\langle y| = \begin{pmatrix} \langle y|x' \rangle & \langle y|y' \rangle \end{pmatrix} = \begin{pmatrix} \sin\theta & \cos\theta \end{pmatrix}$$

Abstracting bra $\langle m|$ state vectors

from

Transformation Matrix

$$T_{m,n'} = \langle m|n' \rangle$$

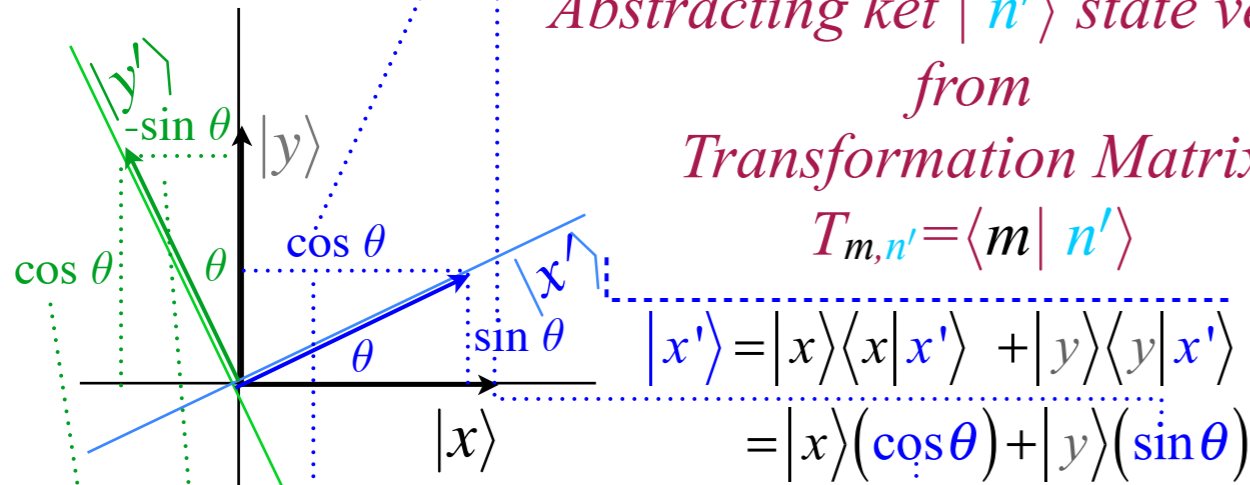
$$|x'\rangle = \begin{pmatrix} \langle x|x' \rangle \\ \langle y|x' \rangle \end{pmatrix} = \begin{pmatrix} \cos\theta \\ \sin\theta \end{pmatrix}, |y'\rangle = \begin{pmatrix} \langle x|y' \rangle \\ \langle y|y' \rangle \end{pmatrix} = \begin{pmatrix} -\sin\theta \\ \cos\theta \end{pmatrix}$$

Abstracting ket $|n'\rangle$ state vectors

from

Transformation Matrix

$$T_{m,n'} = \langle m|n' \rangle$$



$$|x'\rangle = |x\rangle\langle x|x'\rangle + |y\rangle\langle y|x'\rangle$$

$$= |x\rangle(\cos\theta) + |y\rangle(\sin\theta)$$

$$|y'\rangle = |x\rangle\langle x|y'\rangle + |y\rangle\langle y|y'\rangle$$

$$= |x\rangle(-\sin\theta) + |y\rangle(\cos\theta)$$

$$\langle y| = \langle y|x'\rangle\langle x'| + \langle y|y'\rangle\langle y'|$$

$$= (\sin\theta)\langle x'| + (\cos\theta)\langle y'|$$

$$\langle x| = \langle x|x'\rangle\langle x'| + \langle x|y'\rangle\langle y'|$$

$$= (\cos\theta)\langle x'| + (-\sin\theta)\langle y'|$$

$(\theta=+30^\circ)$ -Rotated kets $\{|x'\rangle, |y'\rangle\}$ or $\{\mathbf{x}', \mathbf{y}'\}$ represented in page-aligned $\{|x\rangle, |y\rangle\}$ basis.

$(\theta=-30^\circ)$ -Rotated bras $\{\langle x|, \langle y|\}$ or $\{\mathbf{x}, \mathbf{y}\}$ represented in page-aligned $\{|x'\rangle, |y'\rangle\}$ basis.

Ket vector algebra has the order of $T_{m,n'}$ transposed

Bra vector algebra has the same order as $T_{m,n'}$

$$|x'\rangle = |x\rangle\langle x|x'\rangle + |y\rangle\langle y|x'\rangle = |x\rangle(\cos\theta) + |y\rangle(\sin\theta)$$

$$|y'\rangle = |x\rangle\langle x|y'\rangle + |y\rangle\langle y|y'\rangle = |x\rangle(-\sin\theta) + |y\rangle(\cos\theta)$$

$$\langle x| = \langle x|x'\rangle\langle x'| + \langle x|y'\rangle\langle y'| = (\cos\theta)\langle x'| + (-\sin\theta)\langle y'|$$

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➔ *Review: “Abstraction” of bra and ket vectors from a Transformation Matrix*
Introducing scalar and matrix products

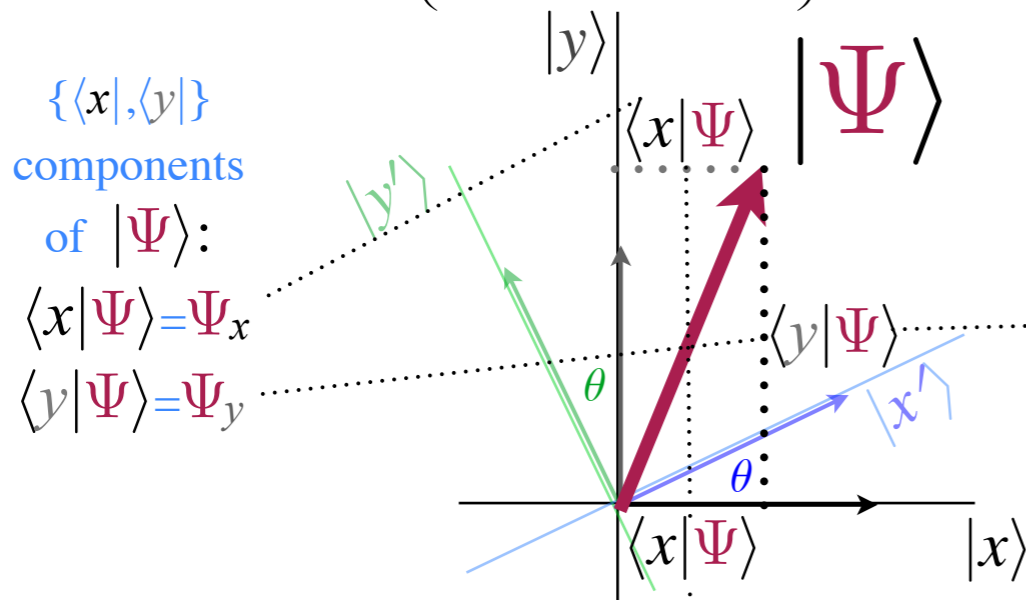
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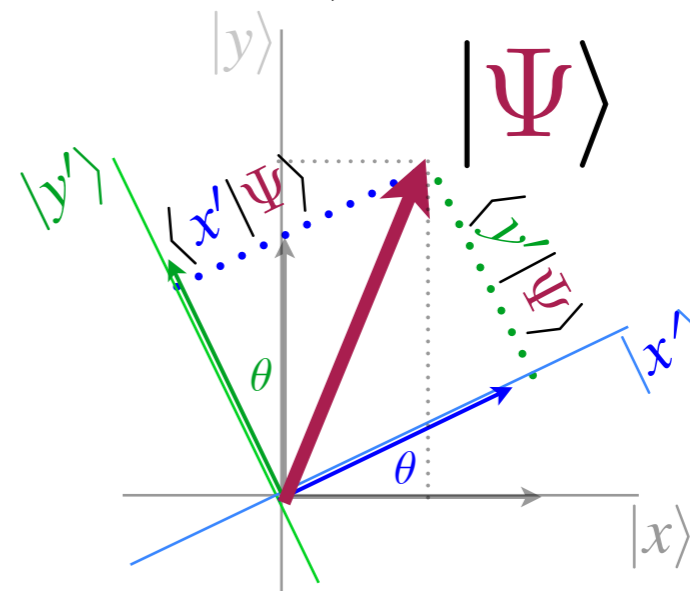
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Transformation matrix $T_{m,n'} = \langle m | n' \rangle$ is array of dot or *scalar products* (dot products) of unit vectors or *direction cosines*.

$$\begin{pmatrix} \langle x | x' \rangle & \langle x | y' \rangle \\ \langle y | x' \rangle & \langle y | y' \rangle \end{pmatrix} = \begin{pmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{pmatrix} = \begin{pmatrix} (\mathbf{x} \cdot \mathbf{x}') & (\mathbf{x} \cdot \mathbf{y}') \\ (\mathbf{y} \cdot \mathbf{x}') & (\mathbf{y} \cdot \mathbf{y}') \end{pmatrix}$$



$\{\langle x |, \langle y | \}$
components
of $|\Psi\rangle$:
 $\langle x | \Psi \rangle = \Psi_x$
 $\langle y | \Psi \rangle = \Psi_y$



$\{\langle x' |, \langle y' | \}$
components
of $|\Psi\rangle$:
 $\langle x' | \Psi \rangle = \Psi_{x'}$
 $\langle y' | \Psi \rangle = \Psi_{y'}$

Any state $|\Psi\rangle$ can be expanded in any basis $\{\langle x |, \langle y | \}$, or $\{\langle x' |, \langle y' | \}$, ...etc.

$$|\Psi\rangle = |x\rangle\langle x | \Psi \rangle + |y\rangle\langle y | \Psi \rangle = |x'\rangle\langle x' | \Psi \rangle + |y'\rangle\langle y' | \Psi \rangle$$

Transformation matrix $T_{m,n'}$ relates $\{\langle x | \Psi \rangle, \langle y | \Psi \rangle\}$ amplitudes to $\{\langle x' | \Psi \rangle, \langle y' | \Psi \rangle\}$.

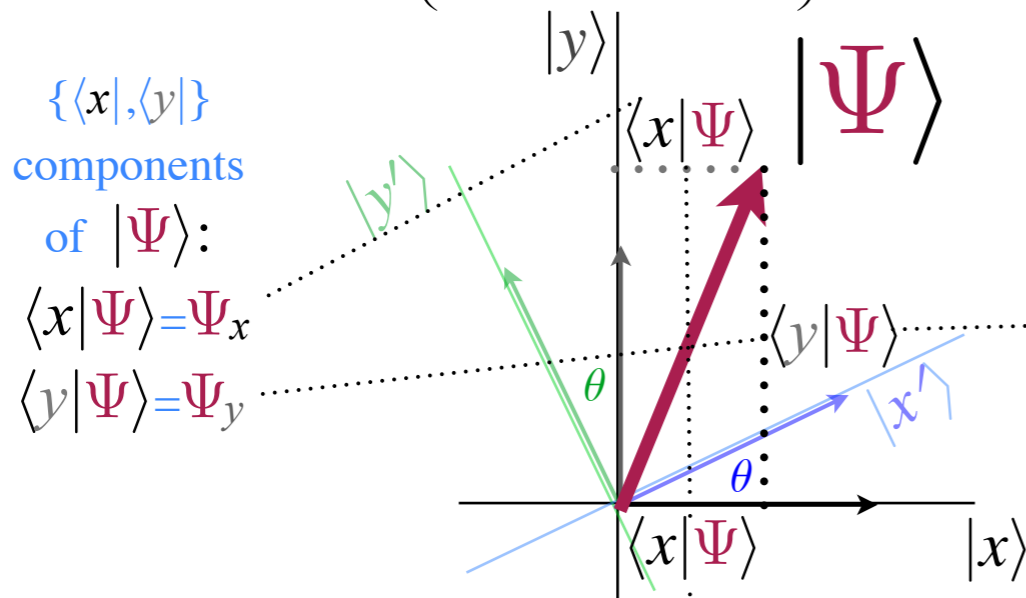
$$\begin{pmatrix} \langle x | \Psi \rangle \\ \langle y | \Psi \rangle \end{pmatrix} = \begin{pmatrix} \langle x | x' \rangle & \langle x | y' \rangle \\ \langle y | x' \rangle & \langle y | y' \rangle \end{pmatrix} \begin{pmatrix} \langle x' | \Psi \rangle \\ \langle y' | \Psi \rangle \end{pmatrix} \quad \text{or:} \quad \begin{pmatrix} \Psi_x \\ \Psi_y \end{pmatrix} = \begin{pmatrix} \langle x | x' \rangle & \langle x | y' \rangle \\ \langle y | x' \rangle & \langle y | y' \rangle \end{pmatrix} \begin{pmatrix} \Psi_{x'} \\ \Psi_{y'} \end{pmatrix}$$

Hybrid
Gibbs-Dirac
notation
(Ug-ly!)

Proof: $\langle x | = \langle x | x' \rangle \langle x' | + \langle x | y' \rangle \langle y' |$ implies: $\langle x | \Psi \rangle = \langle x | x' \rangle \langle x' | \Psi \rangle + \langle x | y' \rangle \langle y' | \Psi \rangle$
 $\langle y | = \langle y | x' \rangle \langle x' | + \langle y | y' \rangle \langle y' |$ implies: $\langle y | \Psi \rangle = \langle y | x' \rangle \langle x' | \Psi \rangle + \langle y | y' \rangle \langle y' | \Psi \rangle$

Transformation matrix $T_{m,n'} = \langle m | n' \rangle$ is array of dot or *scalar products* (dot products) of unit vectors or *direction cosines*.

$$\begin{pmatrix} \langle x | x' \rangle & \langle x | y' \rangle \\ \langle y | x' \rangle & \langle y | y' \rangle \end{pmatrix} = \begin{pmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{pmatrix} = \begin{pmatrix} (\mathbf{x} \cdot \mathbf{x}') & (\mathbf{x} \cdot \mathbf{y}') \\ (\mathbf{y} \cdot \mathbf{x}') & (\mathbf{y} \cdot \mathbf{y}') \end{pmatrix}$$



$\{\langle x |, \langle y | \}$
components
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$$\langle x | \Psi \rangle = \Psi_x$$

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$\{\langle x' |, \langle y' | \}$
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$$\langle x' | \Psi \rangle = \Psi_{x'}$$

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Any state $|\Psi\rangle$ can be expanded in any basis $\{\langle x |, \langle y | \}$, or $\{\langle x' |, \langle y' | \}$, ...etc.

$$|\Psi\rangle = |x\rangle \langle x | \Psi \rangle + |y\rangle \langle y | \Psi \rangle = |x'\rangle \langle x' | \Psi \rangle + |y'\rangle \langle y' | \Psi \rangle$$

Transformation matrix $T_{m,n'}$ relates $\{\langle x | \Psi \rangle, \langle y | \Psi \rangle\}$ amplitudes to $\{\langle x' | \Psi \rangle, \langle y' | \Psi \rangle\}$.

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Hybrid
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(Ug-ly!)

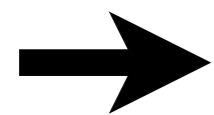
Inverse ($\dagger = T^* = -1$) matrix $T_{n',m}$ relates $\{\langle x' | \Psi \rangle, \langle y' | \Psi \rangle\}$ amplitudes to $\{\langle x | \Psi \rangle, \langle y | \Psi \rangle\}$.

$$\begin{pmatrix} \langle x' | \Psi \rangle \\ \langle y' | \Psi \rangle \end{pmatrix} = \begin{pmatrix} \langle x' | x \rangle & \langle x' | y \rangle \\ \langle y' | x \rangle & \langle y' | y \rangle \end{pmatrix} \begin{pmatrix} \langle x | \Psi \rangle \\ \langle y | \Psi \rangle \end{pmatrix} \quad \text{or:} \quad \begin{pmatrix} \Psi_{x'} \\ \Psi_{y'} \end{pmatrix} = \begin{pmatrix} \langle x' | x \rangle & \langle x' | y \rangle \\ \langle y' | x \rangle & \langle y' | y \rangle \end{pmatrix} \begin{pmatrix} \Psi_x \\ \Psi_y \end{pmatrix}$$

Hybrid
Gibbs-Dirac
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Planck's energy and N -quanta (*Cavity/Beam of volume V with wave mode of frequency $\omega=2\pi\nu$*)

Planck axiom: E -field energy density U in cavity/beam mode- ω is: $U=N\hbar\omega/V =Nh\nu/V$ (N "photons")

$$h=2\pi\hbar=6.6310^{-34}Js$$

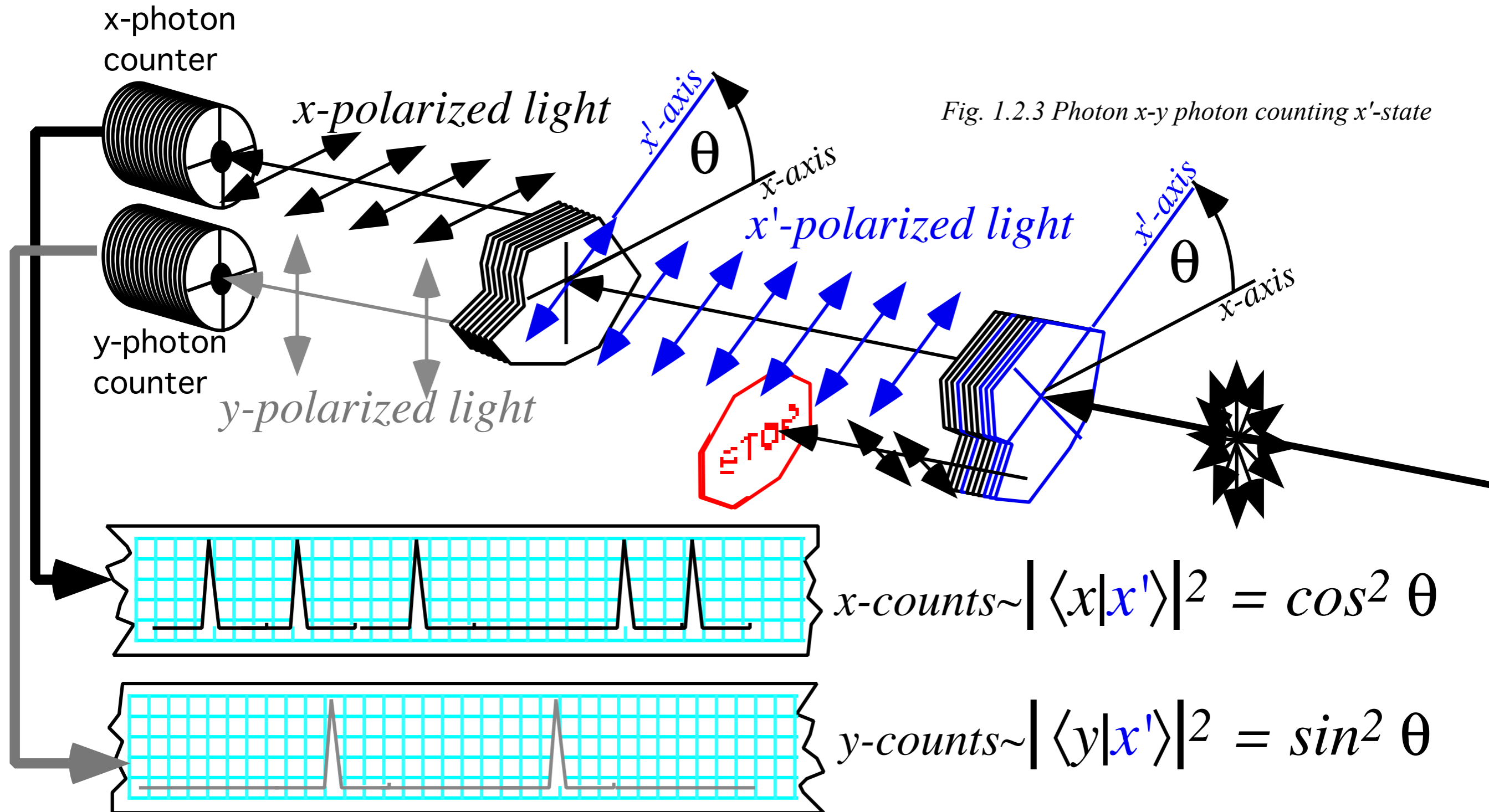
Planck constant

Planck's energy and N -quanta (Cavity/Beam of volume V with wave mode of frequency $\omega=2\pi\nu$)

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E -field vector $(E_x, E_y) = f \cdot (\Psi_x, \Psi_y)$ where *quantum field proportionality constant* is f .

$$h=2\pi\hbar=6.63 \cdot 10^{-34} \text{Js}$$

Planck constant

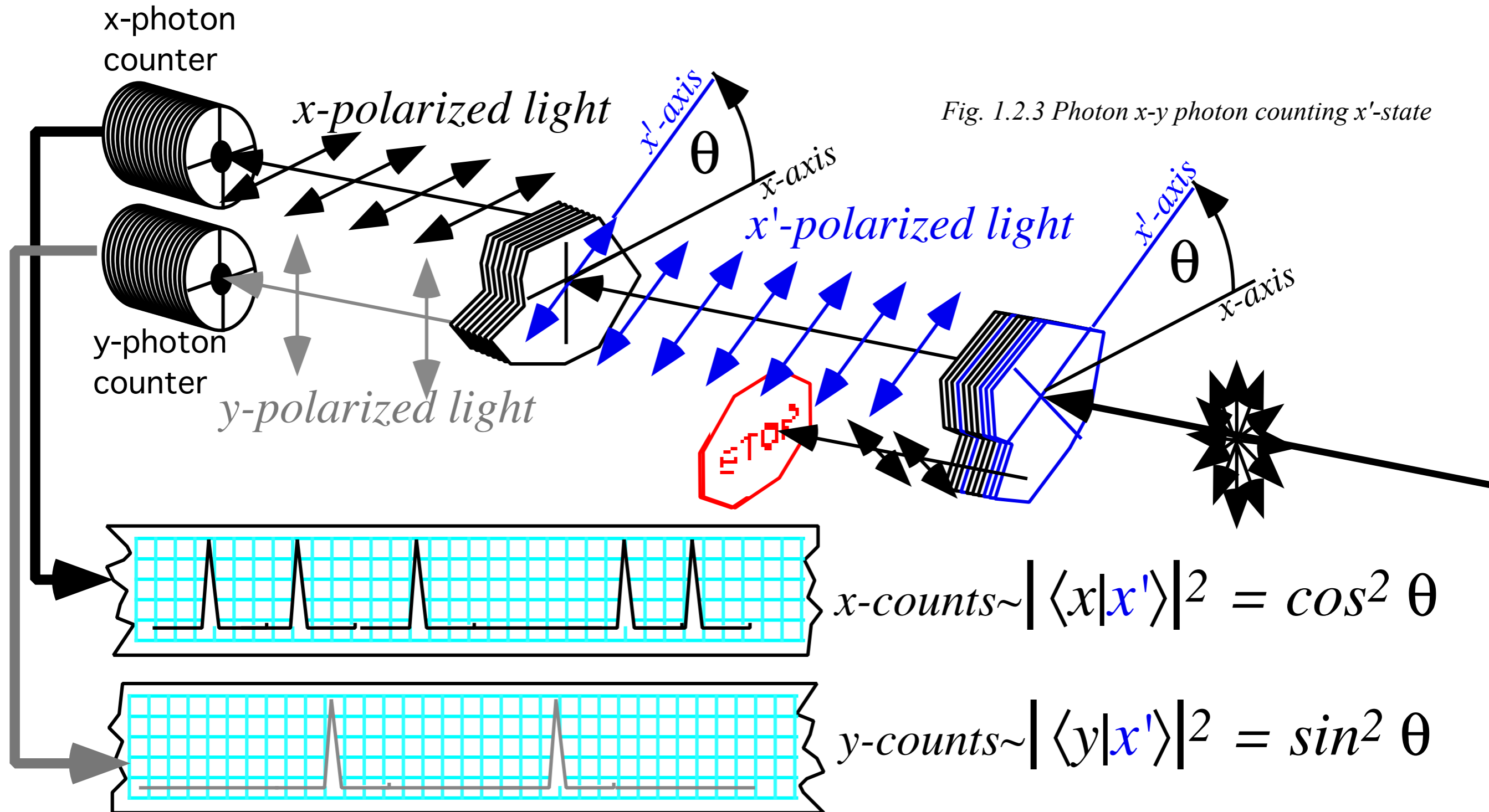


Fig. 1.2.3 Photon x-y photon counting x' -state

Planck's energy and N -quanta (Cavity/Beam of volume V with wave mode of frequency $\omega=2\pi\nu$)

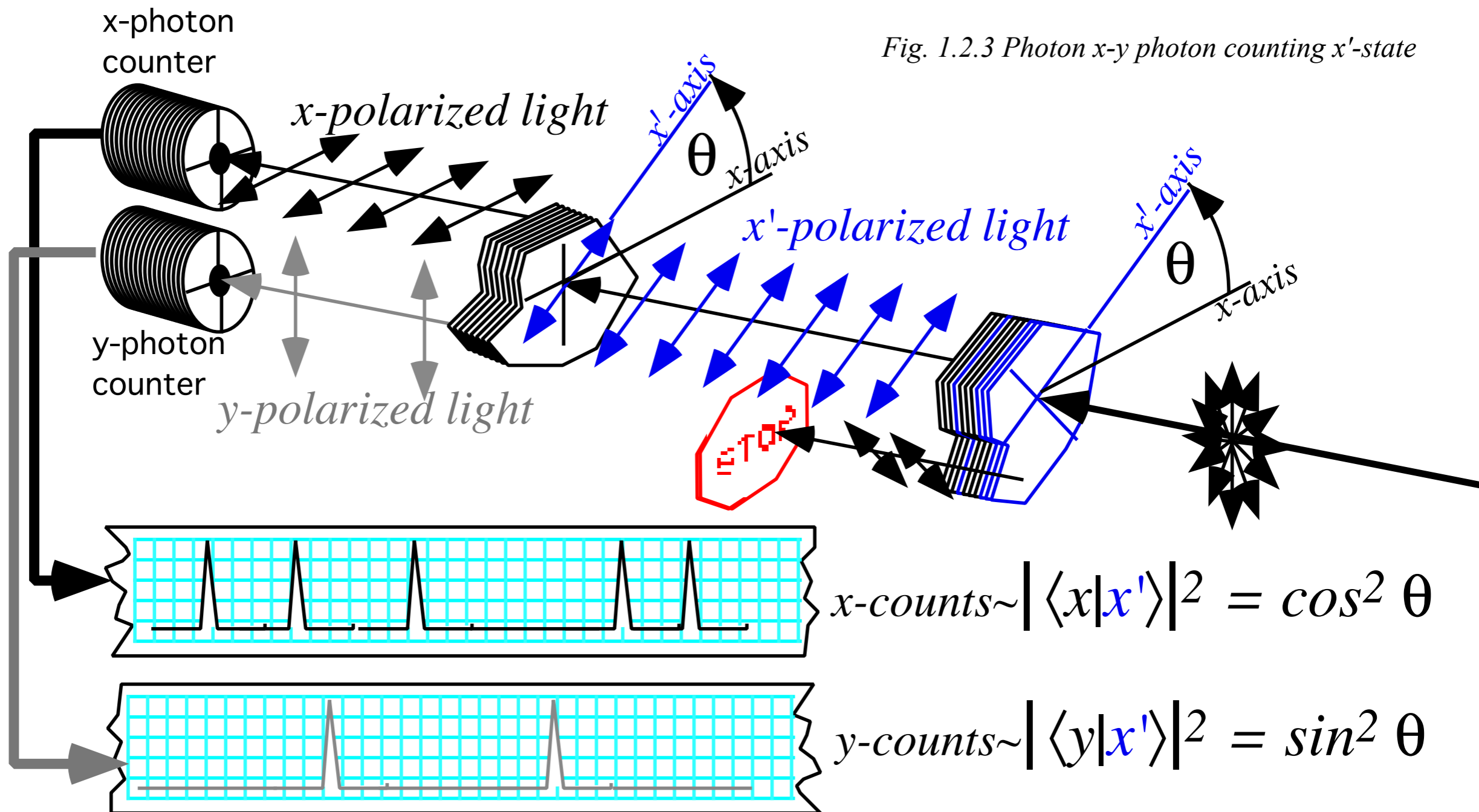
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$$\begin{pmatrix} E_x \\ E_y \end{pmatrix} = \begin{pmatrix} E_x(0)e^{ikz-i\omega t} \\ E_y(0)e^{ikz-i\omega t} \end{pmatrix} \cong \begin{pmatrix} E_x(0)e^{-i\omega t} \\ E_y(0)e^{-i\omega t} \end{pmatrix} = f \begin{pmatrix} \Psi_x \\ \Psi_y \end{pmatrix}$$

$h=2\pi\hbar=6.63 \cdot 10^{-34} \text{Js}$

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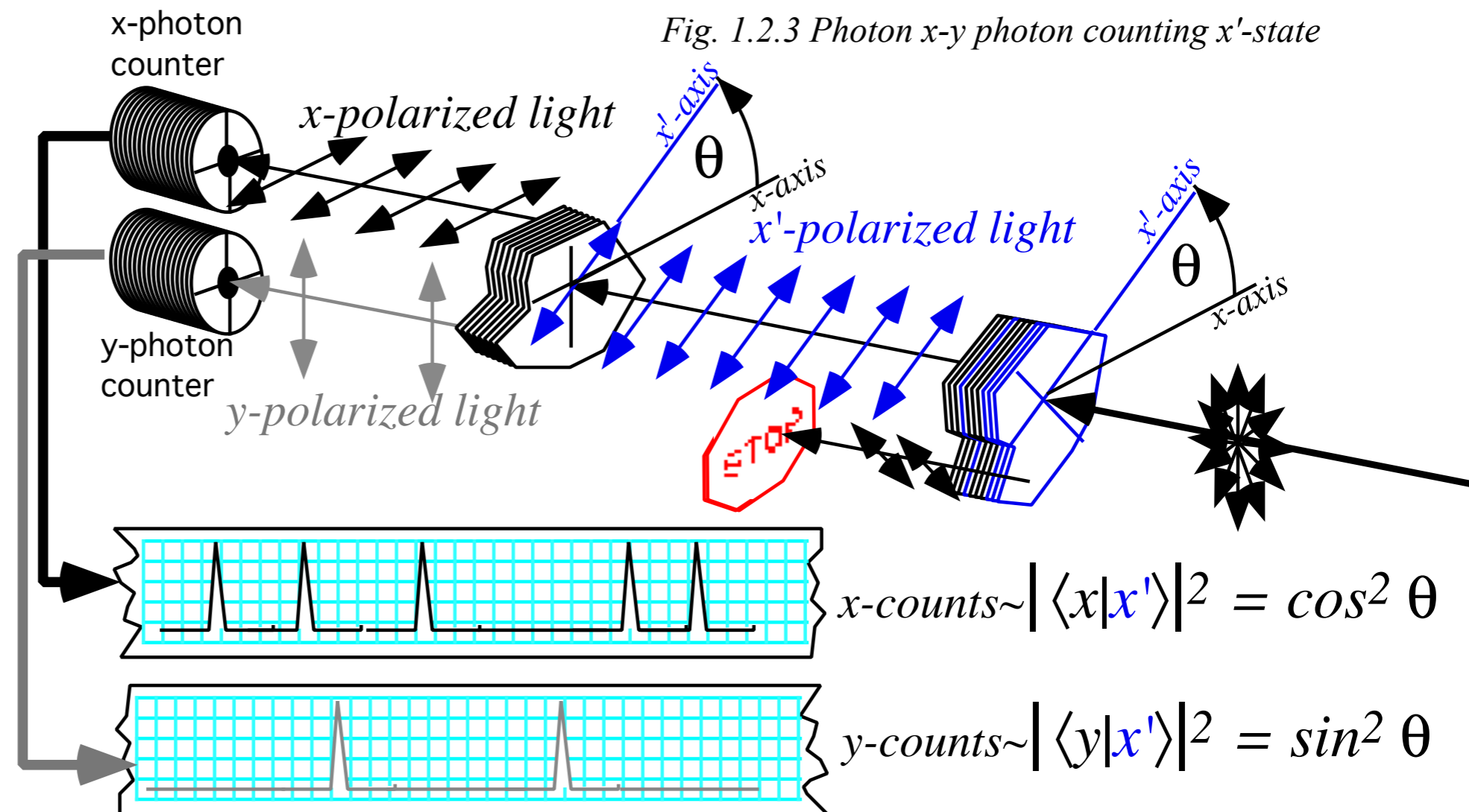
$h=2\pi\hbar=6.626\ 075 \cdot 10^{-34} \text{ Js}$ Planck constant
 $c = 2.997\ 924\ 58 \cdot 10^8 \text{ ms}^{-1}$ Light speed
 $\epsilon_0 = 8.854 \cdot 10^{-12} \text{ C}^2 \text{ N}^{-1} \text{ m}^{-2}$ Electrostatic constant

Coulomb constant: $k = 1/4\pi\epsilon_0 = 9 \cdot 10^9 \text{ J/C}$

Ψ -amplitude-squares sum to *exciton-number* N . (...or *photon-number* N)

$$N = |\Psi_x|^2 + |\Psi_y|^2 = \langle x | \Psi \rangle^* \langle x | \Psi \rangle + \langle y | \Psi \rangle^* \langle y | \Psi \rangle$$

Quantum
em wave
theory



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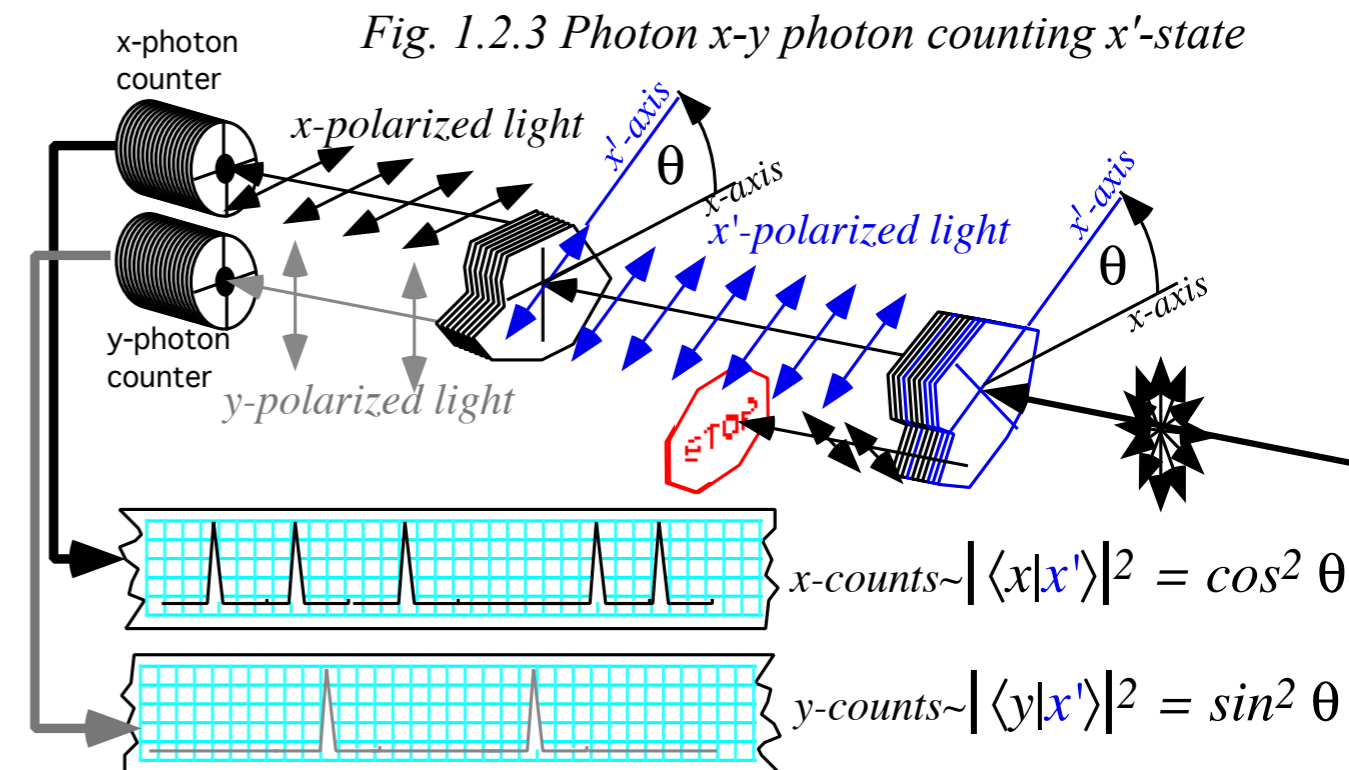
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Quantum
em wave
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Poynting energy flux $S(\text{J/m}^2\text{s})$ or *energy density* $U(\text{J/m}^3)$ of light beam.

$$S = cU, \text{ where: } U = \epsilon_0 \left(|E_x|^2 + |E_y|^2 \right) = \epsilon_0 \left(E_x^* E_x + E_y^* E_y \right) = \epsilon_0 \left(E_x(0)^2 + E_y(0)^2 \right)$$

Classical
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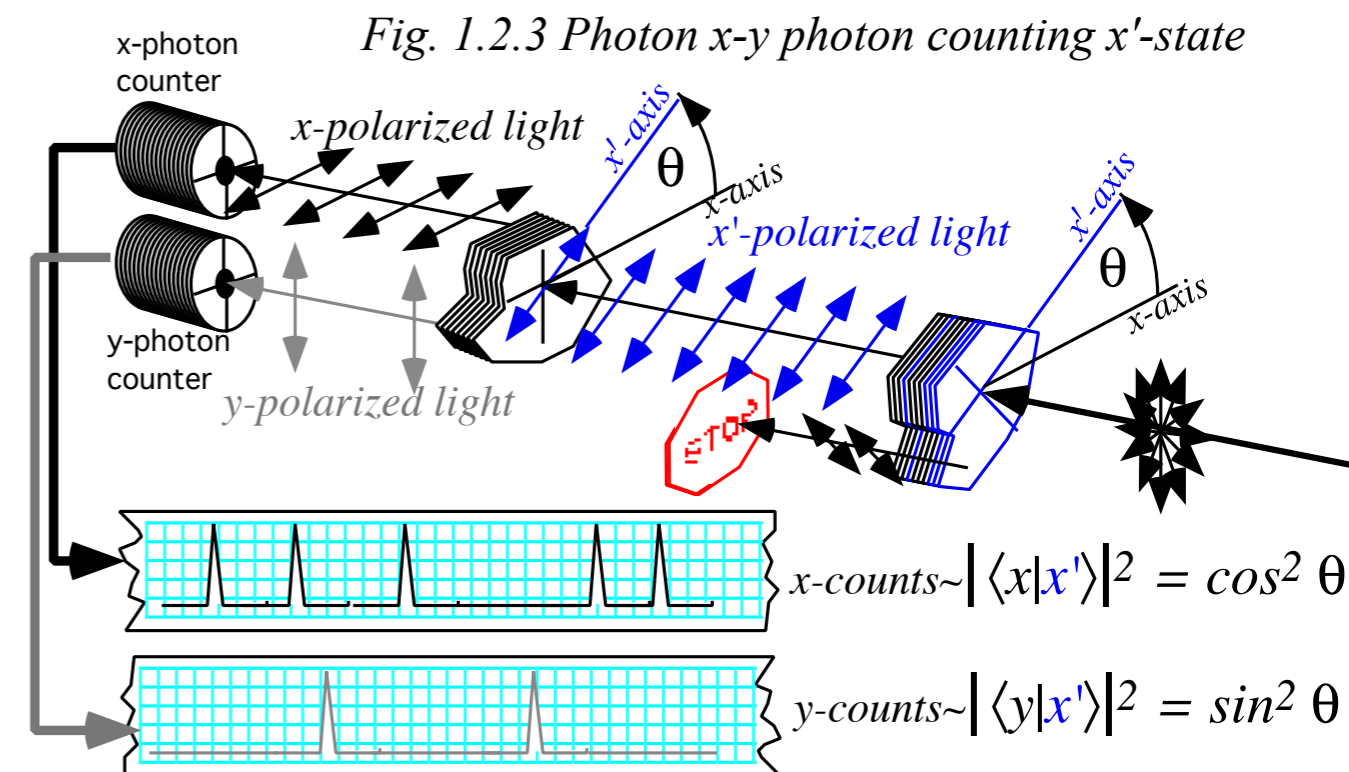
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Classical
em wave
theory

Equate U to Planck's N -*"photon"* quantum energy density $N\hbar\omega/V$



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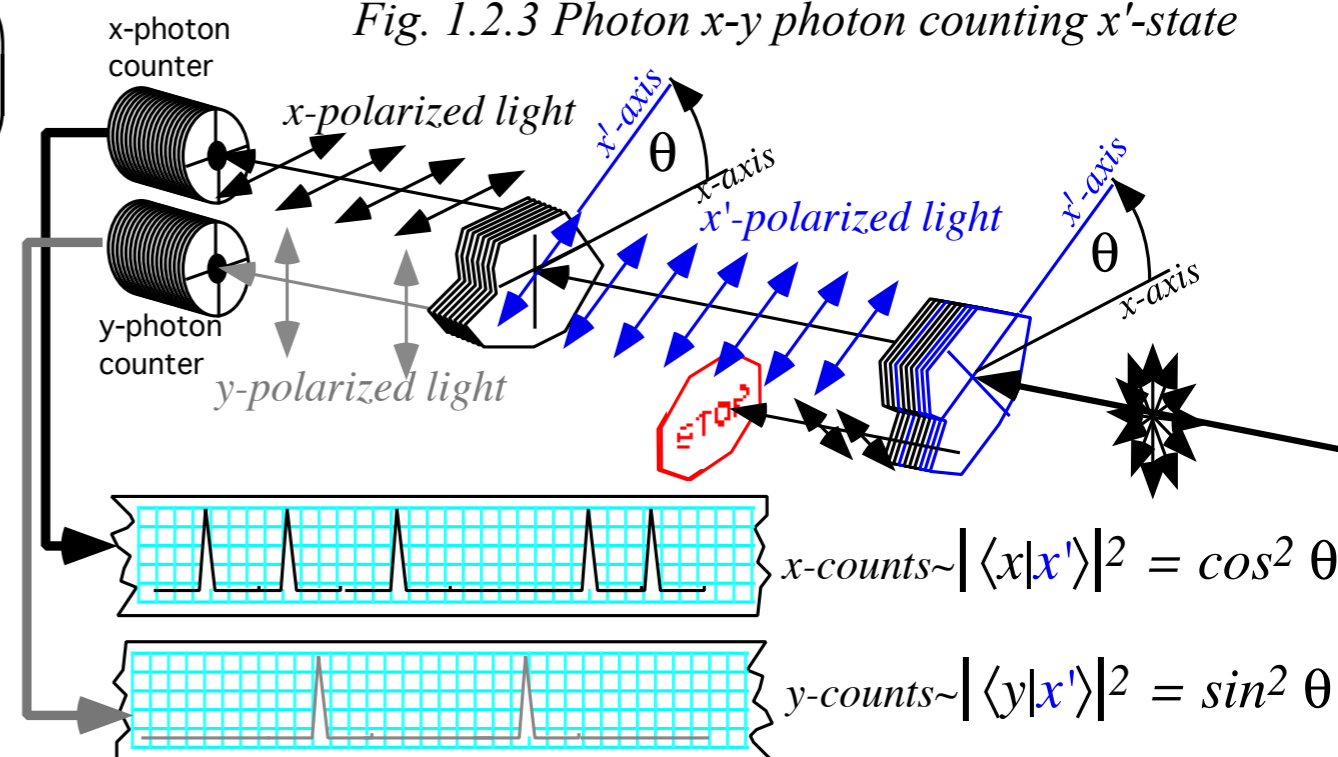
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Classical
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Equate U to Planck's N -*"photon"* quantum energy density $N\hbar\omega/V$

$$\frac{N\hbar\omega}{V} = U = \epsilon_0 \left(|E_x|^2 + |E_y|^2 \right) = \epsilon_0 f^2 \underbrace{\left(|\Psi_x|^2 + |\Psi_y|^2 \right)}_N$$

Fig. 1.2.3 Photon x - y photon counting x' -state



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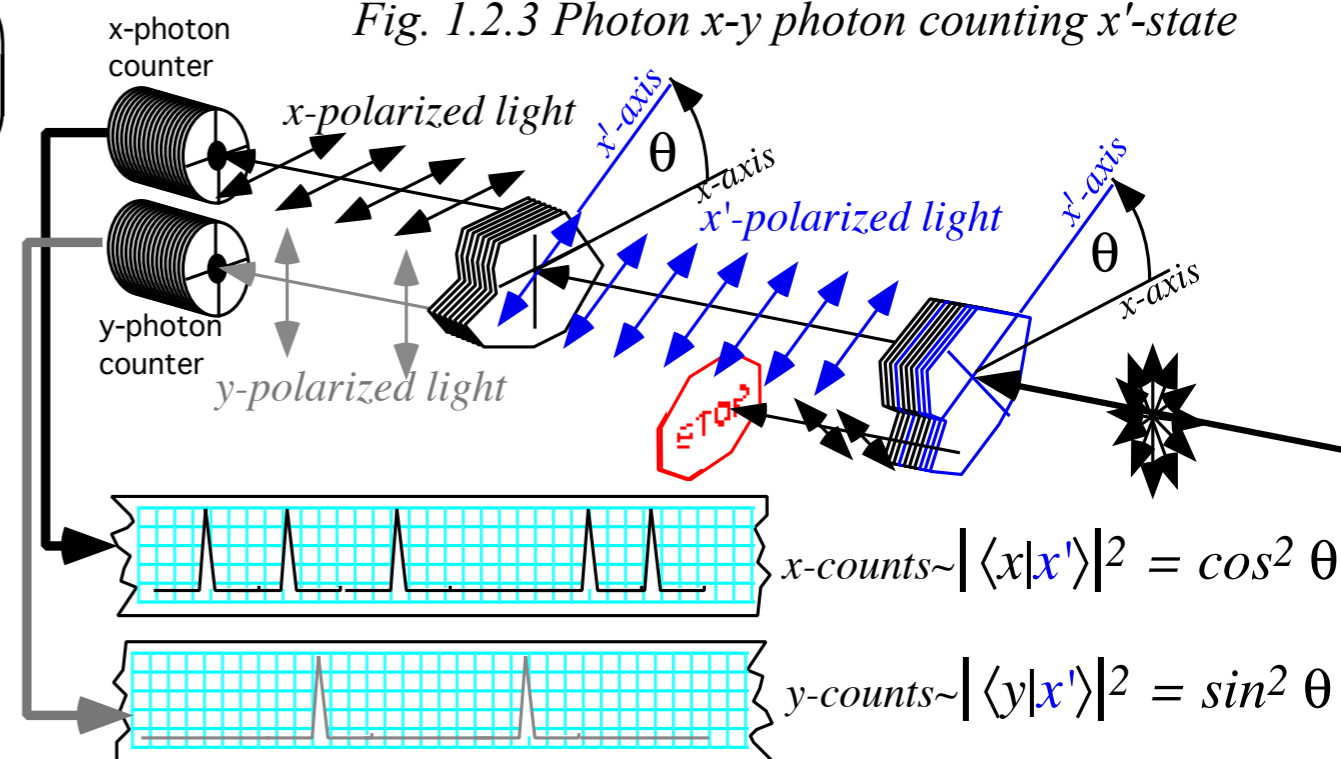
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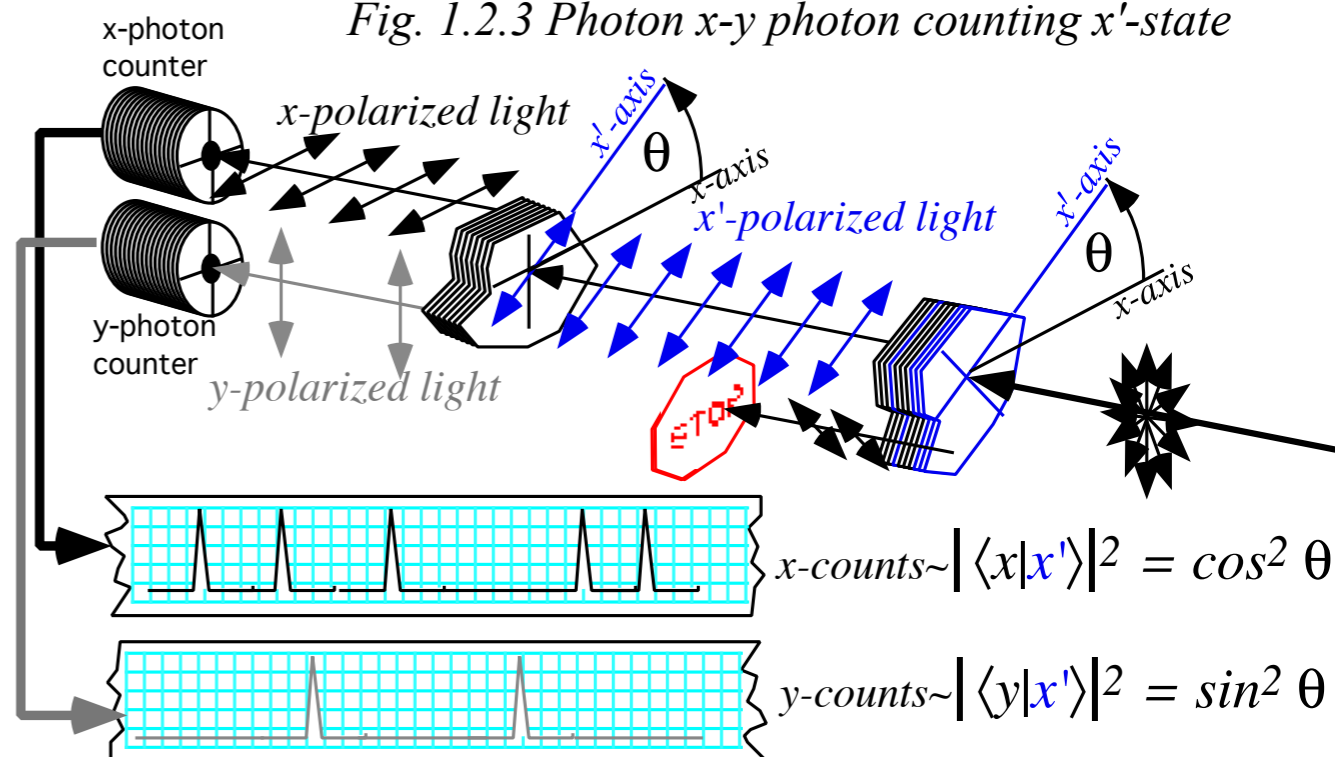
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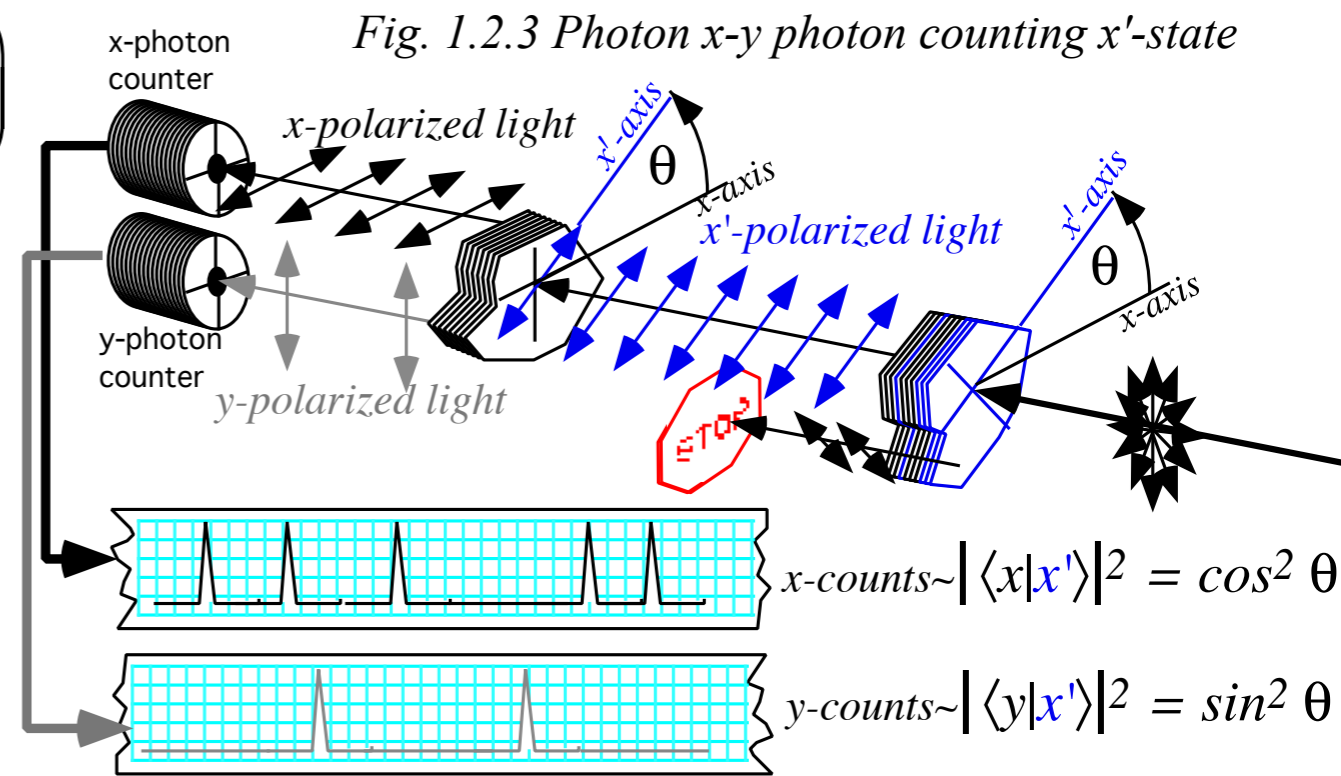
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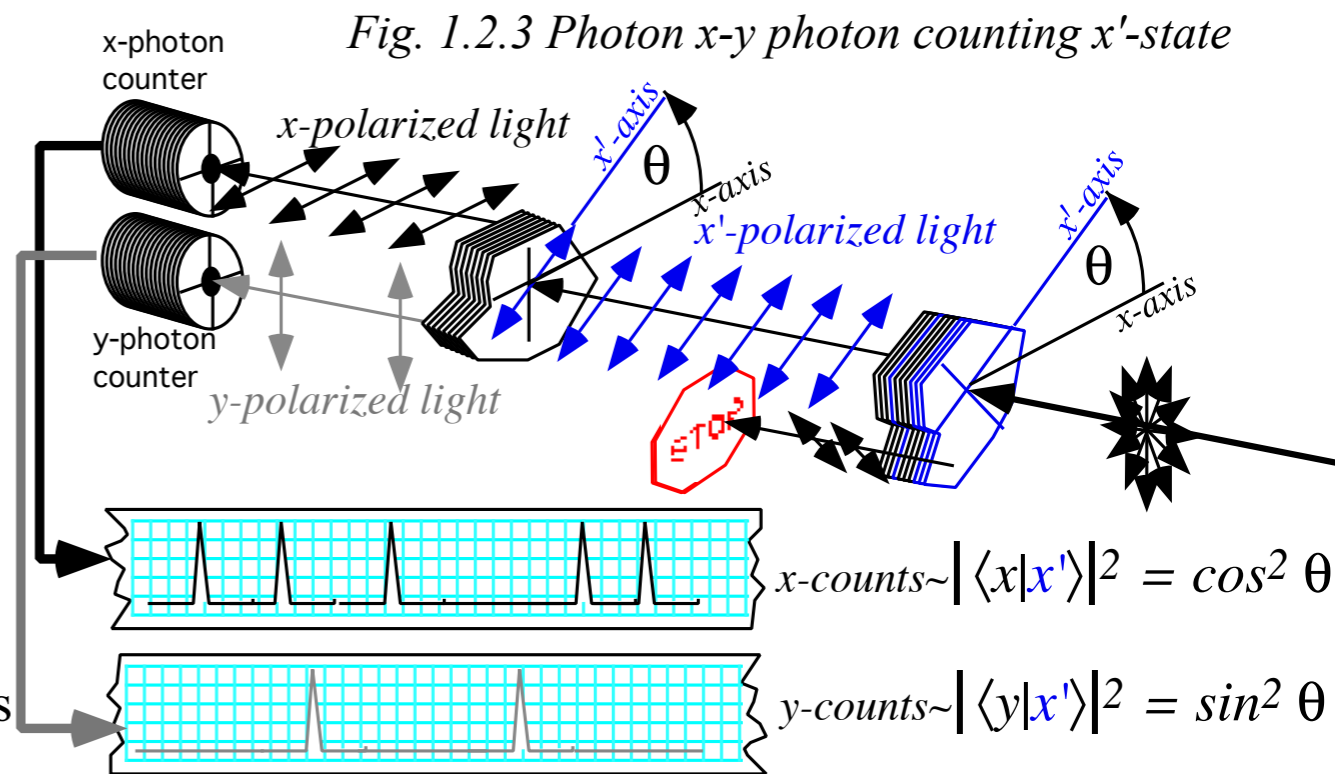


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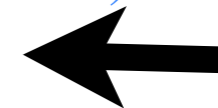
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$$\epsilon = (\text{const.}) \omega^2 A^2 = 1/2 m \omega^2 A^2$$

Energy U for classical electromagnetic cavity mode is **quadratic** in frequency ω and vector potential \mathbf{A} .

$$U = \epsilon_0 (|E_x|^2 + |E_y|^2) = \epsilon_0 |\mathbf{E}|^2 = \epsilon_0 \omega^2 |\mathbf{A}|^2 \quad \text{where: } \mathbf{E} = -\partial_t \mathbf{A} = i\omega \mathbf{A}$$

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Was Planck's linear-in-frequency- ν energy axiom a Goof?

Not at all. It's really not linear-in-frequency after all.

The $\epsilon_0 |\mathbf{E}|^2 = h N \cdot \nu$ axiom IS in fact a product of TWO frequencies!

The 2nd "frequency" is COUNT RATE N .

As shown in Unit 2: Both frequencies transform by Relativistic Doppler factor $e^{\pm\phi}$:

Coherent frequency is light quality: $\nu' = e^{\pm\phi} \nu$. Incoherent frequency is light quantity: $N' = e^{\pm\phi} N$.

This would imply that the $|\mathbf{E}|$ -field also transforms like a frequency: $|\mathbf{E}'| = e^{\pm\phi} |\mathbf{E}|$.

Indeed, E-field amplitude is a frequency, too!

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Feynman amplitude axiom 1

(1) The probability axiom

The first axiom deals with physical interpretation of amplitudes $\langle j|k' \rangle$.

Axiom 1: The absolute square $|\langle j|k' \rangle|^2 = \langle j|k' \rangle^ \langle j|k' \rangle$ gives probability for occurrence in state- j of a system that started in state- $k'=1',2',\dots,$ or n' from one sorter and then was forced to choose between states $j=1,2,\dots,n$ by another sorter.*

*Feynman-Dirac
Interpretation of*

$$\langle j|k' \rangle$$

*=Amplitude of state- j after
state- k' is forced to choose
from available m -type states*

Amplitude axioms apply to all intensity-^{probability}-conserving systems

This includes, first of all, spin-1/2 electron, proton, ..., ^{13}C , ... particles (Fermions)

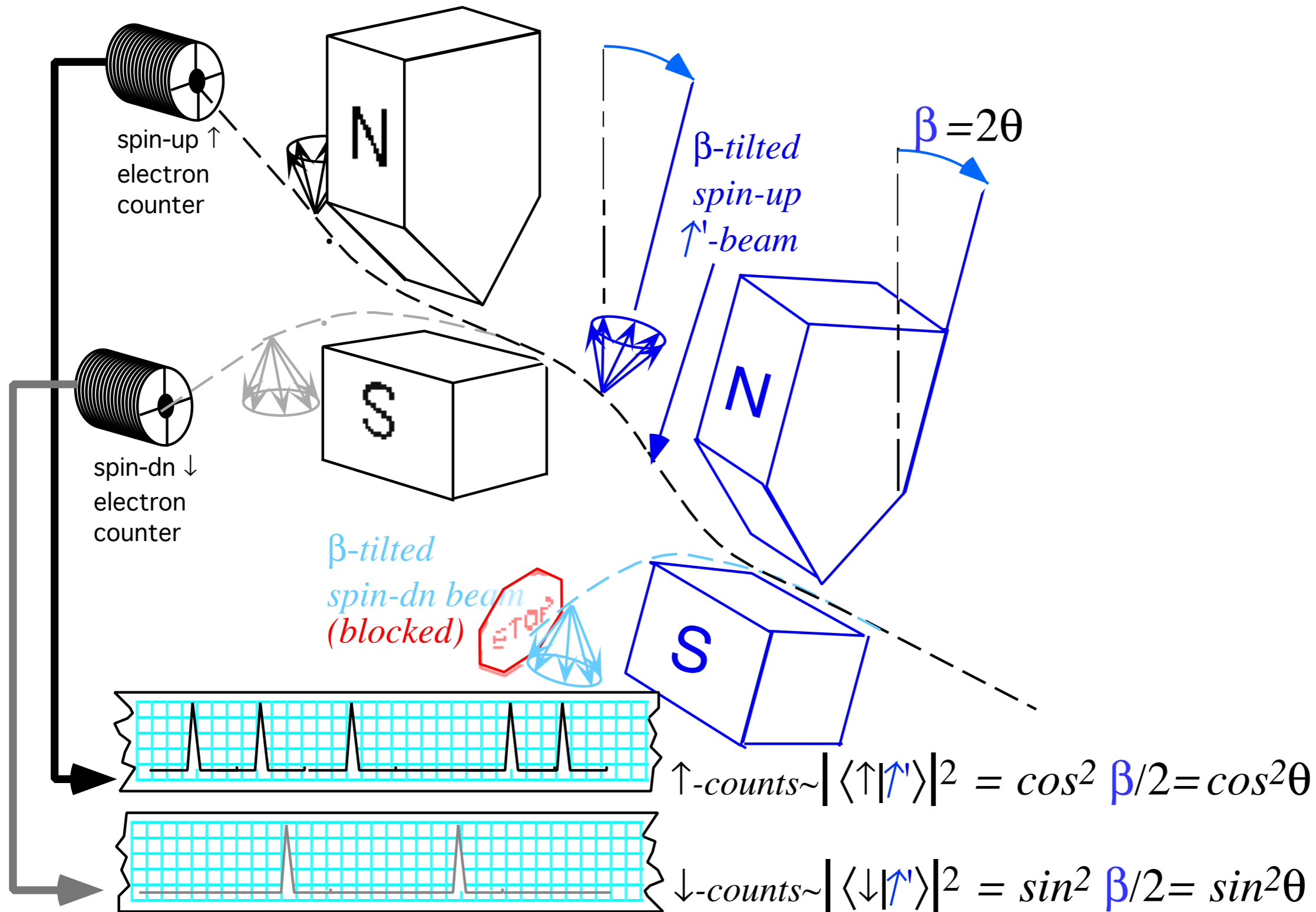


Fig. 1.2.4 Electron up-dn-spin counting of a tilted spin-up (\uparrow')-state

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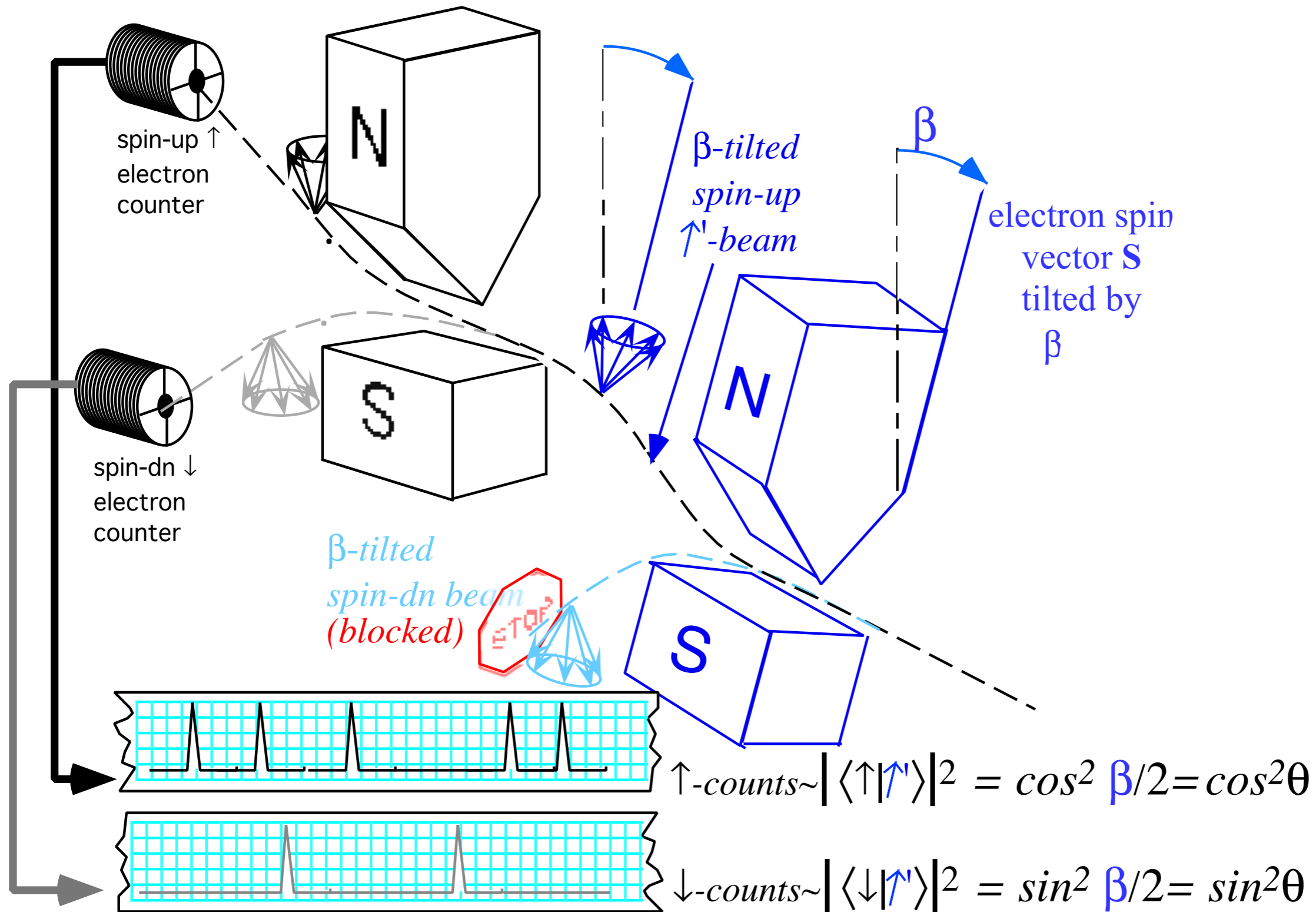


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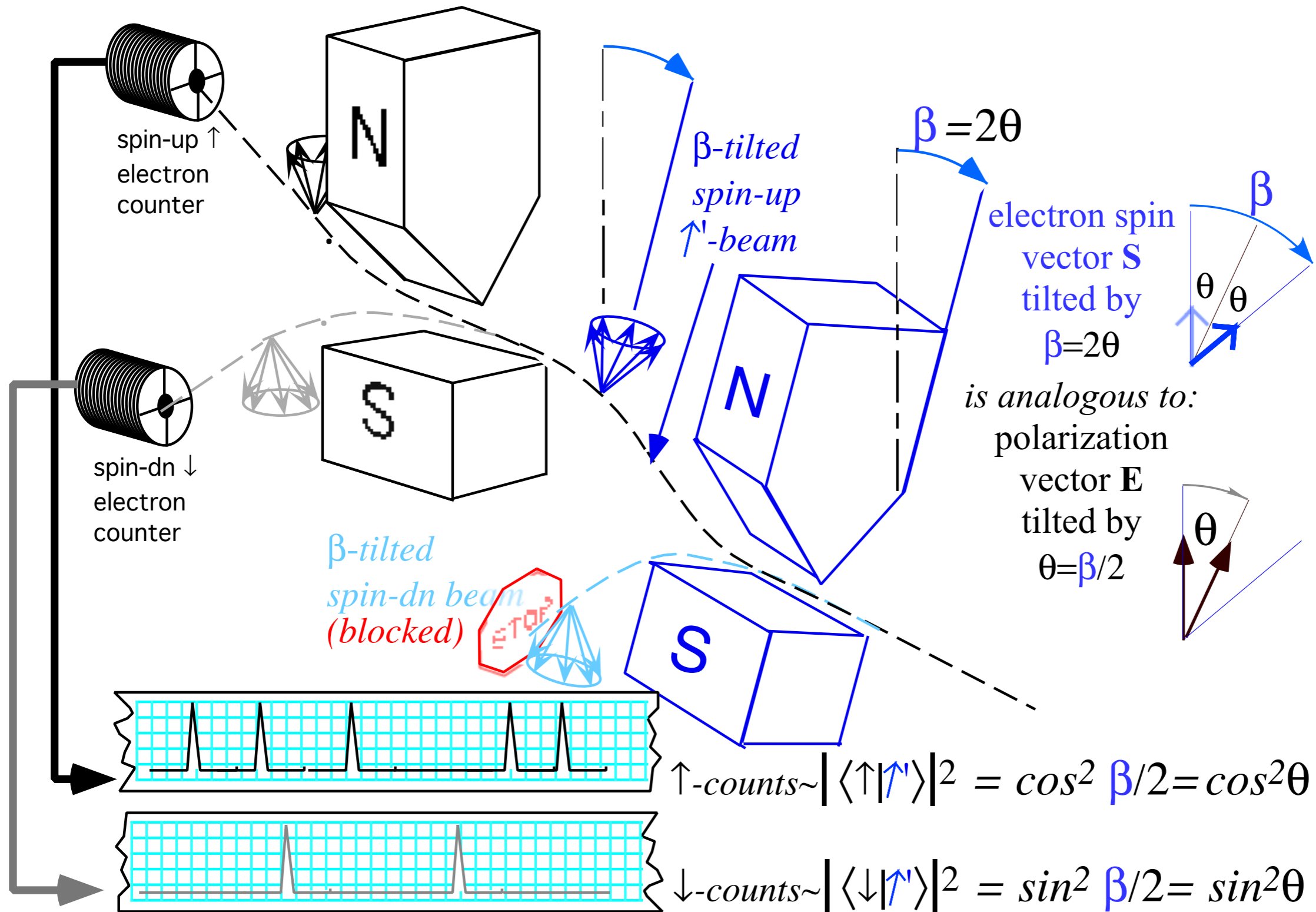


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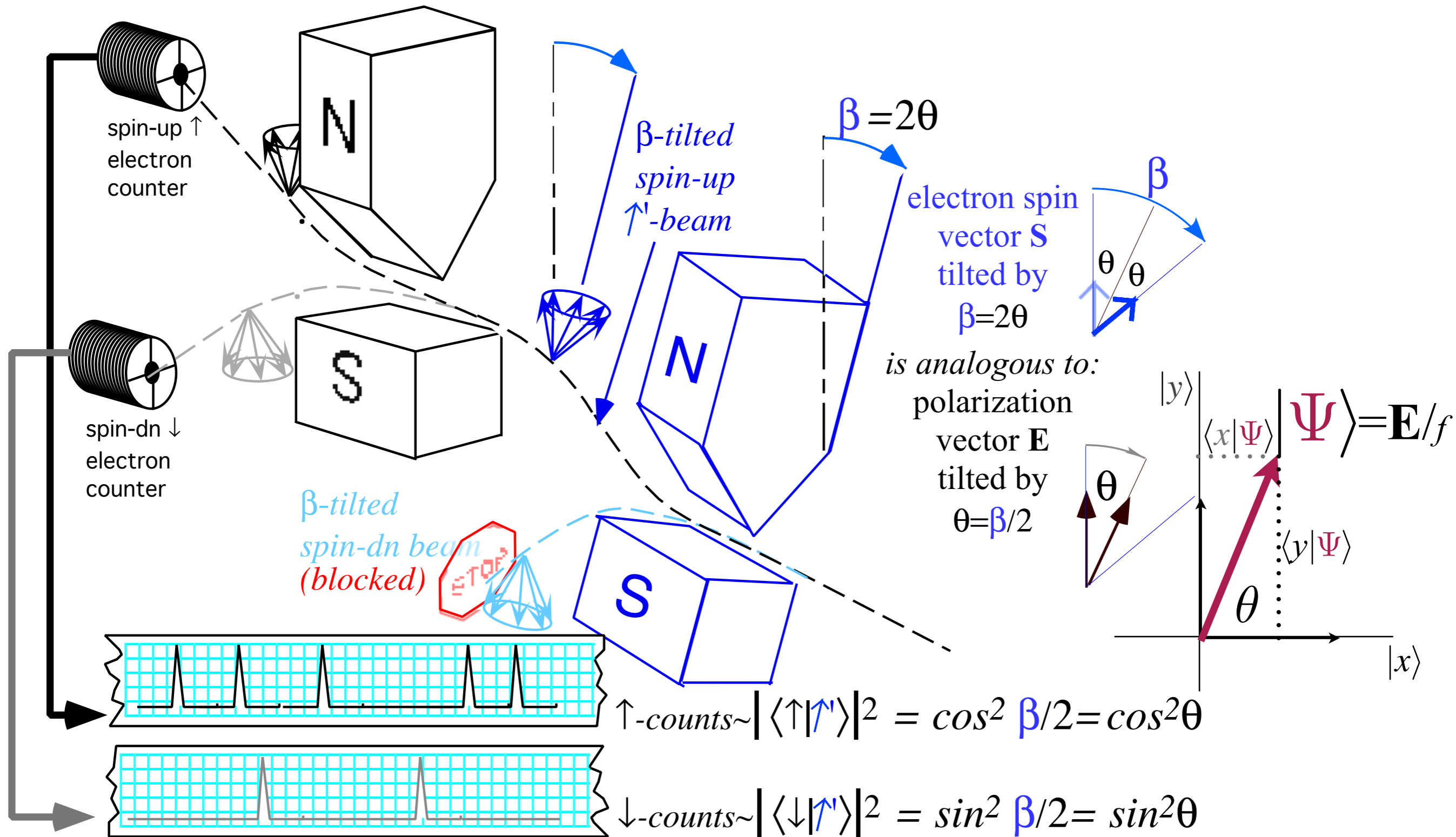


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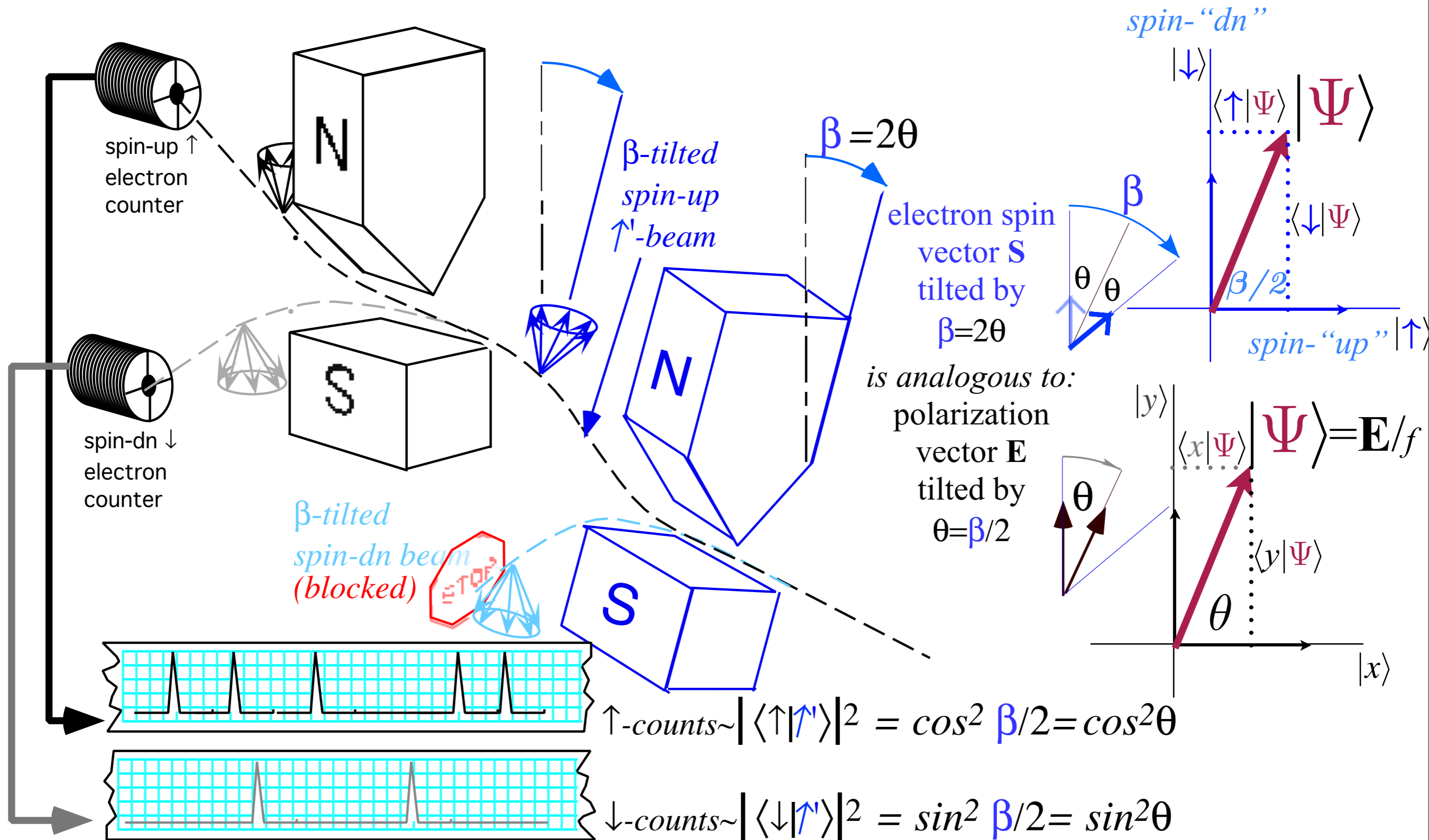


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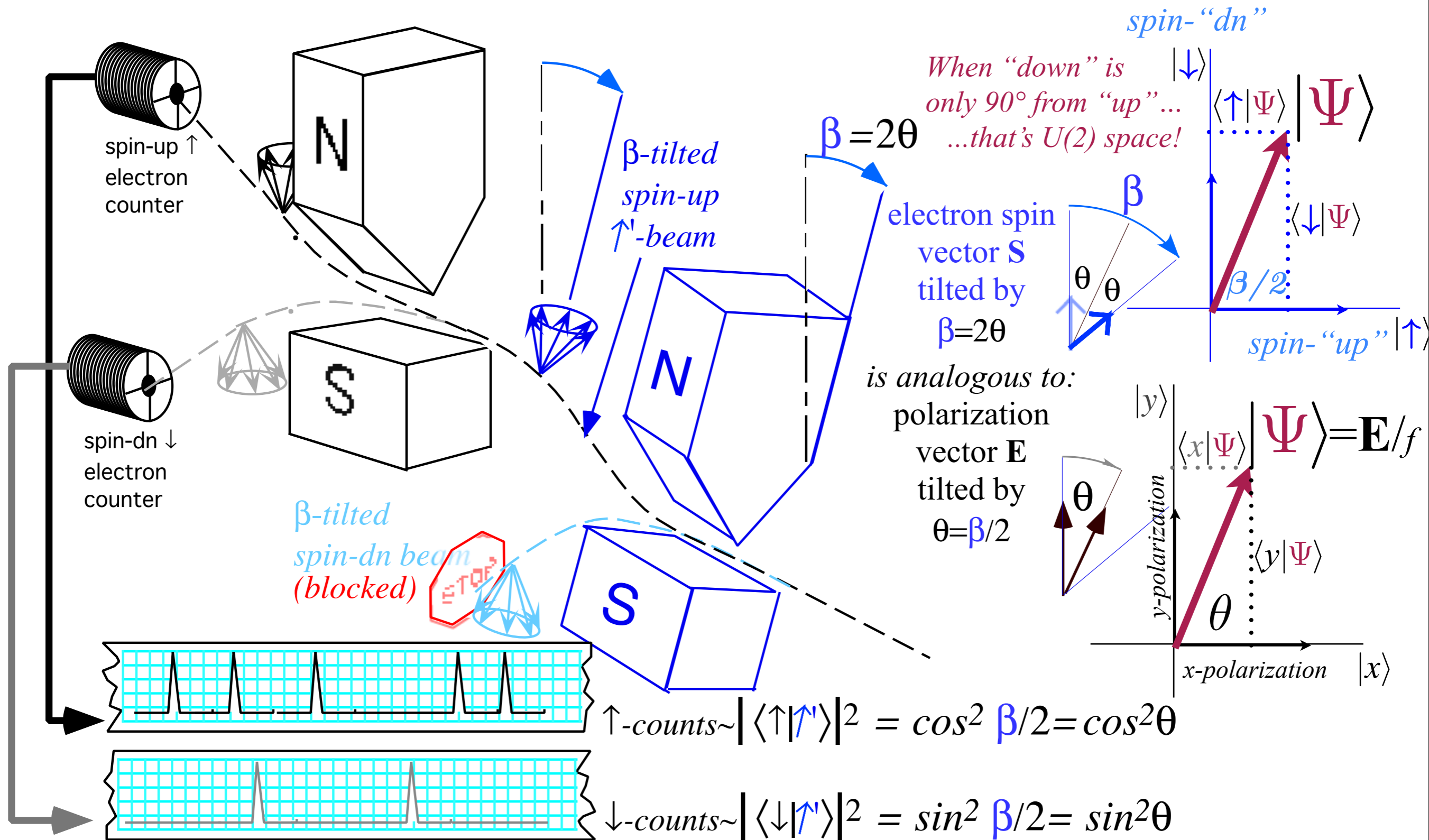


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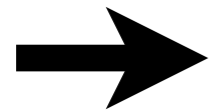
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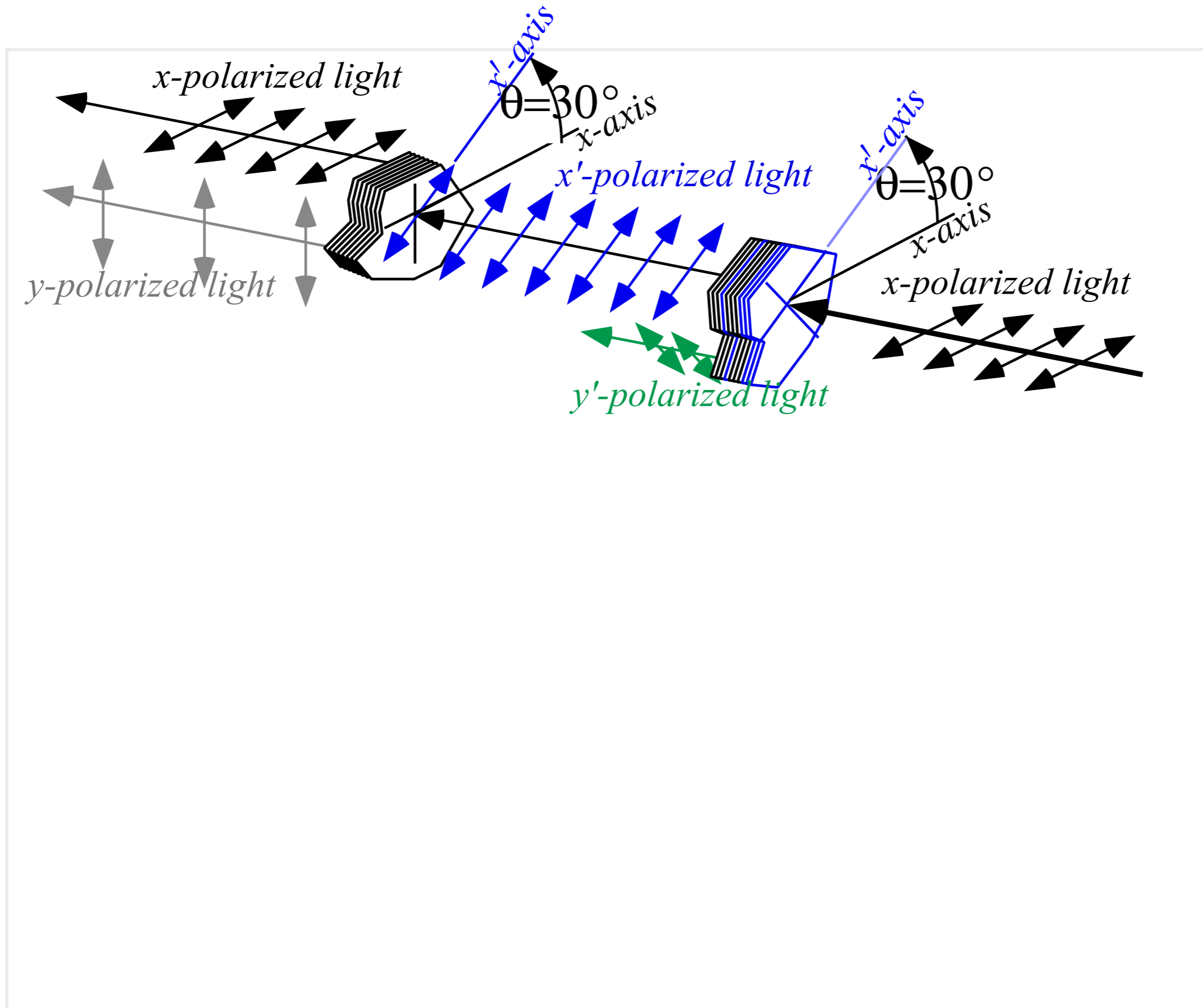
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A: Want to determine or calculate:

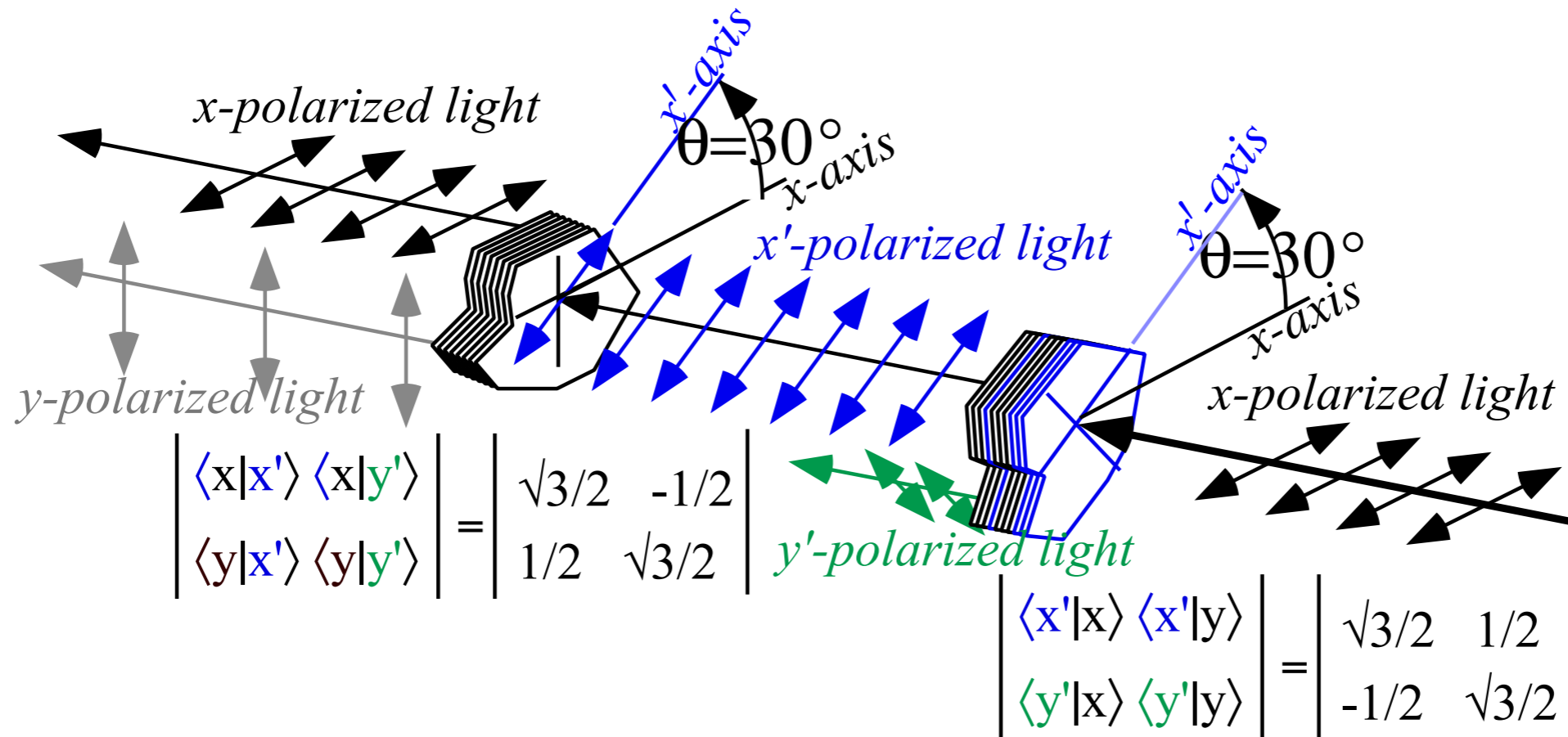
its *base state* $|b\rangle$, its *amplitude* $\langle b|\Psi\rangle$, and its *probability* $|\langle b|\Psi\rangle|^2$ using T -matrices



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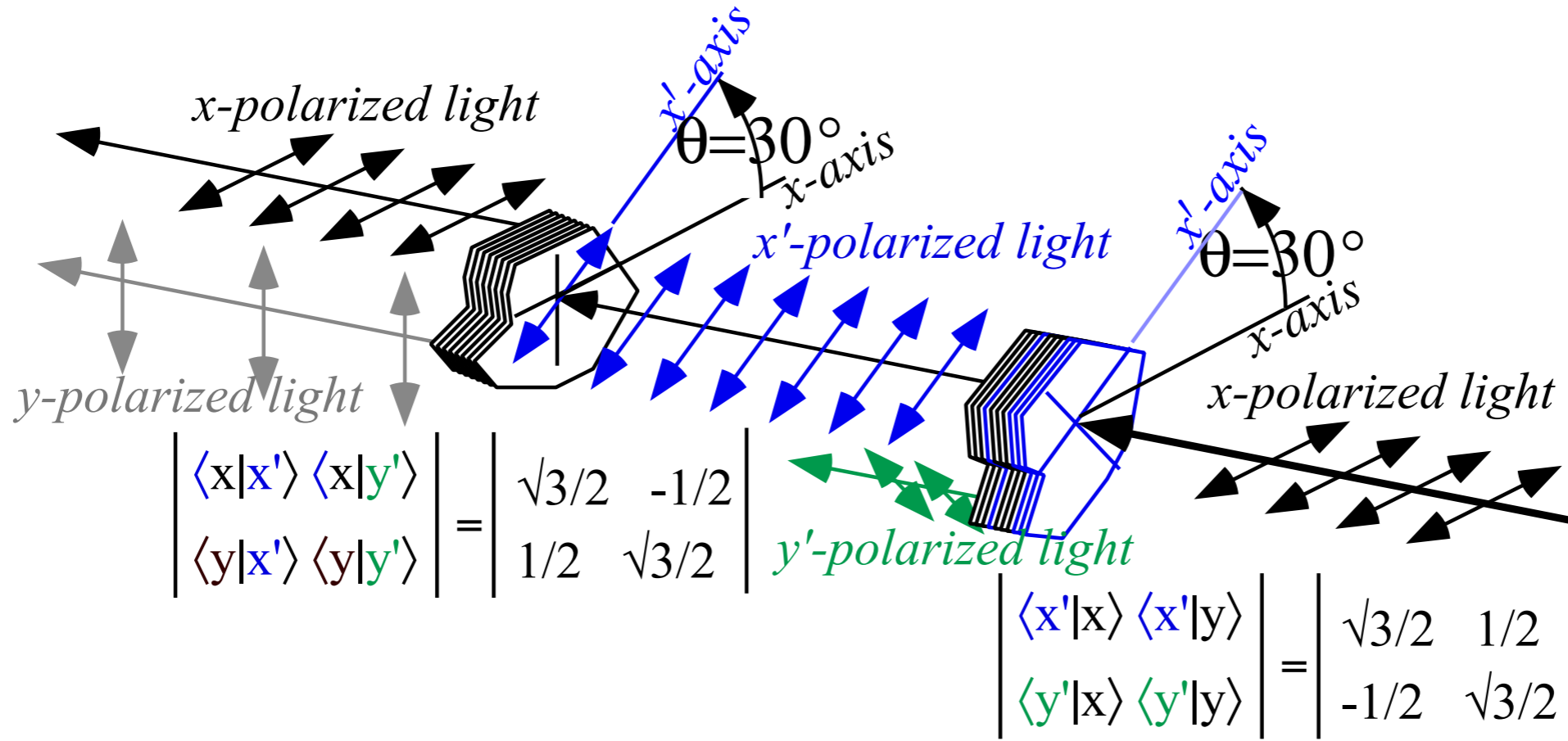
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$$\begin{vmatrix} \langle x|x'\rangle & \langle x|y'\rangle \\ \langle y|x'\rangle & \langle y|y'\rangle \end{vmatrix} = \begin{vmatrix} \sqrt{3}/2 & -1/2 \\ 1/2 & \sqrt{3}/2 \end{vmatrix}$$

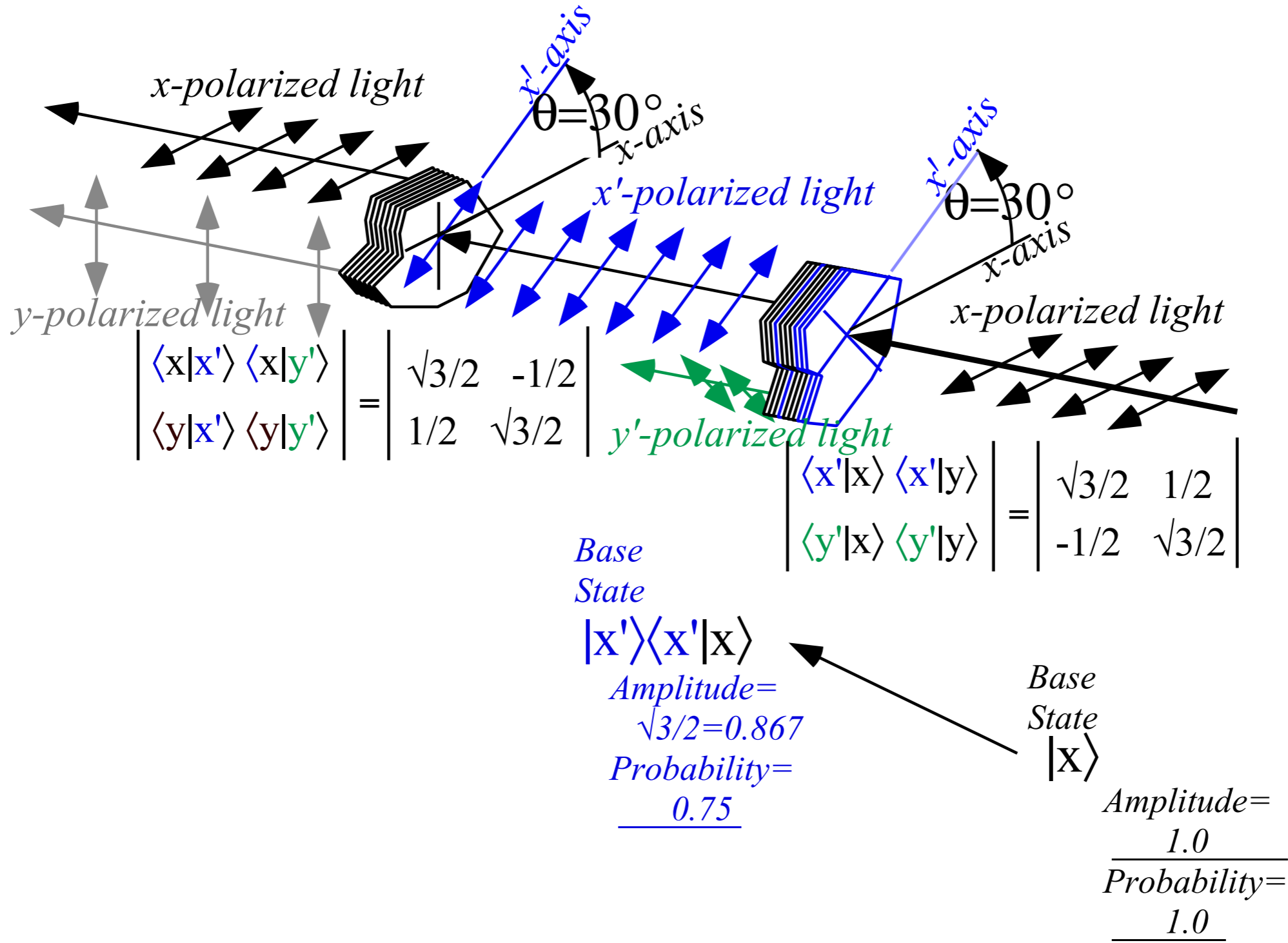
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Base State $|x\rangle$
 Amplitude = $\frac{1.0}{\quad}$
 Probability = $\frac{1.0}{\quad}$

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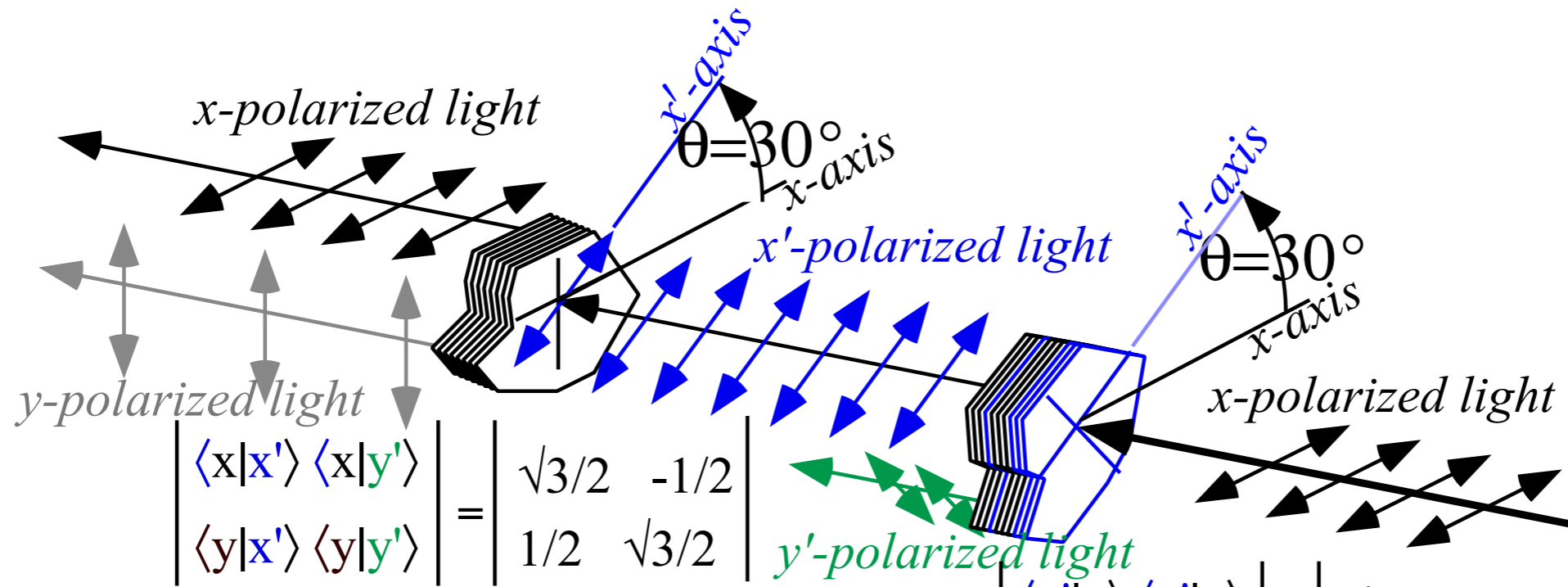
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Base State

$$|x'\rangle\langle x'|x\rangle$$

Amplitude = $\sqrt{3}/2 = 0.867$
 Probability = 0.75

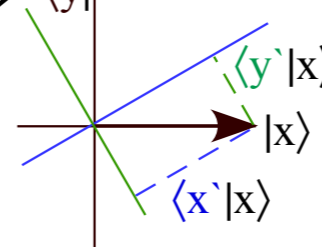
Base State

$$|y'\rangle\langle y'|x\rangle$$

Amplitude = $-1/2 = -0.500$
 Probability = 0.25

Base State $|x\rangle$

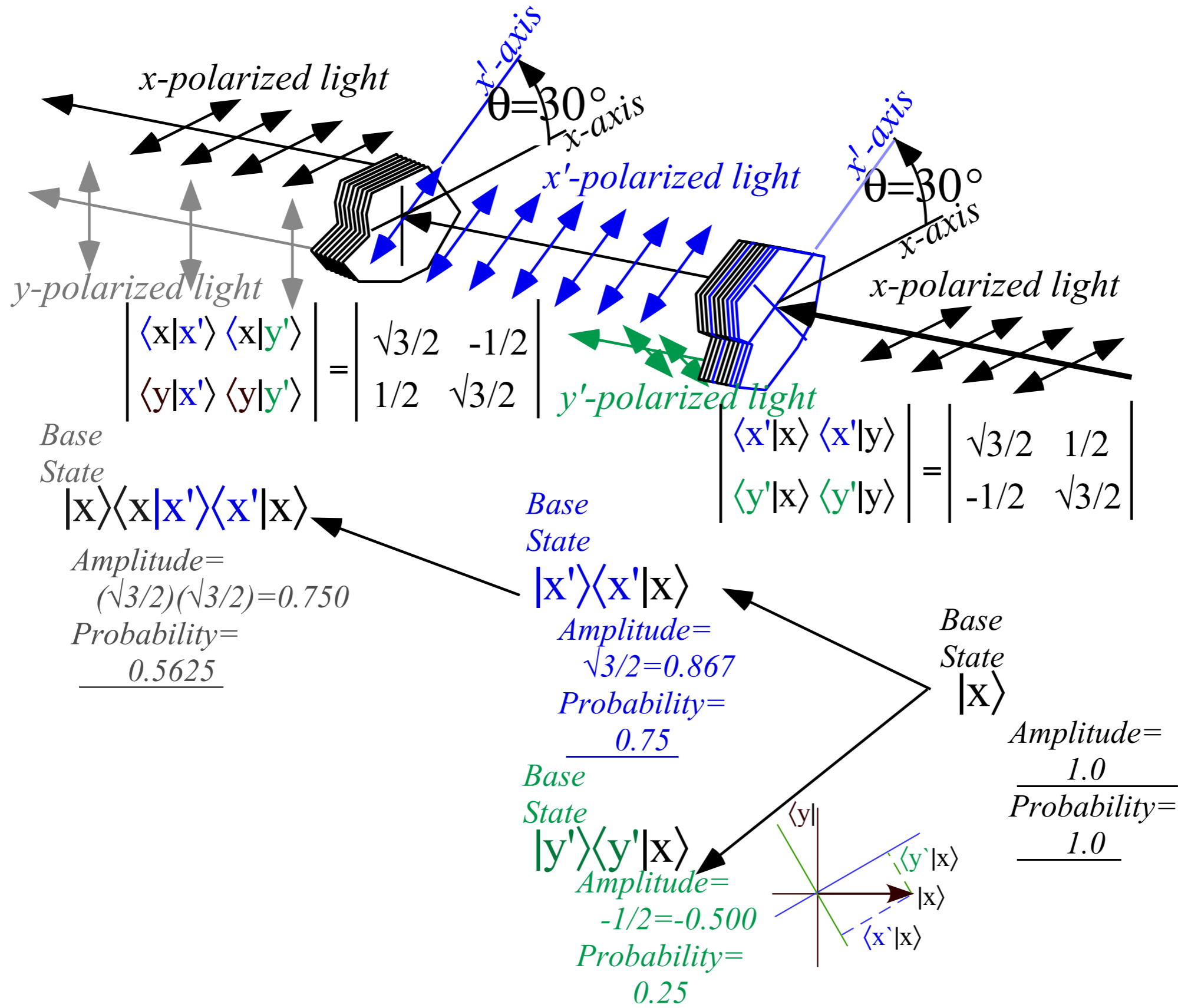
Amplitude = 1.0
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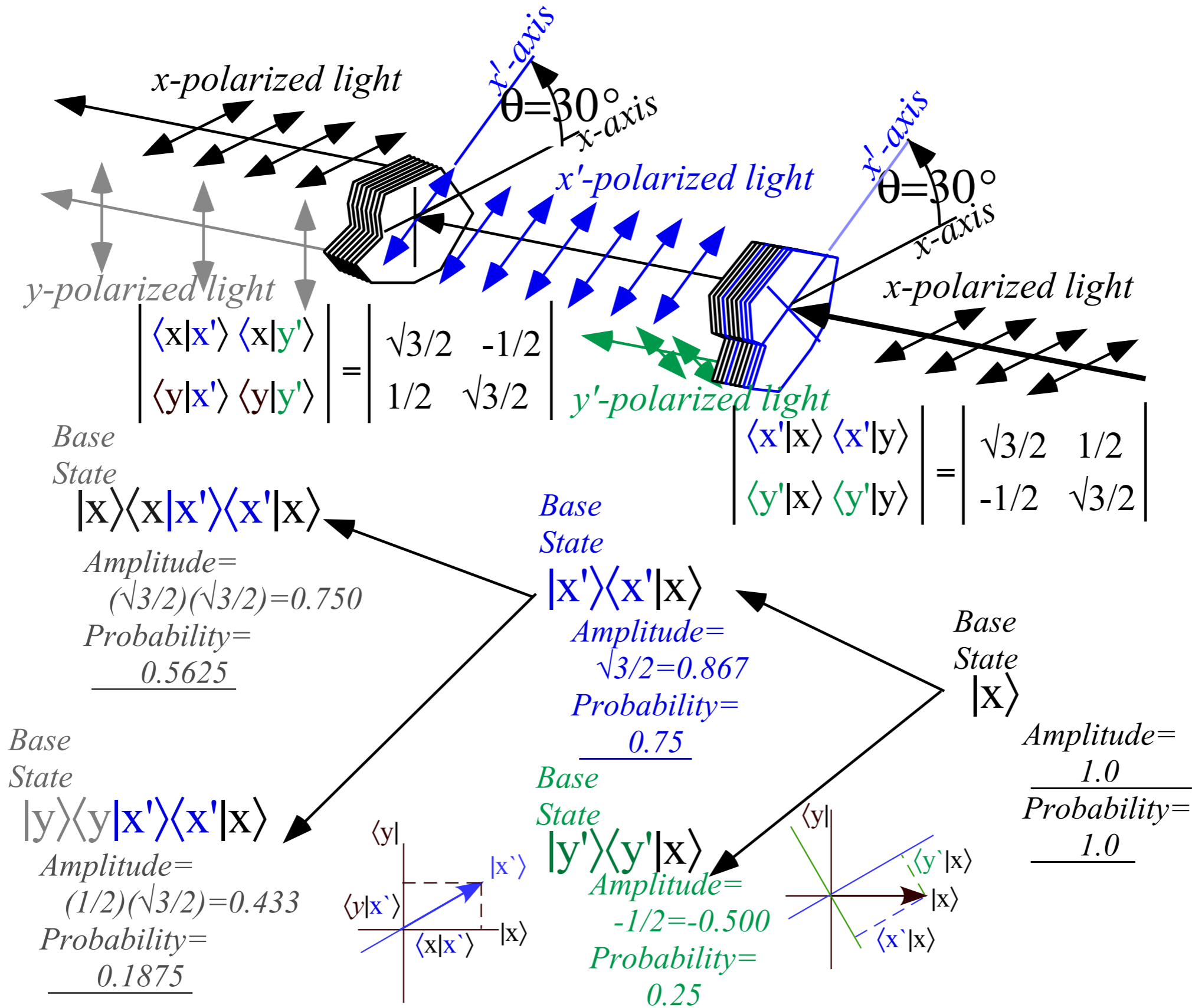
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$$\langle e_p | D_o C_n B_m | \Psi \rangle = \langle e_p | d_o \rangle \langle d_o | c_n \rangle \langle c_n | b_m \rangle \langle b_m | \Psi \rangle$$

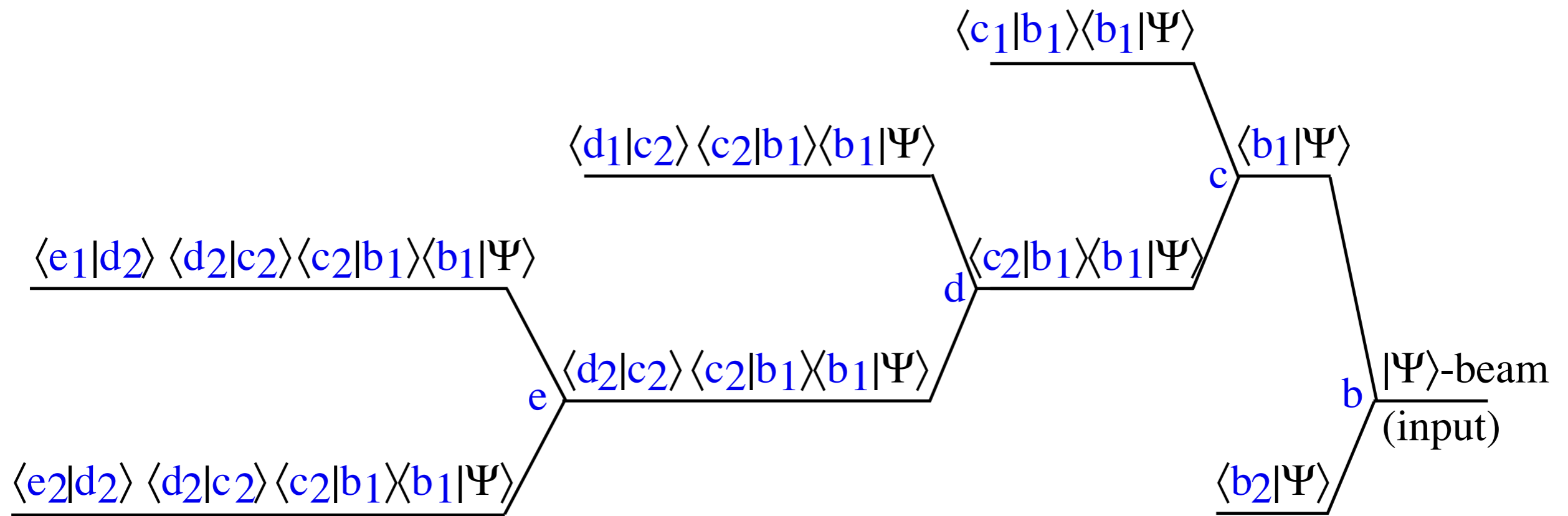


Fig. 1.3.10 Beams-amplitude products for successive beam sorting

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Sorter-counter, Filter, 1/2-wave plate, 1/4-wave plate

Feynman amplitude axiom 1 (Given above p.35)

(1) The probability axiom

The first axiom deals with physical interpretation of amplitudes $\langle j|k' \rangle$.

Axiom 1: The absolute square $|\langle j|k' \rangle|^2 = \langle j|k' \rangle^ \langle j|k' \rangle$ gives probability for occurrence in state- j of a system that started in state- $k'=1',2',\dots,$ or n' from one sorter and then was forced to choose between states $j=1,2,\dots,n$ by another sorter.*

*Feynman-Dirac
Interpretation of*

$$\langle j|k' \rangle$$

*=Amplitude of state- j after
state- k' is forced to choose
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Feynman amplitude axioms 1-2

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(2) The conjugation or inversion axiom (time reversal symmetry)

The second axiom concerns going backwards through a sorter or the reversal of amplitudes.

Axiom 2: The complex conjugate $\langle j|k' \rangle^$ of an amplitude $\langle j|k' \rangle$ equals its reverse: $\langle j|k' \rangle^* = \langle k'|j \rangle$*

Feynman amplitude axioms 1-3

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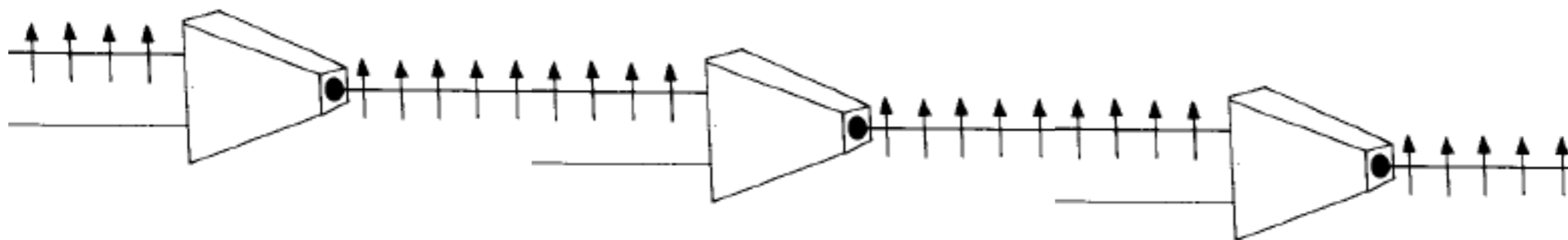
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Review: “Abstraction” of bra and ket vectors from a Transformation Matrix

Introducing scalar and matrix products

Planck's energy and N-quanta (Cavity/Beam wave mode)

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Feynman amplitude axiom 1

What comes out of a beam sorter channel or branch-b?

Sample calculations

Feynman amplitude axioms 2-3

 *Beam analyzers: Sorter-unsorters*

The “Do-Nothing” analyzer

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Some “Do-Something” analyzers

Sorter-counter, Filter, 1/2-wave plate, 1/4-wave plate

Beam Analyzers: Fundamental Quantum Processes (Sorter-Unsorter pairs)

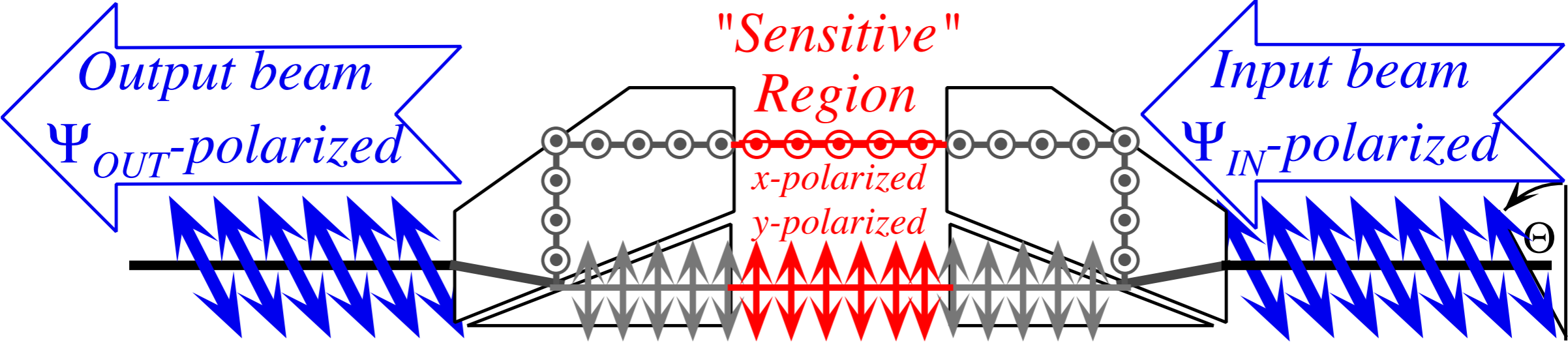


Fig. 1.3.1 Anatomy of ideal optical polarization analyzer in its "Do nothing" mode

Beam Analyzers: Fundamental Quantum Processes (Sorter-Unsorter pairs)

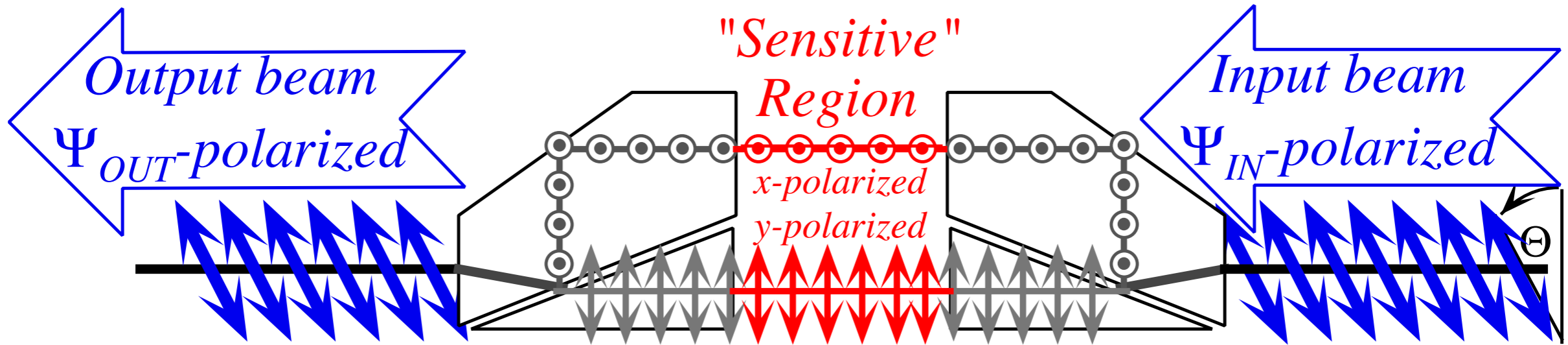


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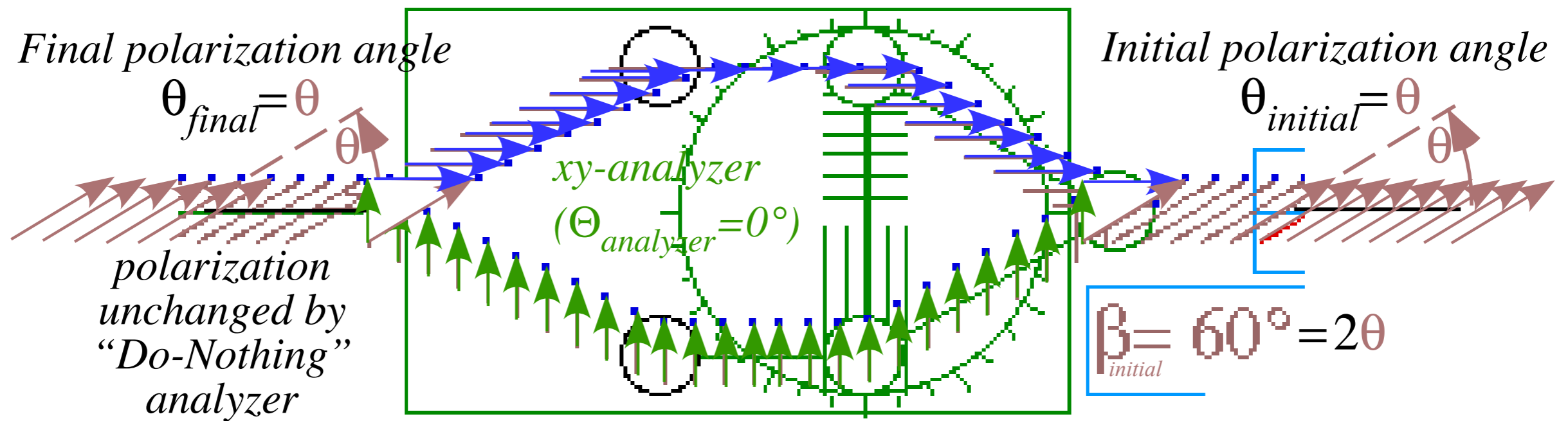


Fig. 1.3.2 Computer sketch of simulated polarization analyzer in "do-nothing" mode

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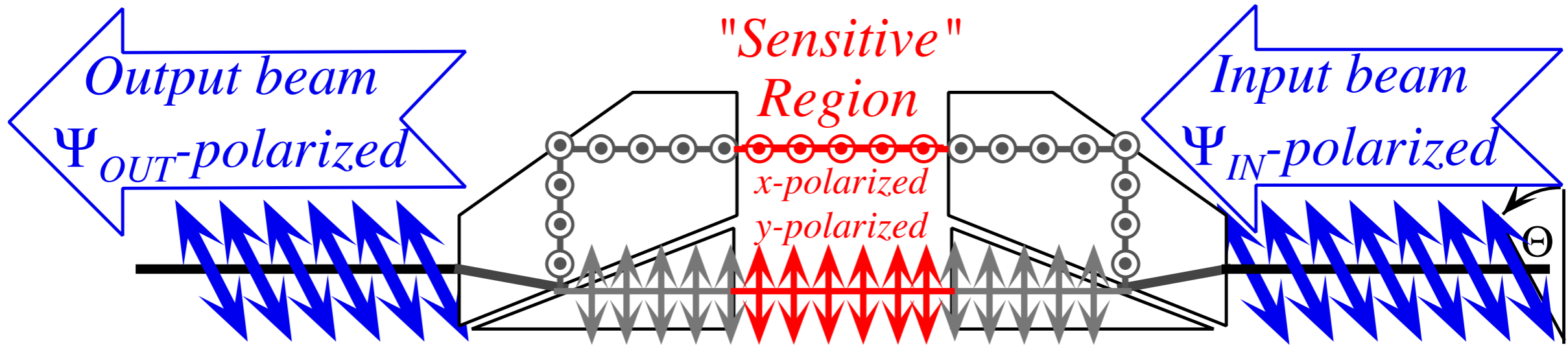


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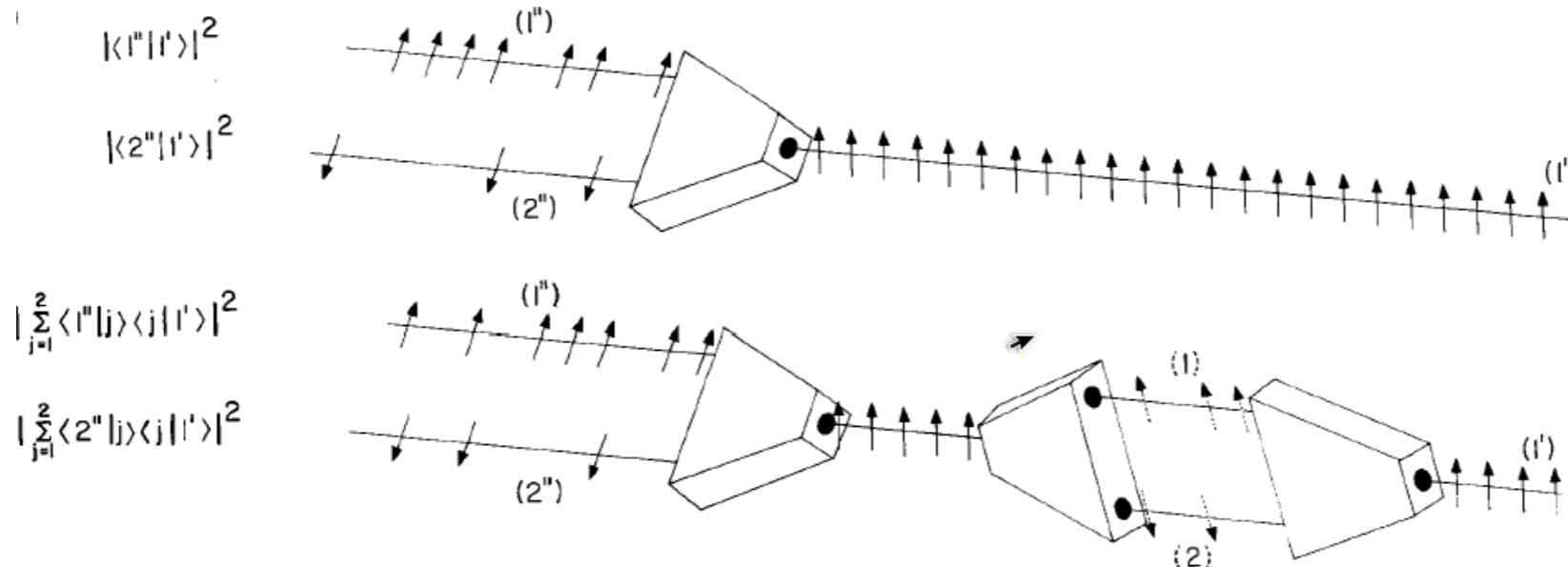
Feynman amplitude axiom 4

(4) The completeness or closure axiom

The fourth axiom concerns the "Do-nothing" property of an ideal analyzer, that is, a sorter followed by an "unsorter" or "put-back-togetherer" as sketched above.

Axiom 4. Ideal sorting followed by ideal recombination of amplitudes has no effect:

$$\langle j'' | m' \rangle = \sum_{k=1}^n \langle j'' | k \rangle \langle k | m' \rangle$$



Beam Analyzers: Fundamental Quantum Processes (Sorter-Unsorter pairs)

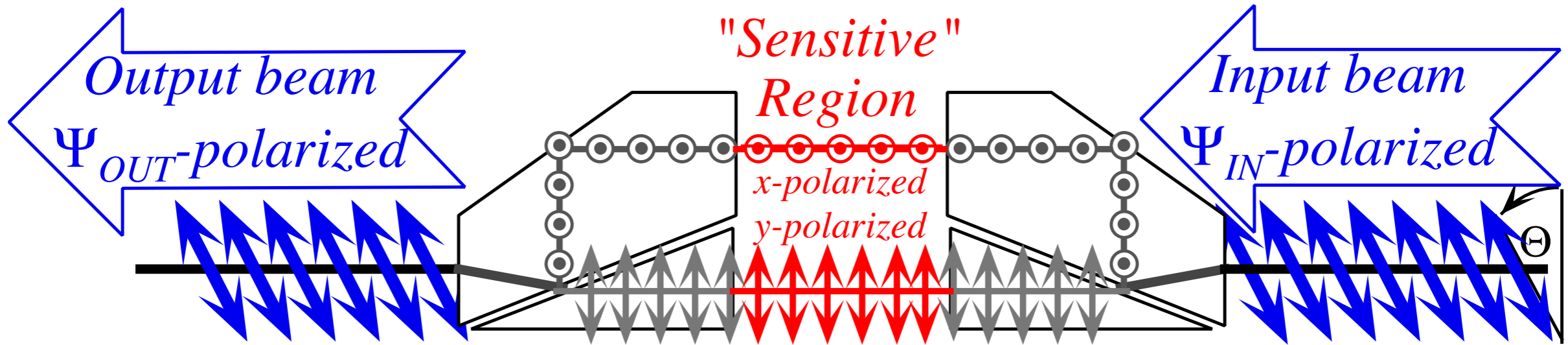


Fig. 1.3.1 Anatomy of ideal optical polarization analyzer in its "Do nothing" mode

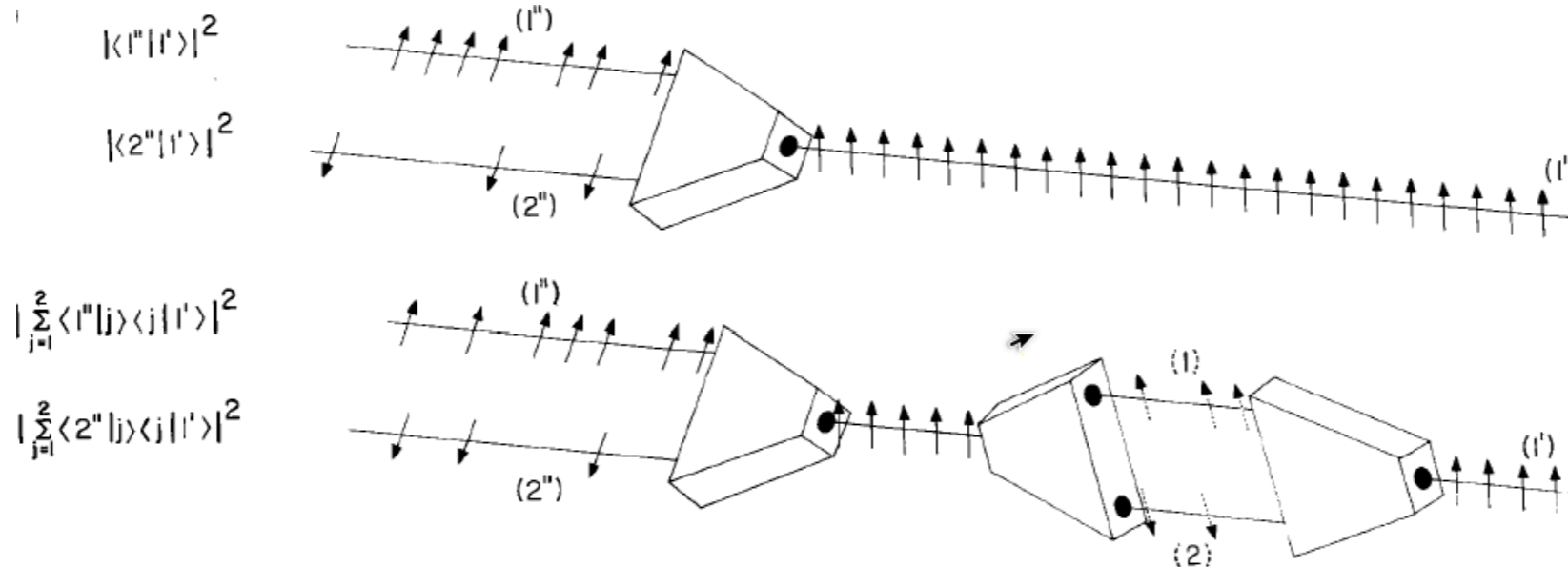
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May use axioms 1-3 to prove special case:

$$1 = \langle m' | m' \rangle \text{ by 3}$$

\Downarrow by 1

$$1 = \sum_{k=1}^n |\langle k | m' \rangle|^2$$

\Downarrow by 1

$$1 = \sum_{k=1}^n \langle k | m' \rangle^* \langle k | m' \rangle$$

\Downarrow by 2

$$\sum_{k=1}^n \langle m' | k \rangle \langle k | m' \rangle = \langle m' | m' \rangle$$

Feynman amplitude axioms 1-4

Feynman-Dirac
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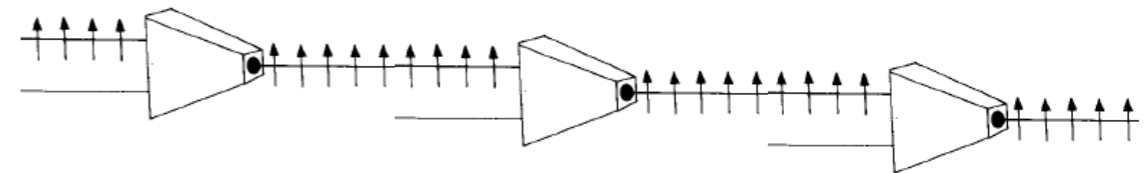
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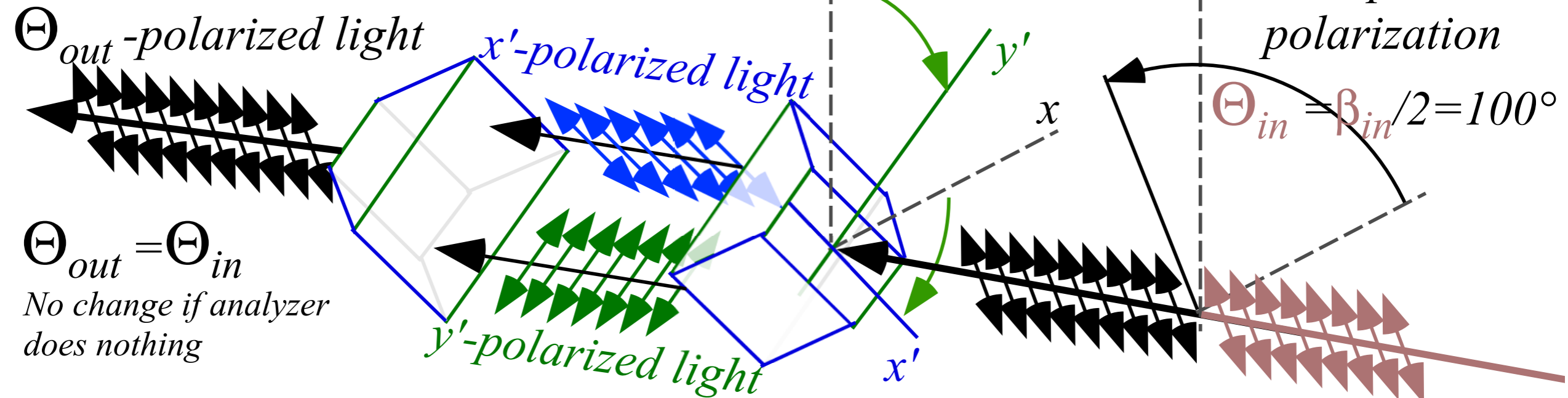
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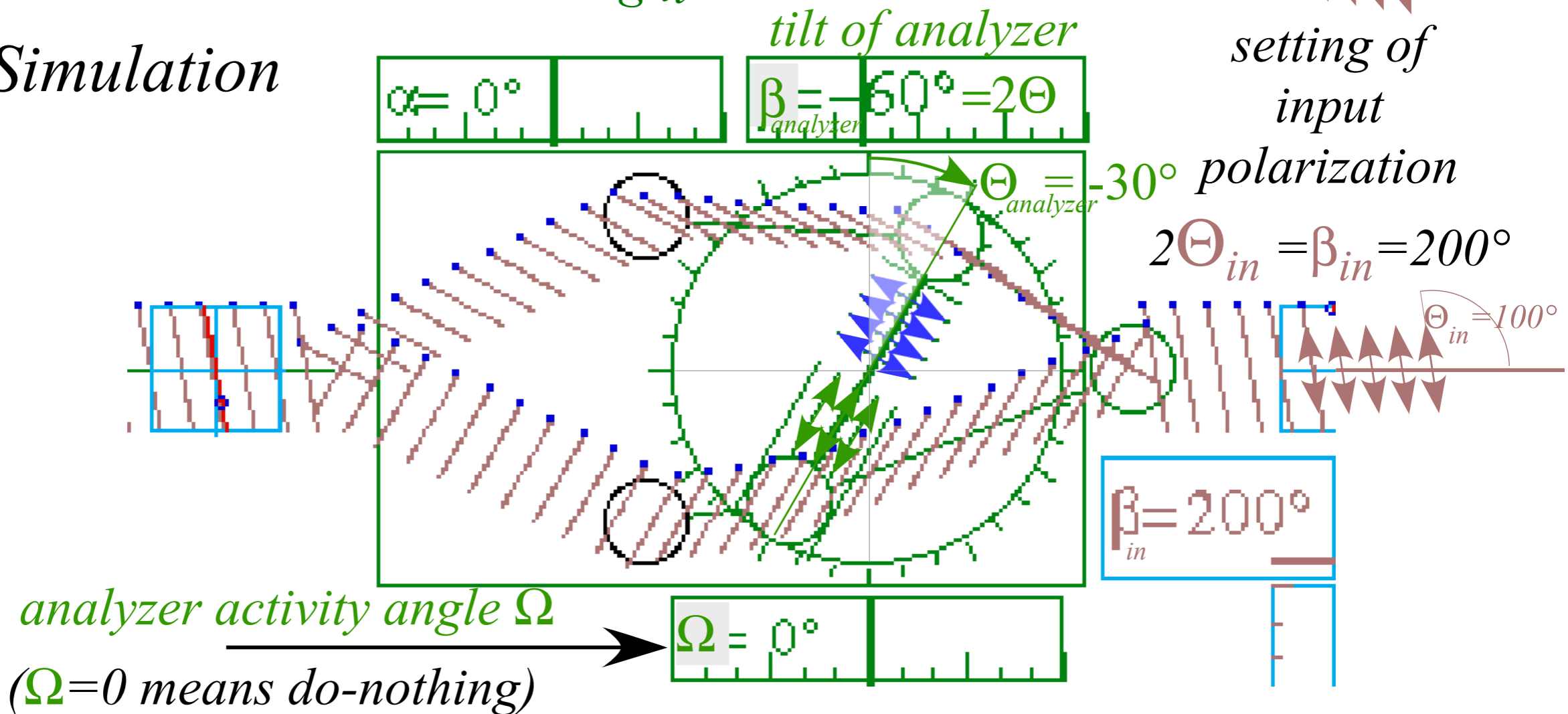
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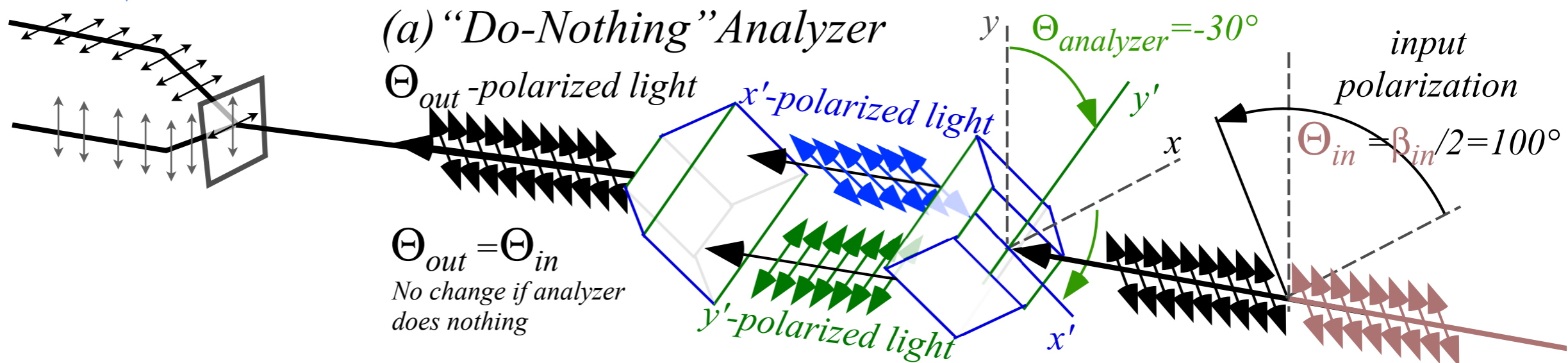
(a) "Do-Nothing" Analyzer



(b) Simulation



Imagine final xy -sorter analyzes output beam into x and y -components.



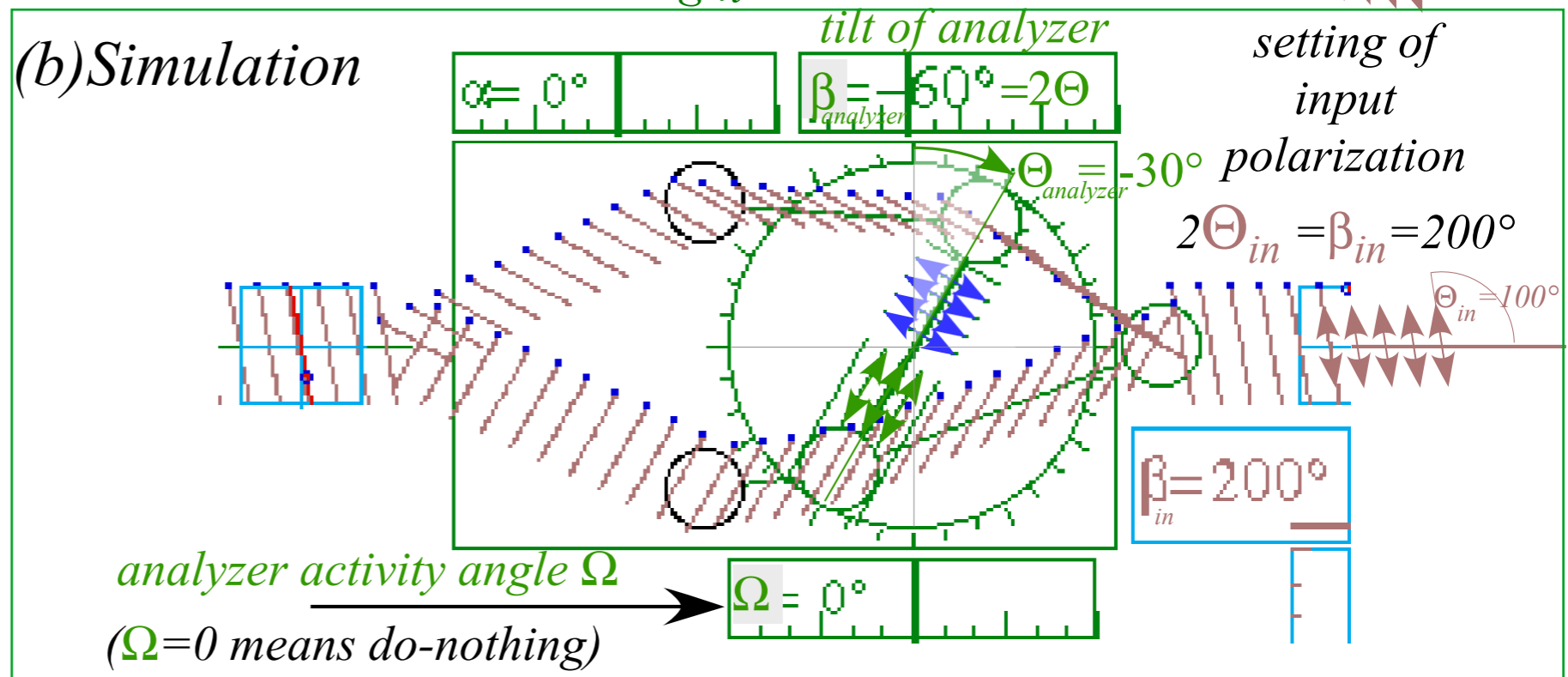
Amplitude in x or y -channel is sum over x' and y' -amplitudes

$\langle x' | \Theta_{in} \rangle = \cos(\Theta_{in} - \Theta)$

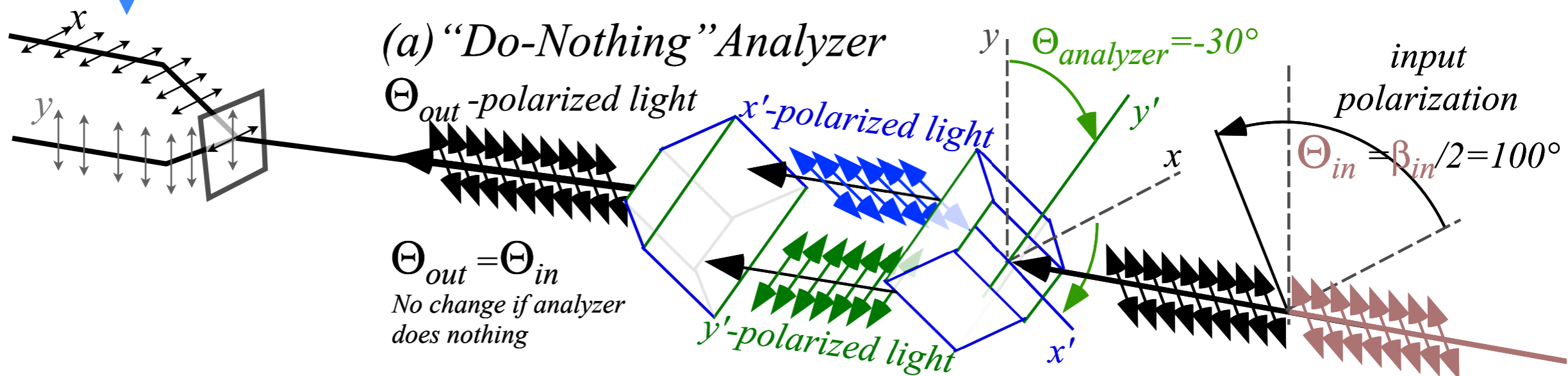
$\langle y' | \Theta_{in} \rangle = \sin(\Theta_{in} - \Theta)$

with relative angle $\Theta_{in} - \Theta$

of Θ_{in} to Θ -analyzer axes- (x', y')



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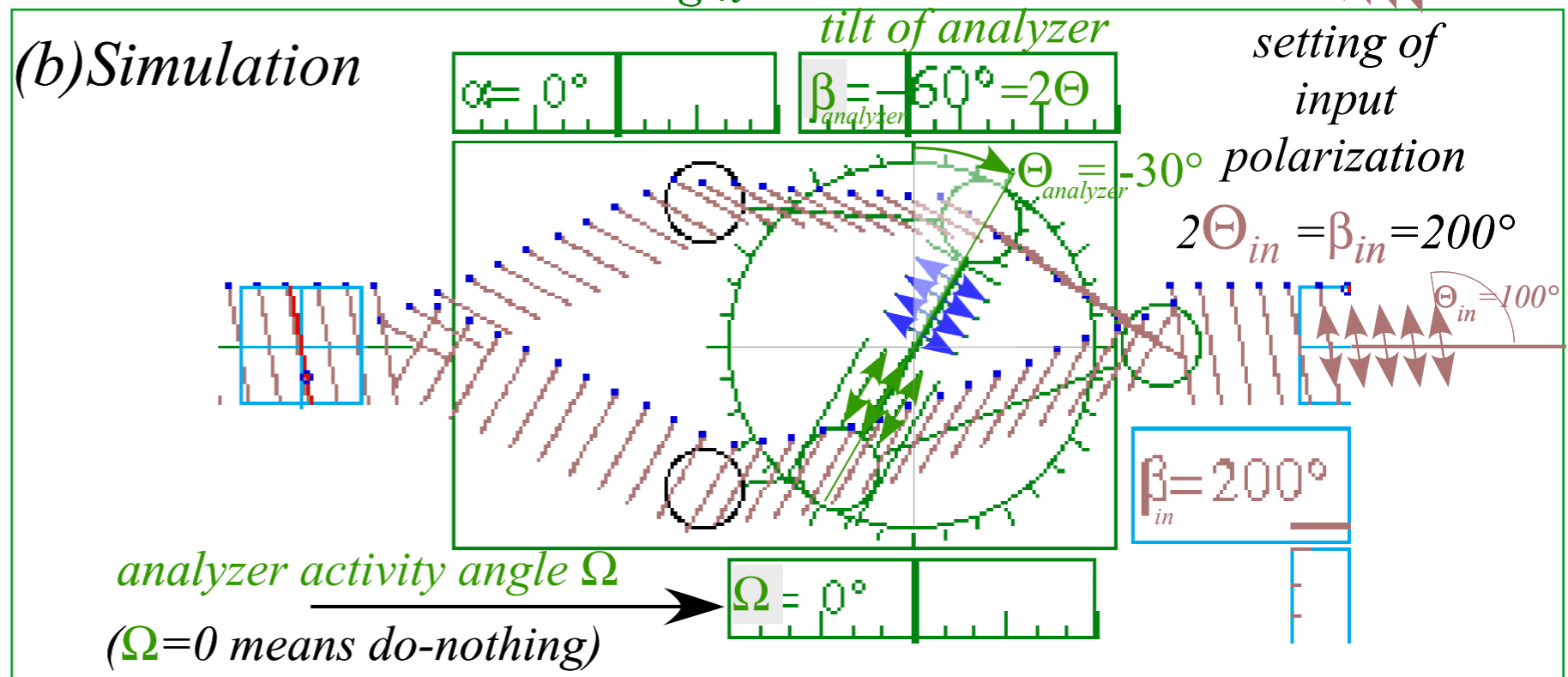
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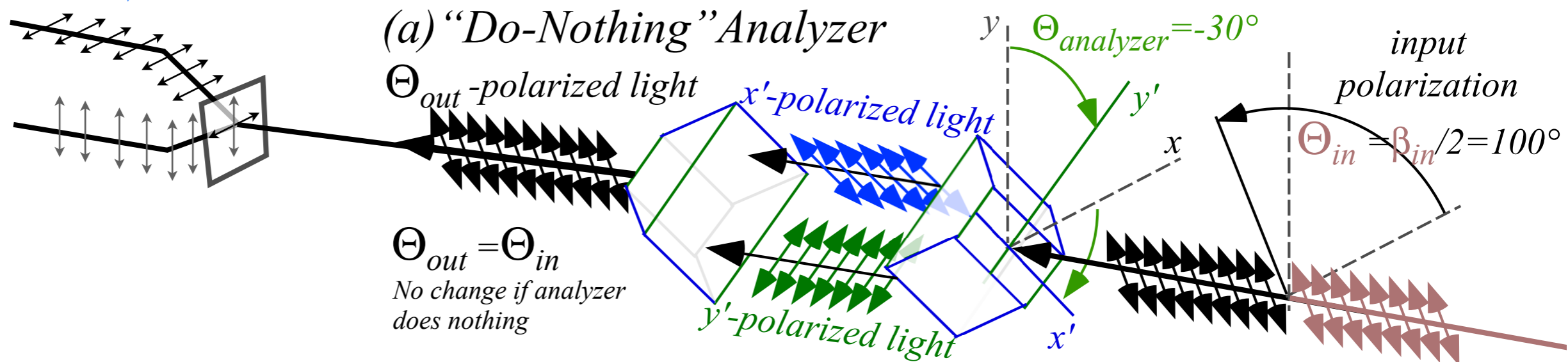
in products with final xy -sorter:

lab x -axis: $\langle x | x' \rangle = \cos \Theta = \langle y | y' \rangle$

y -axis: $\langle y | x' \rangle = \sin \Theta = -\langle x | y' \rangle$.



Imagine final xy -sorter analyzes output beam into x and y -components.



Amplitude in x or y -channel is sum over x' and y' -amplitudes

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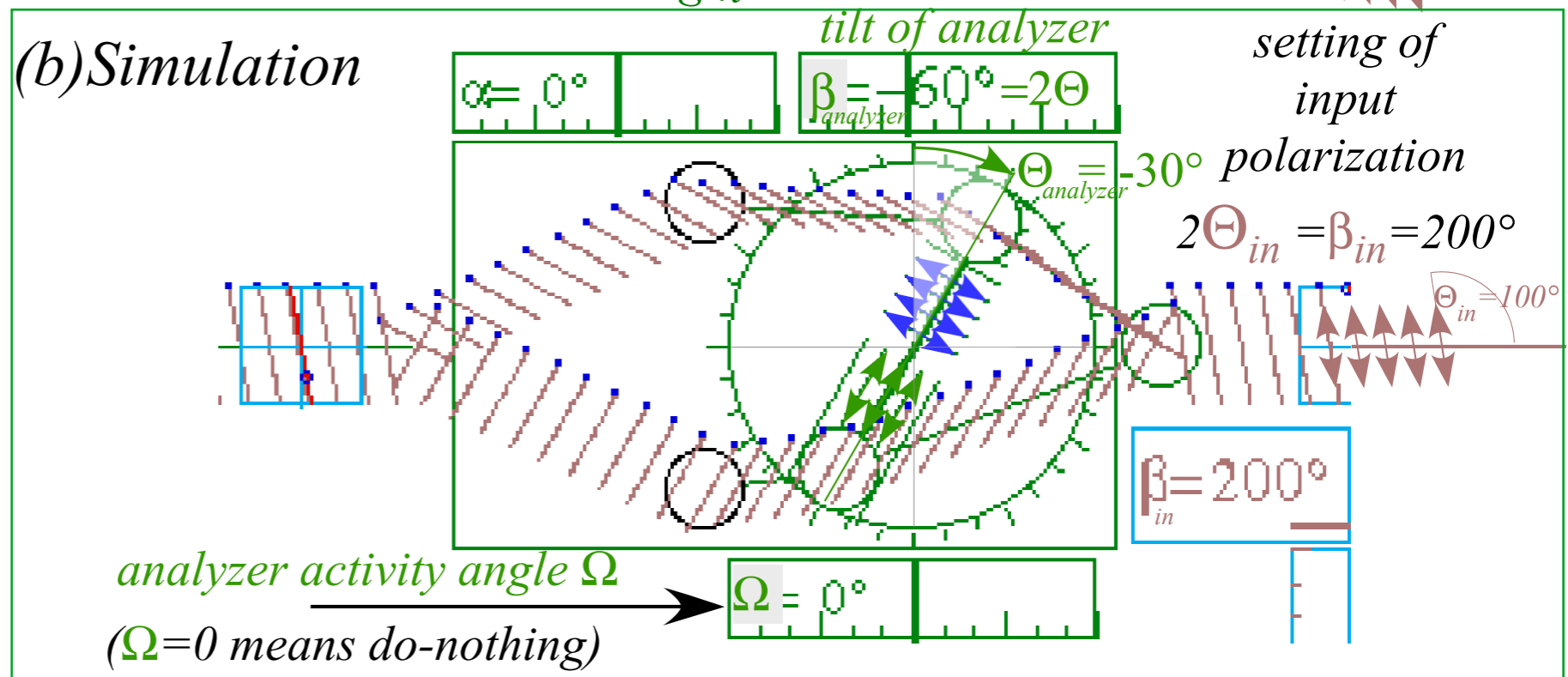
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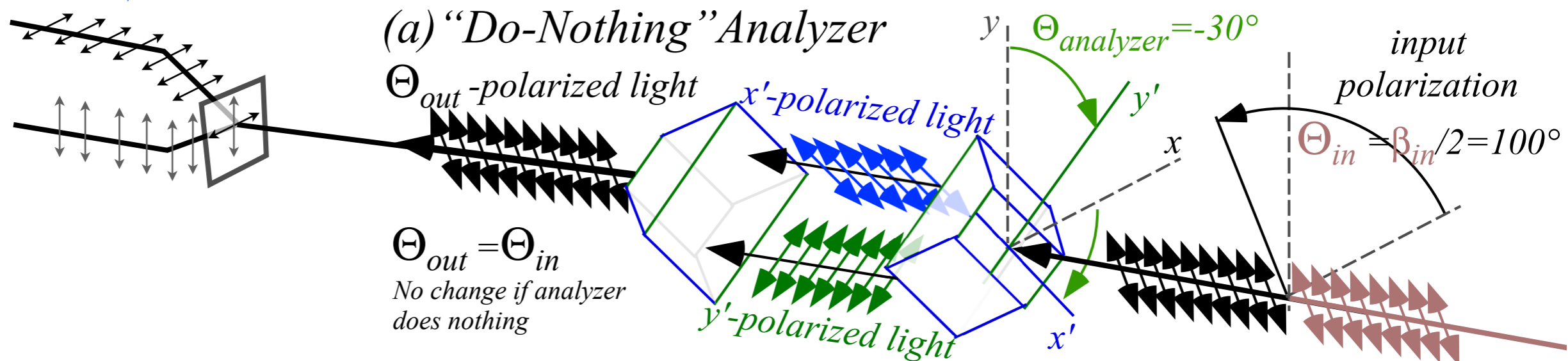


x -Output is: $\langle x | \Theta_{out} \rangle = \langle x | x' \rangle \langle x' | \Theta_{in} \rangle + \langle x | y' \rangle \langle y' | \Theta_{in} \rangle = \cos \Theta \cos(\Theta_{in} - \Theta) - \sin \Theta \sin(\Theta_{in} - \Theta) = \cos \Theta_{in}$

y -Output is: $\langle y | \Theta_{out} \rangle = \langle y | x' \rangle \langle x' | \Theta_{in} \rangle + \langle y | y' \rangle \langle y' | \Theta_{in} \rangle = \sin \Theta \cos(\Theta_{in} - \Theta) - \cos \Theta \sin(\Theta_{in} - \Theta) = \sin \Theta_{in}$.

(Recall $\cos(a+b) = \cos a \cos b - \sin a \sin b$ and $\sin(a+b) = \sin a \cos b + \cos a \sin b$)

Imagine final xy -sorter analyzes output beam into x and y -components.



Amplitude in x or y -channel is sum over x' and y' -amplitudes

$$\langle x' | \Theta_{in} \rangle = \cos(\Theta_{in} - \Theta)$$

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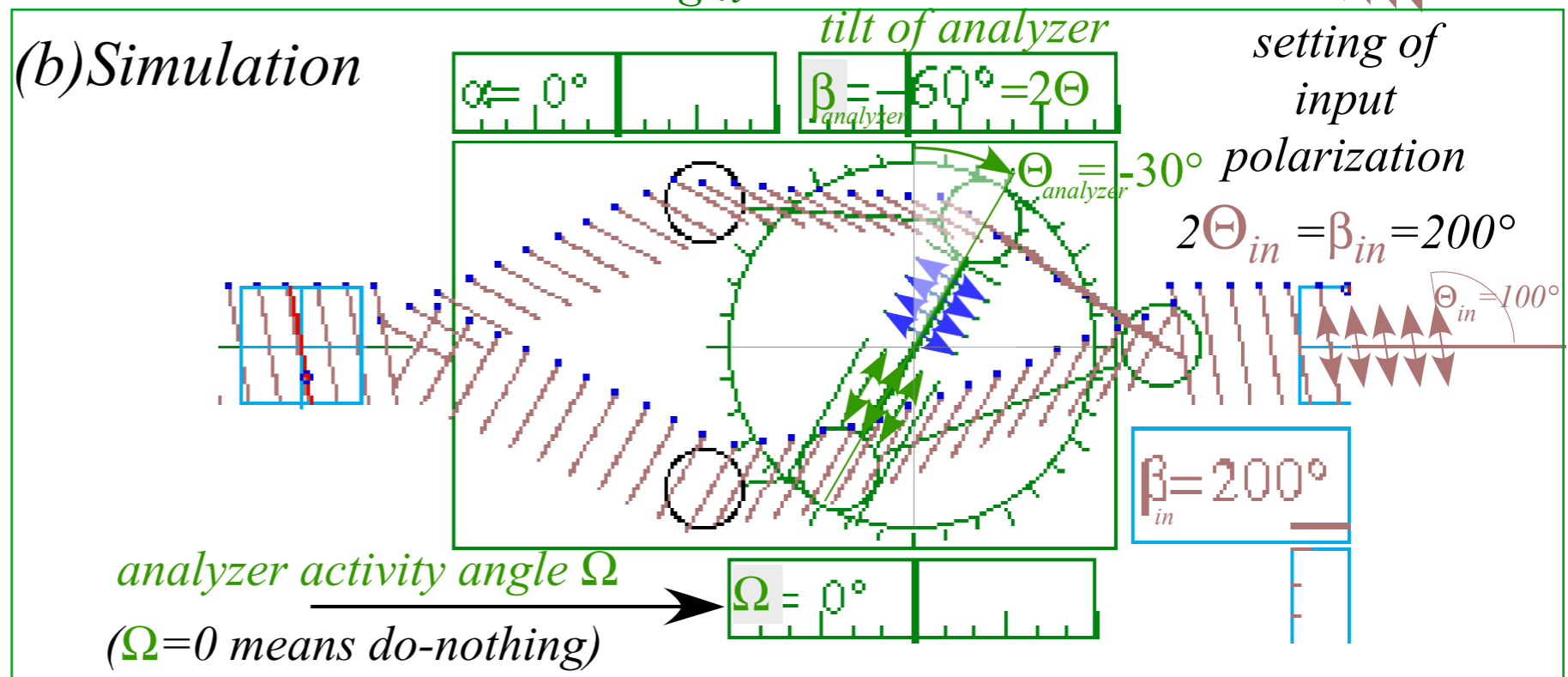
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$$y\text{-Output is: } \langle y | \Theta_{out} \rangle = \langle y | x' \rangle \langle x' | \Theta_{in} \rangle + \langle y | y' \rangle \langle y' | \Theta_{in} \rangle = \sin \Theta \cos(\Theta_{in} - \Theta) - \cos \Theta \sin(\Theta_{in} - \Theta) = \sin \Theta_{in}.$$

(Recall $\cos(a+b) = \cos a \cos b - \sin a \sin b$ and $\sin(a+b) = \sin a \cos b + \cos a \sin b$)

Conclusion:

$$\langle x | \Theta_{out} \rangle = \cos \Theta_{out} = \cos \Theta_{in} \text{ or: } \Theta_{out} = \Theta_{in} \text{ so "Do-Nothing" Analyzer in fact does nothing.}$$

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The “Do-Nothing” analyzer

Feynman amplitude axiom 4

Some “Do-Something” analyzers

 *Sorter-counter, Filter, 1/2-wave plate, 1/4-wave plate*

(1) Optical analyzer in sorter-counter configuration

Analyzer reduced to a simple sorter-counter by blocking output of x -high-road and y -low-road with counters

$$x\text{-counts} \sim |\langle x|x' \rangle|^2$$

$$= \cos^2 \theta = 0.75$$

$$y\text{-counts} \sim |\langle y|x' \rangle|^2$$

$$= \sin^2 \theta = 0.25$$

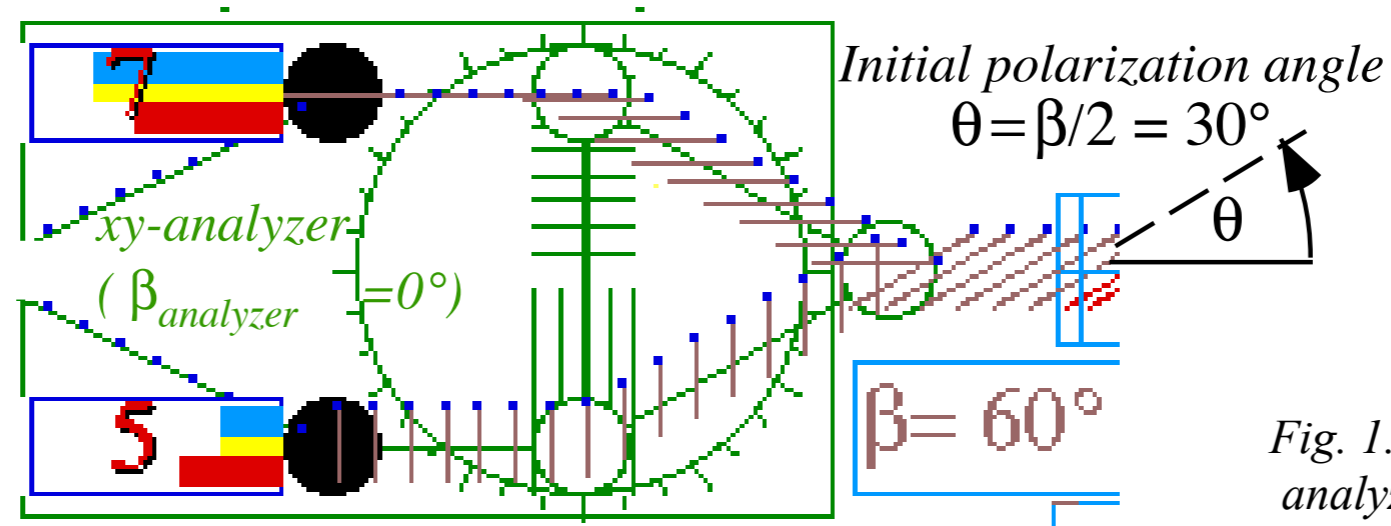


Fig. 1.3.3 Simulated polarization analyzer set up as a sorter-counter

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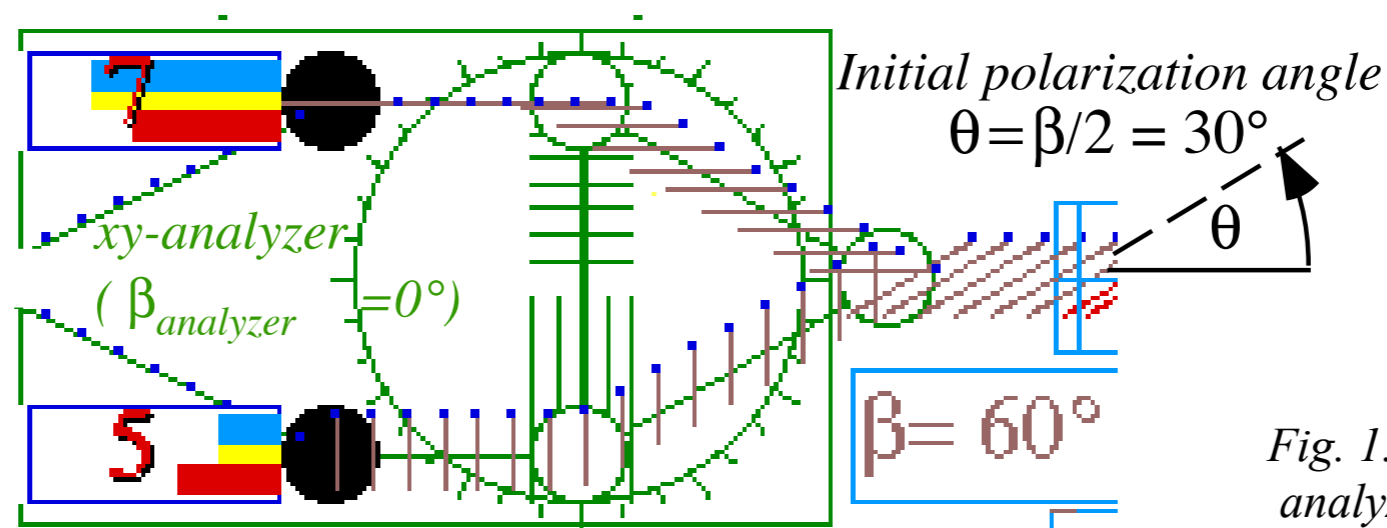


Fig. 1.3.3 Simulated polarization analyzer set up as a sorter-counter

(2) Optical analyzer in a filter configuration (Polaroid© sunglasses)

Analyzer blocks one path which may have photon counter without affecting function.

$$x\text{-counts} \sim |\langle y|x' \rangle|^2 = 0.75$$

(Blocked and filtered out)

$$y\text{-output} \sim |\langle y|x' \rangle|^2 = \sin^2 \theta = 0.25$$

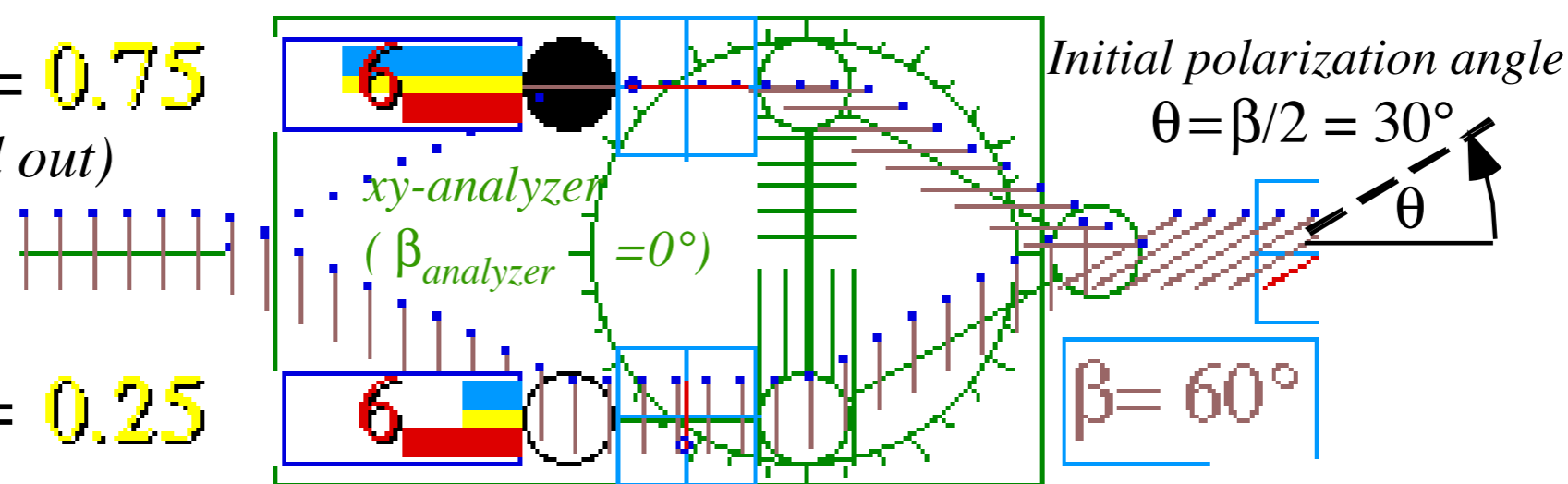


Fig. 1.3.4 Simulated polarization analyzer set up to filter out the x -polarized photons

Review: “Abstraction” of bra and ket vectors from a Transformation Matrix

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Planck's energy and N-quanta (Cavity/Beam wave mode)

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Sample calculations

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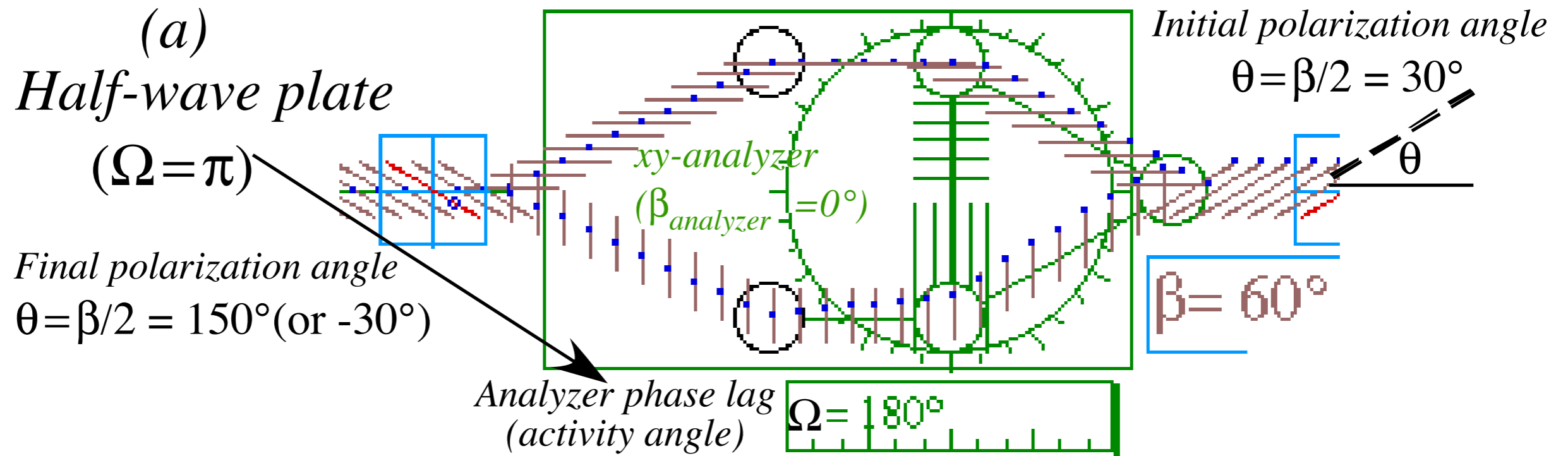
Feynman amplitude axiom 4

Some “Do-Something” analyzers

Sorter-counter, Filter, 1/2-wave plate, 1/4-wave plate



(3) Optical analyzers in the "control" configuration: Half or Quarter wave plates



(3) Optical analyzers in the "control" configuration: Half or Quarter wave plates

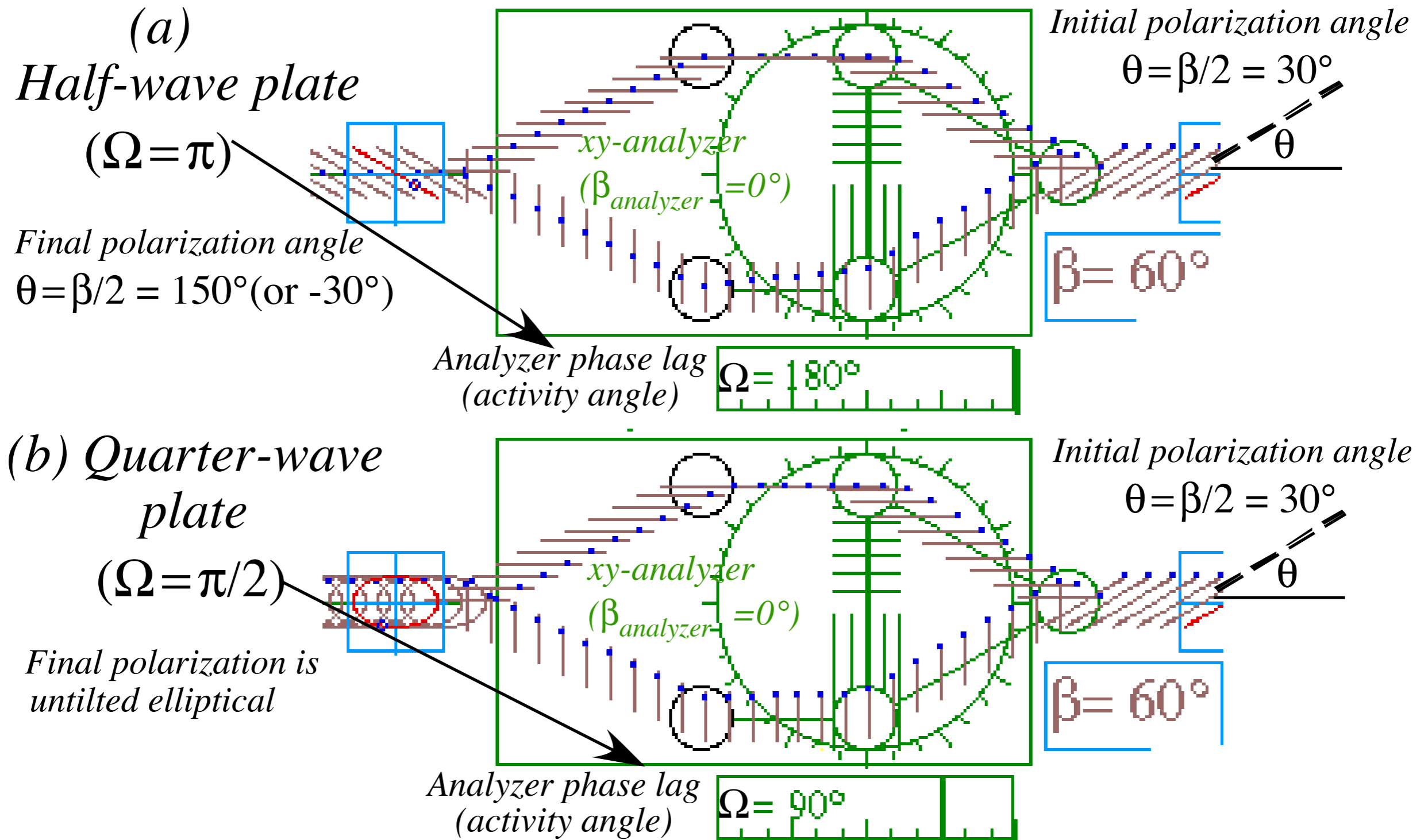
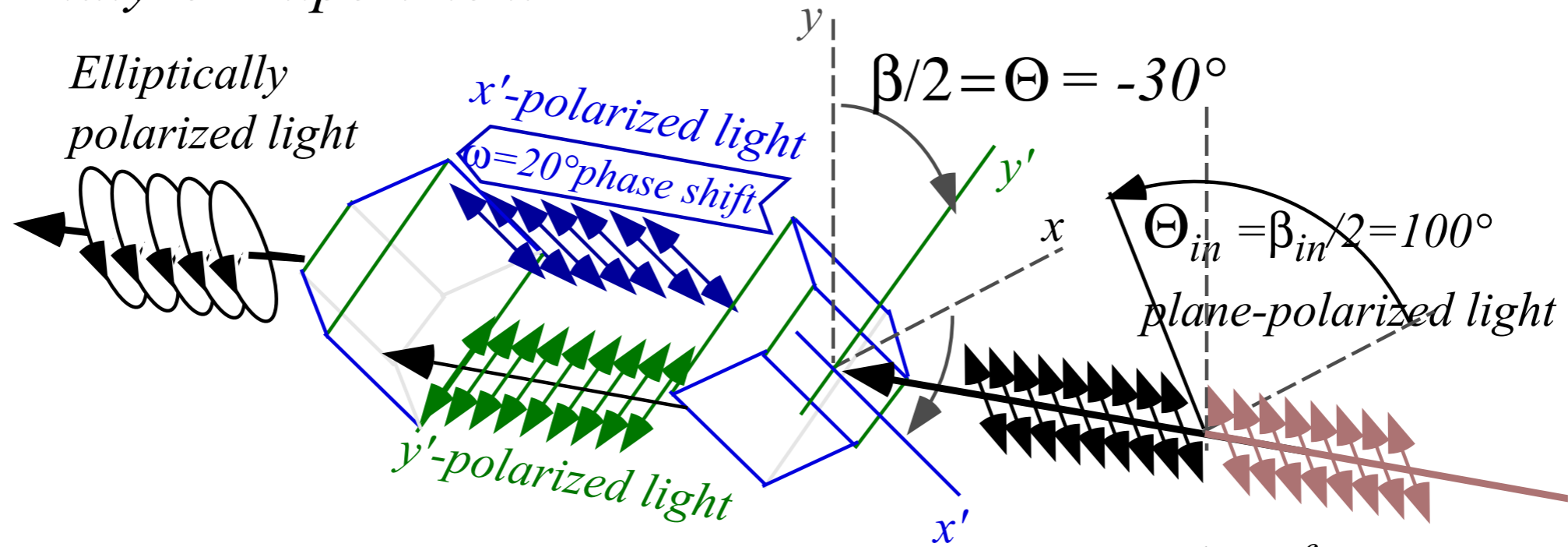
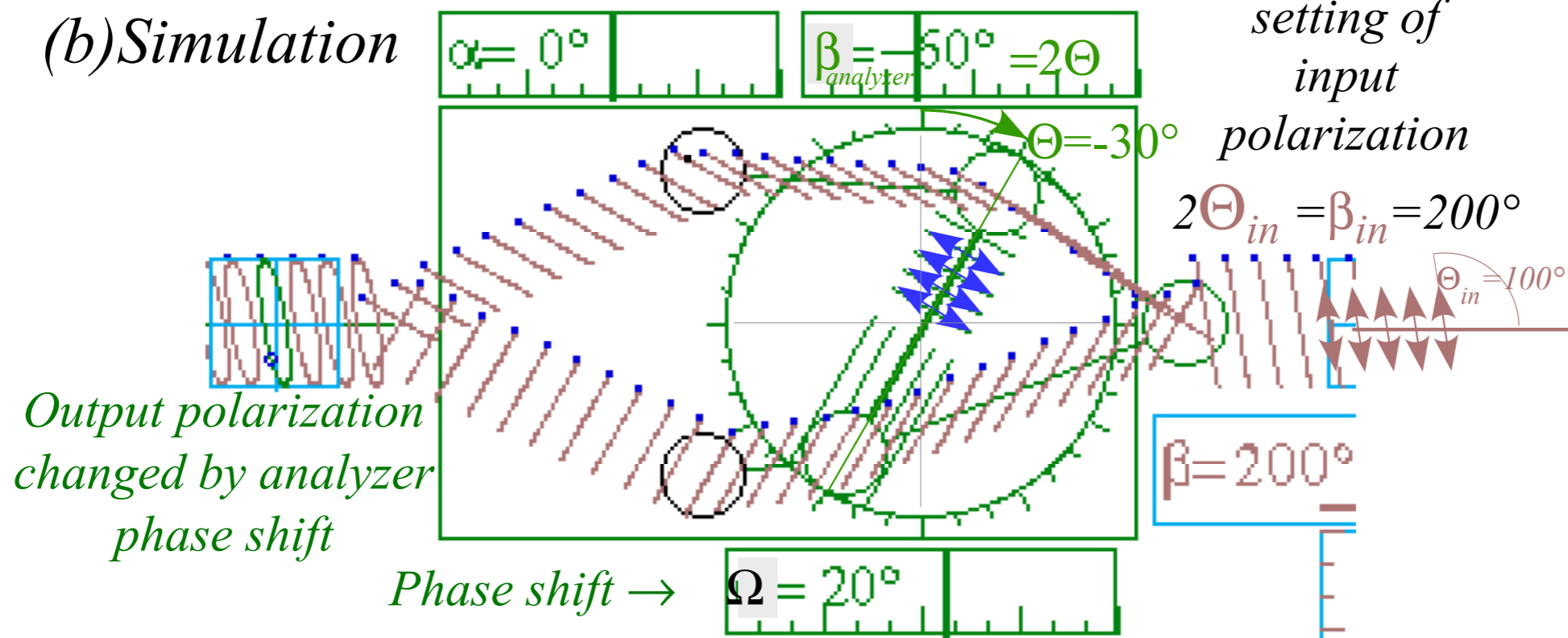


Fig. 1.3.5 Polarization control set to shift phase by (a) Half-wave ($\Omega = \pi$) , (b) Quarter wave ($\Omega = \pi/2$)

(a) Analyzer Experiment



(b) Simulation



Similar to "do-nothing" analyzer but has extra phase factor $e^{-i\Omega} = 0.94 - i 0.34$ on the x' -path .

$$x\text{-output: } \langle x | \Psi_{out} \rangle = \langle x | x' \rangle e^{-i\Omega} \langle x' | \Psi_{in} \rangle + \langle x | y' \rangle \langle y' | \Psi_{in} \rangle = e^{-i\Omega} \cos \Theta \cos(\Theta_{in} - \Theta) - \sin \Theta \sin(\Theta_{in} - \Theta)$$

$$y\text{-output: } \langle y | \Psi_{out} \rangle = \langle y | x' \rangle e^{-i\Omega} \langle x' | \Psi_{in} \rangle + \langle y | y' \rangle \langle y' | \Psi_{in} \rangle = e^{-i\Omega} \sin \Theta \cos(\Theta_{in} - \Theta) + \cos \Theta \sin(\Theta_{in} - \Theta)$$

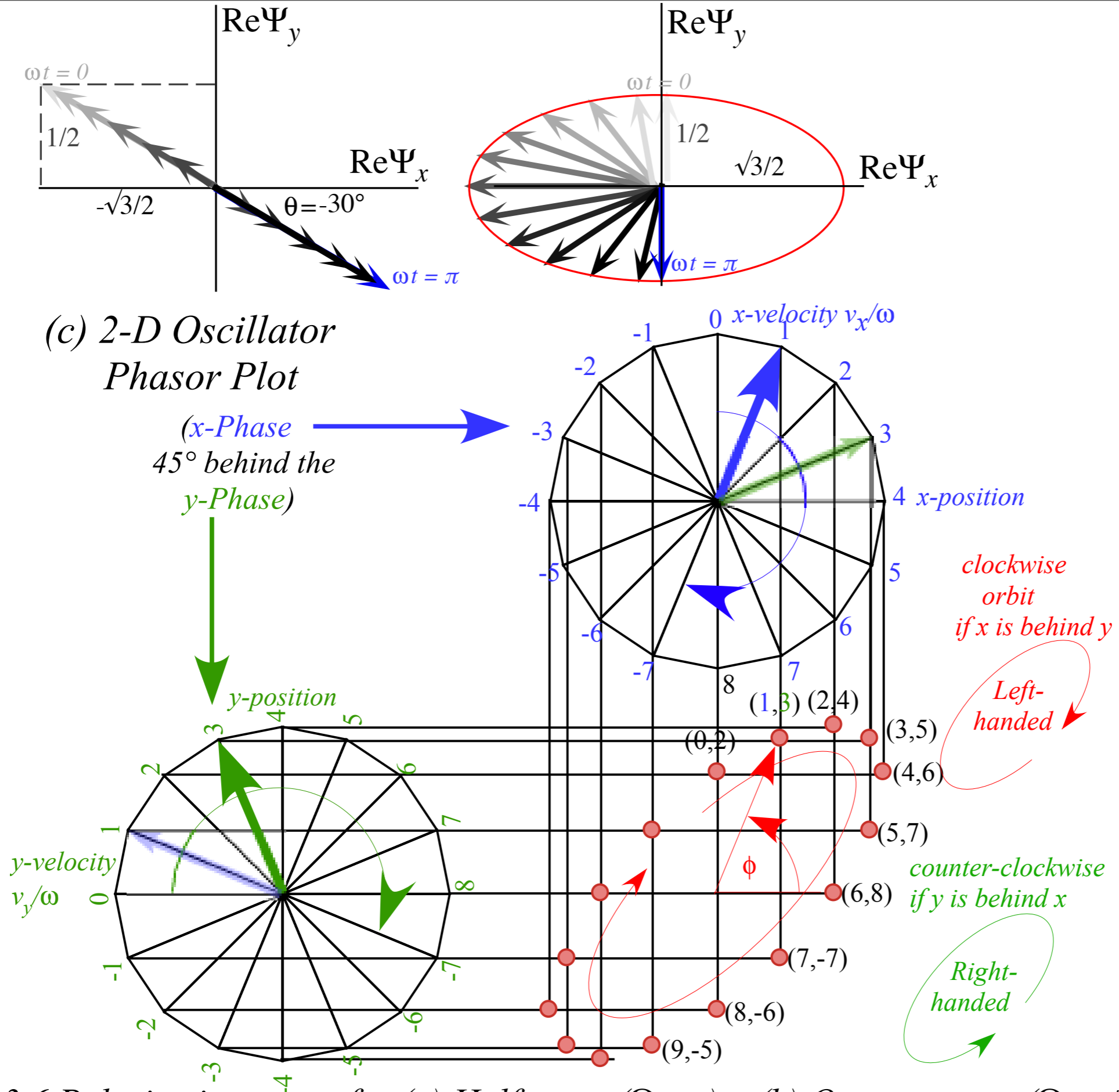


Fig. 1.3.6 Polarization states for (a) Half-wave ($\Omega=\pi$) , (b) Quarter wave ($\Omega=\pi/2$) (c) ($\Omega=-\pi/4$)

