

AMOP Lectures 9.0
Tue. 3.4-Thur. 3.6 2014

Relativity of interfering and galloping waves: Amplitude and SWR.

(Ch. 2-5 of CMwBang-Unit 8 Ch. 6 of QTforCA Unit 2)

Unmatched amplitudes giving galloping waves

Standing Wave Ratio (SWR) and Standing Wave Quotient (SWQ)

Analogy with group and phase

Galloping waves

Analogy between wave galloping, Keplerian IHO orbits, and optical polarization

Galloping dynamics algebra

Waves that go back in time - The Feynman-Wheeler Switchback

The Ship-Barn-and-Butler saga of confused causality

1st Quantization: Quantizing phase variables ω and k

Understanding how quantum transitions require “mixed-up” states

Closed cavity vs ring cavity

Relativistic effects on charge, current, and Maxwell Fields

*Current density changes by Lorentz **asynchrony***

Magnetic B-field is relativistic $\sinh\rho$ 1st order-effect



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$$A_{\rightarrow} e^{i(k_{\rightarrow} x - \omega_{\rightarrow} t)} + A_{\leftarrow} e^{i(k_{\leftarrow} x - \omega_{\leftarrow} t)} = e^{i(k_{\Sigma} x - \omega_{\Sigma} t)} [A_{\rightarrow} e^{i(k_{\Delta} x - \omega_{\Delta} t)} + A_{\leftarrow} e^{-i(k_{\Delta} x - \omega_{\Delta} t)}]$$

Waves have half-sum mean-phase rates $(k_{\Sigma}, \omega_{\Sigma})$ and half-difference group rates $(k_{\Delta}, \omega_{\Delta})$.

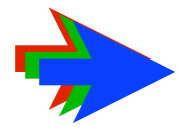
$$k_{\Sigma} = (k_{\rightarrow} + k_{\leftarrow}) / 2$$

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Detailed wave motion depends on standing-wave-ratio *SWR* or the inverse standing-wave-quotient *SWQ*.

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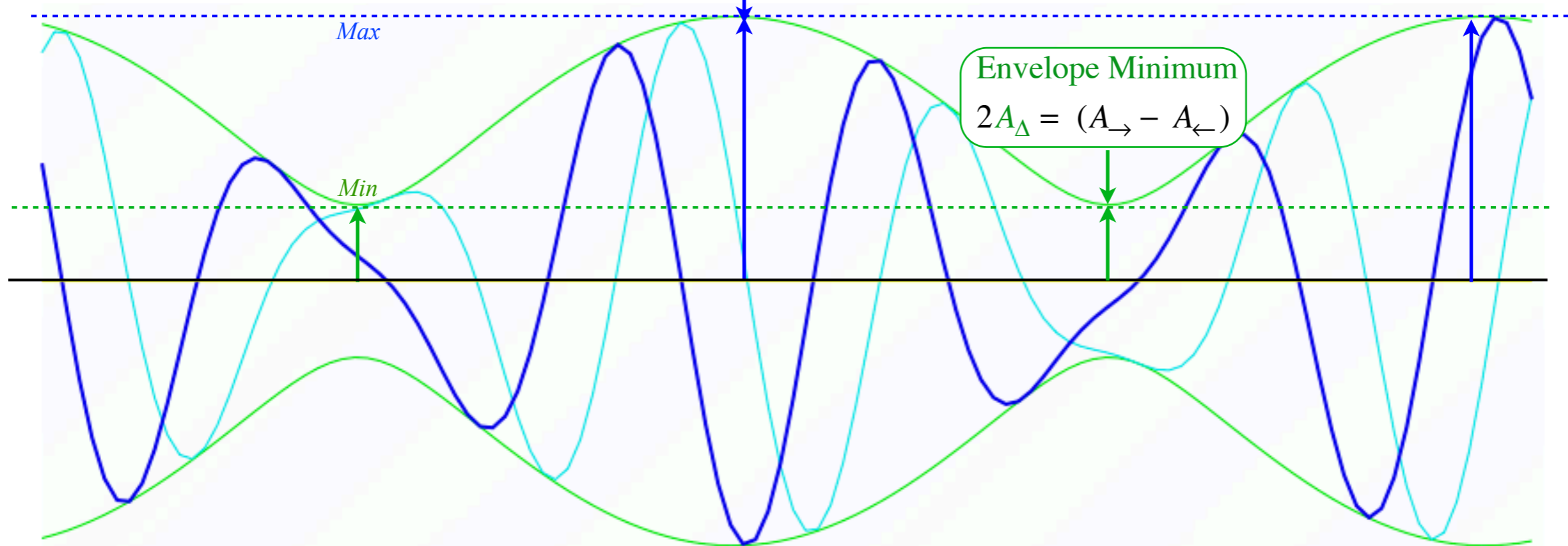
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$$\text{Envelope Maximum} \\ 2A_{\Sigma} = (A_{\rightarrow} + A_{\leftarrow})$$

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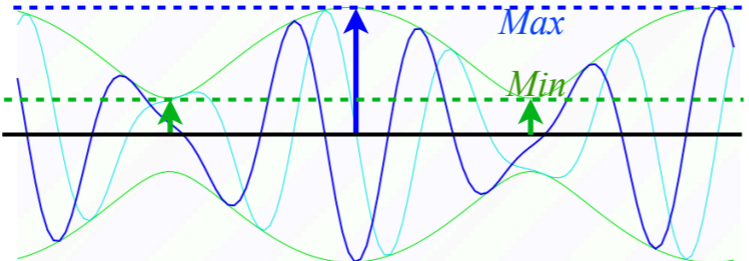
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These are analogous to frequency ratios for *group velocity* $V_{group} < c$ and its inverse that is *phase velocity* $V_{phase} > c$.

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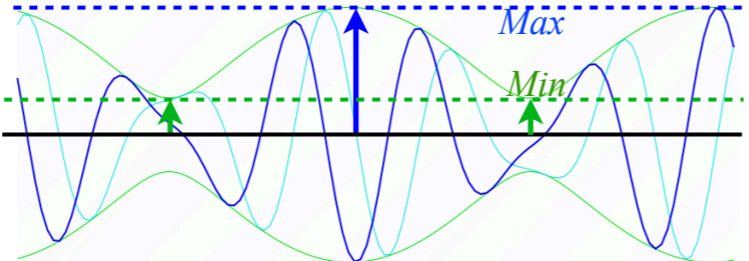
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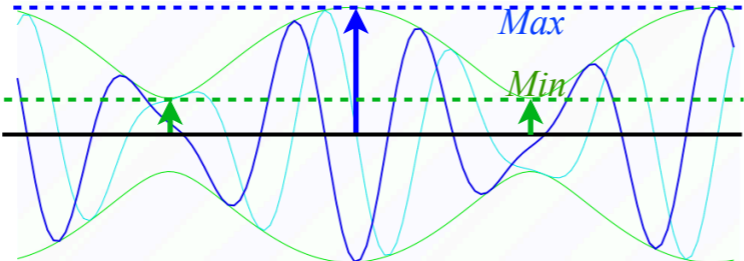
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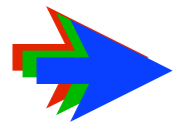
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$$\frac{V_{group}}{c} = \frac{c}{V_{phase}} \quad \text{is analogous to:} \quad \text{SWR} = \frac{1}{\text{SWQ}}$$

Unmatched amplitudes giving galloping waves

Standing Wave Ratio (SWR) and Standing Wave Quotient (SWQ)

Analogy with group and phase

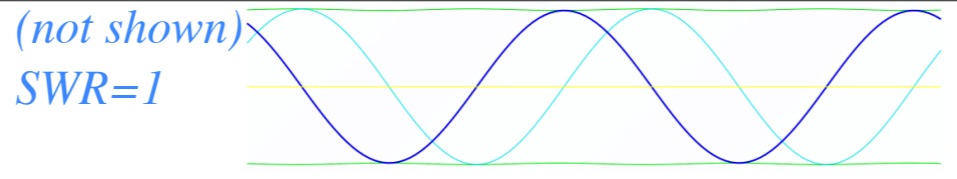
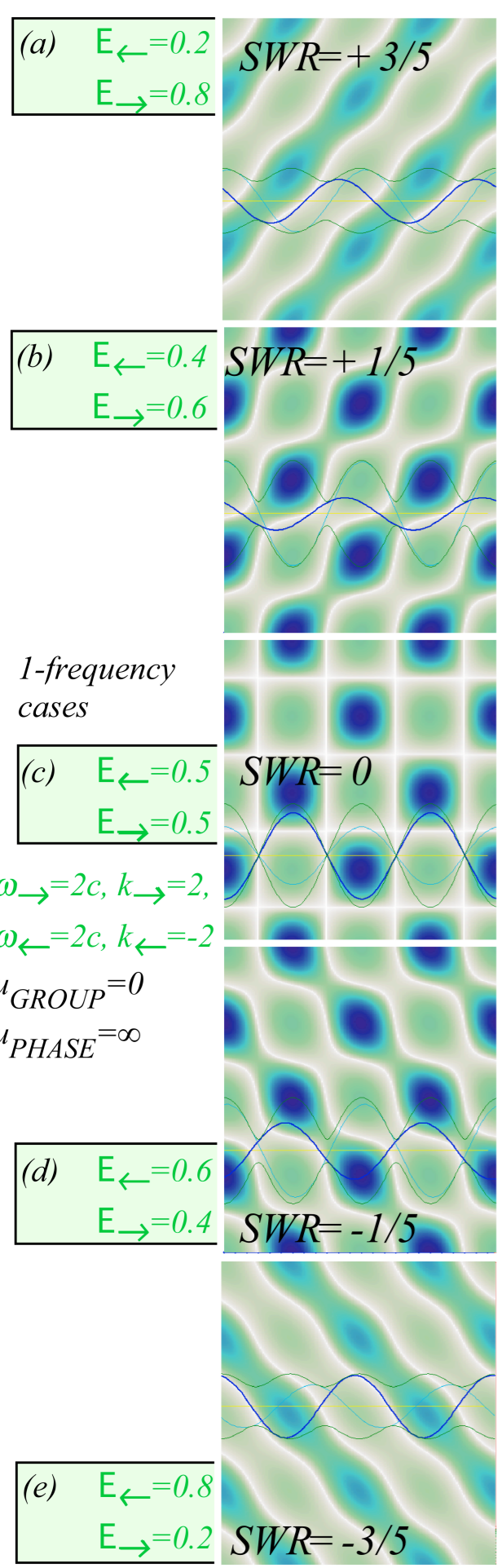


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Waves that go back in time - The Feynman-Wheeler Switchback



$SWR=+3/5$

$SWR=+1/5$

$SWR=0$

$SWR=-1/5$

$SWR=-3/5$

(not shown in (x,ct) plots)

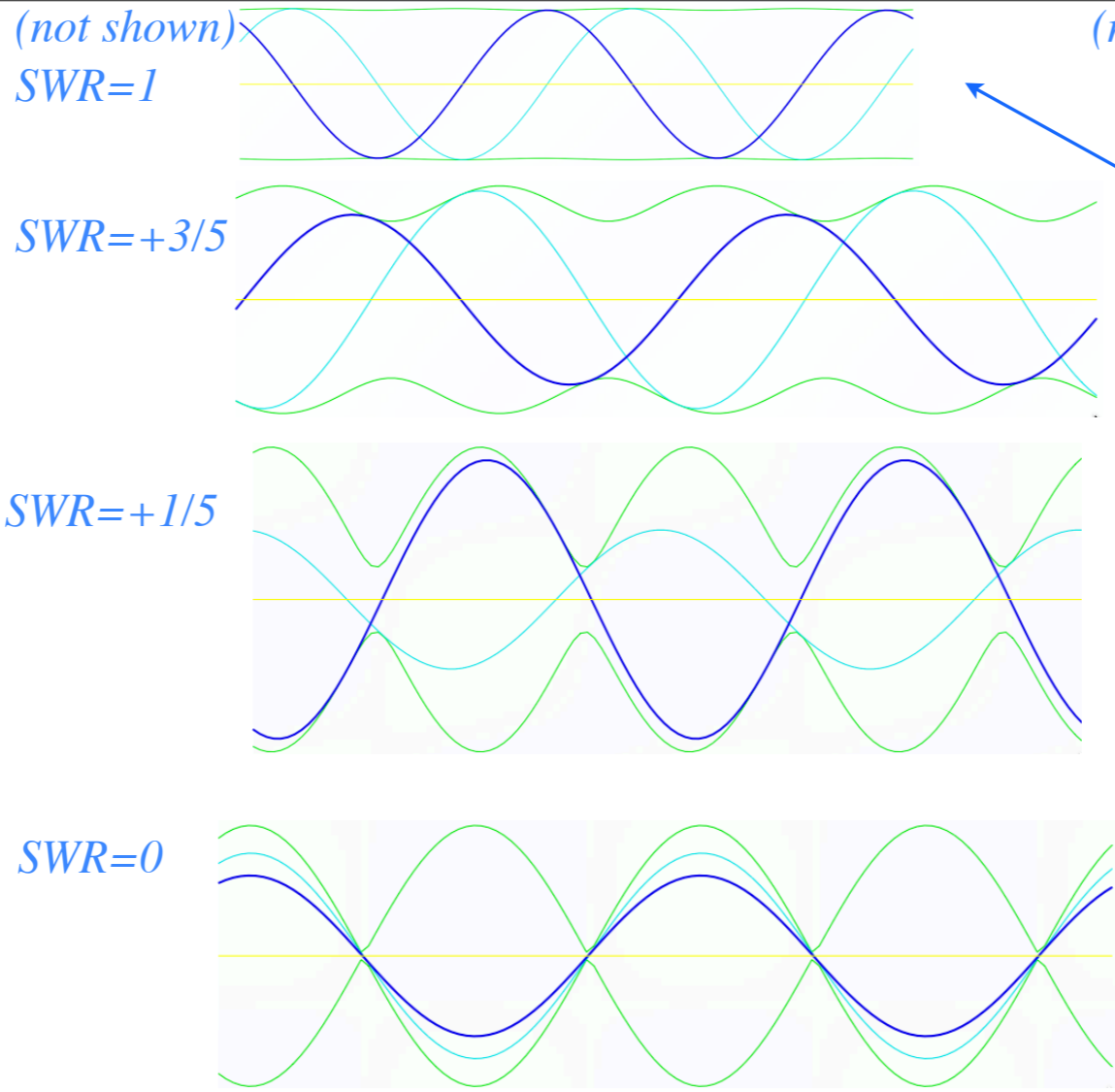
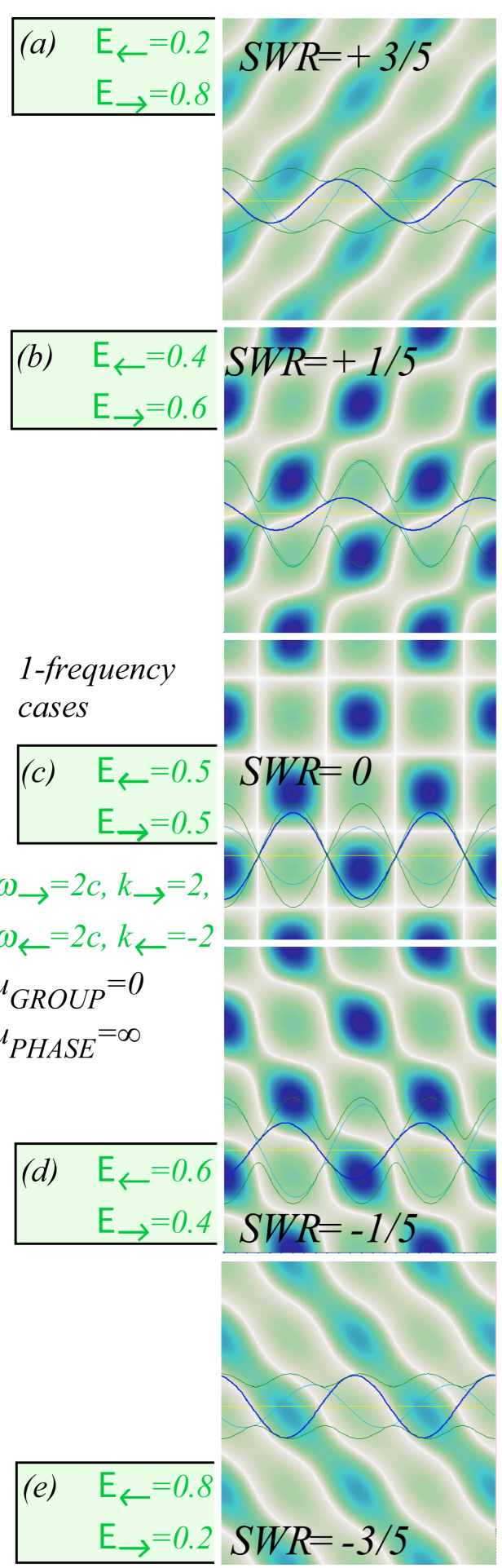
$SWR=1$

Two extremes for Standing Wave Ratio

$SWR=0$

Fig. 6.1 Monochromatic (1-frequency) 2-CW wave space-time patterns.

from: Fig. 4.5.2 QTforCA Unit 2 Ch.4
from: Fig. 8.6.3 CMwBang! Unit 8 Ch.6

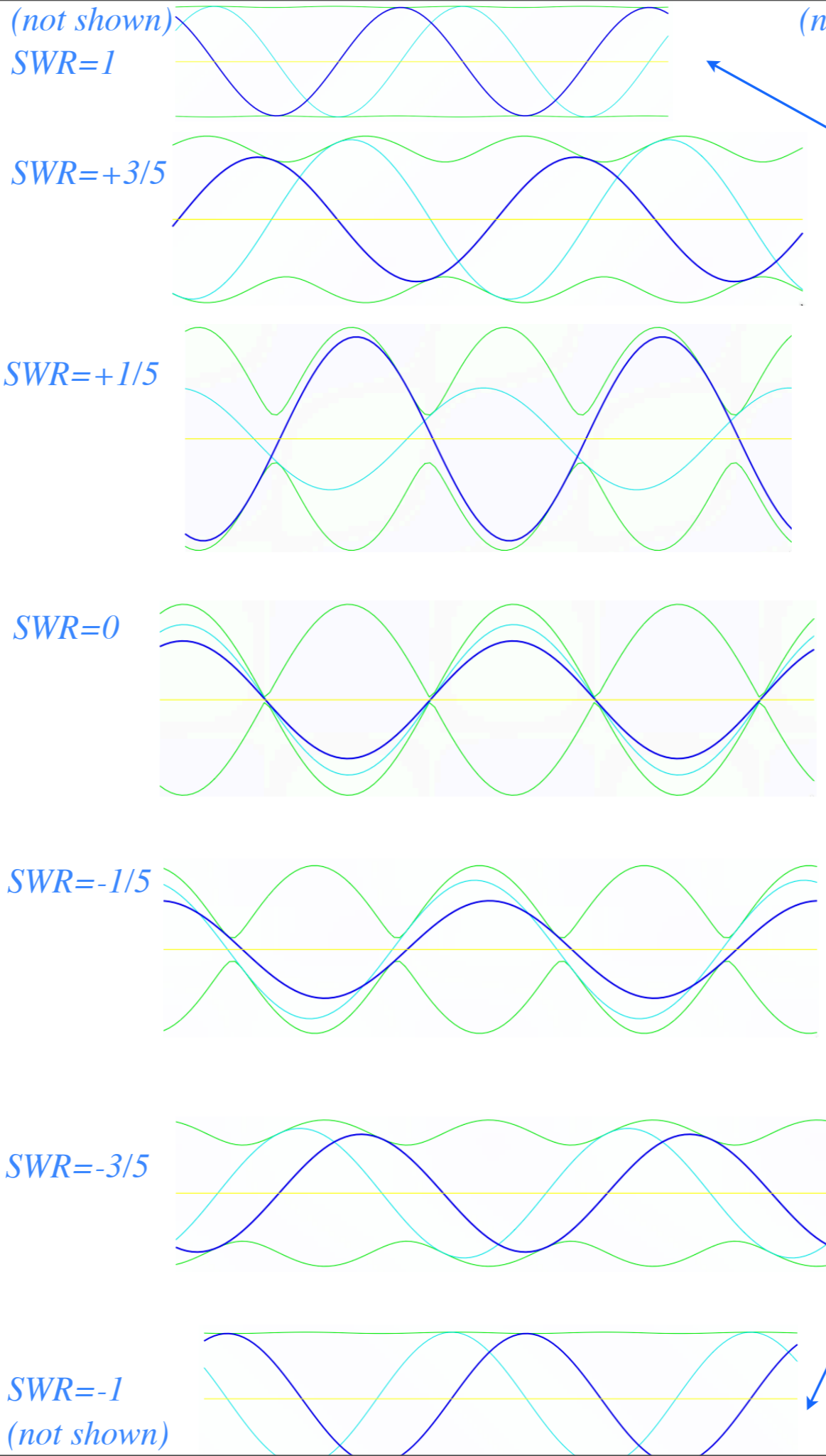
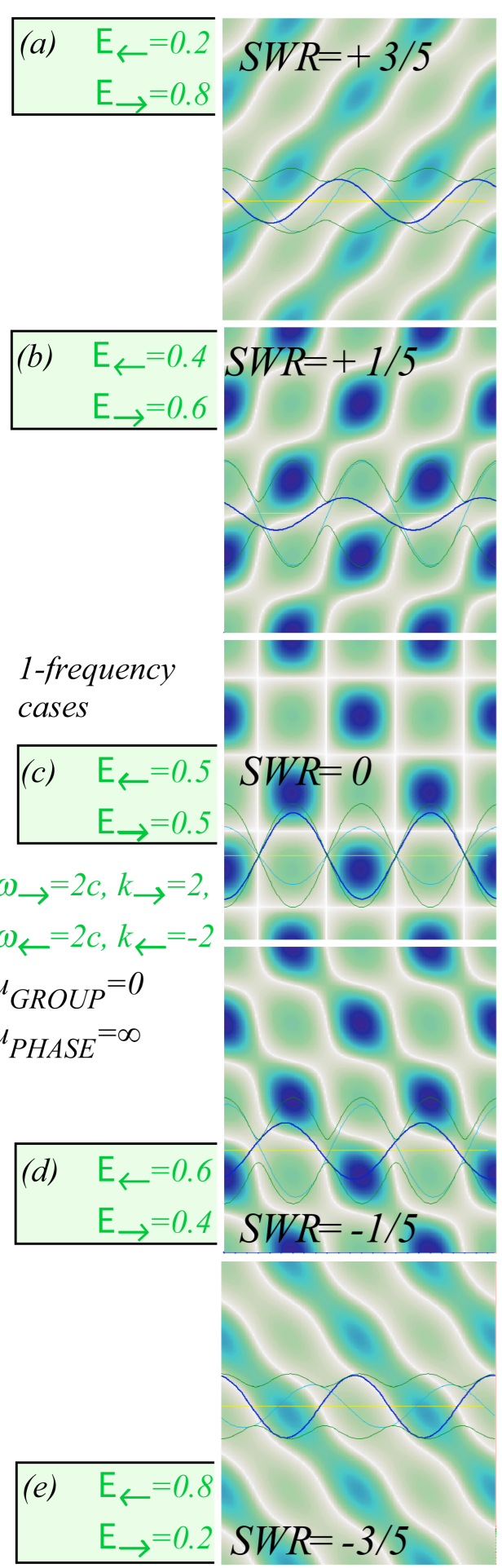


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Fig. 6.1 Monochromatic (1-frequency) 2-CW wave space-time patterns.



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...and
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from: Fig. 4.5.2
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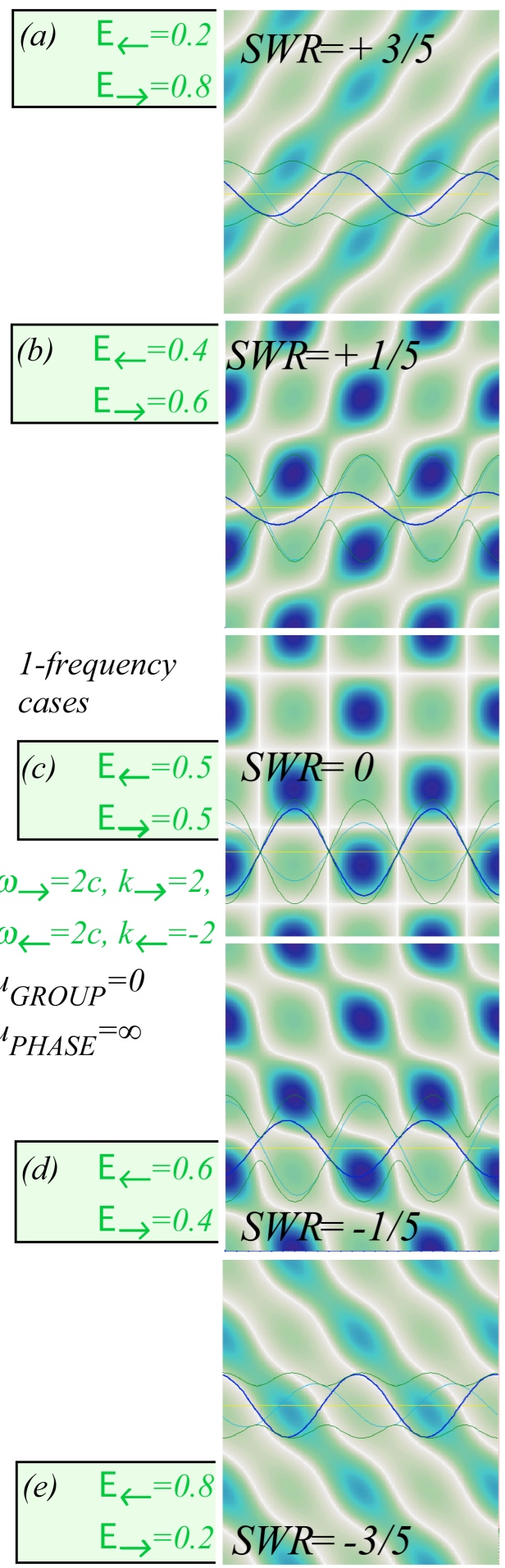


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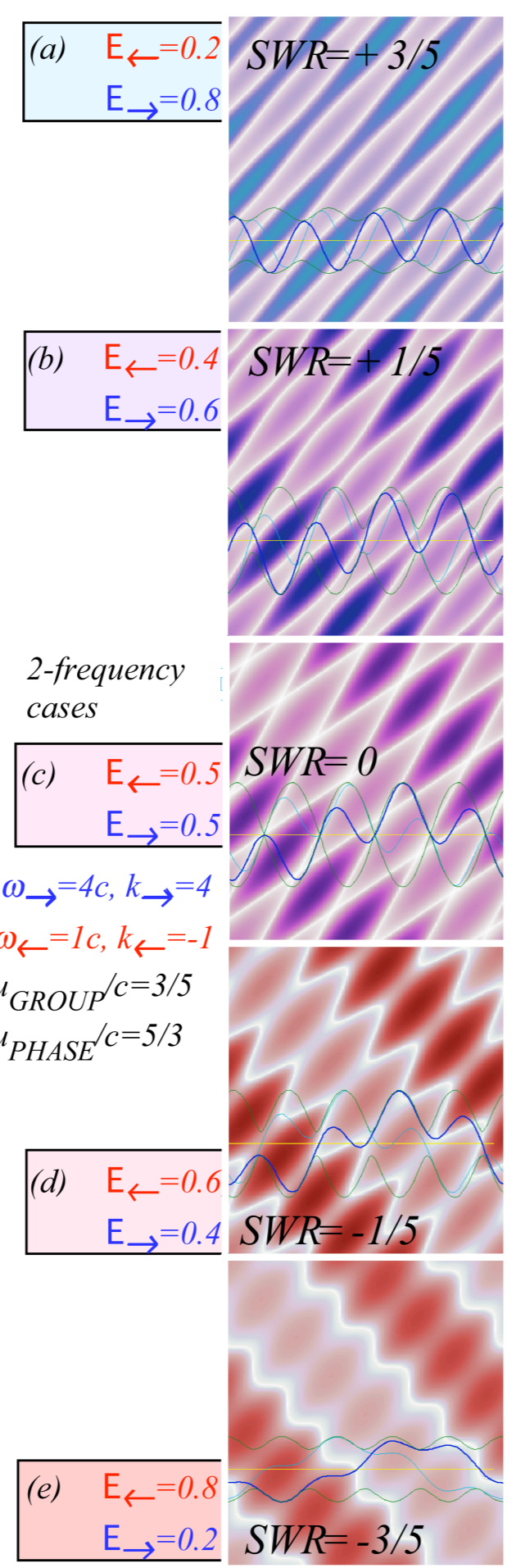
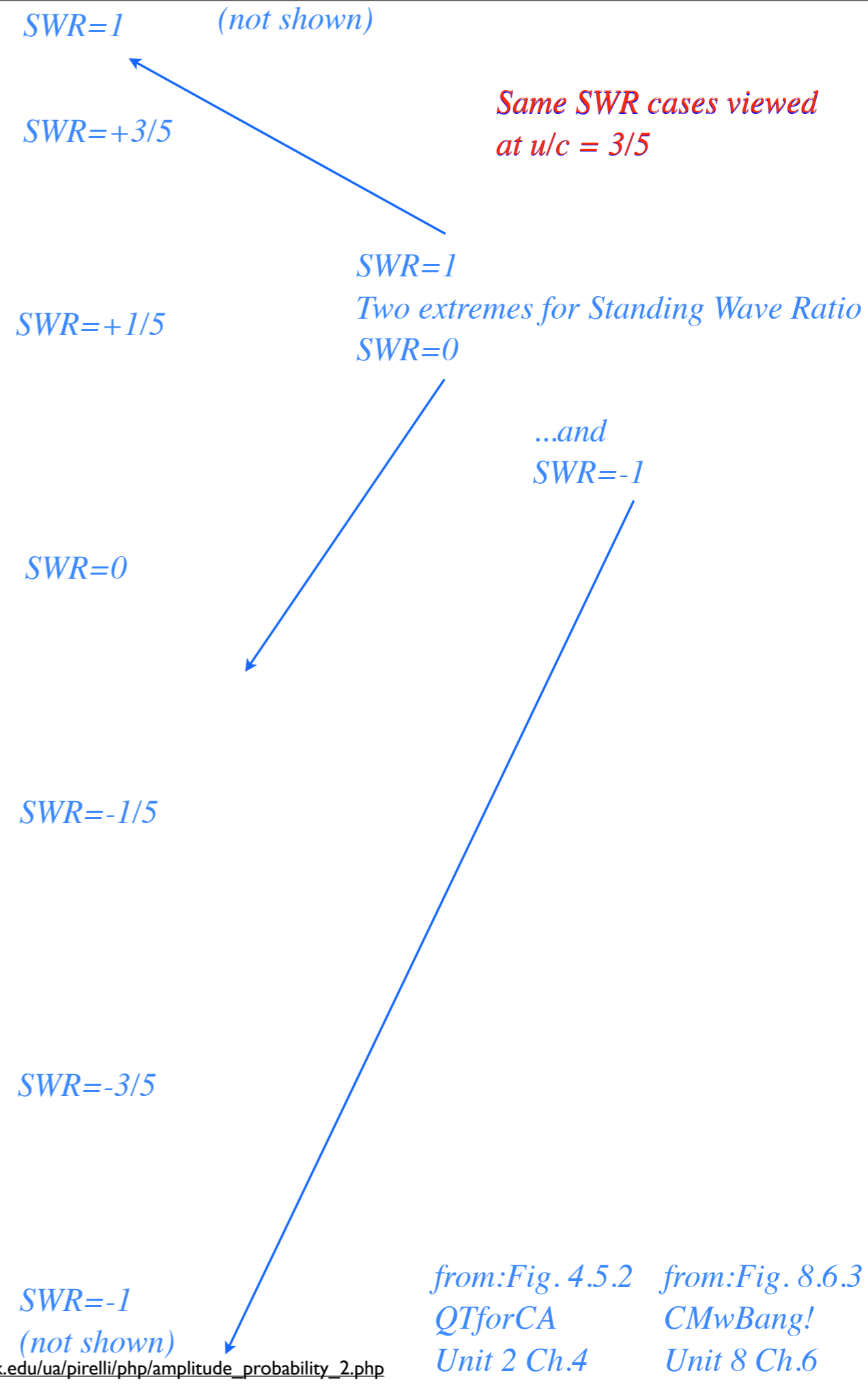


Fig. 6.2 Dichromatic (2-frequency) 2-CW wave space-time patterns.

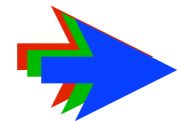


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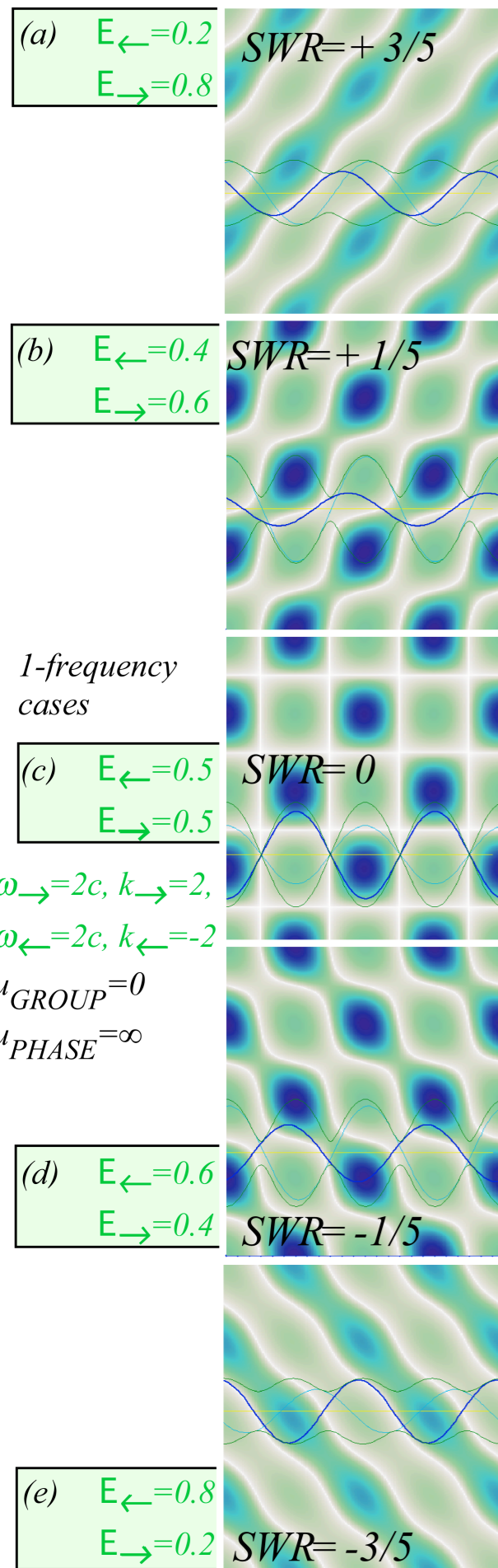


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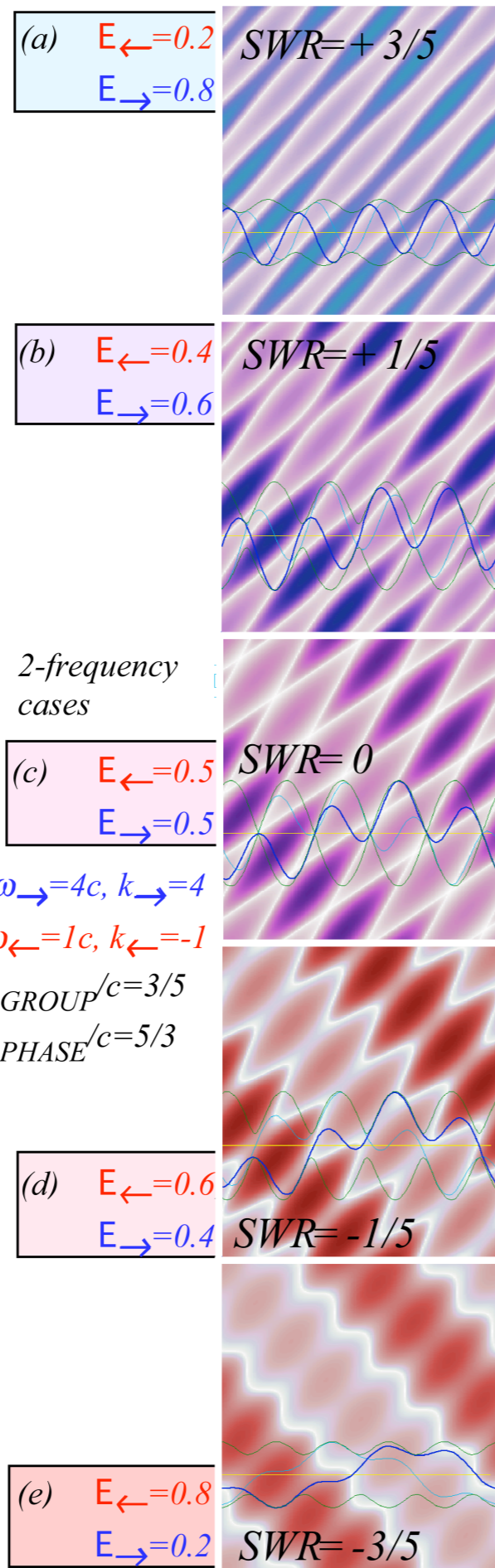


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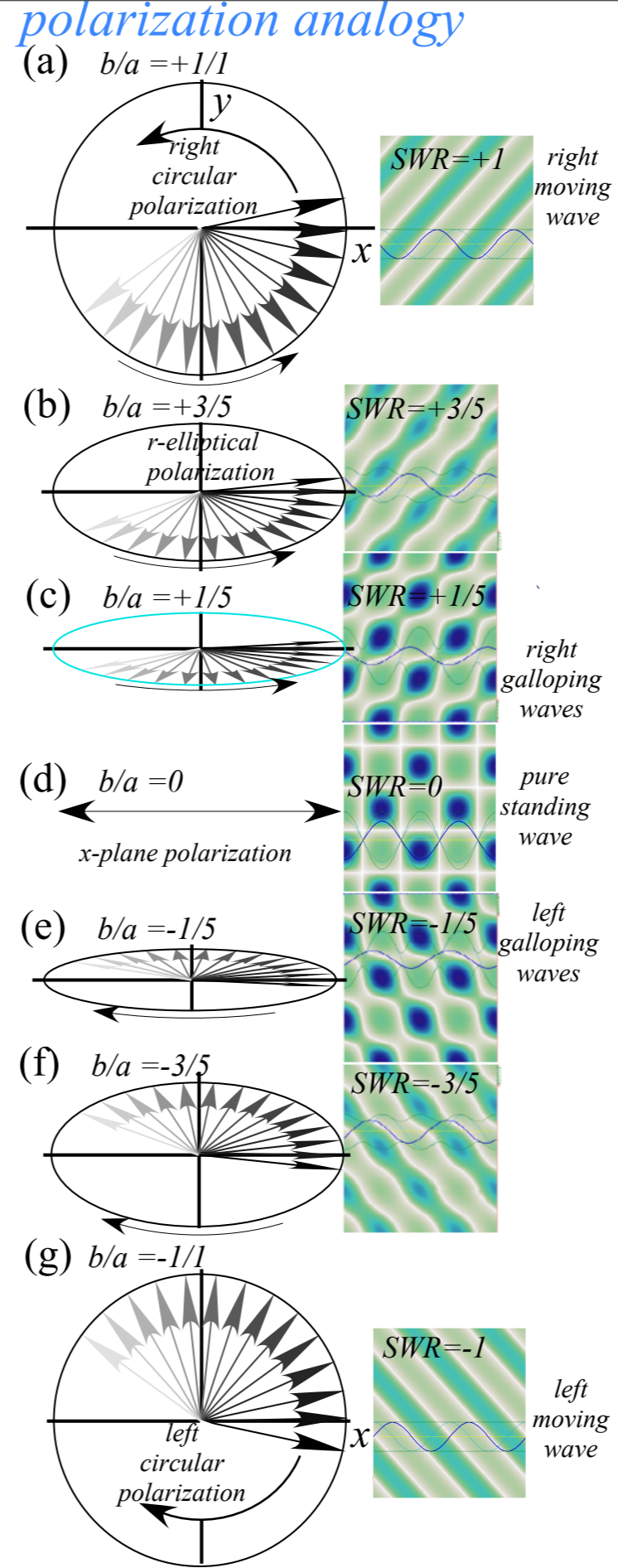


Fig. 6.3 (a-g) Elliptic polarization ellipses relate to galloping waves in Fig. 6.1. (h-i) Kepler anomalies.

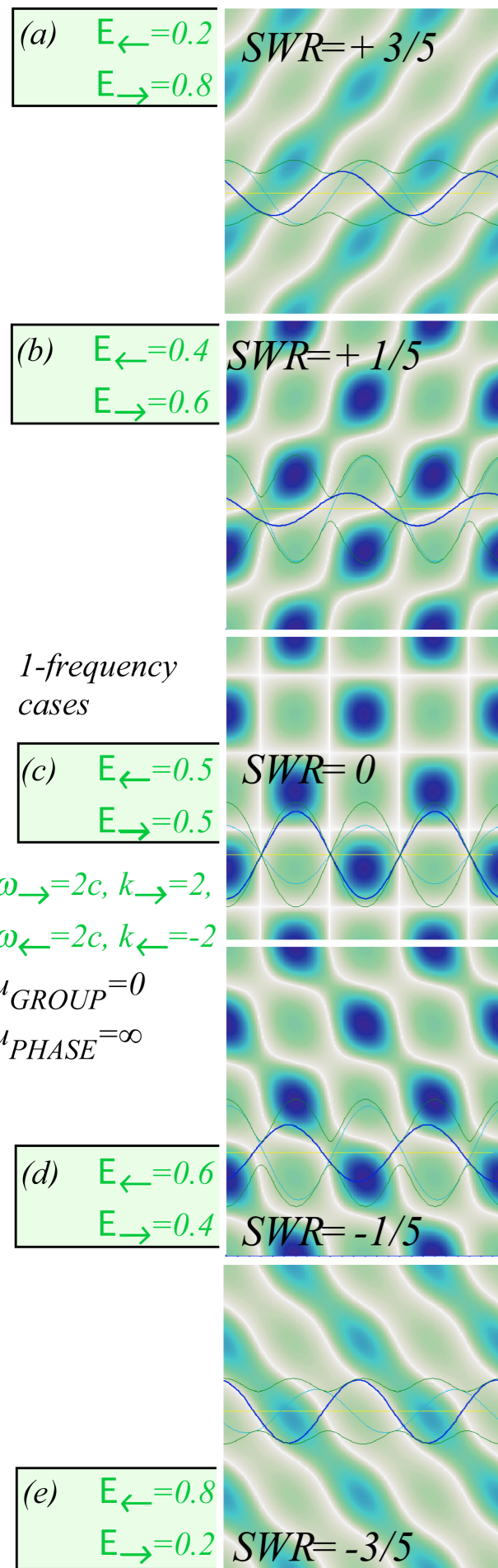


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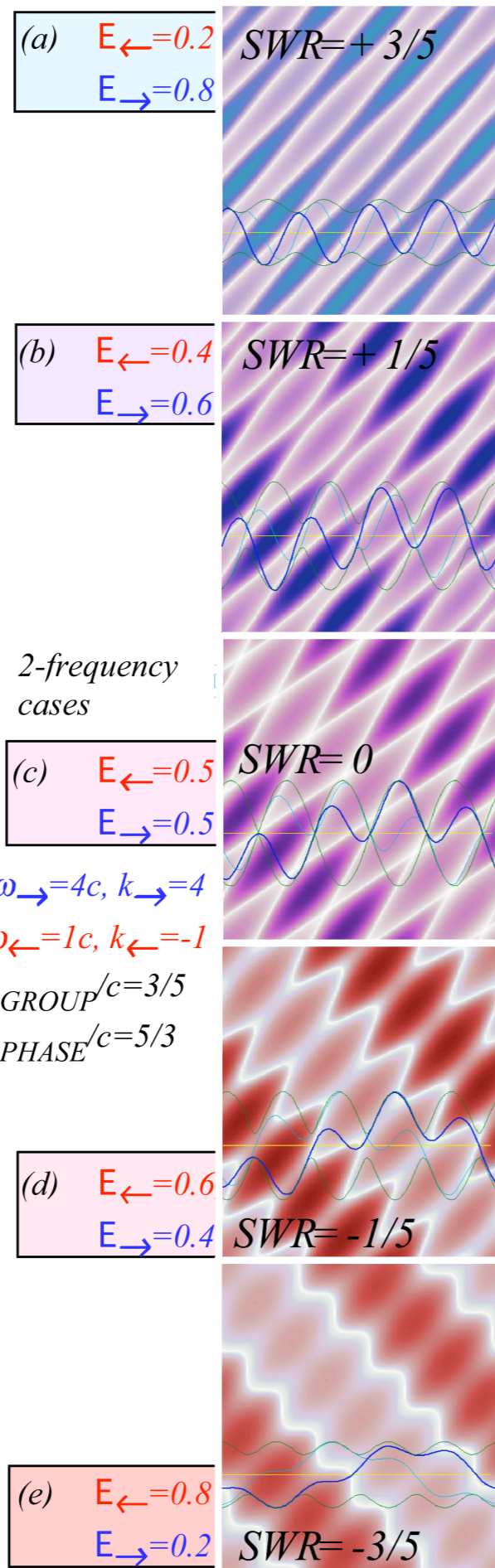


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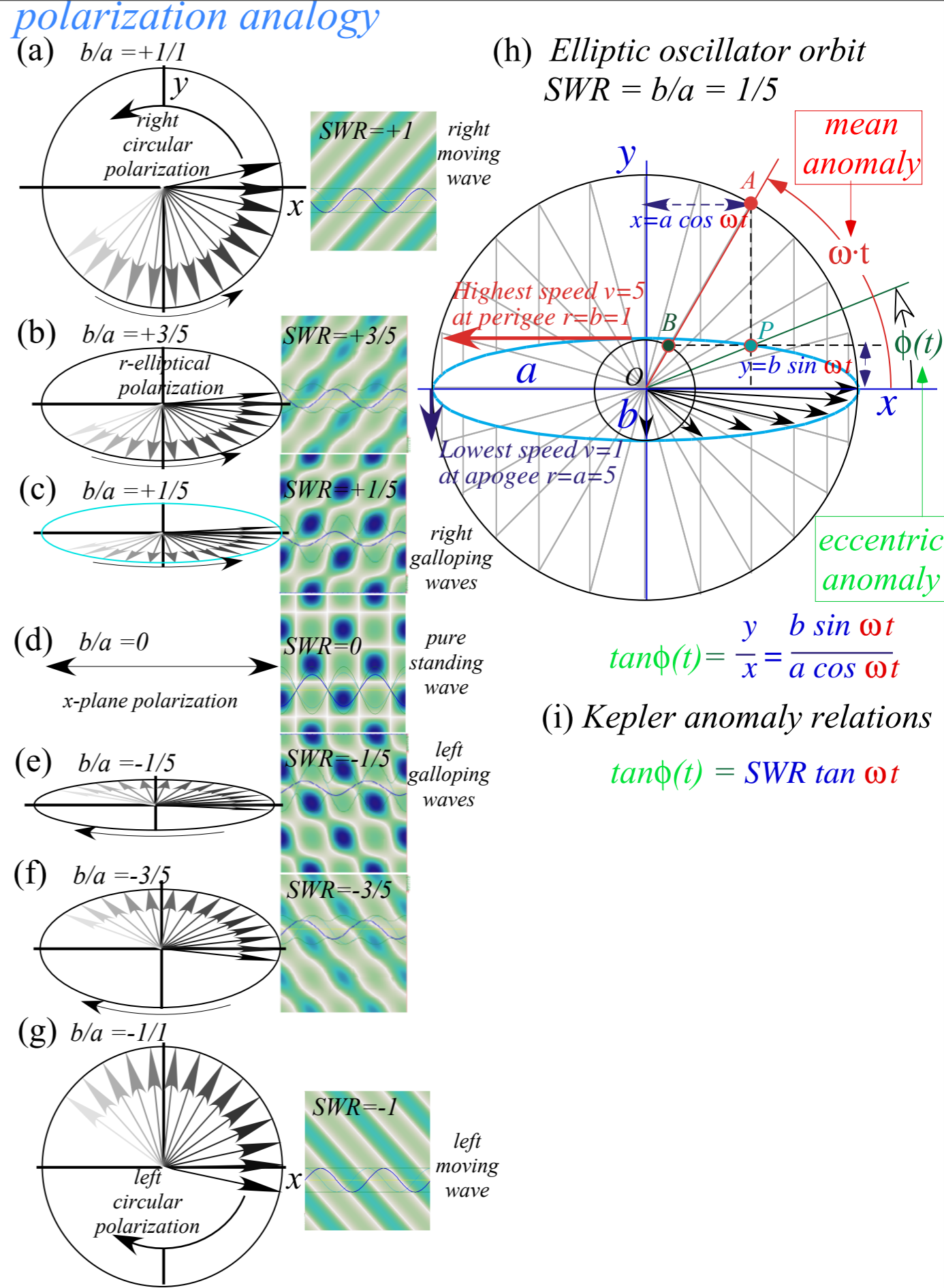


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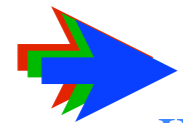
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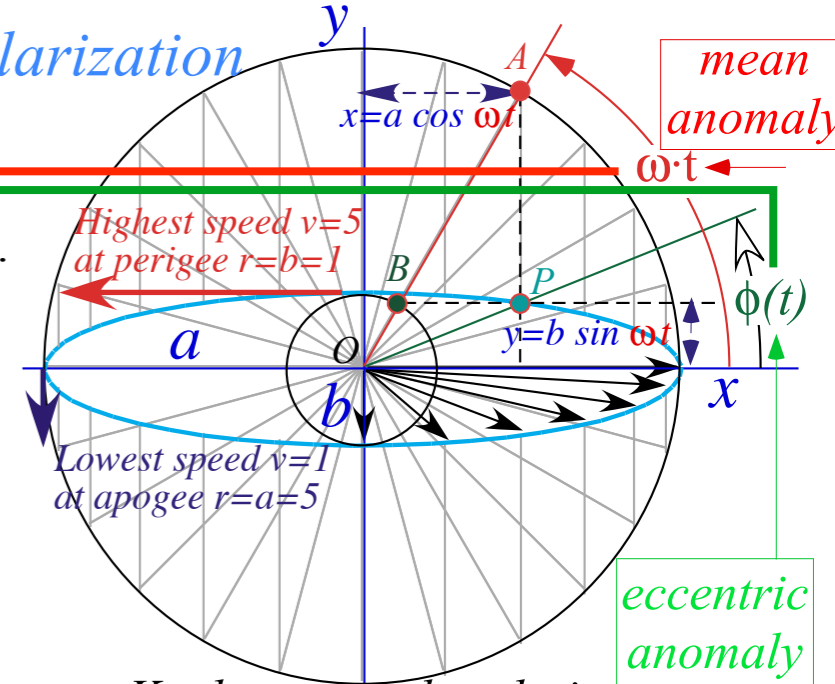
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Analogy between wave galloping, Keplerian IHO orbits, and optical polarization

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$$\tan \phi(t) = \frac{b}{a} \tan \omega \cdot t$$



Kepler anomaly relations

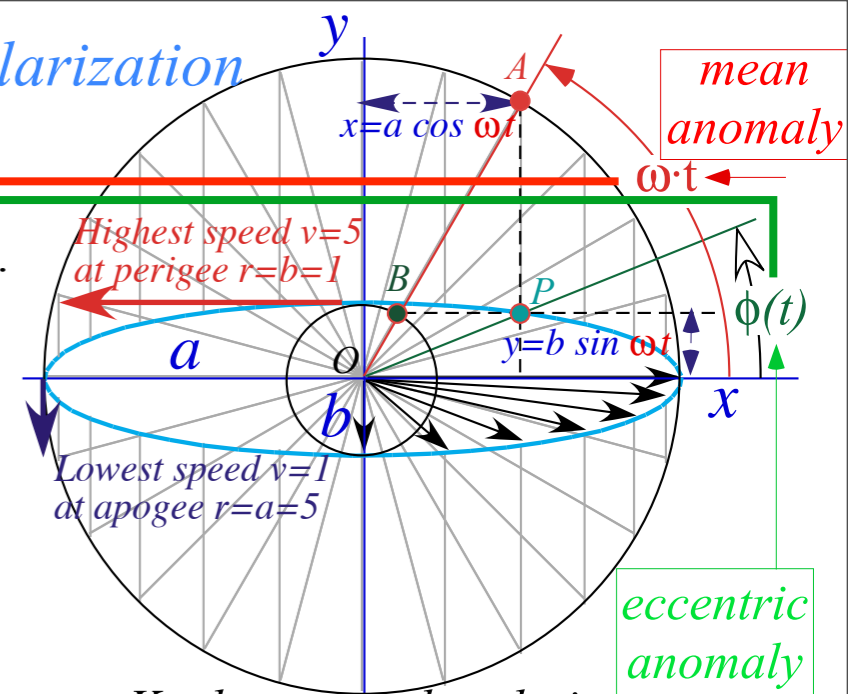
$$\tan \phi(t) = \frac{y}{x} = \frac{b \sin \omega t}{a \cos \omega t} = SWR \cdot \tan \omega t$$

from: Fig. 4.5.2 from: Fig. 8.6.3
 QTforCA CMwBang!
 Unit 2 Ch.4 Unit 8 Ch.6

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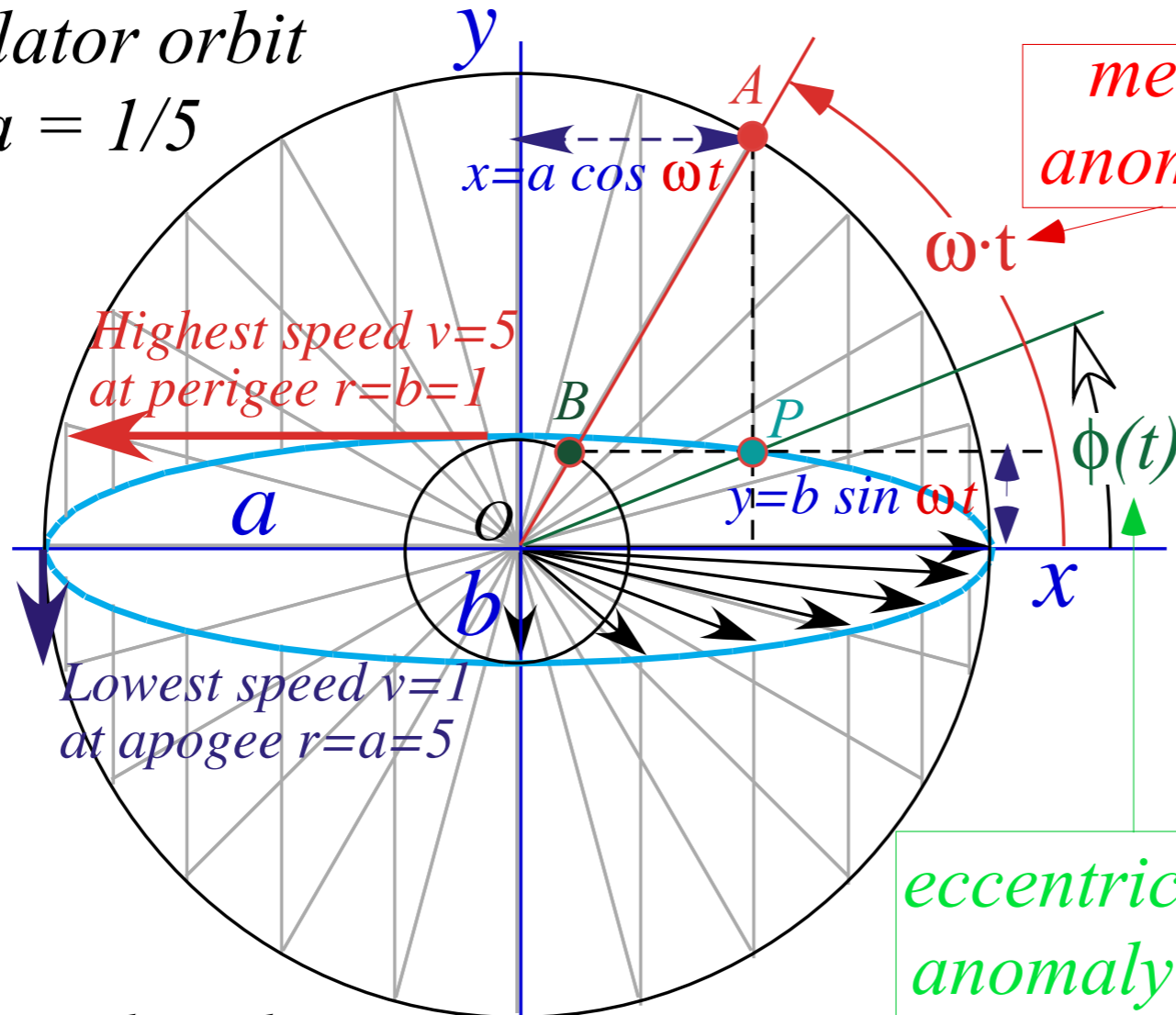


Kepler anomaly relations

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Elliptic oscillator orbit

$$SWR = b/a = 1/5$$



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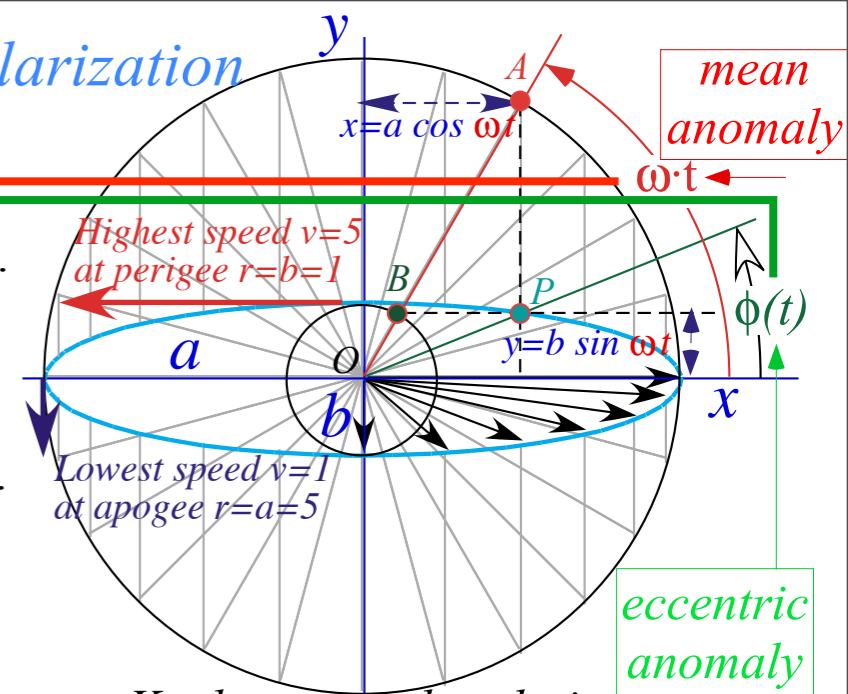
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The eccentric anomaly time derivative of ϕ (angular velocity) gallops between $\omega \cdot b/a$ and $\omega \cdot a/b$.

$$\dot{\phi} = \frac{d\phi}{dt} = \omega \cdot \frac{b \sec^2 \omega t}{a \sec^2 \phi} = \omega \cdot \frac{b \sec^2 \omega t}{a (1 + \tan^2 \phi)} = \frac{\omega \cdot b/a}{\cos^2 \omega t + (b/a)^2 \cdot \sin^2 \omega t} = \begin{cases} \omega \cdot b/a & \text{for: } \omega t = 0, \pi, 2\pi, \dots \\ \omega \cdot a/b & \omega t = \pi/2, 3\pi/2, \dots \end{cases}$$

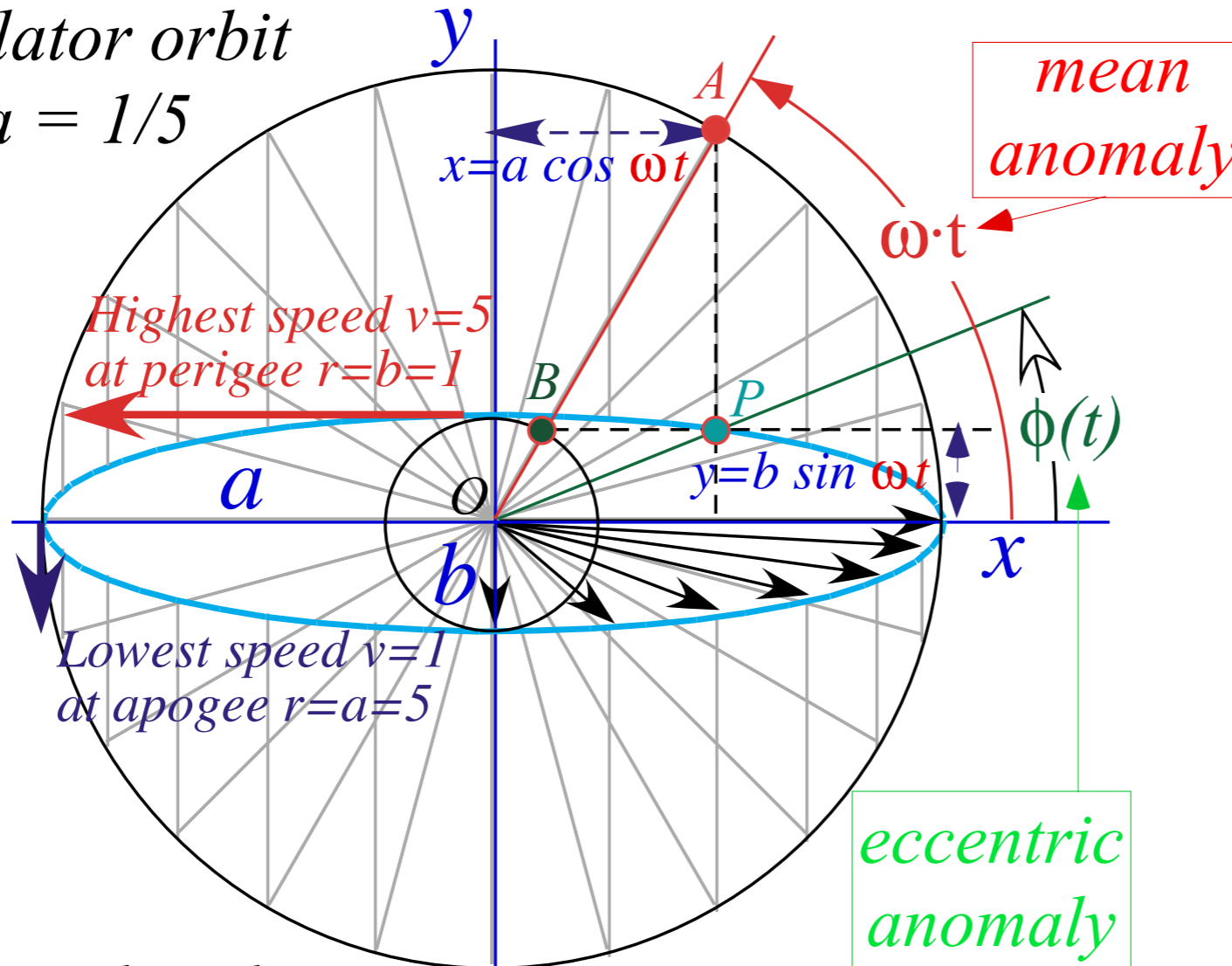


Kepler anomaly relations

$$\tan \phi(t) = \frac{y}{x} = \frac{b \sin \omega t}{a \cos \omega t} = SWR \cdot \tan \omega t$$

Elliptic oscillator orbit

$$SWR = b/a = 1/5$$



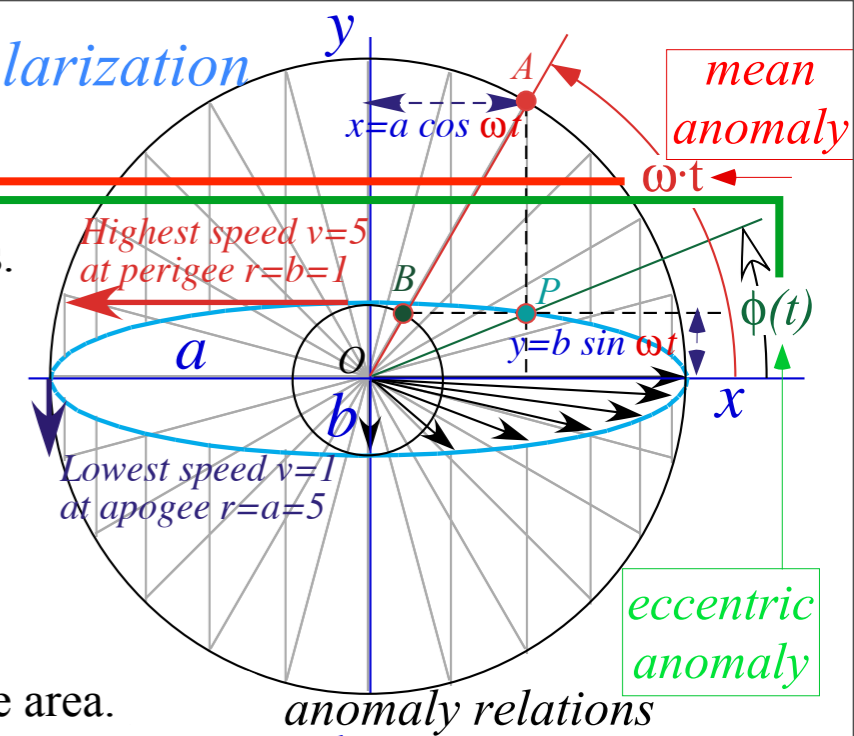
Kepler anomaly relations

$$\tan \phi(t) = SWR \cdot \tan \omega t$$

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from: Fig. 4.5.2 QTforCA Unit 2 Ch.4
 from: Fig. 8.6.3 CMwBang! Unit 8 Ch.6

Analogy between wave galloping, Keplerian IHO orbits, and optical polarization



We'll show wave galloping is analogous to Keplerian orbital motion of angles $\omega \cdot t$ and ϕ of orbits.

$$\tan \phi(t) = \frac{b}{a} \tan \omega \cdot t$$

The eccentric anomaly time derivative of ϕ (angular velocity) gallops between $\omega \cdot b/a$ and $\omega \cdot a/b$.

$$\dot{\phi} = \frac{d\phi}{dt} = \omega \cdot \frac{b \sec^2 \omega t}{a \sec^2 \phi} = \omega \cdot \frac{b \sec^2 \omega t}{a (1 + \tan^2 \phi)} = \frac{\omega \cdot b/a}{\cos^2 \omega t + (b/a)^2 \cdot \sin^2 \omega t} = \begin{cases} \omega \cdot b/a & \text{for: } \omega t = 0, \pi, 2\pi, \dots \\ \omega \cdot a/b & \omega t = \pi/2, 3\pi/2, \dots \end{cases}$$

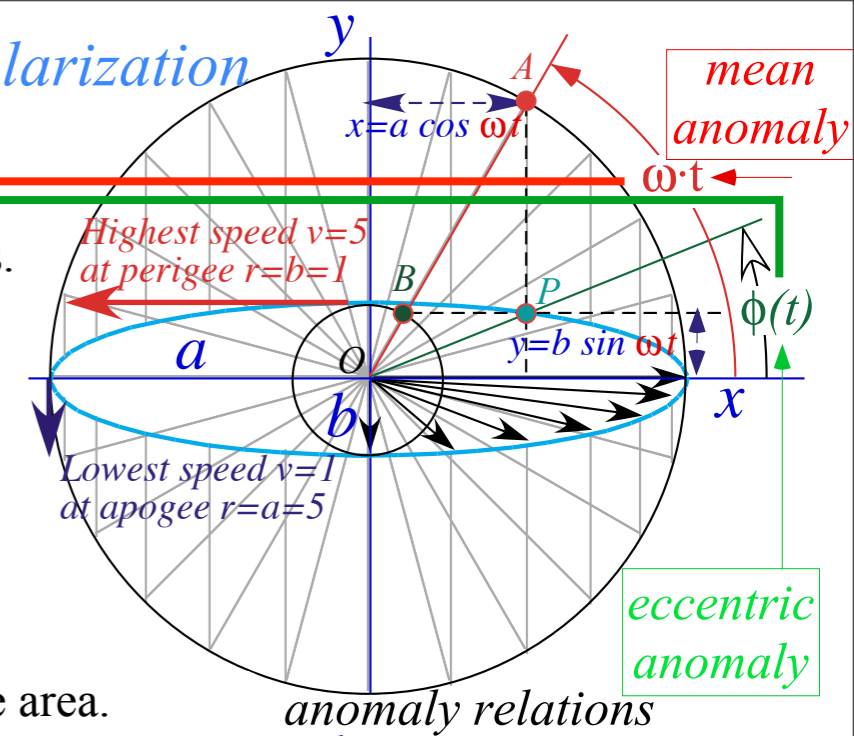
The product of angular momentum r^2 and $\dot{\phi}$ is orbital momentum, a constant proportional to ellipse area.

$$r^2 \frac{d\phi}{dt} = \text{constant} = (a^2 \cos^2 \omega t + b^2 \cdot \sin^2 \omega t) \frac{d\phi}{dt} = \omega \cdot ab$$

anomaly relations

$$\tan \phi(t) = \frac{y}{x} = \frac{b \sin \omega t}{a \cos \omega t} = \text{SWR} \cdot \tan \omega t$$

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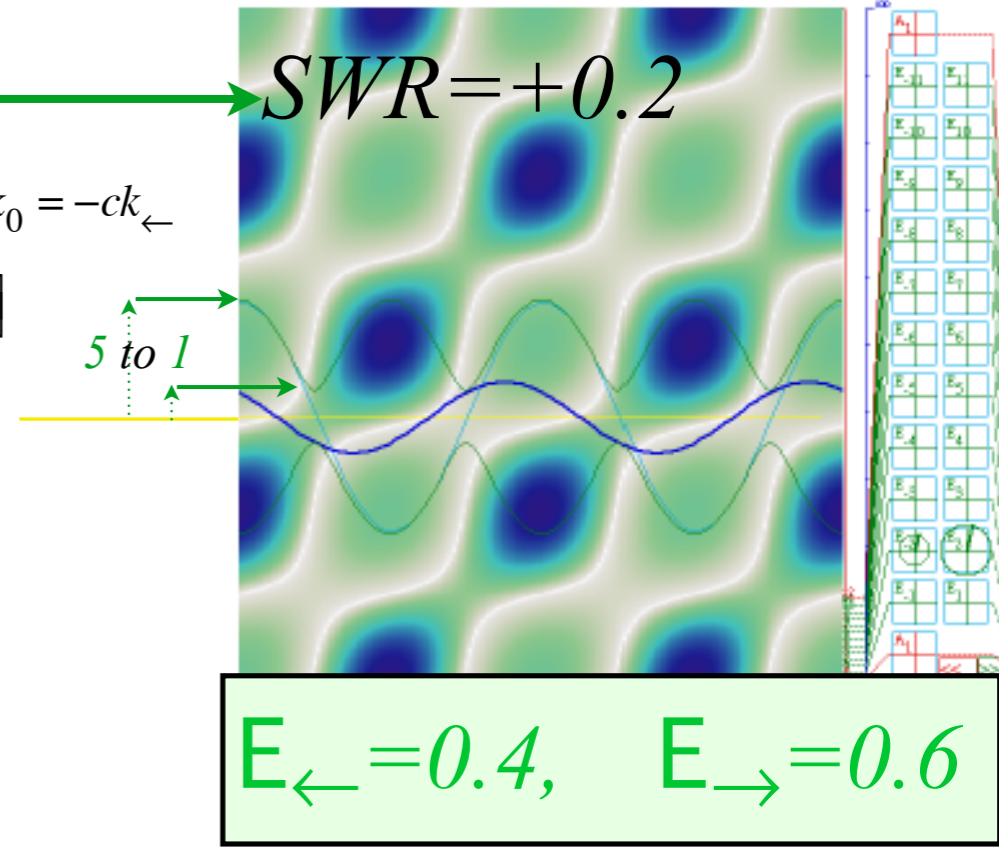
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Consider galloping wave zeros of a monochromatic wave having $SWR = 1/5$.

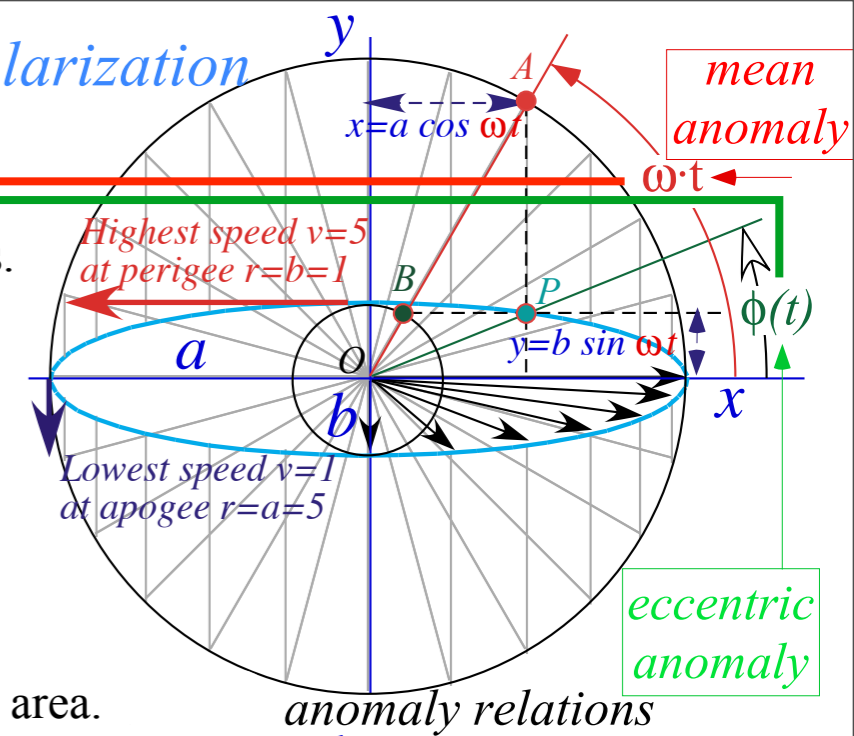
$$0 = \text{Re } \Psi(x, t) = \text{Re} \left[A_{\rightarrow} e^{i(k_0 x - \omega_0 t)} + A_{\leftarrow} e^{i(-k_0 x - \omega_0 t)} \right] \text{ where: } \omega_{\rightarrow} = \omega_0 = \omega_{\leftarrow} = ck_0 = -ck_{\leftarrow}$$

$$0 = A_{\rightarrow} [\cos k_0 x \cos \omega_0 t + \sin k_0 x \sin \omega_0 t] + A_{\leftarrow} [\cos k_0 x \cos \omega_0 t - \sin k_0 x \sin \omega_0 t]$$

$$(A_{\rightarrow} + A_{\leftarrow}) [\cos k_0 x \cos \omega_0 t] = -(A_{\rightarrow} - A_{\leftarrow}) [\sin k_0 x \sin \omega_0 t]$$



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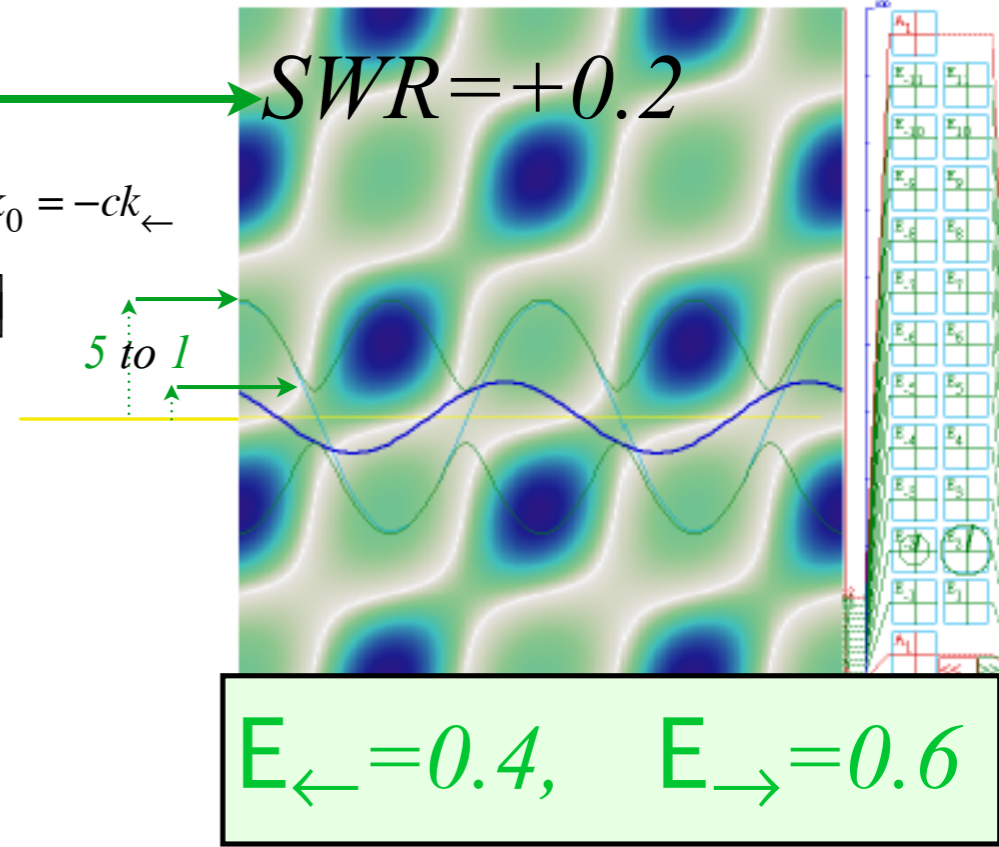
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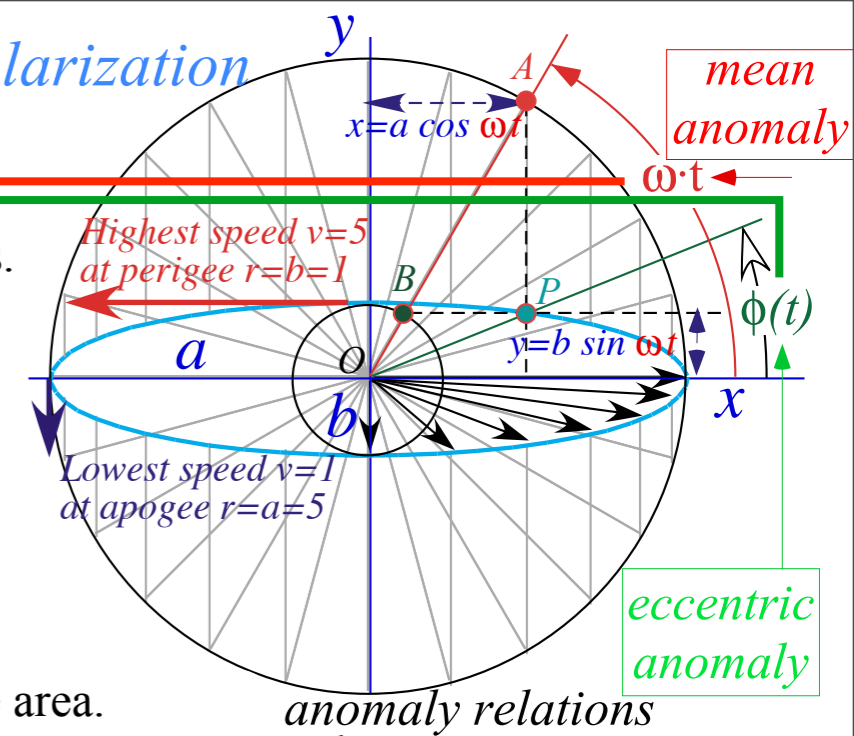
$$(A_{\rightarrow} + A_{\leftarrow}) [\cos k_0 x \cos \omega_0 t] = -(A_{\rightarrow} - A_{\leftarrow}) [\sin k_0 x \sin \omega_0 t]$$

Space $k_0 x$ varies with time $\omega_0 t$ in the same way that eccentric anomaly ϕ varies with $\omega \cdot t$.

$$\tan k_0 x = -SWR \cdot \cot \omega_0 t = SWR \cdot \tan \omega_0 \bar{t} \text{ where: } \omega_0 \bar{t} = \omega_0 t - \pi/2$$



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$$0 = A_{\rightarrow} [\cos k_0 x \cos \omega_0 t + \sin k_0 x \sin \omega_0 t] + A_{\leftarrow} [\cos k_0 x \cos \omega_0 t - \sin k_0 x \sin \omega_0 t]$$

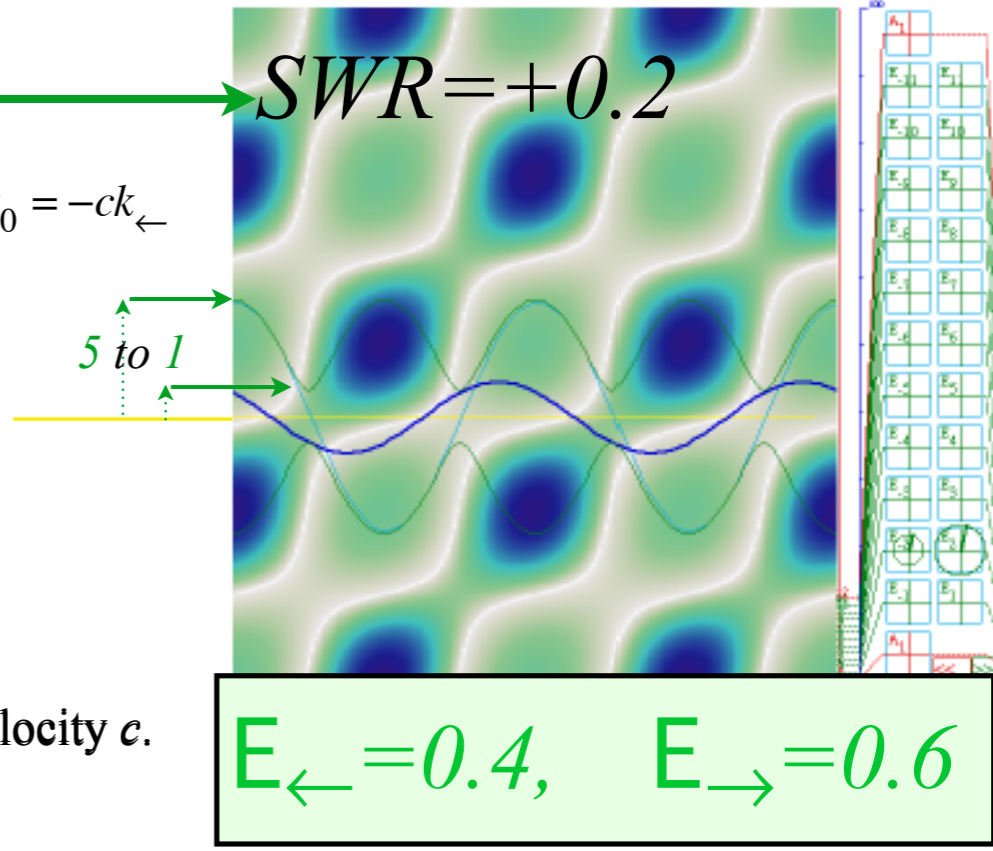
$$(A_{\rightarrow} + A_{\leftarrow}) [\cos k_0 x \cos \omega_0 t] = -(A_{\rightarrow} - A_{\leftarrow}) [\sin k_0 x \sin \omega_0 t]$$

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$$\tan k_0 x = -SWR \cdot \cot \omega_0 t = SWR \cdot \tan \omega_0 \bar{t} \text{ where: } \omega_0 \bar{t} = \omega_0 t - \pi/2$$

Speed of galloping wave zeros is the time derivative of root location x in units of light velocity c .

$$\frac{dx}{dt} = c \cdot SWR \frac{\sec^2 \omega_0 \bar{t}}{\sec^2 k_0 x} = \frac{c \cdot SWR}{\cos^2 \omega_0 \bar{t} + SWR^2 \cdot \sin^2 \omega_0 \bar{t}} = \begin{cases} c \cdot SWR & \text{for: } \bar{t} = 0, \pi, 2\pi, \dots \\ c \cdot SWQ & \bar{t} = \pi/2, 3\pi/2, \dots \end{cases}$$



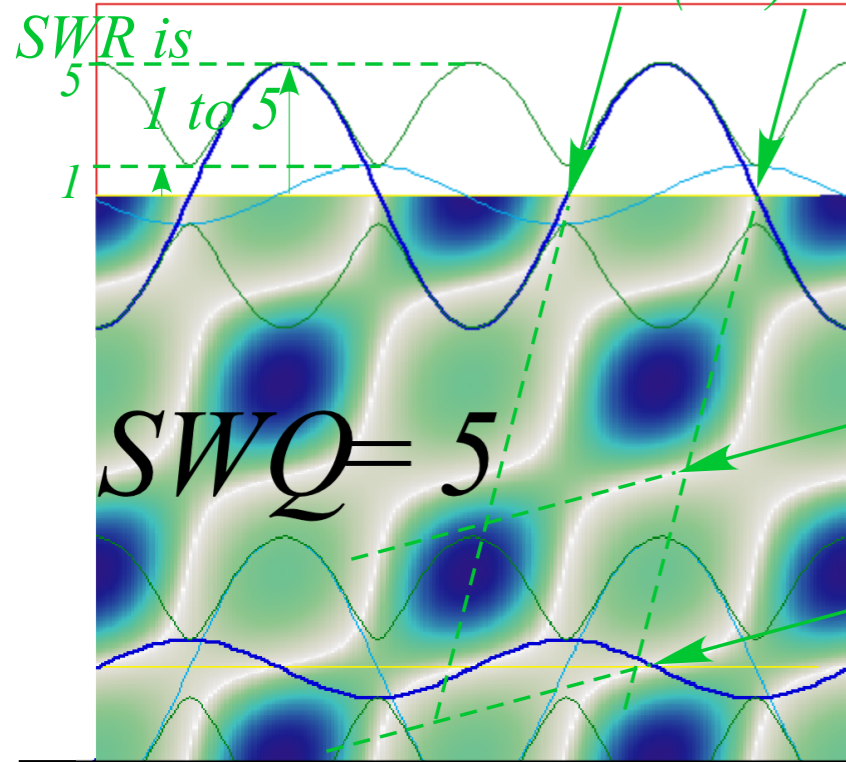
Wave-Zero Speed-Limits

Standing Wave Ratio SWR and Quotient SWQ

$$SWR = (E_{\rightarrow} - E_{\leftarrow}) / (E_{\rightarrow} + E_{\leftarrow}) = 1 / SWQ$$

Wave zeros
"resting"
at $(1/5)c$

$$SWR = 1/5$$



$$\begin{aligned} \omega_{\rightarrow} &= 2c & \omega_{\leftarrow} &= 2c \\ k_{\rightarrow} &= 2 & k_{\leftarrow} &= -2 \\ u_{GROUP} &= 0 & u_{PHASE} &= \infty \end{aligned}$$

Wave zeros
"galloping"
at $5c$

Wave zeros
"standing"
at 0-speed

$$SWR = 0$$

$$SWQ = \infty$$

$$E_{\leftarrow} = 0.5, E_{\rightarrow} = 0.5$$

$$E_{\leftarrow} = 0.4, E_{\rightarrow} = 0.6$$

Speed of galloping wave zeros is the time derivative of root location x in units of light velocity c .

$$\frac{dx}{dt} = c \cdot SWR \frac{\sec^2 \omega_0 \bar{t}}{\sec^2 k_0 x} = \frac{c \cdot SWR}{\cos^2 \omega_0 \bar{t} + SWR^2 \cdot \sin^2 \omega_0 \bar{t}} = \begin{cases} c \cdot SWR & \text{for: } \bar{t} = 0, \pi, 2\pi, \dots \\ c \cdot SWQ & \bar{t} = \pi/2, 3\pi/2, \dots \end{cases}$$

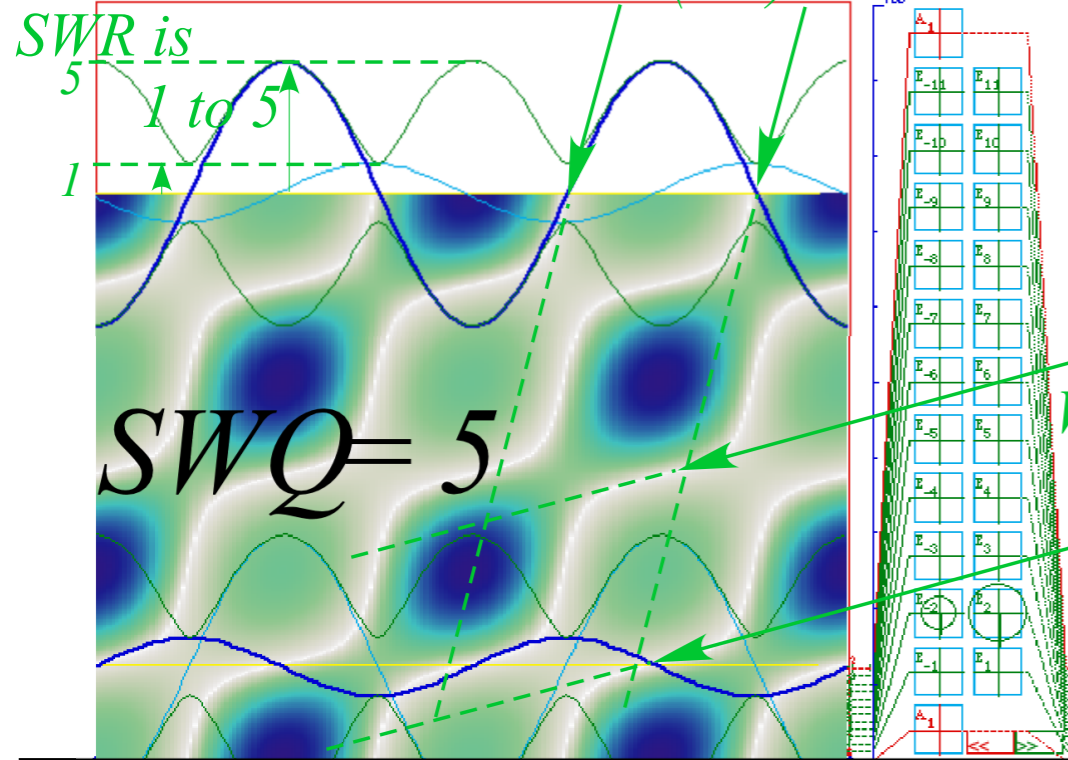
Wave-Zero Speed-Limits

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Wave zeros
"resting"
at $(1/5)c$

$$SWR = 1/5$$



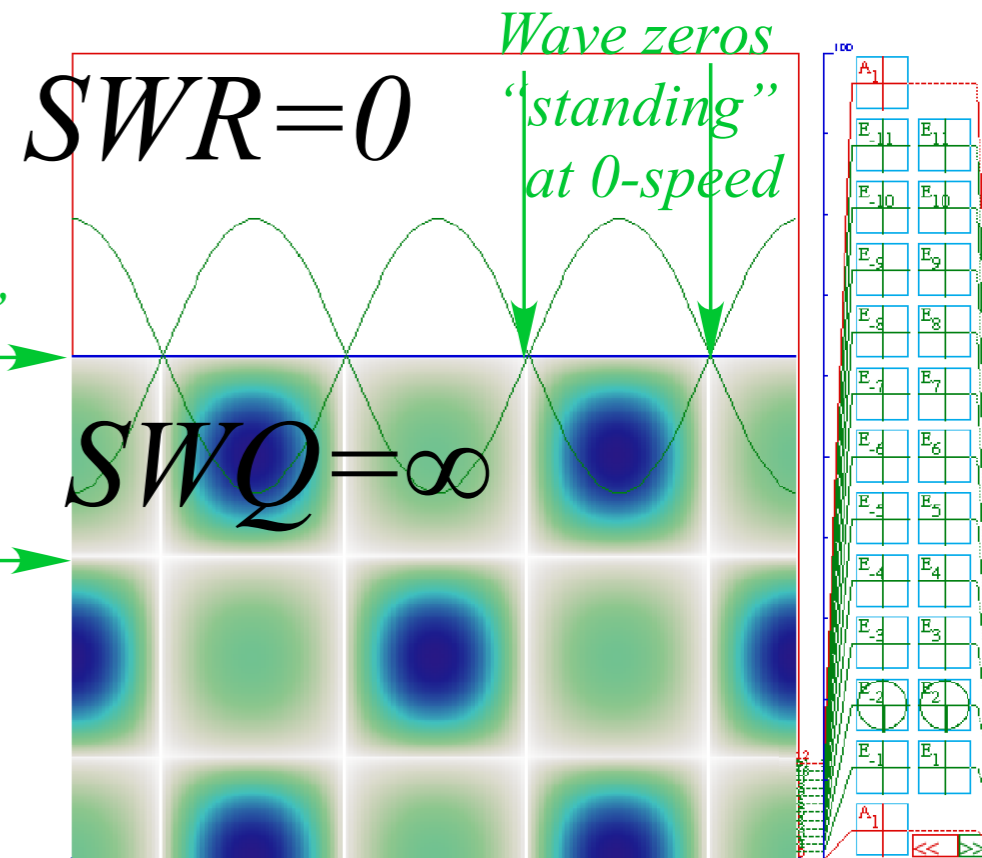
$$E_{\leftarrow} = 0.4, E_{\rightarrow} = 0.6$$

$SWR=1/5$ is analogous to (5-to-1)
Right Elliptic Polarization



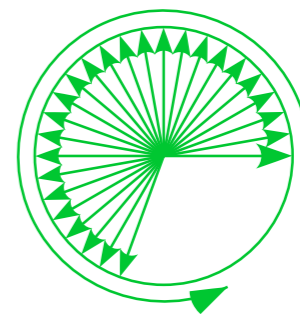
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Wave zeros
"galloping"
at $5c$



$$E_{\leftarrow} = 0.5, E_{\rightarrow} = 0.5$$

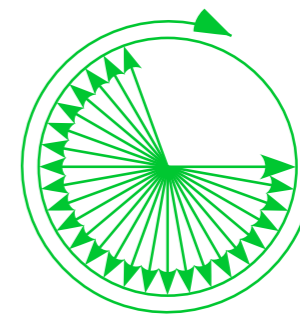
$SWR=1$ is analogous to (1,i)
Right Circular Polarization



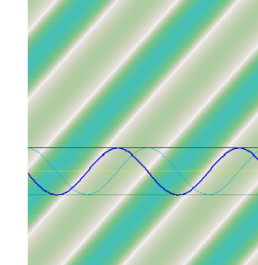
$SWR=0$ is analogous to (1,0)
x-Plane Linear Polarization



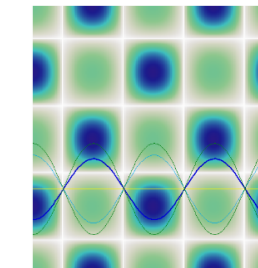
$SWR=-1$ is analogous to (1,-i)
Left Circular Polarization



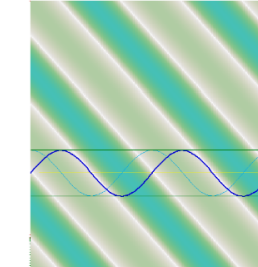
$SWR=+1$

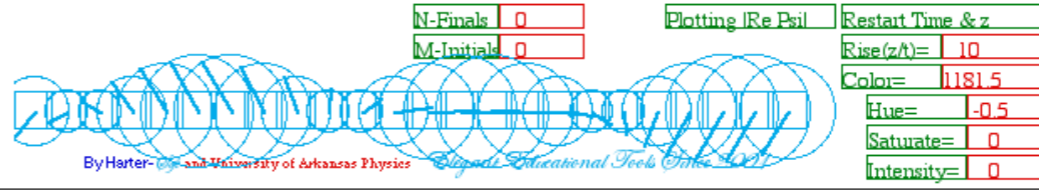
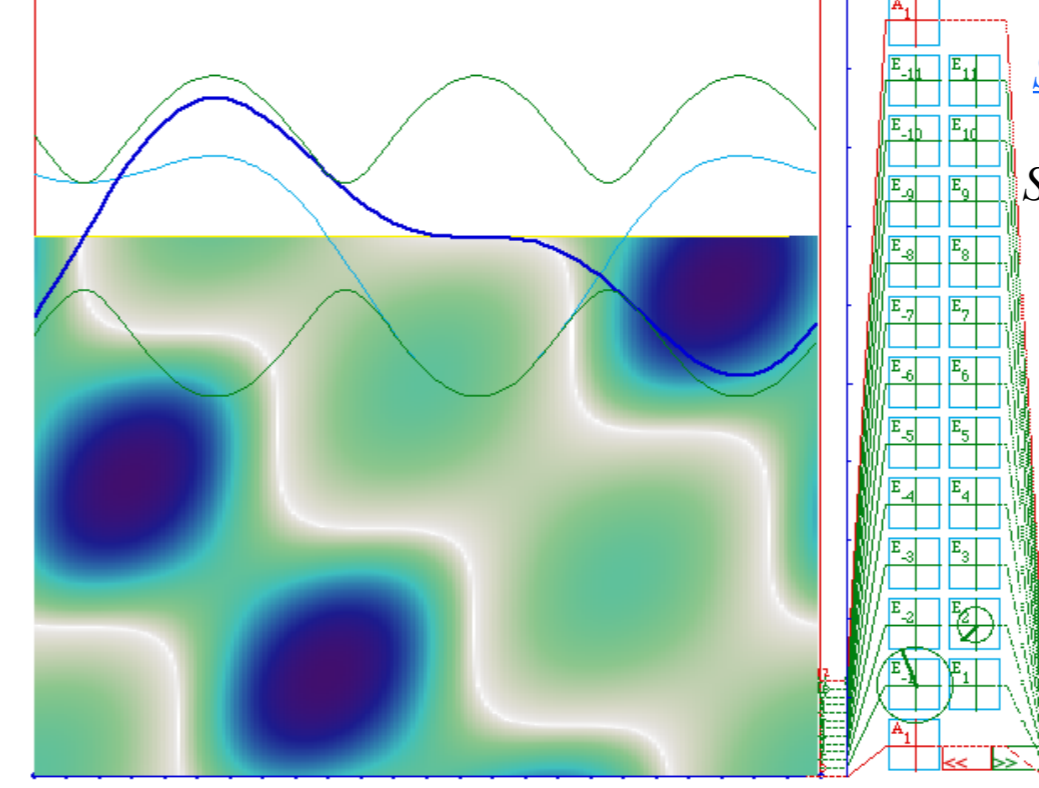
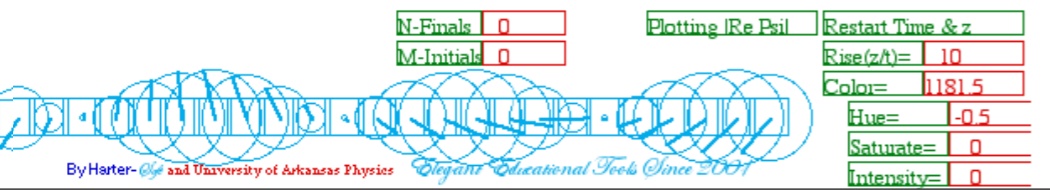
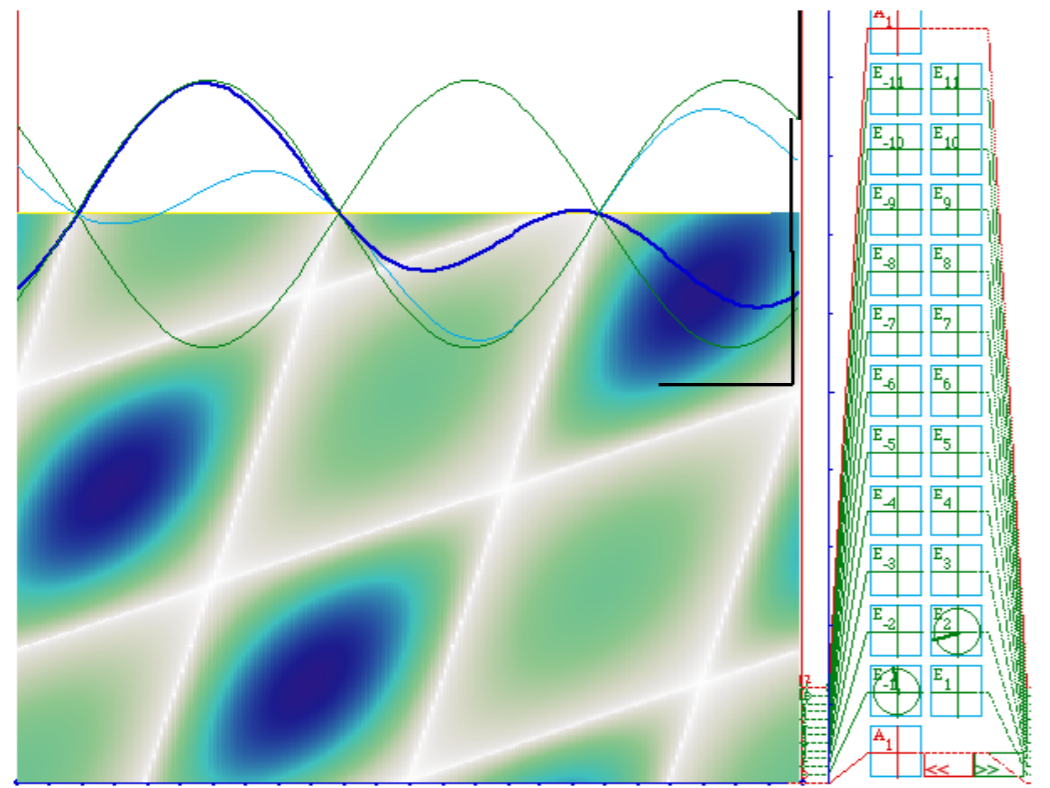
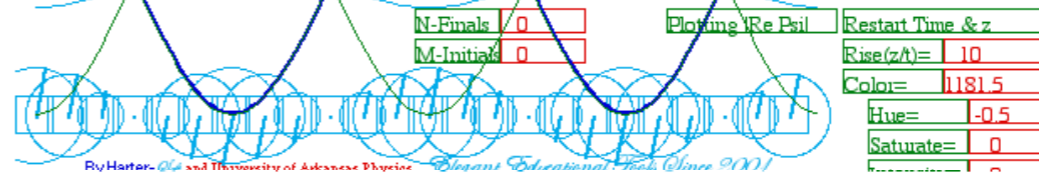
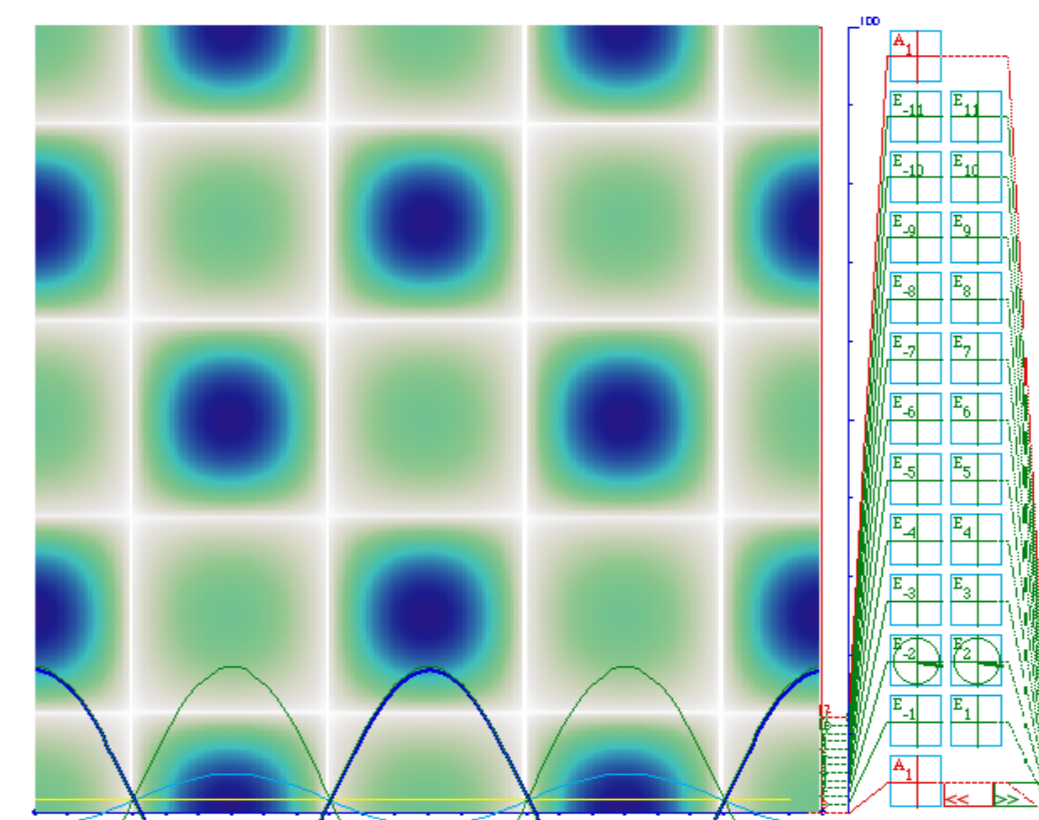
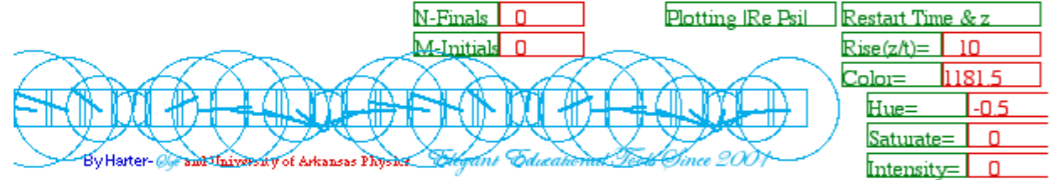
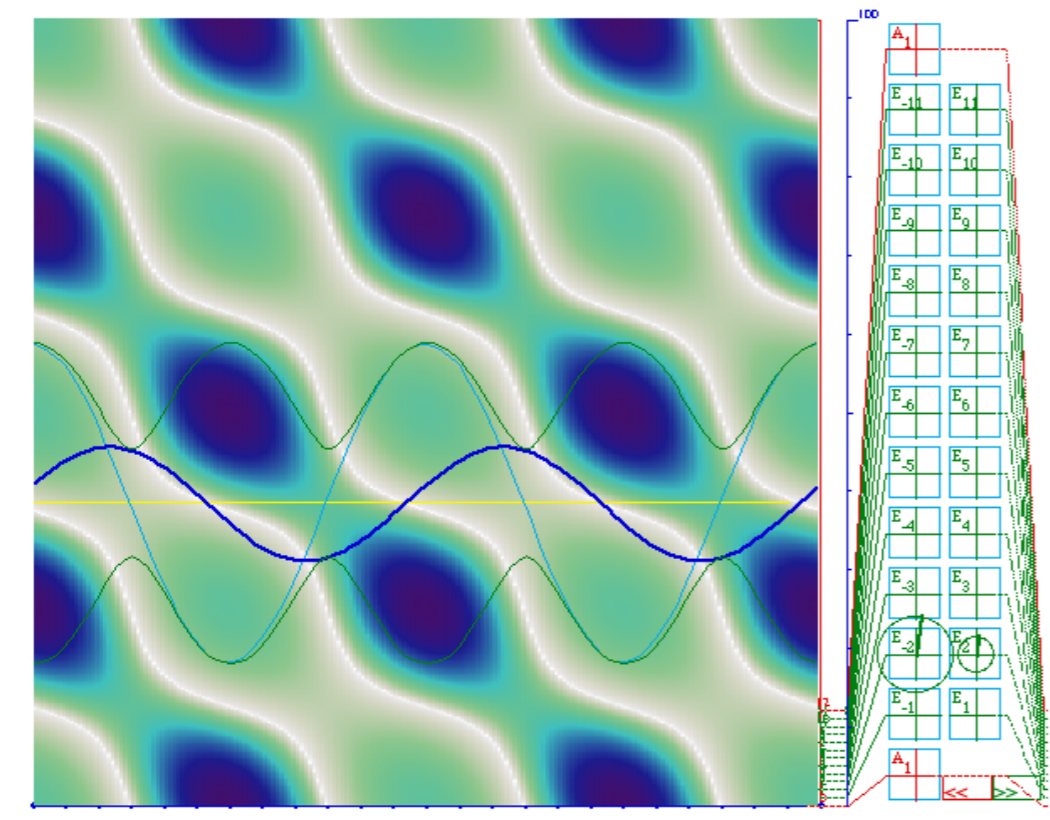


$SWR=0$



$SWR=-1$





Staircase Galloping
 Speed of galloping
 $SWR = +1/2$ cancelled
 by group velocity
 $u_{GROUP}/c = -1/2$.

Unmatched amplitudes giving galloping waves

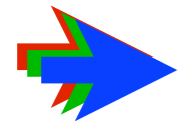
Standing Wave Ratio (SWR) and Standing Wave Quotient (SWQ)

Analogy with group and phase

Galloping waves

Analogy between wave galloping, Keplarian IHO orbits, and optical polarization

Galloping dynamics algebra



Waves that go back in time - The Feynman-Wheeler Switchback

The Ship-Barn-and-Butler saga of confused causality

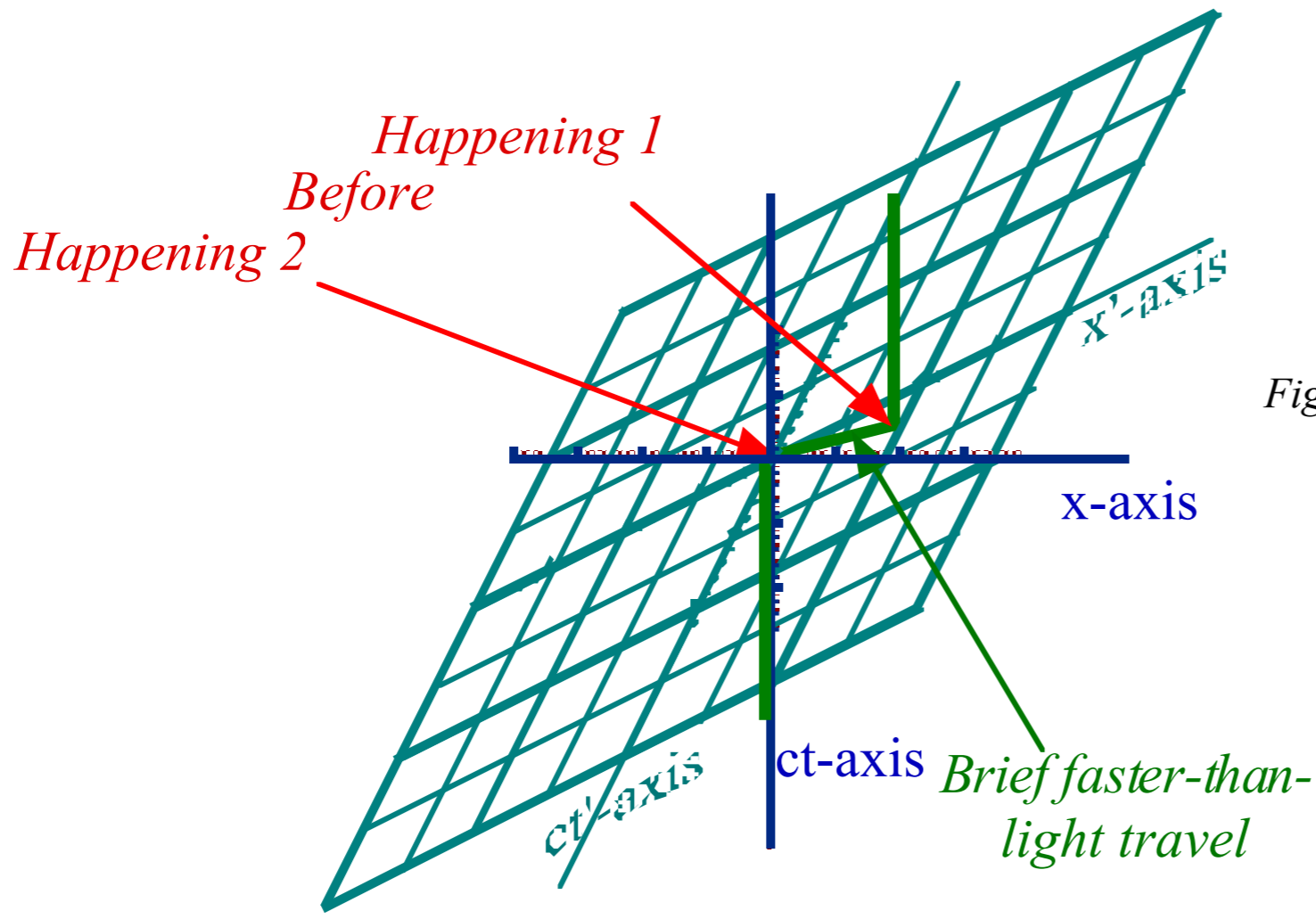


Fig. 2.B.10 Lighthouse plot of two Happenings

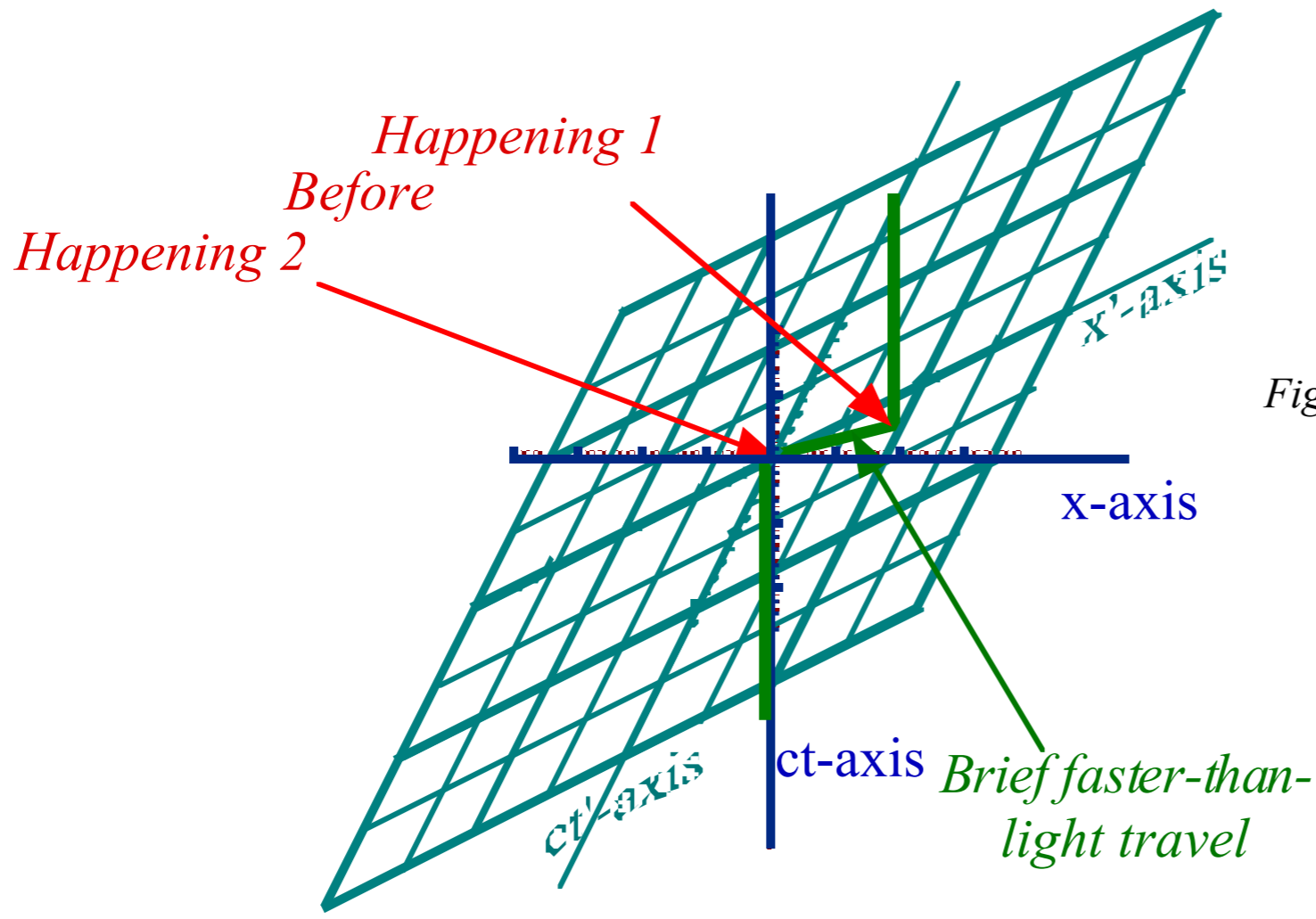


Fig. 2.B.10 Lighthouse plot of two Happenings

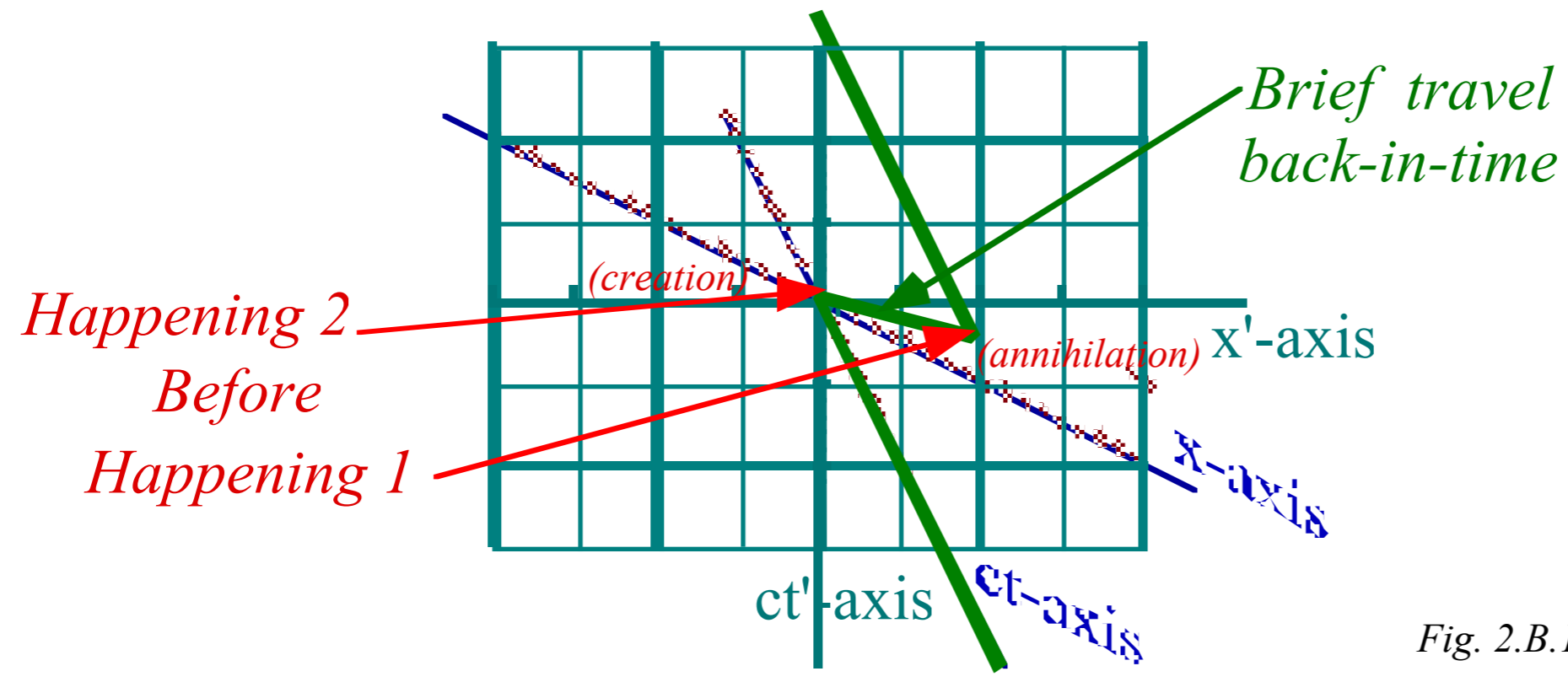


Fig. 2.B.11 Ship plot of two Happenings

Waves that go back in time - The Feynman-Wheeler Switchback

Minkowski Zero-Grids are Spacetime Switchbacks for $-u_{GROUP} < SWR < 0$

$\omega_{\rightarrow} = 4c$	$\omega_{\leftarrow} = 1c$
$k_{\rightarrow} = 4$	$k_{\leftarrow} = -1$
$u_{GROUP} = c3/5$	$u_{PHASE} = c5/3$

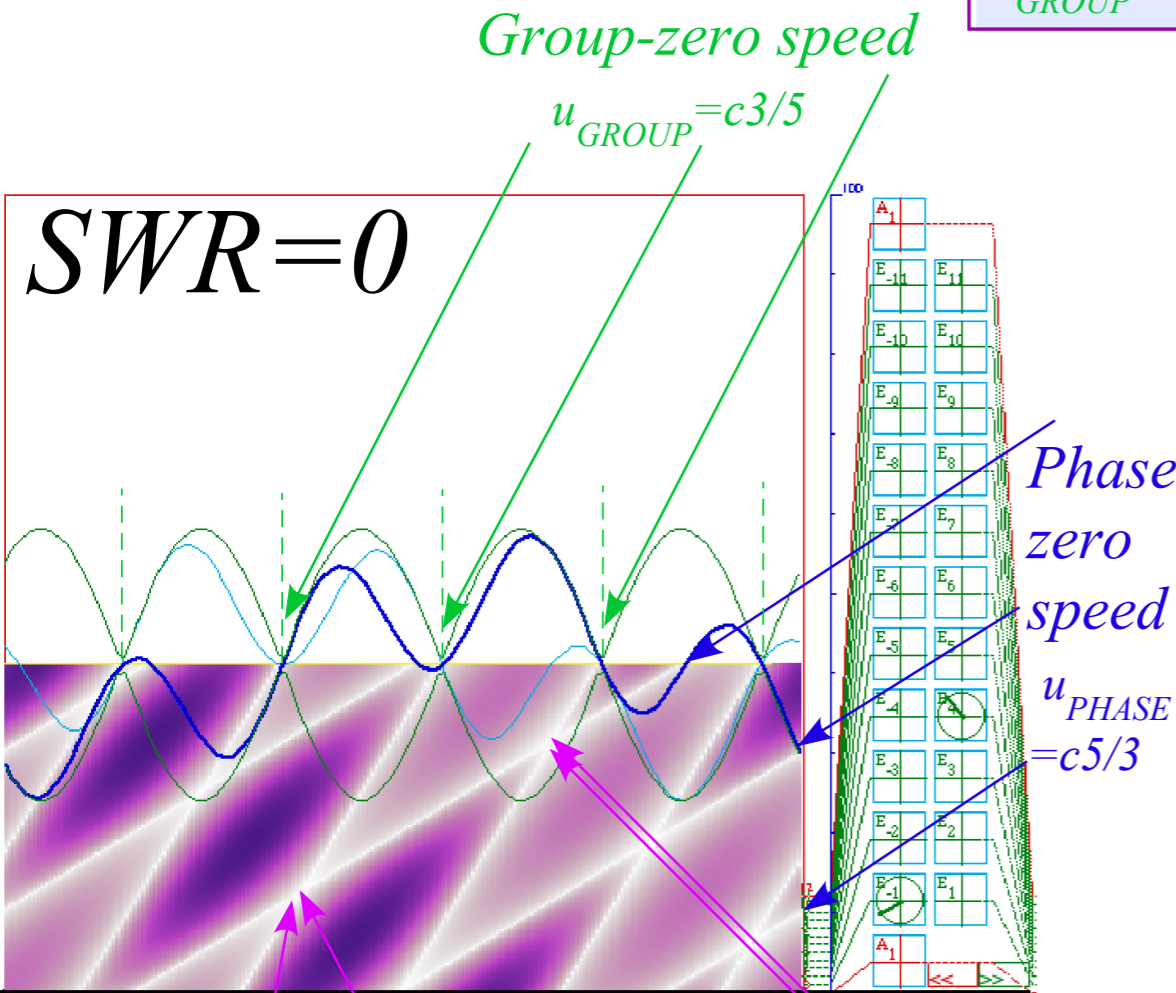
Group zero speed limit
 $u_{GROUP} + SWR$
 $\frac{1 + u_{GROUP} \cdot SWR}{c^2} = 5c/11$

$$\frac{\frac{3}{5} + \frac{-1}{5}}{1 + \frac{3-1}{5 \cdot 5}} = \frac{\frac{2}{5}}{\frac{22}{25}} = \frac{5}{11}$$

Phase "anti-zero" going "back-in-time"

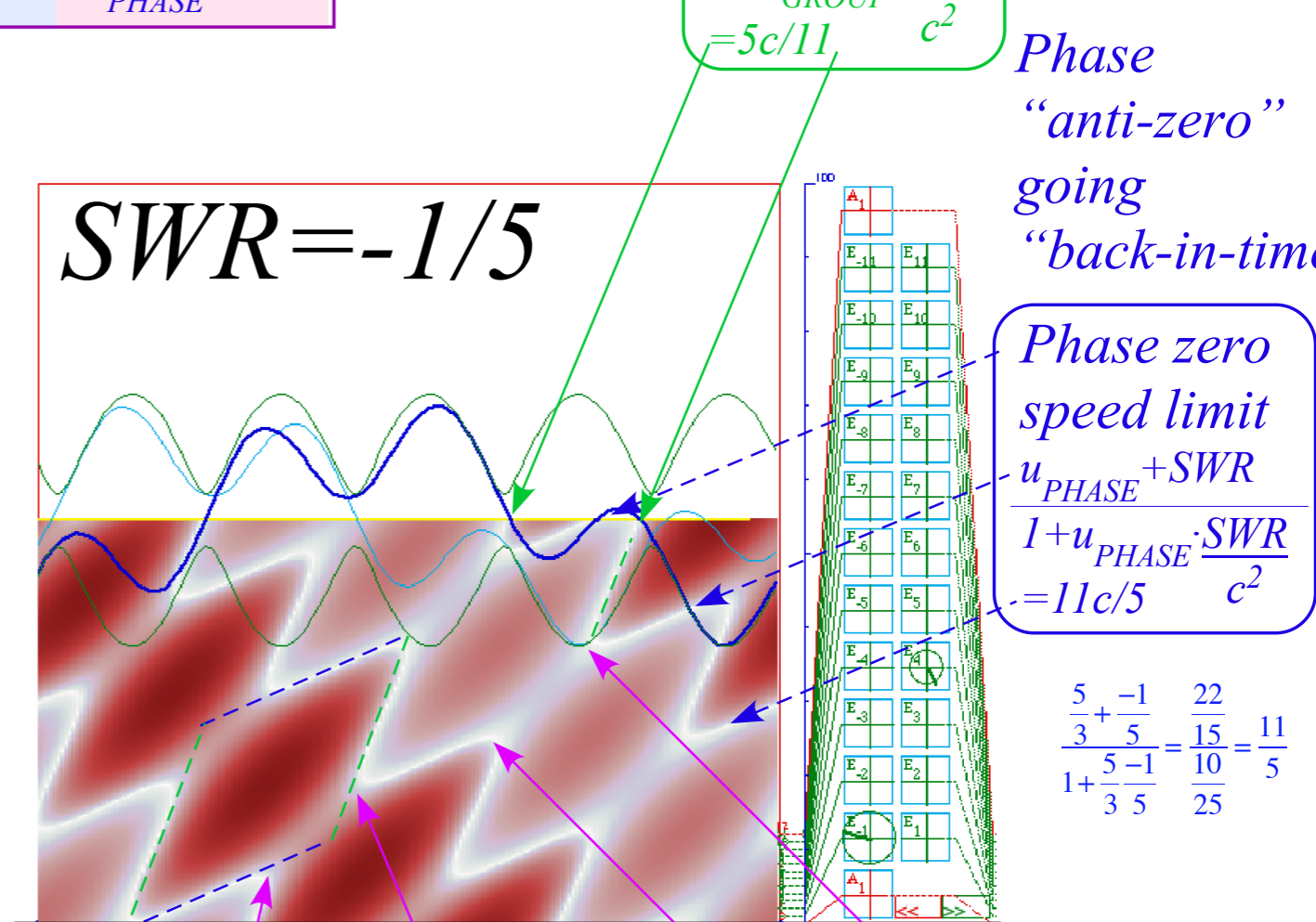
Phase zero speed limit
 $u_{PHASE} + SWR$
 $\frac{1 + u_{PHASE} \cdot SWR}{c^2} = 11c/5$

$$\frac{\frac{5}{3} + \frac{-1}{5}}{1 + \frac{5-1}{3 \cdot 5}} = \frac{\frac{22}{15}}{\frac{10}{25}} = \frac{11}{5}$$



$E_{\leftarrow} = 0.5, E_{\rightarrow} = 0.5$

Wave zero-anti-zero annihilation and creation occur together at the same spacetime point for $SWR=0$



$E_{\leftarrow} = 0.6, E_{\rightarrow} = 0.4$

Wave zero-anti-zero annihilation and creation occur separately at different spacetime points for $-u_{GROUP} < SWR < 0$

At High Speed 2-CW Modes Look More Like 1-CW Beams

$$\psi = E \sqrt{\frac{\epsilon_0}{h\nu}}$$

Various combinations of *opposite-k* 1-CW beams occur with open boundaries.

E-wave: $\mathbf{E} = \mathbf{E}_{\rightarrow} e^{i(k_{\rightarrow}x - \omega_{\rightarrow}t)} + \mathbf{E}_{\leftarrow} e^{i(k_{\leftarrow}x - \omega_{\leftarrow}t)}$ is related to Ψ -wave: $\Psi = \psi_{\rightarrow} e^{i(k_{\rightarrow}x - \omega_{\rightarrow}t)} + \psi_{\leftarrow} e^{i(k_{\leftarrow}x - \omega_{\leftarrow}t)}$

Standing Wave Ratio (or Quotient)

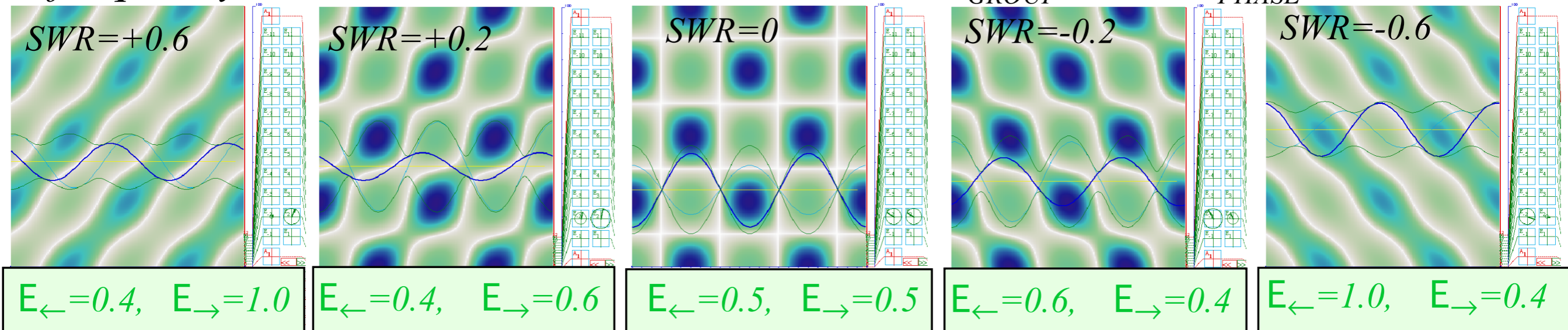
$$SWR = (\mathbf{E}_{\rightarrow} - \mathbf{E}_{\leftarrow}) / (\mathbf{E}_{\rightarrow} + \mathbf{E}_{\leftarrow}) = 1/SWQ$$

key numbers

Wave Group (or Phase) Velocity

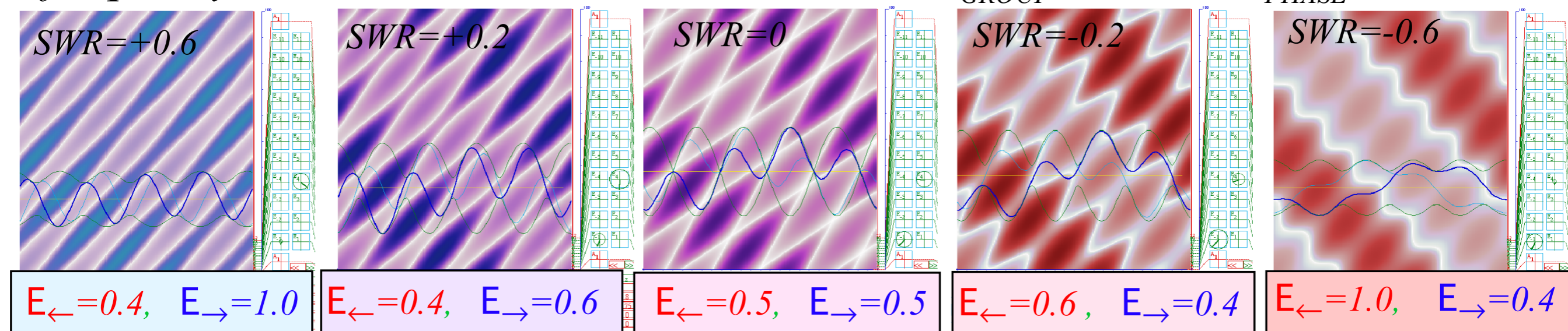
$$u_{GROUP}/c = (\omega_{\rightarrow} - \omega_{\leftarrow}) / (\omega_{\rightarrow} + \omega_{\leftarrow}) = c/u_{PHASE}$$


1-frequency case : $\omega_{\rightarrow} = 2c, k_{\rightarrow} = 2, \omega_{\leftarrow} = 2c, k_{\leftarrow} = -2$ gives: $u_{GROUP} = 0$ and $u_{PHASE} = \infty$



2-frequency case : $\omega_{\rightarrow} = 4c, k_{\rightarrow} = 4, \omega_{\leftarrow} = 1c, k_{\leftarrow} = -1$ gives: $u_{GROUP}/c = 3/5$ and $u_{PHASE}/c = 5/3$

Staircase Galloping





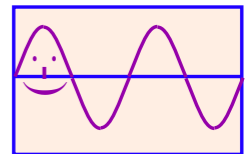
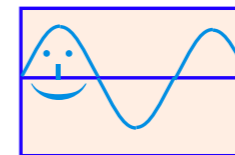
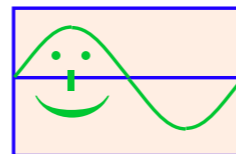
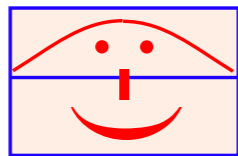
1st Quantization: Quantizing phase variables ω and k
Understanding how quantum transitions require “mixed-up” states
Closed cavity vs ring cavity

Quantized ω and k Counting wave kink numbers

If everything is made of waves then we expect *quantization* of everything because waves only thrive if *integral* numbers n of their “kinks” fit into whatever structure (box, ring, etc.) they’re supposed to live. The numbers n are called *quantum numbers*.

OK box quantum numbers: $n=1$ $n=2$ $n=3$ $n=4$

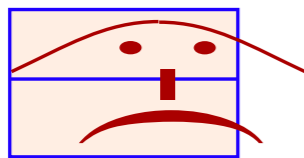
(+ integers only)



Some

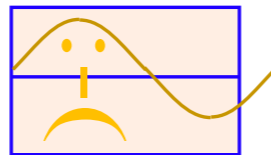
NOT OK numbers: $n=0.67$

too fat!



$n=1.7$

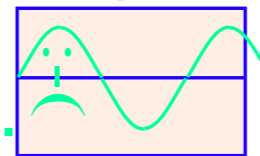
too thin!



$n=2.59$

wrong color again!

misfits...



$n=4$

...not tolerated!

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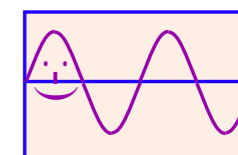
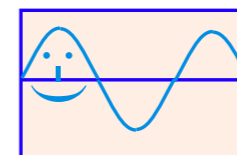
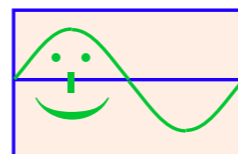
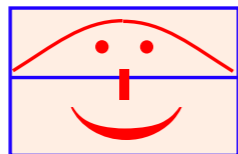
This doesn't mean a system's energy can't vary continuously between "OK" values $E_1, E_2, E_3, E_4, \dots$

Quantized ω and k Counting wave kink numbers

If everything is made of waves then we expect *quantization* of everything because waves only thrive if *integral* numbers n of their “kinks” fit into whatever structure (box, ring, etc.) they’re supposed to live. The numbers n are called *quantum numbers*.

OK box quantum numbers: $n=1$ $n=2$ $n=3$ $n=4$

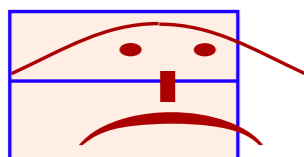
(+ integers only)



Some

NOT OK numbers: $n=0.67$

too fat!



$n=1.7$

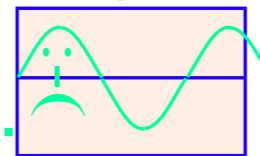
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misfits...



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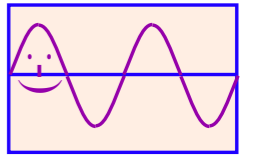
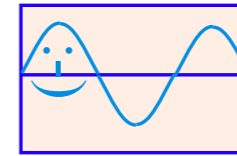
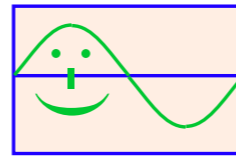
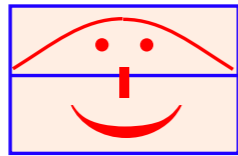
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Understanding how quantum transitions require “mixed-up” states
Closed cavity vs ring cavity

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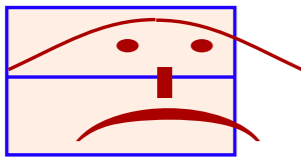
(+ integers only)



Some

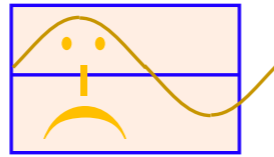
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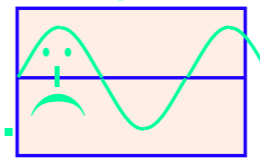
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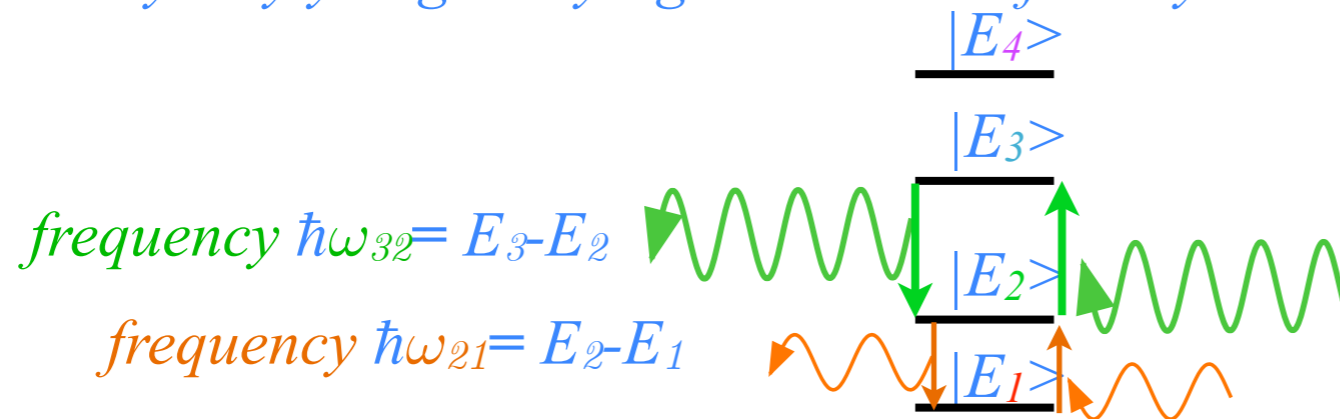


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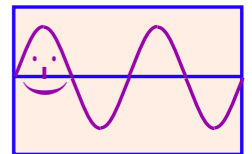
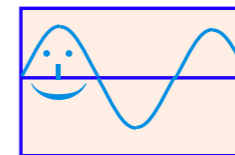
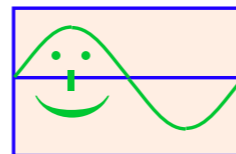
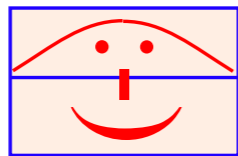


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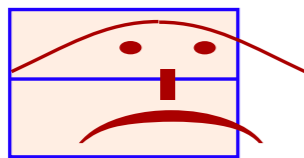
(+ integers only)



Some

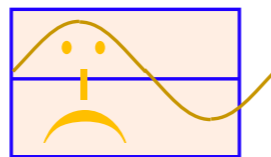
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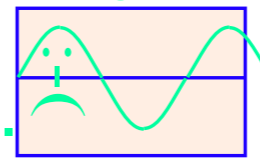
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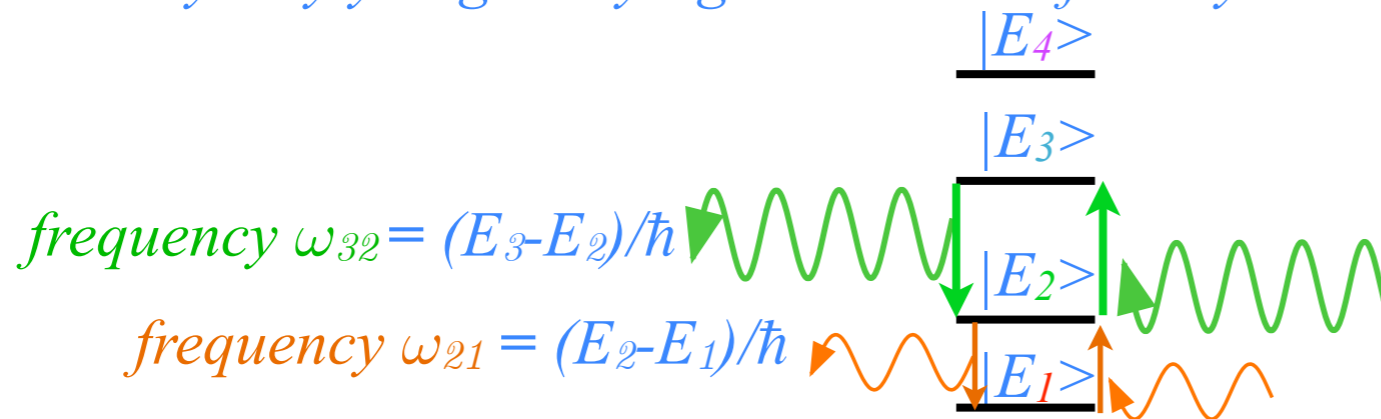


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...not tolerated!

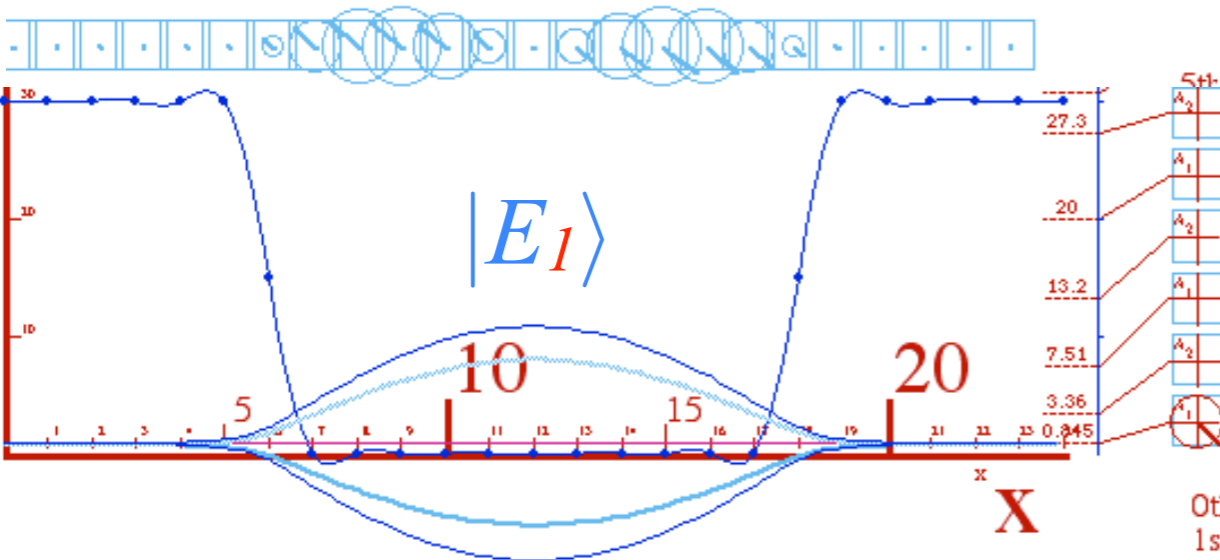
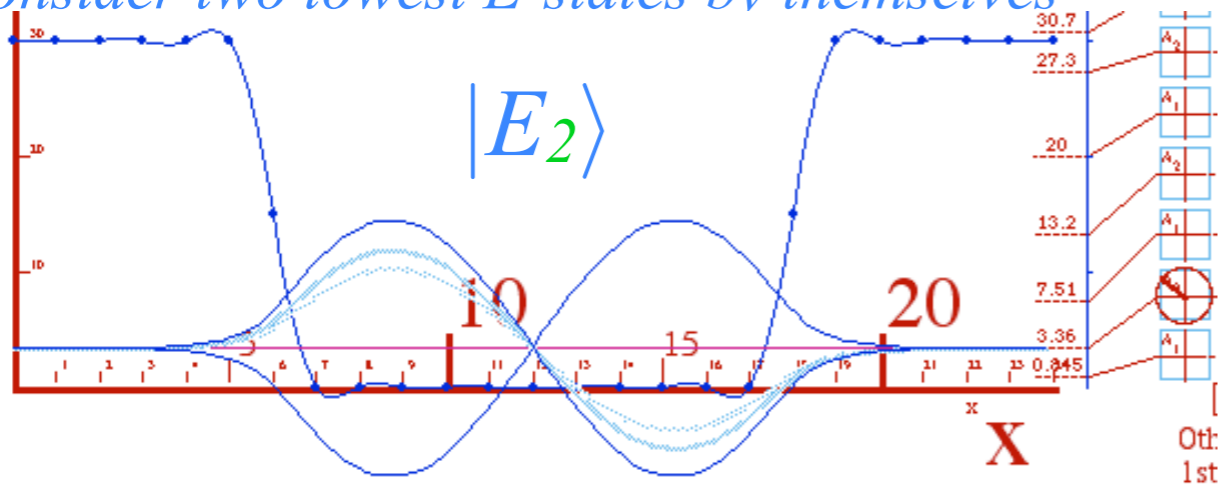
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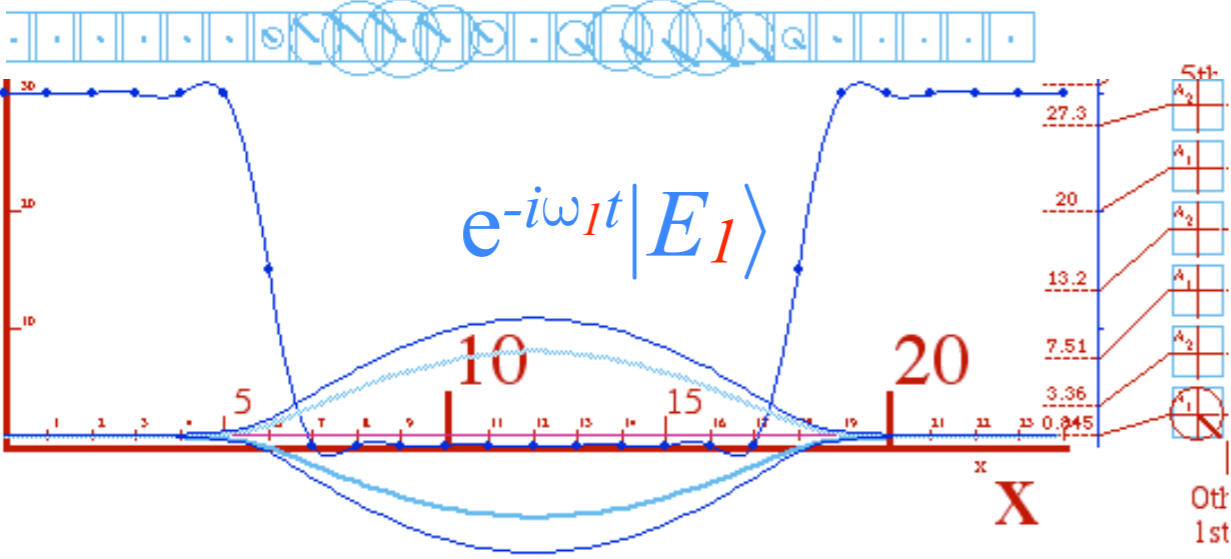
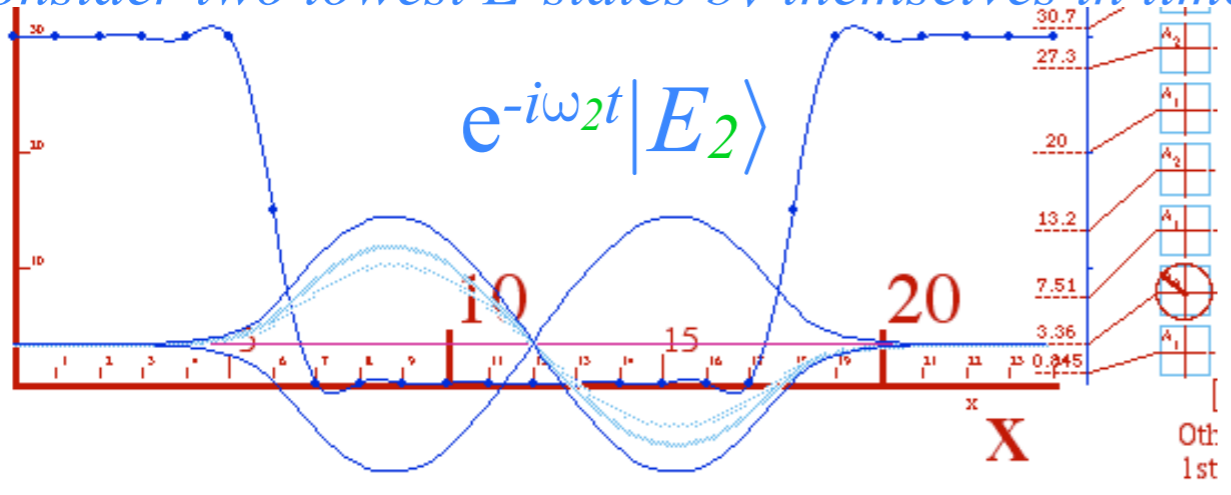
These *eigenstates* are the only ways the system can “play dead” ...
 ... “sleep with the fishes” ...

Consider two lowest E -states by themselves



By Harter-*et al* and University of Arkansas Physics *Elegant Educational Tools Since 2001*

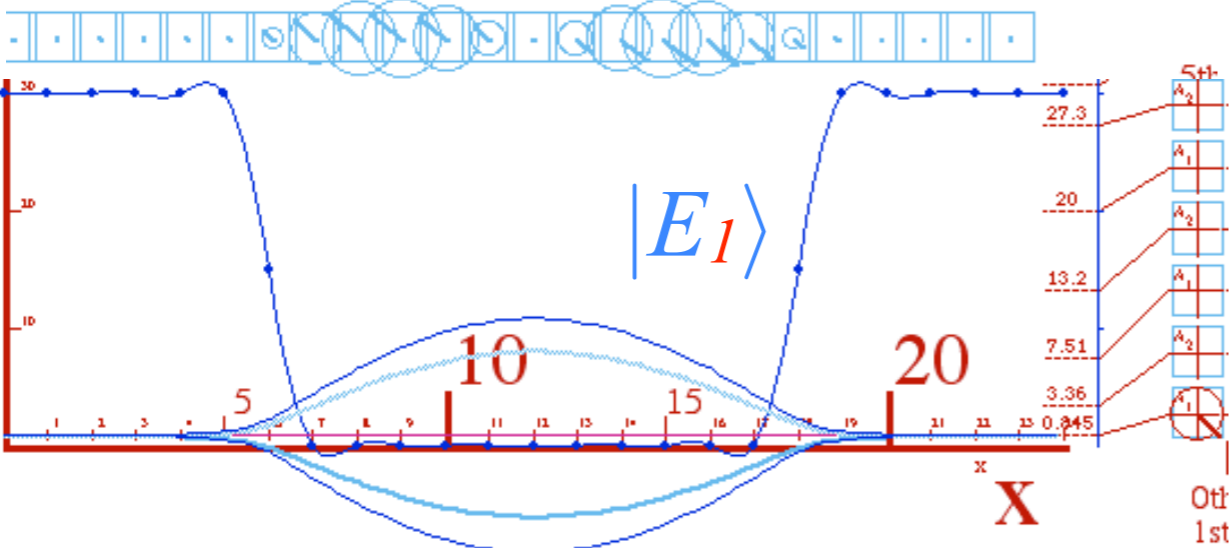
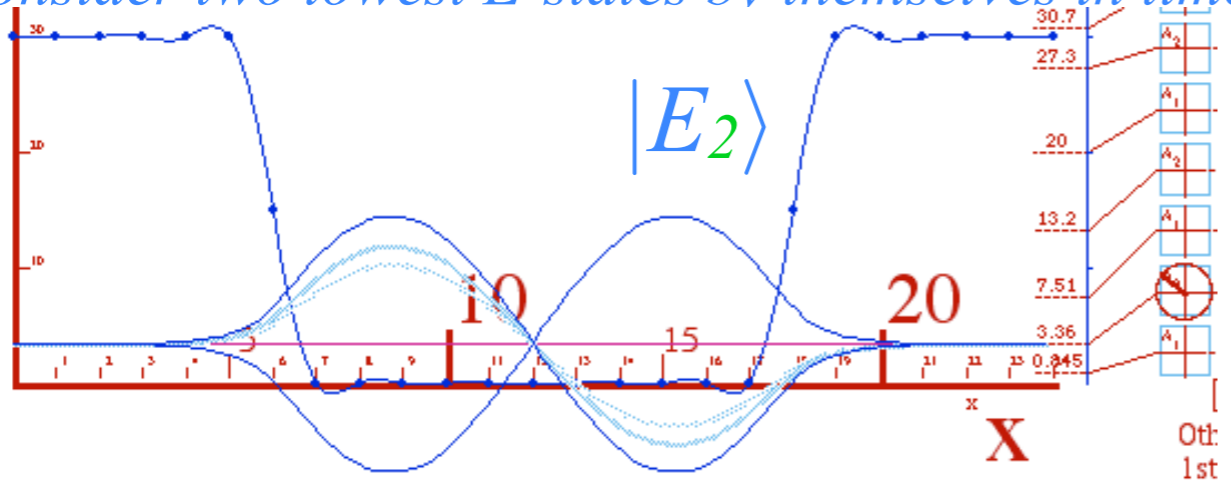
Consider two lowest E-states by themselves in time



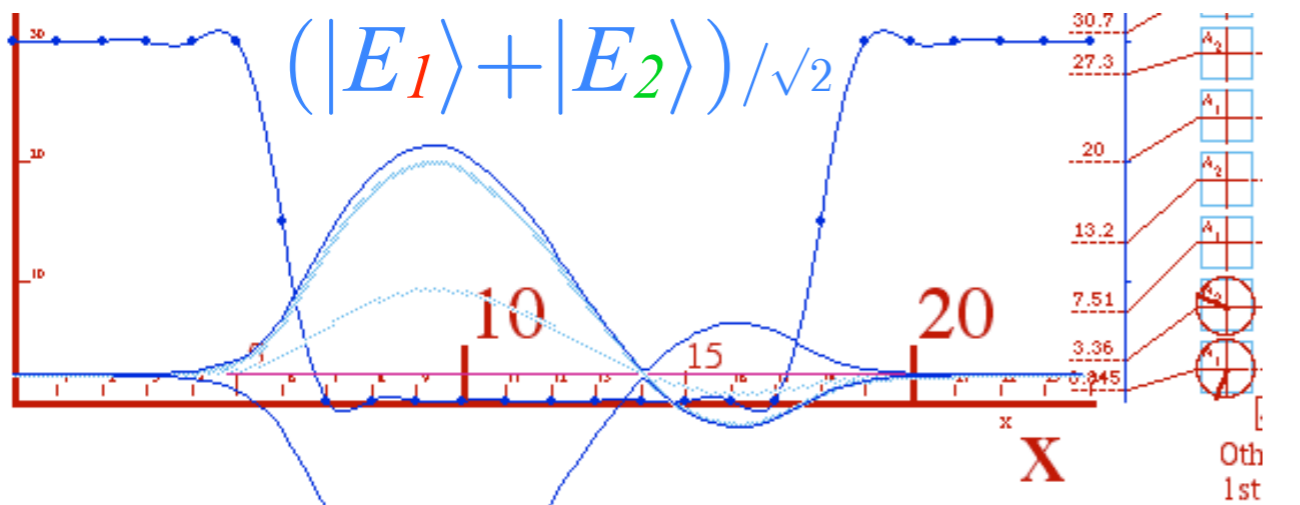
By Harter-*et al* and University of Arkansas Physics *Elegant Educational Tools Since 2001*

Consider two lowest E -states by themselves in time

Now combine (add) them

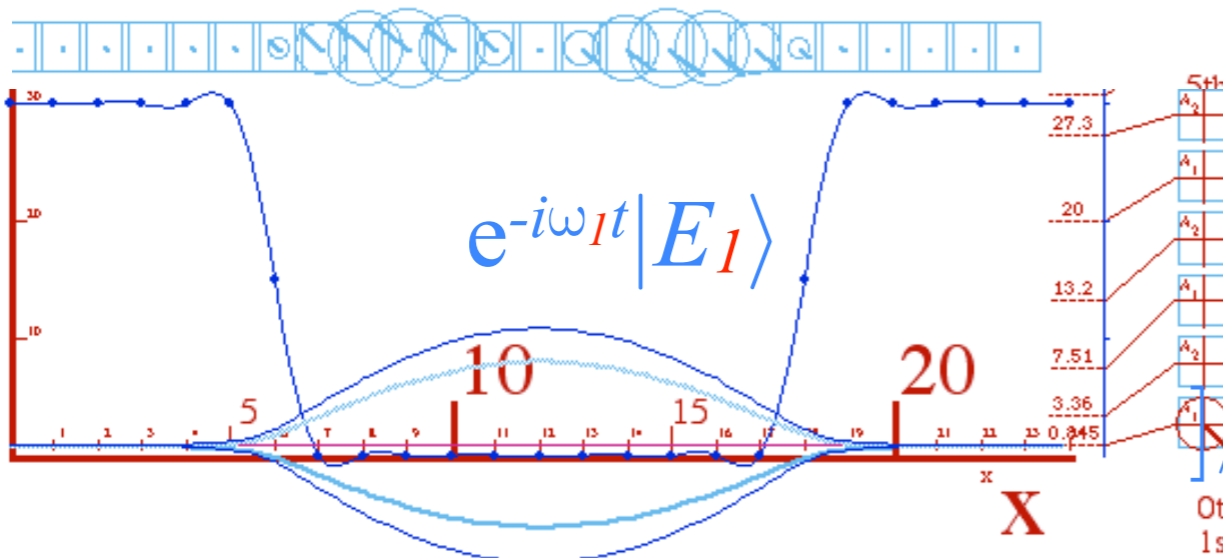
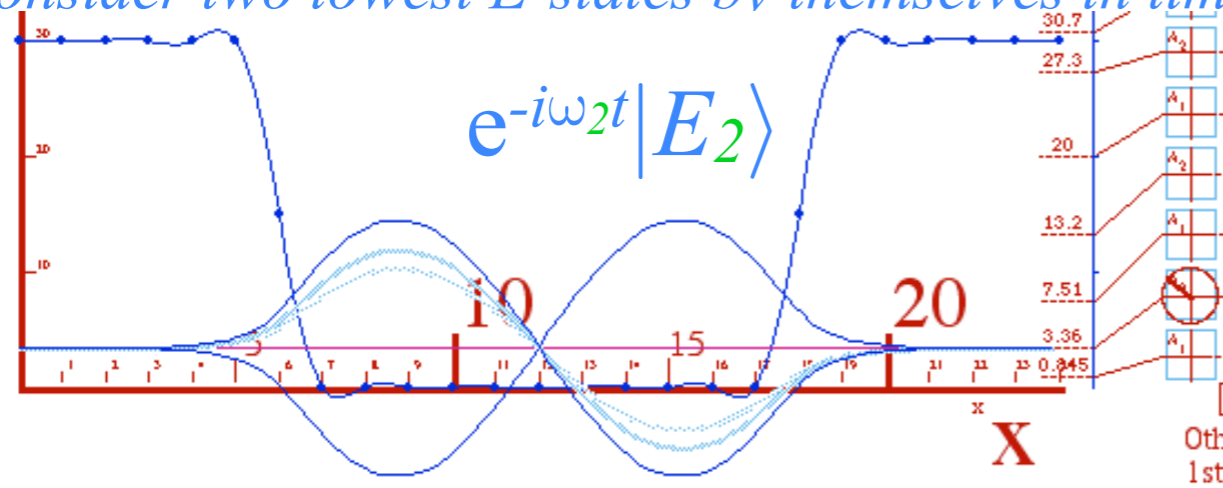


By Harter- and University of Arkansas Physics *Elegant Educational Tools Since 2001*



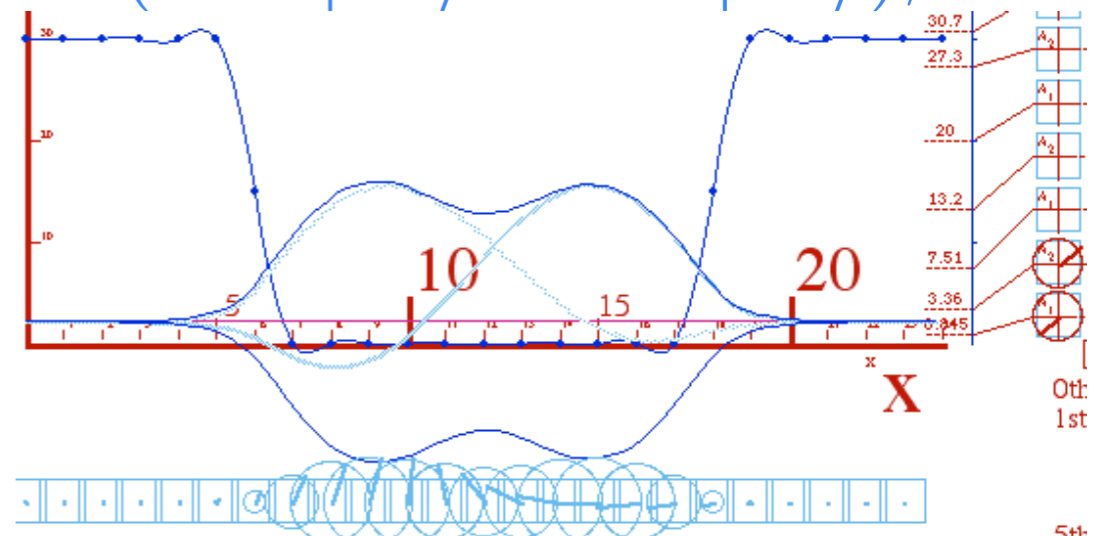
By Harter- and University of Arkansas Physics *Elegant Educational Tools Since 2001*

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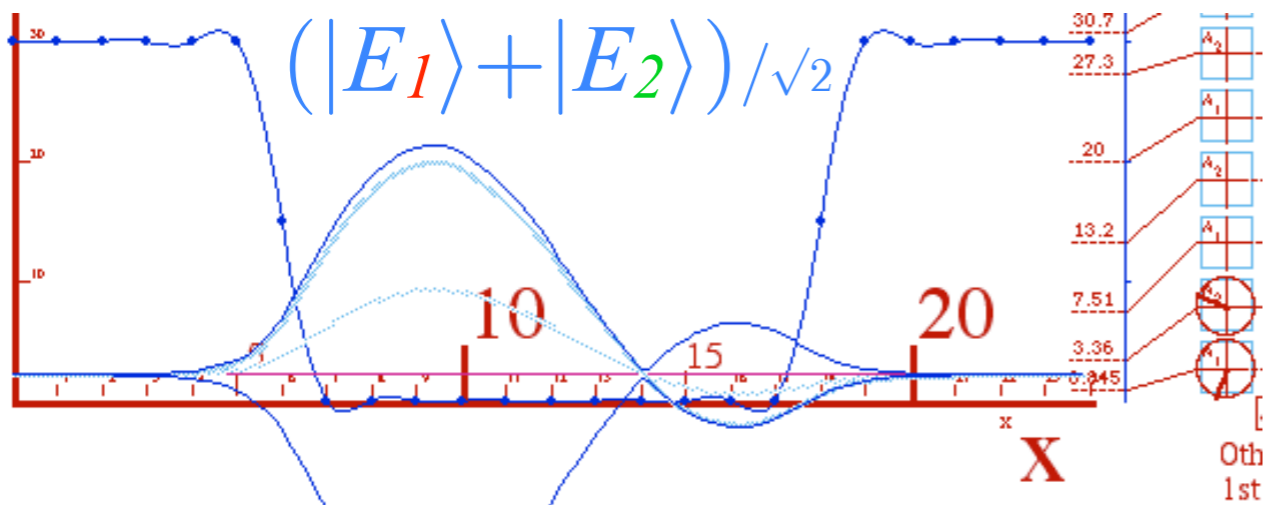
Now combine (add) them and let time roll!

$$(e^{-i\omega_1 t} |E_1\rangle + e^{-i\omega_2 t} |E_2\rangle) / \sqrt{2}$$



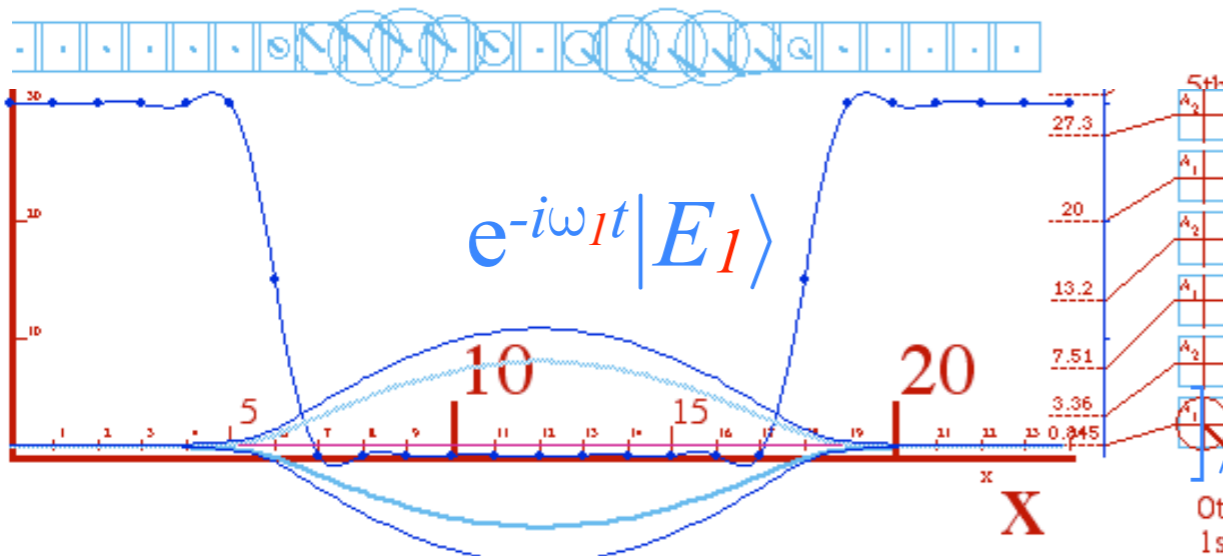
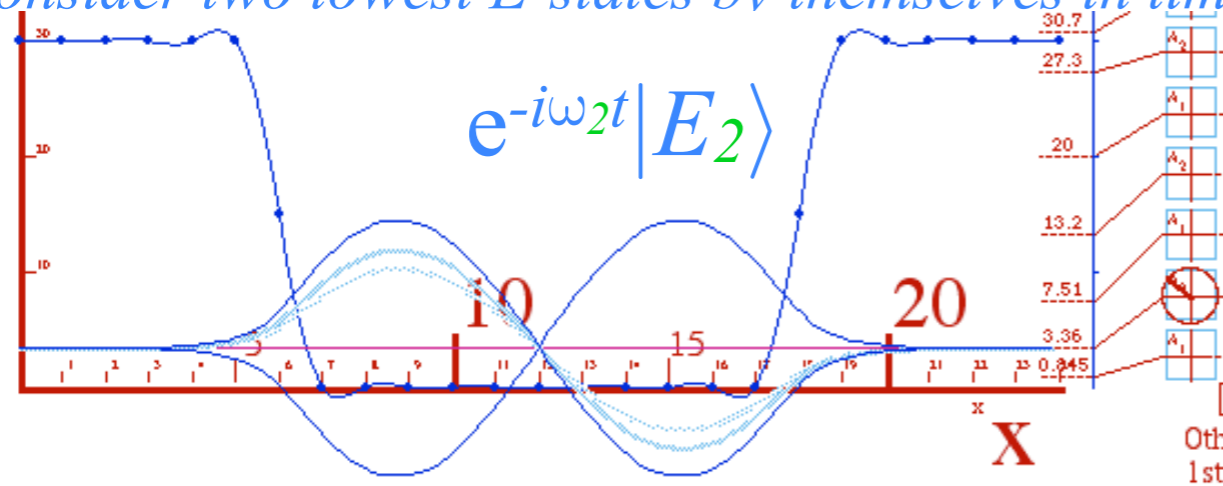
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$$(|E_1\rangle + |E_2\rangle) / \sqrt{2}$$



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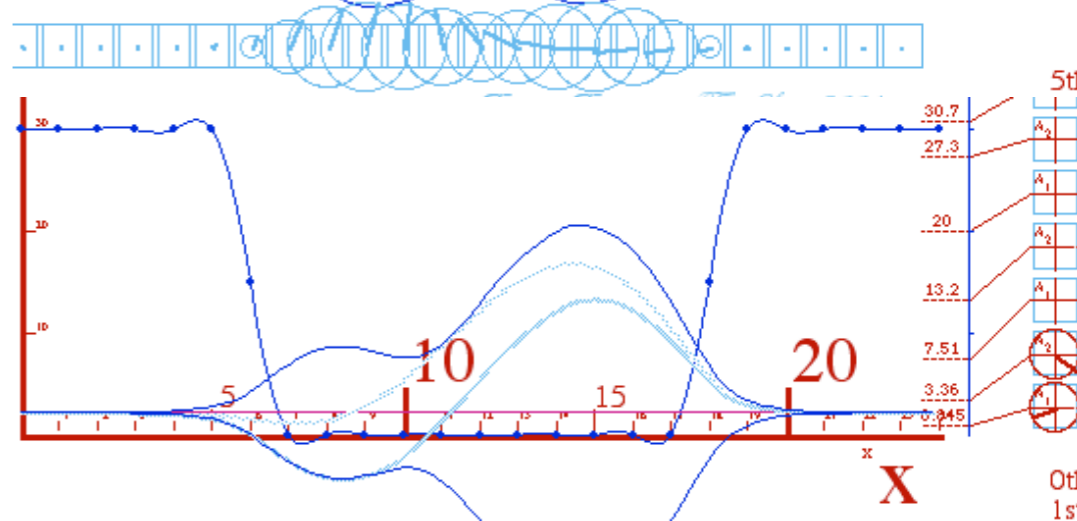
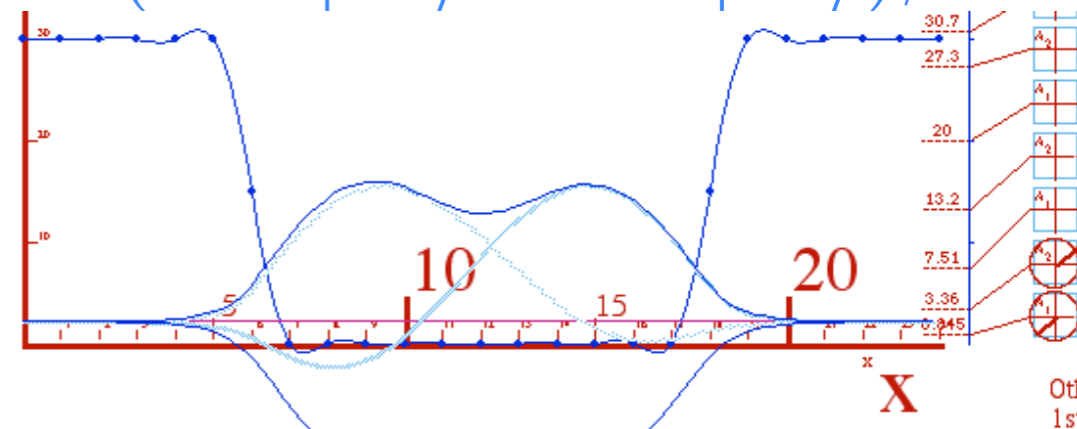
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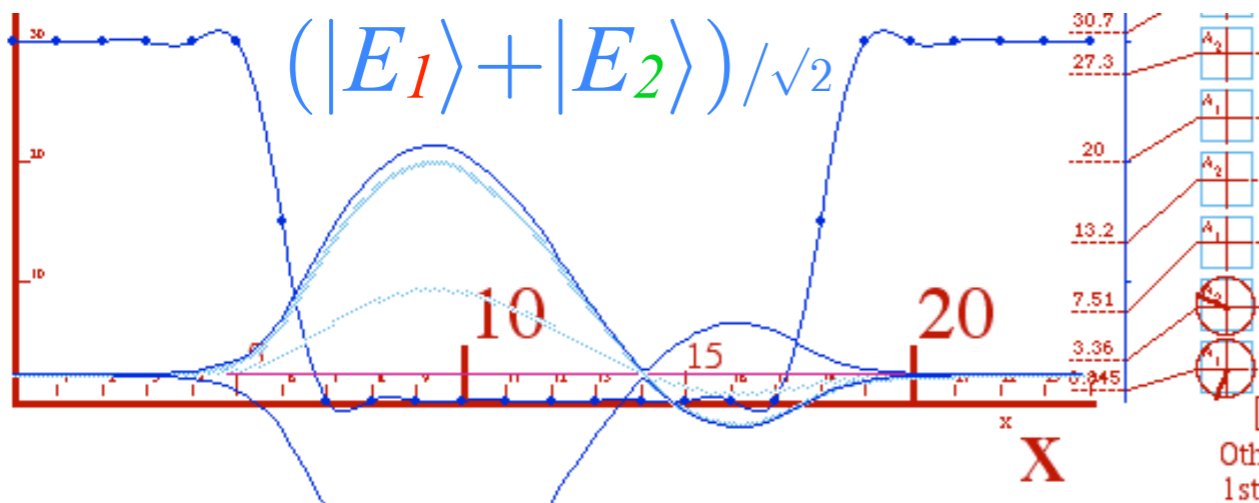
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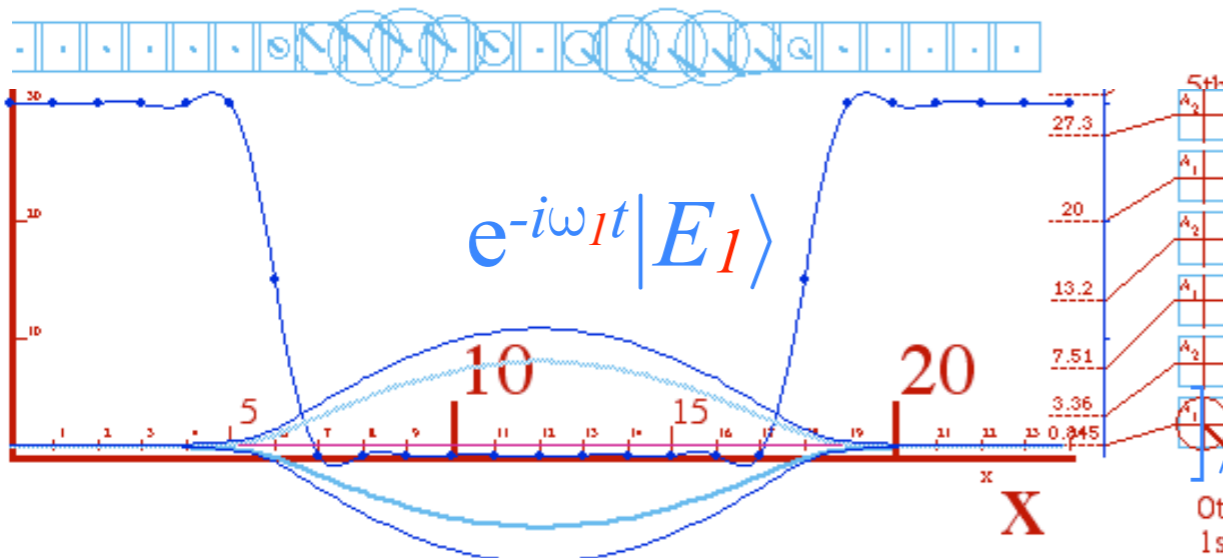
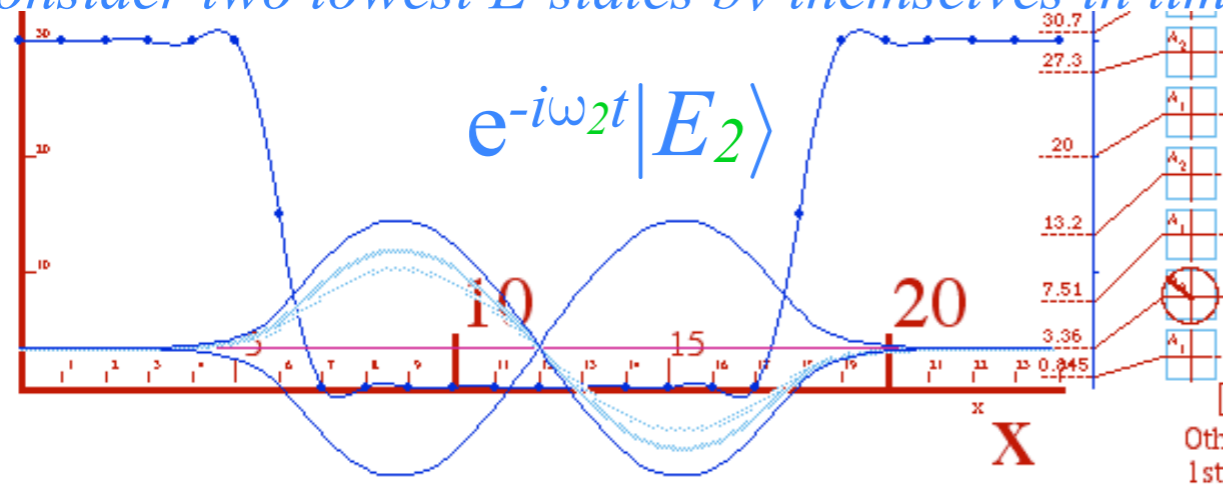


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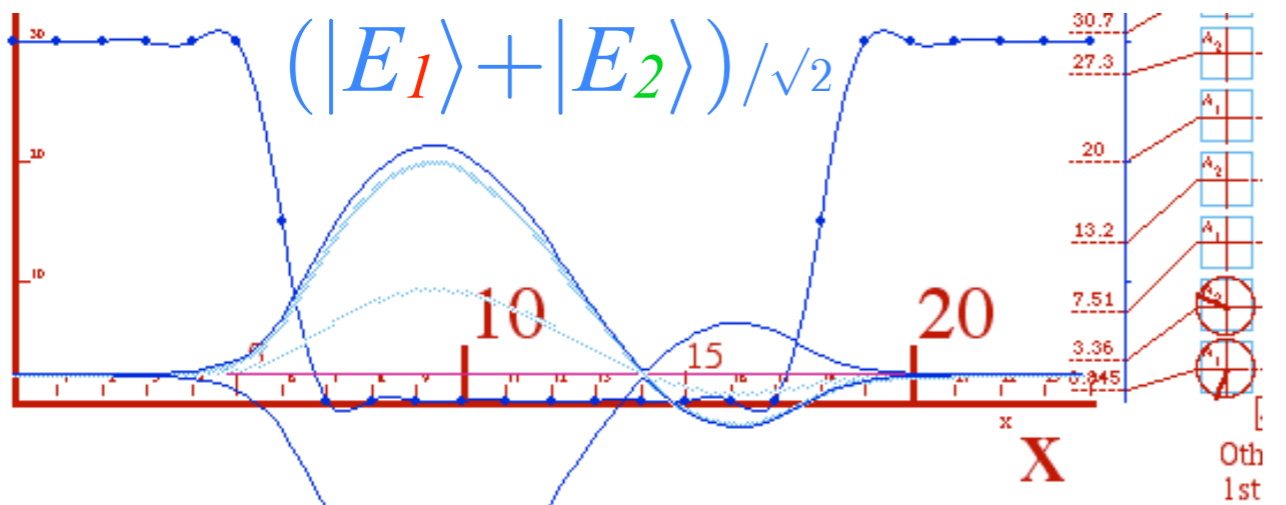


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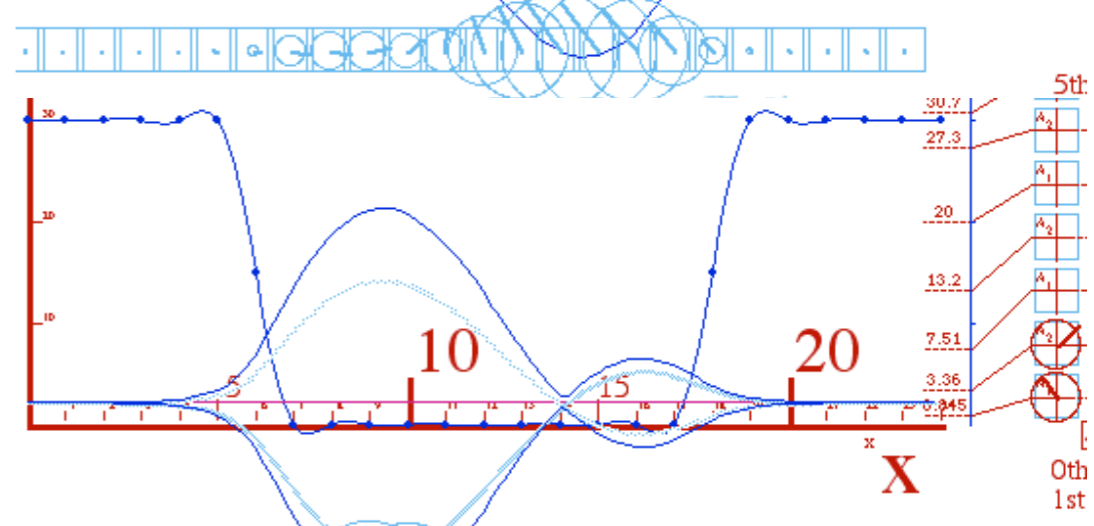
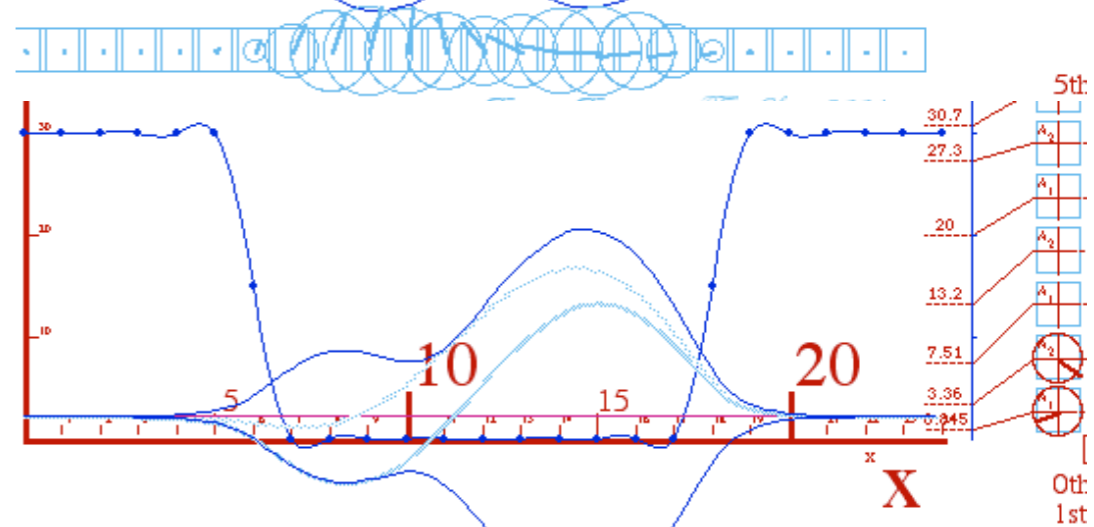
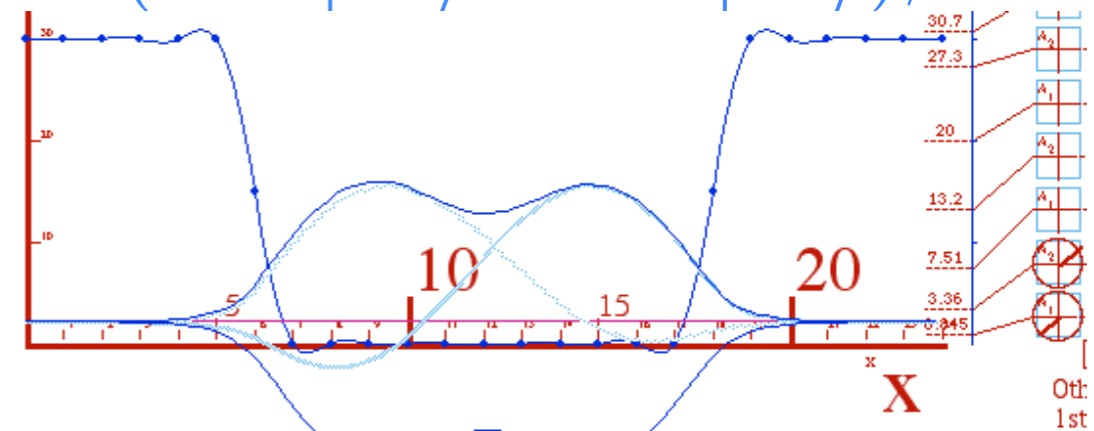
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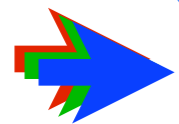
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1st Quantization: Quantizing phase variables ω and k

Understanding how quantum transitions require “mixed-up” states



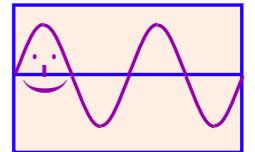
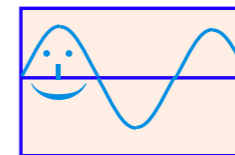
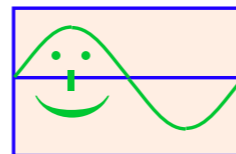
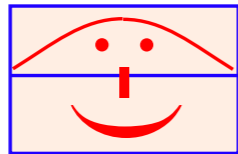
Closed cavity vs ring cavity

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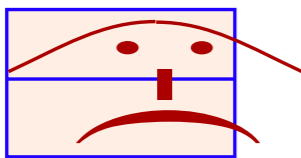
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Some

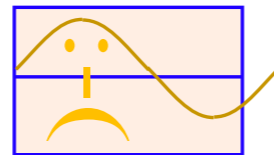
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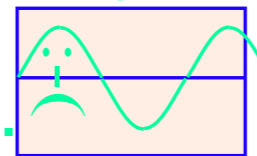
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misfits...



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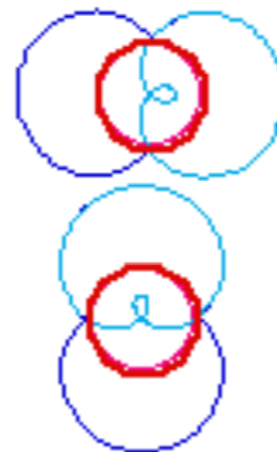
Rings tolerate a *zero* (kinkless) quantum wave but require \pm integral wave number.

OK ring quantum numbers: $m=0$

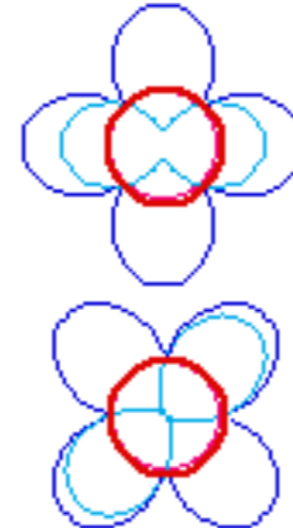
(\pm integral number of wavelengths)



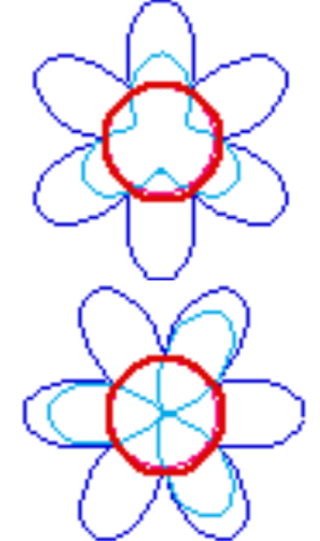
$m=\pm 1$



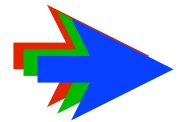
$m=\pm 2$



$m=3$



Bohr’s models of *atomic spectra* (1913-1923) are beginnings of *quantum wave mechanics* built on *Planck-Einstein* (1900-1905) relation $E=h\nu$. *DeBroglie* relation $p=h/\lambda$ comes around 1923.



2nd Quantization: Quantizing amplitudes (“photons”, “vibrons”, and “what-ever-ons”)

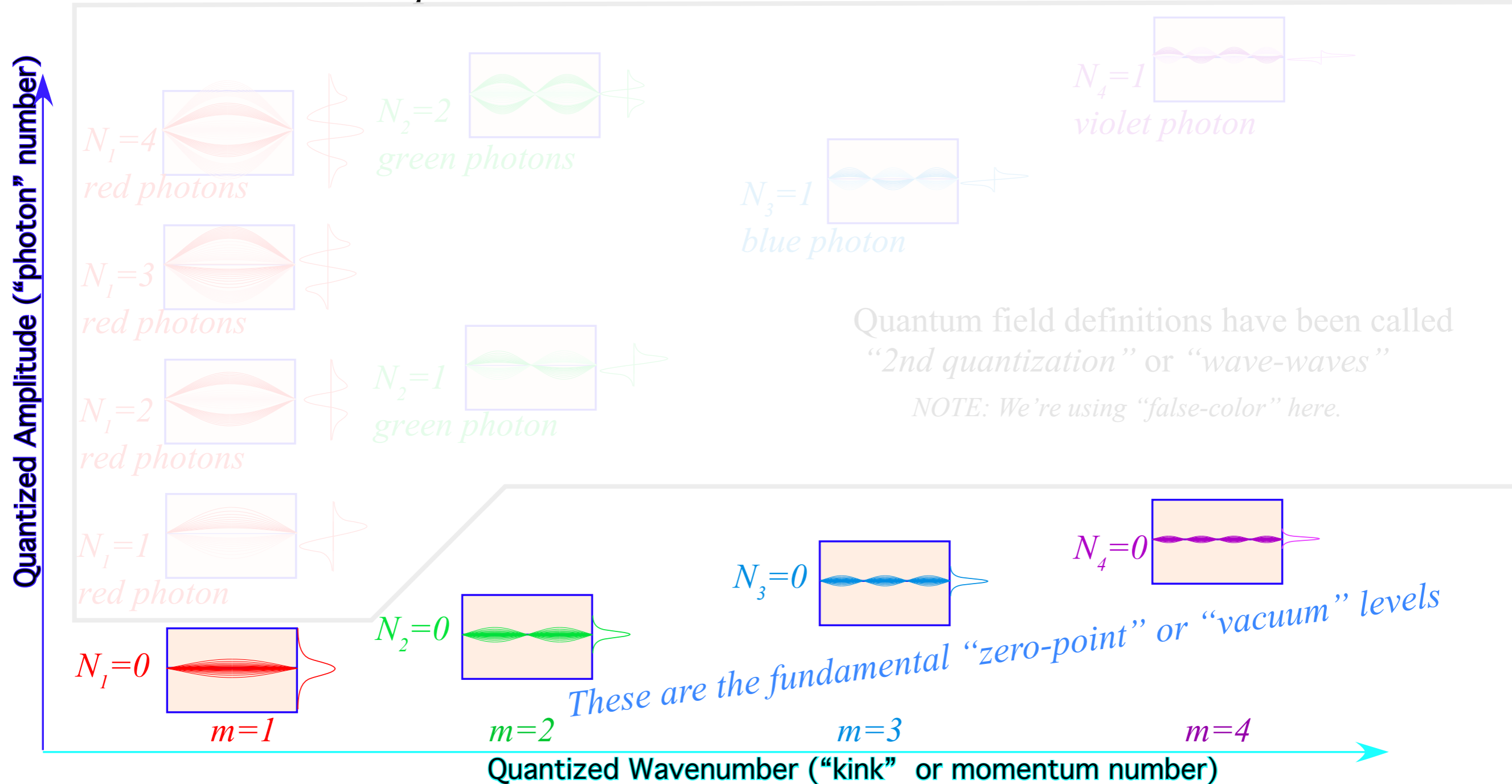
Introducing coherent states (What lasers use)

Analogy with (ω, k) wave packets

Wave coordinates need coherence

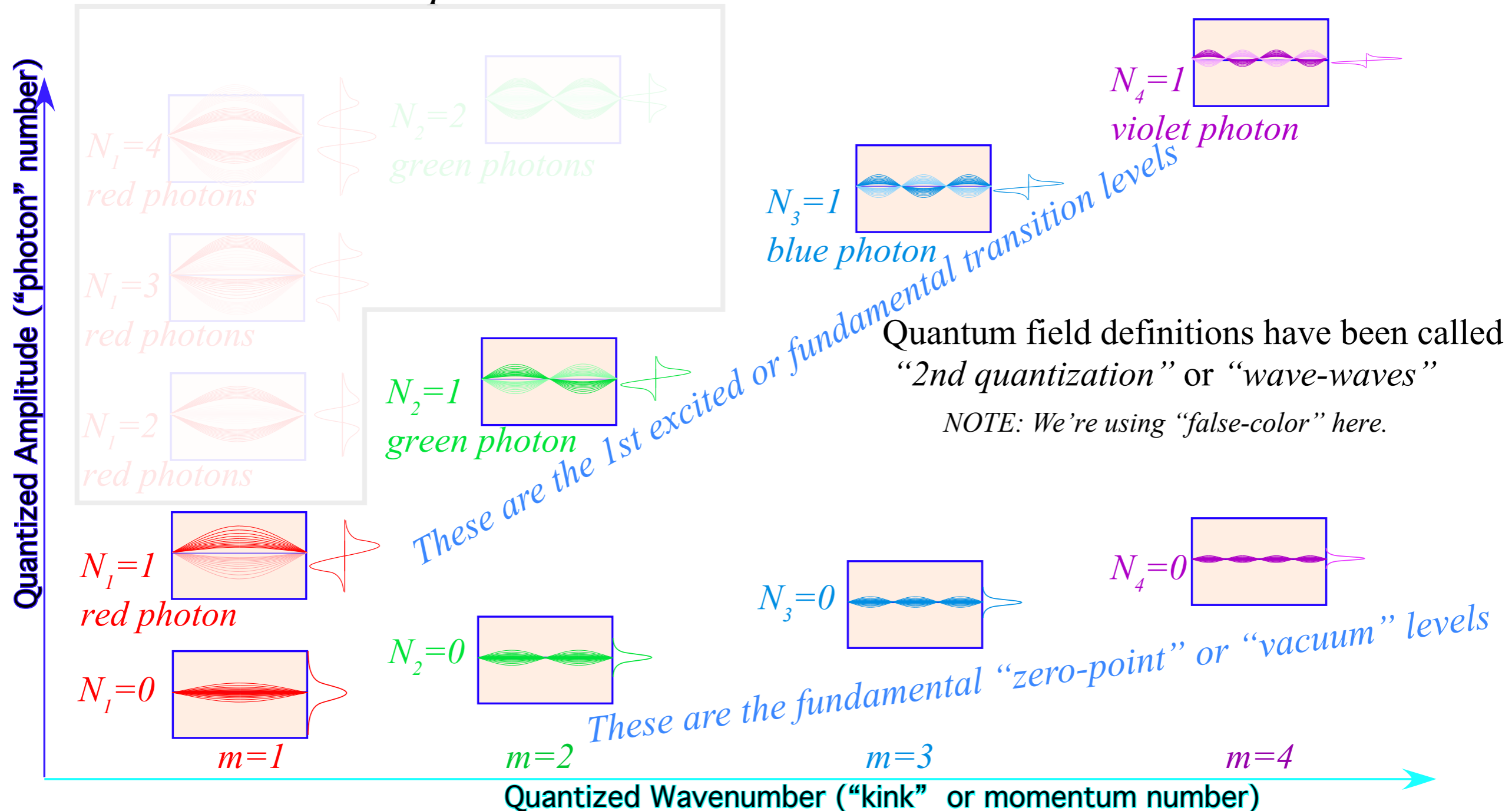
Quantized *Amplitude* Counting “photon” number

Planck’s relation $E=Nh\nu$ began as a tentative axiom to explain low-T light. Then he tried to disavow it! Einstein picked it up in his 1905 paper. Since then its use has grown enormously and continues to amaze, amuse (or bewilder) all who study it. A current view is that it represents the *quantization* of optical field *amplitude*. We picture this below as N -*photon* wave states for each box-mode of m wave kinks.



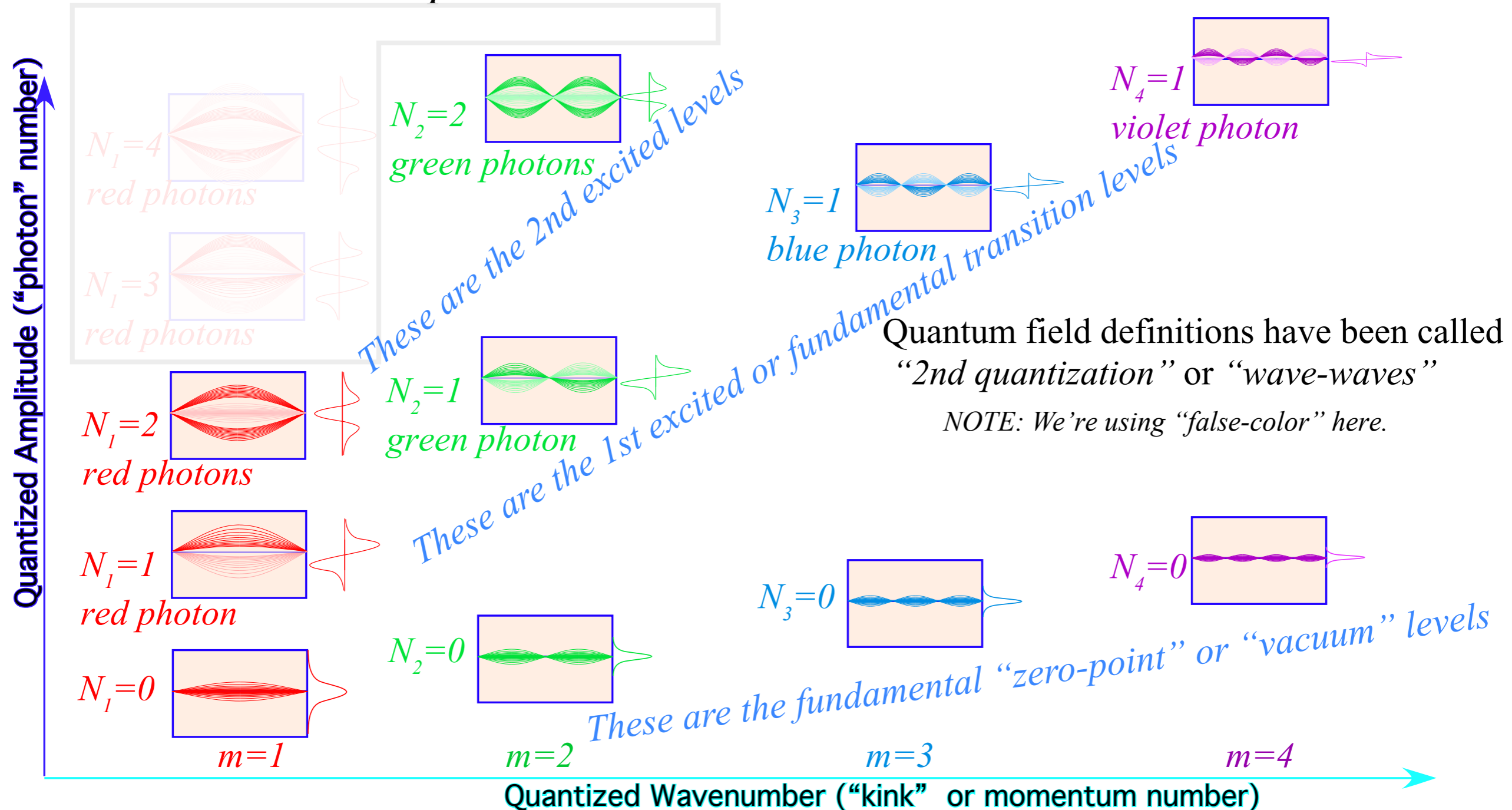
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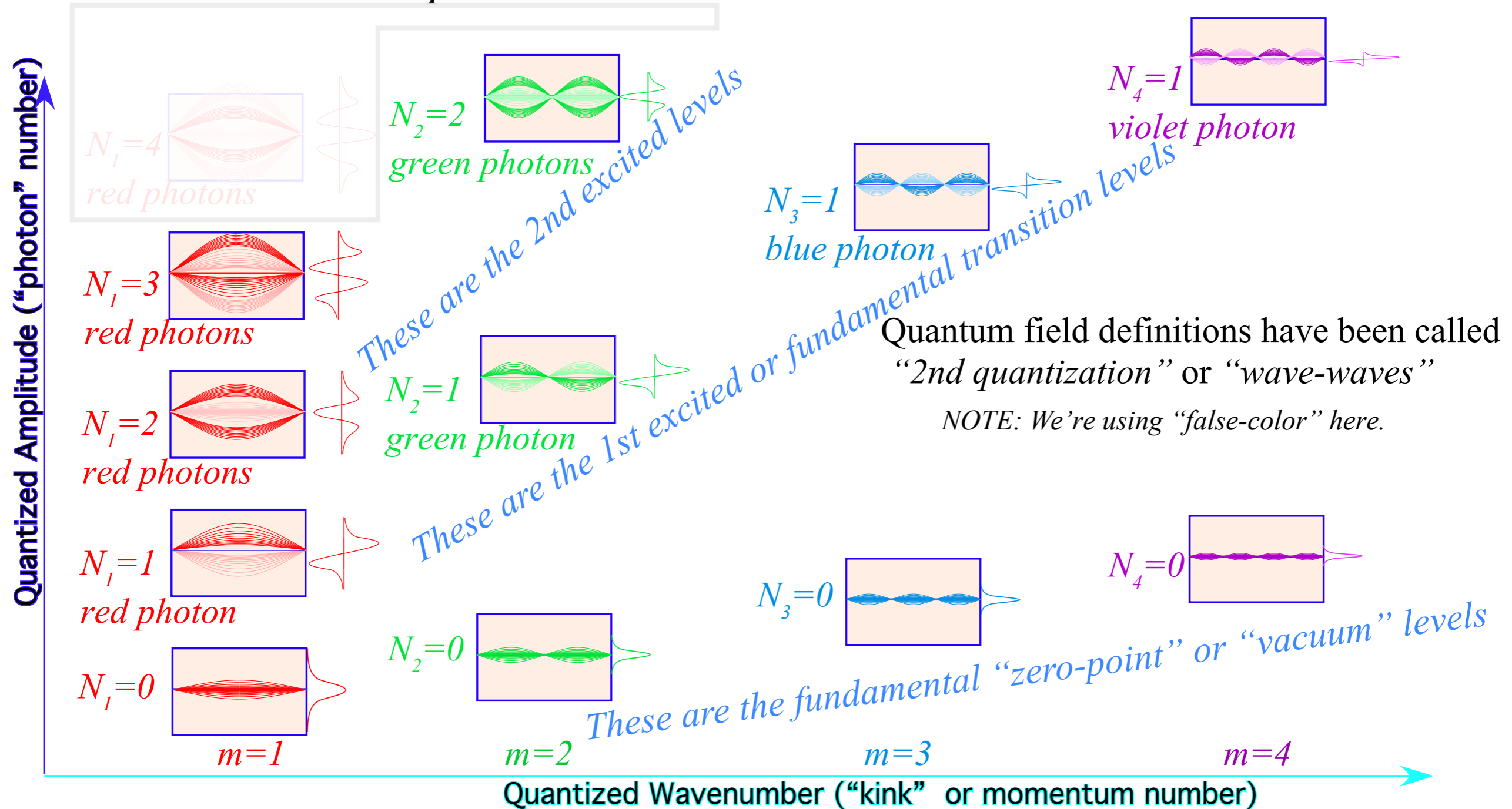
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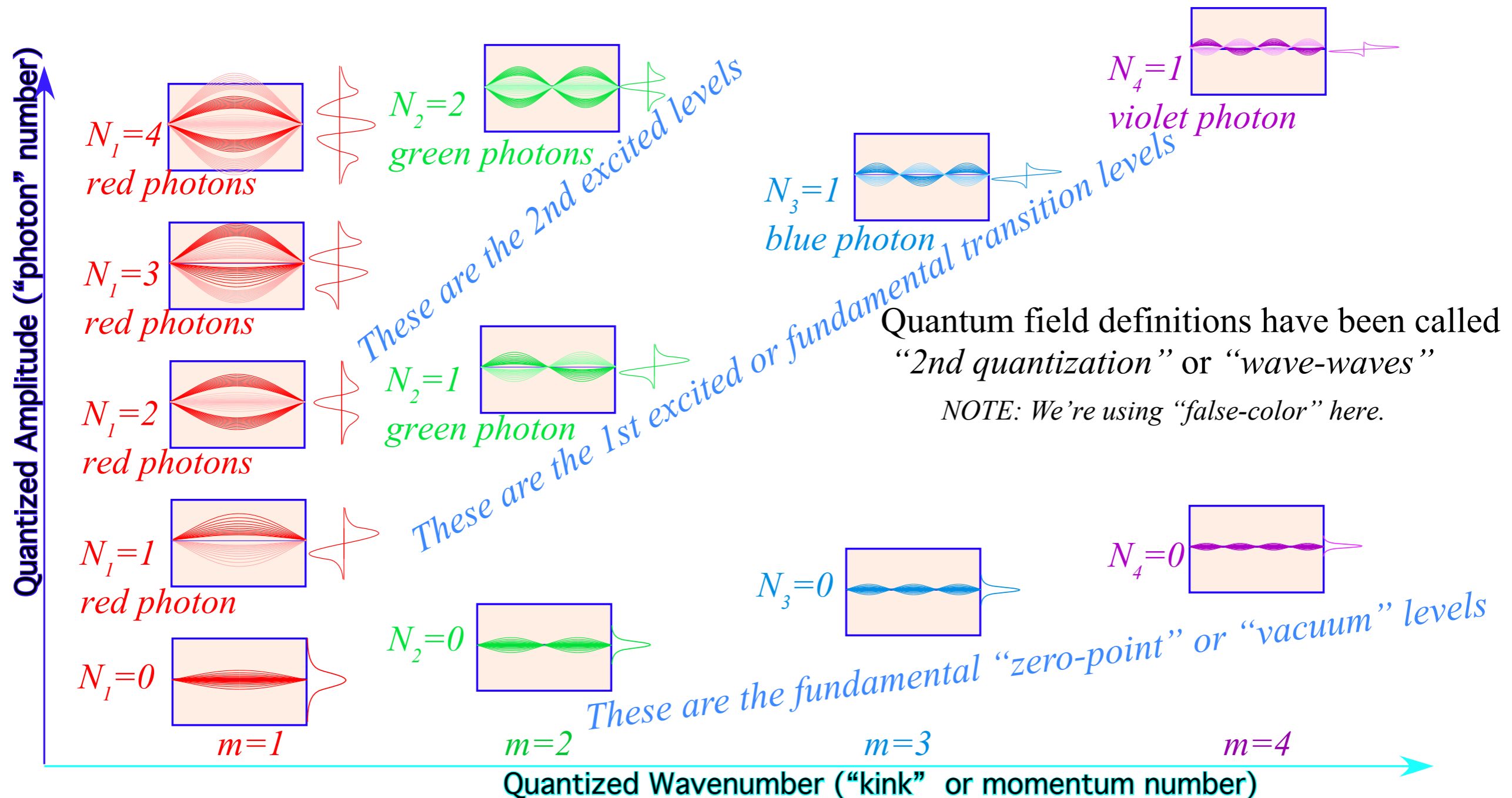
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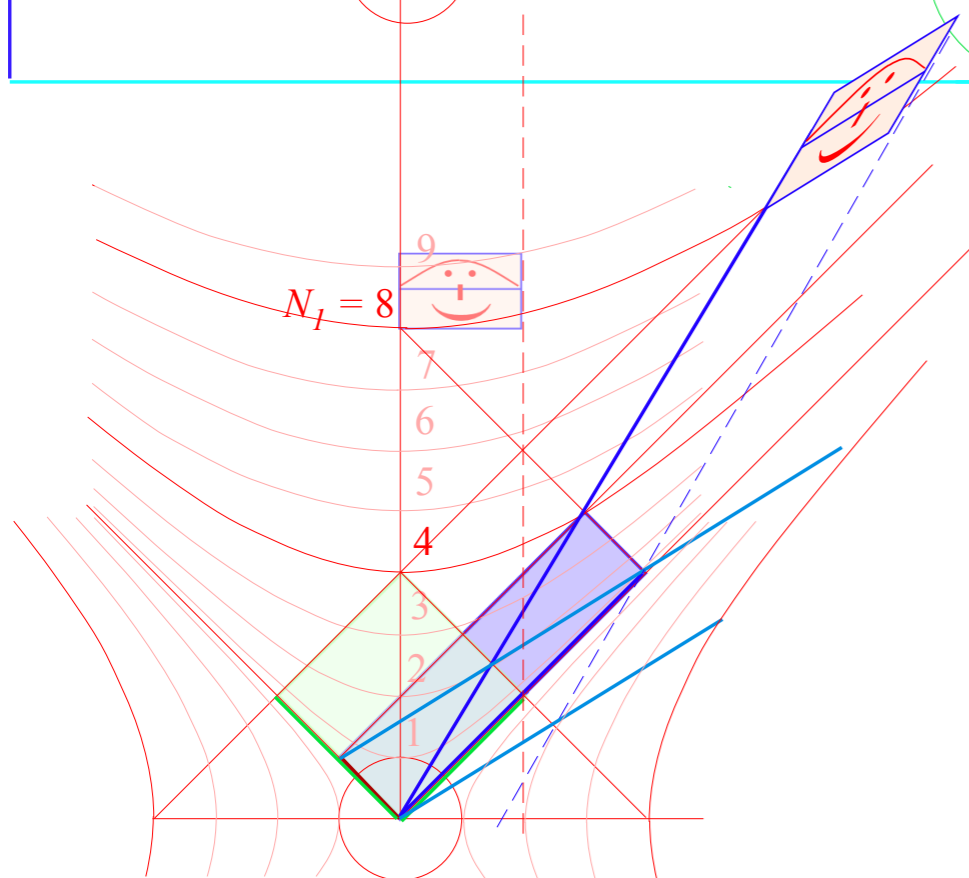
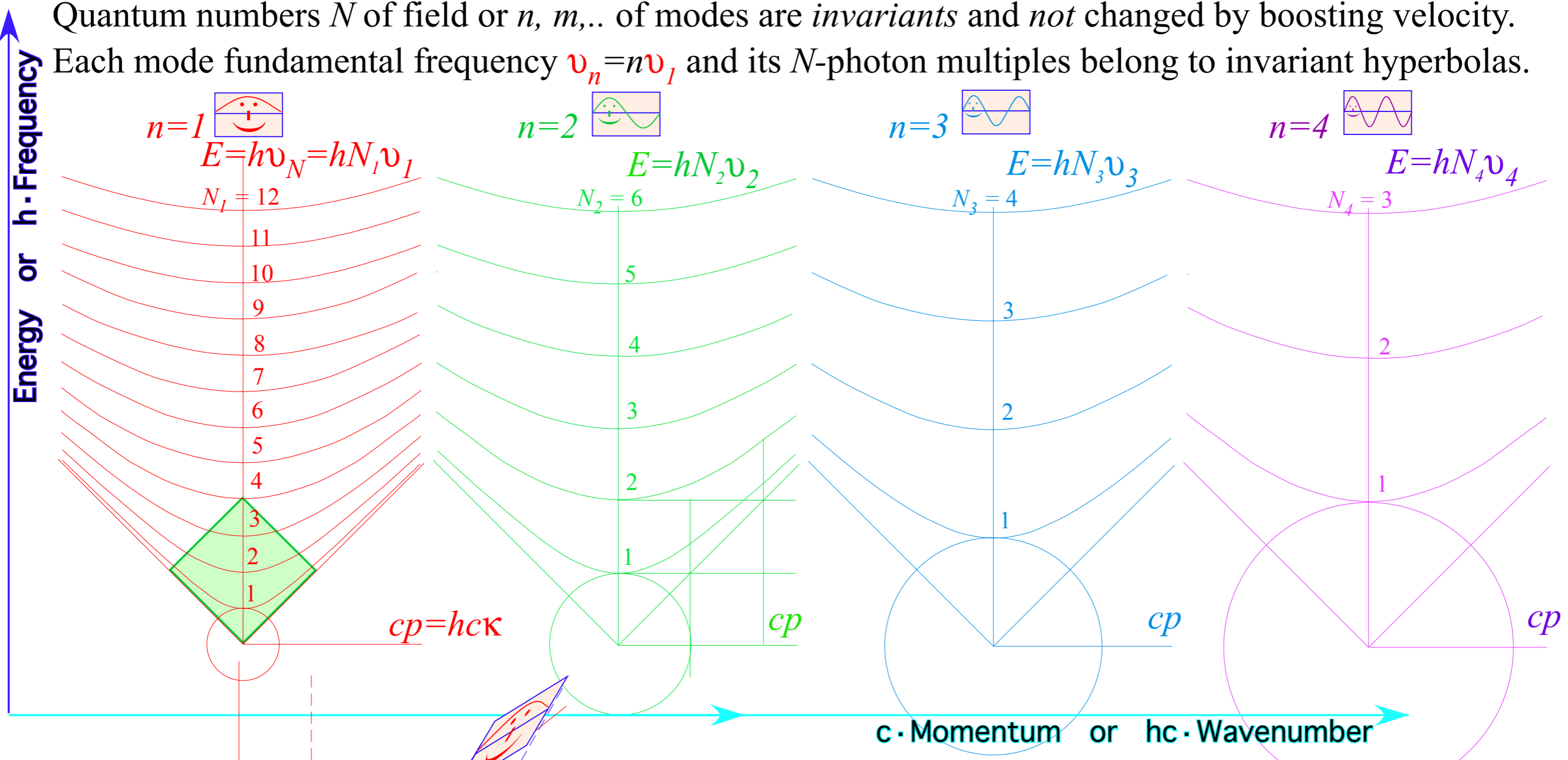


Quantized *Amplitude* Counting “photon” number

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Quantum numbers N of field or n, m, \dots of modes are *invariants* and *not* changed by boosting velocity. Each mode fundamental frequency $\nu_n = n\nu_1$ and its N -photon multiples belong to invariant hyperbolas.



Boosted observers see distorted frequencies and lengths, but will agree on the *numbers* n and N of mode *nodes* and *photons*.

This is how light waves can “fake” some of the properties of classical “things” such as *invariance* or *object permanence*.

It takes at least *TWO CW*’s to achieve such invariance. One CW is not enough and cannot have non-zero invariant N . Invariance is an *interference* effect that needs at least *two-to-tango*!

Lecture 30 ended here

2nd Quantization: Quantizing amplitudes (“photons”, “vibrons”, and “what-ever-ons”)



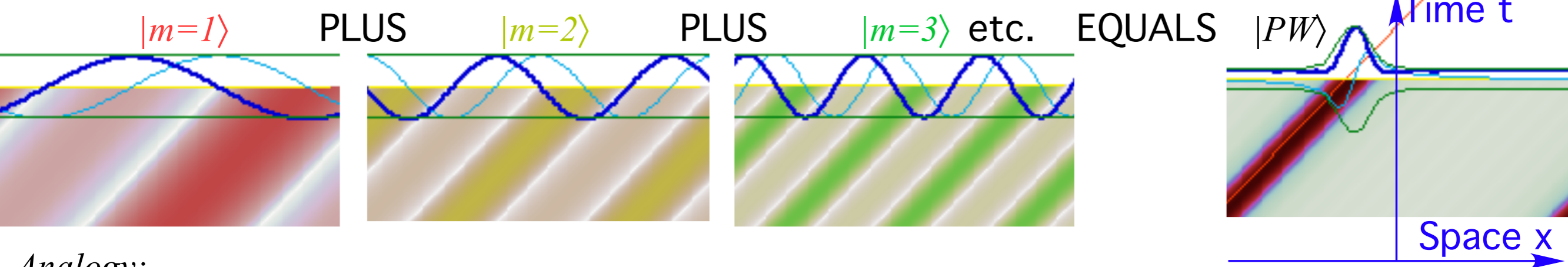
Introducing coherent states (What lasers use)

Analogy with (ω, k) wave packets

Wave coordinates need coherence

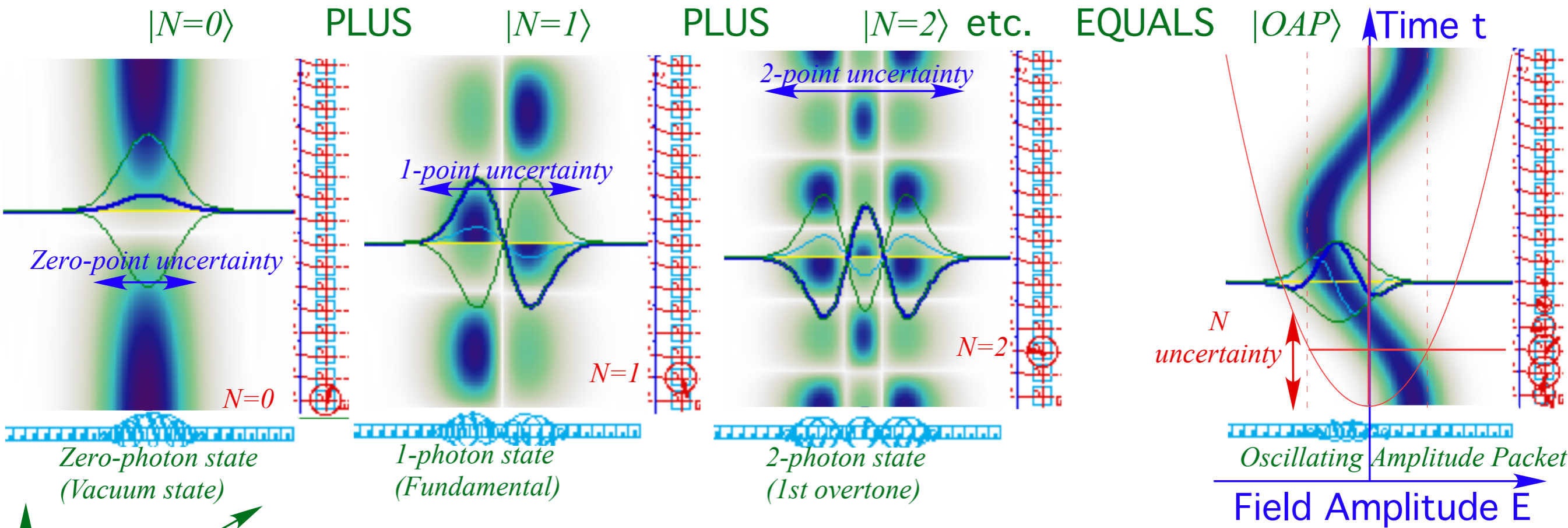
Coherent States: Oscillator Amplitude Packets analogous to Wave Packets

We saw how adding CW's (Continuous Waves $m=1,2,3\dots$) can make PW (Pulse Wave) or WP (Wave Packet) that is more like a classical "thing" with more localization in space x and time t .



Analogy:

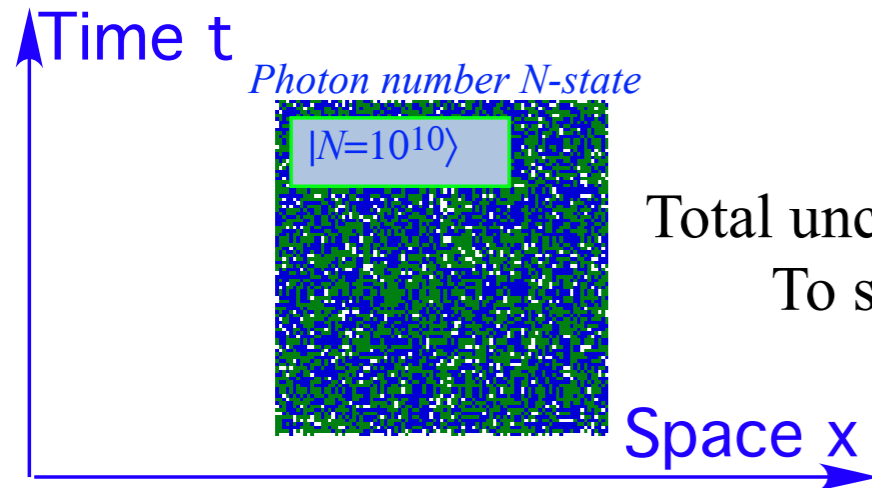
Adding photons (Quantized amplitude $N=0,1,2\dots$) can make a CS (Coherent State) or OAP (Oscillator Amplitude Packet) that is more like a classical wave oscillation with more localization in field amplitude.



Pure photon states have localized (certain) N but delocalized (uncertain) amplitude and phase.
 OAP states have delocalized (uncertain) N but more localized (certain) amplitude and phase.

Coherent States(contd.) Spacetime wave grid is impossible without coherent states

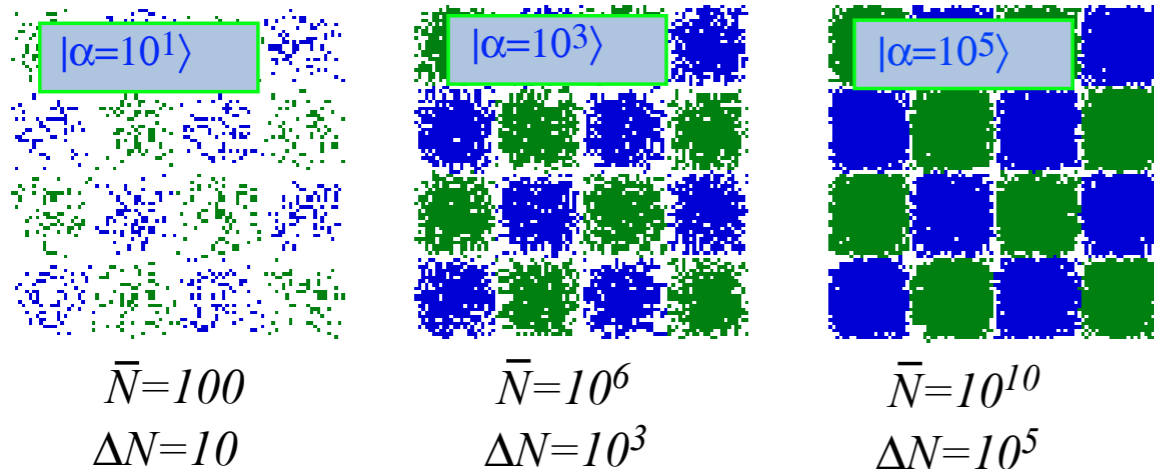
Pure photon number N -states would make useless spacetime coordinates



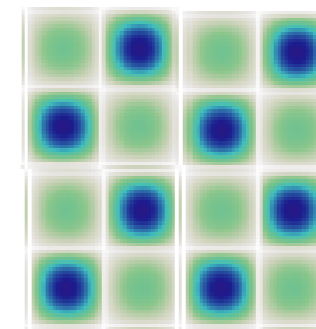
Total uncertainty of amplitude and phase makes the count pattern a wash.
To see grids *some N -uncertainty is necessary!*

Coherent- α -states are defined by continuous amplitude-packet parameter α whose square is average photon number $\bar{N}=|\alpha|^2$. Coherent-states make better spacetime coordinates for larger $\bar{N}=|\alpha|^2$.

Quantum field coherent α -states

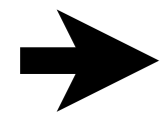


Classical limit



Coherent-state uncertainty in photon number (and mass) varies with amplitude parameter $\Delta N \sim \alpha \sim \sqrt{N}$ so a coherent state with $\bar{N}=|\alpha|^2=10^6$ only has a 1-in-1000 uncertainty $\Delta N \sim \alpha \sim \sqrt{N}=1000$.

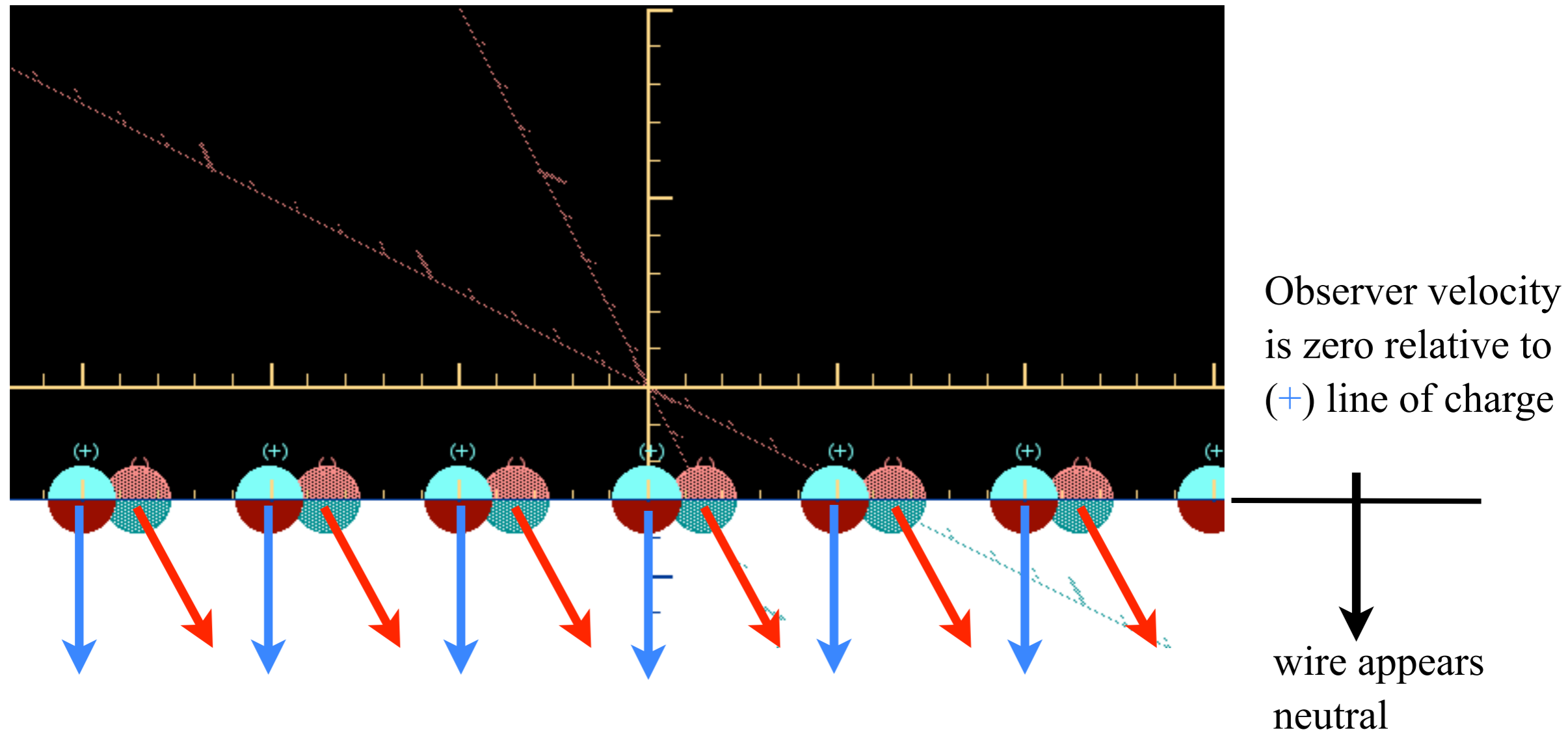
Relativistic effects on charge, current, and Maxwell Fields



*Current density changes by Lorentz **asynchrony***

Magnetic B-field is relativistic $\sinh\rho$ 1st order-effect

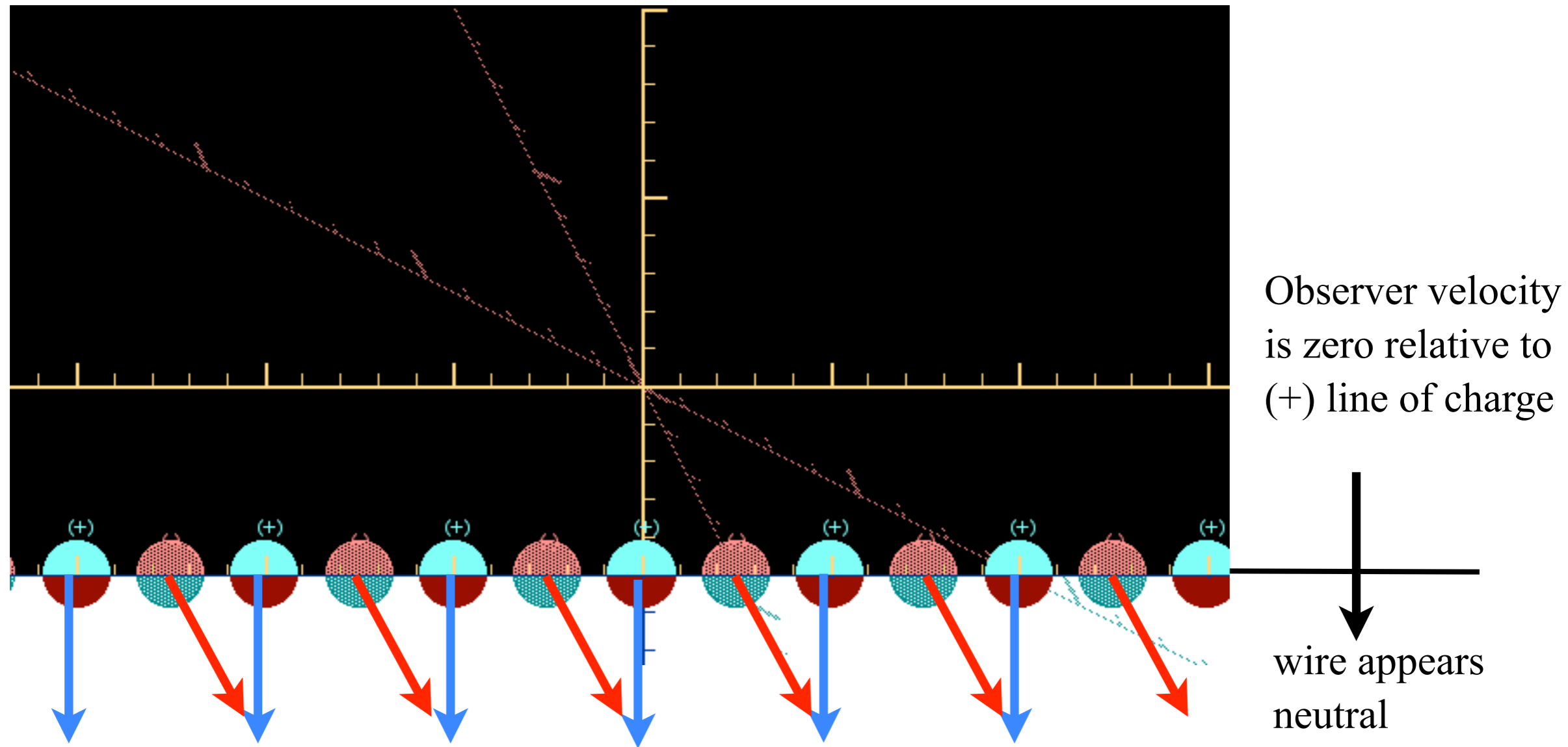
Relativistic effects on charge, current, and Maxwell Fields



(+) Charge fixed (-) Charge moving to right (*Negative current density*)

(+) Charge density is Equal to the (-) Charge density

Relativistic effects on charge, current, and Maxwell Fields



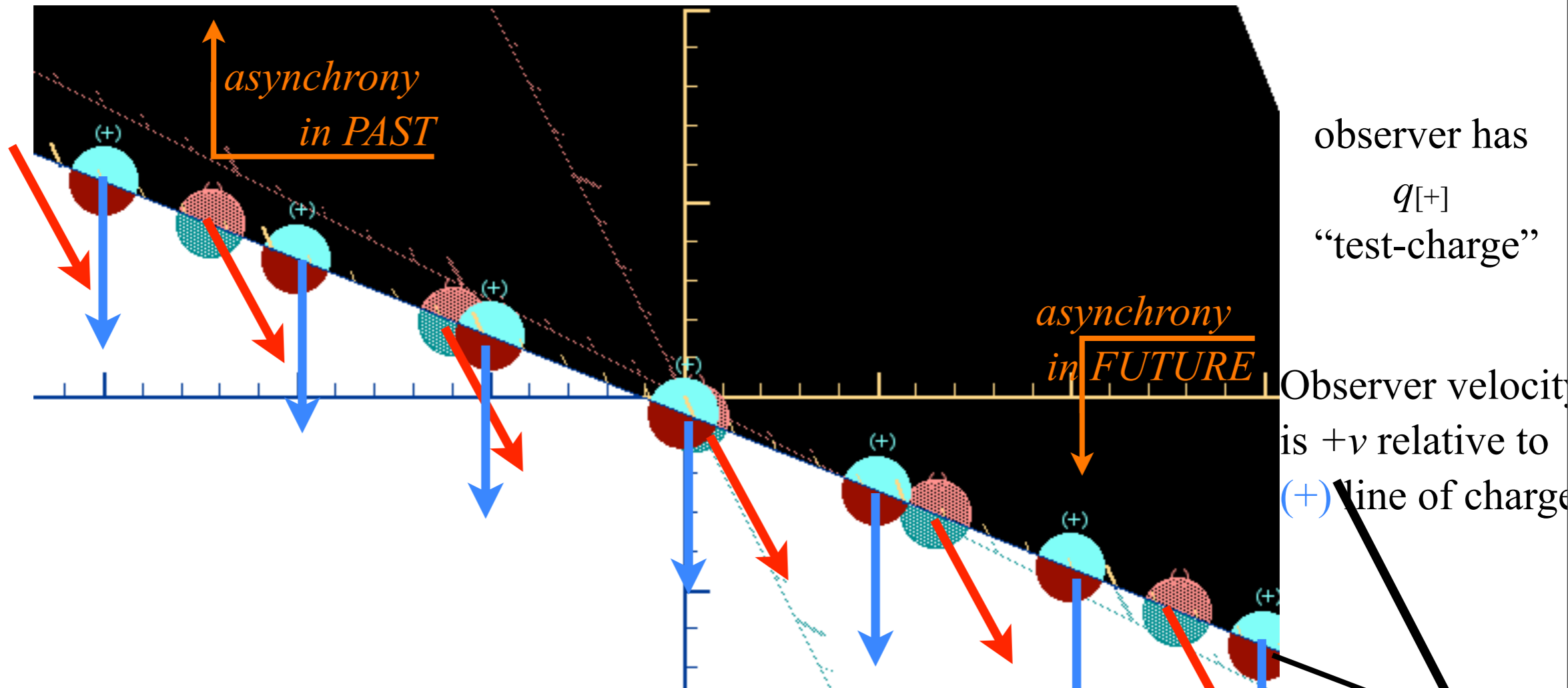
- (+) Charge fixed (-) Charge moving to right (*Negative current density* $\vec{j}(x,t)$)
- (+) Charge density is Equal to the (-) Charge density (*Zero* $\rho(x,t)=0$)

Relativistic effects on charge, current, and Maxwell Fields

Current density changes by Lorentz *asynchrony*

Asynchrony due to off-diagonal $\sinh \rho$ (a 1st-order effect)

in Lorentz transform:
$$\begin{pmatrix} \cosh \rho & \sinh \rho \\ \sinh \rho & \cosh \rho \end{pmatrix} \sim \begin{pmatrix} 1 & v/c \\ v/c & 1 \end{pmatrix}$$



(+) Charge fixed (-) Charge moving to right (Negative current density $\vec{j}(x,t)$)
 (+) Charge density is Greater than (-) Charge density (Positive $\rho(x,t) > 0$)

observer has $q_{[+]}$ "test-charge"

Observer velocity is $+v$ relative to (+) line of charge

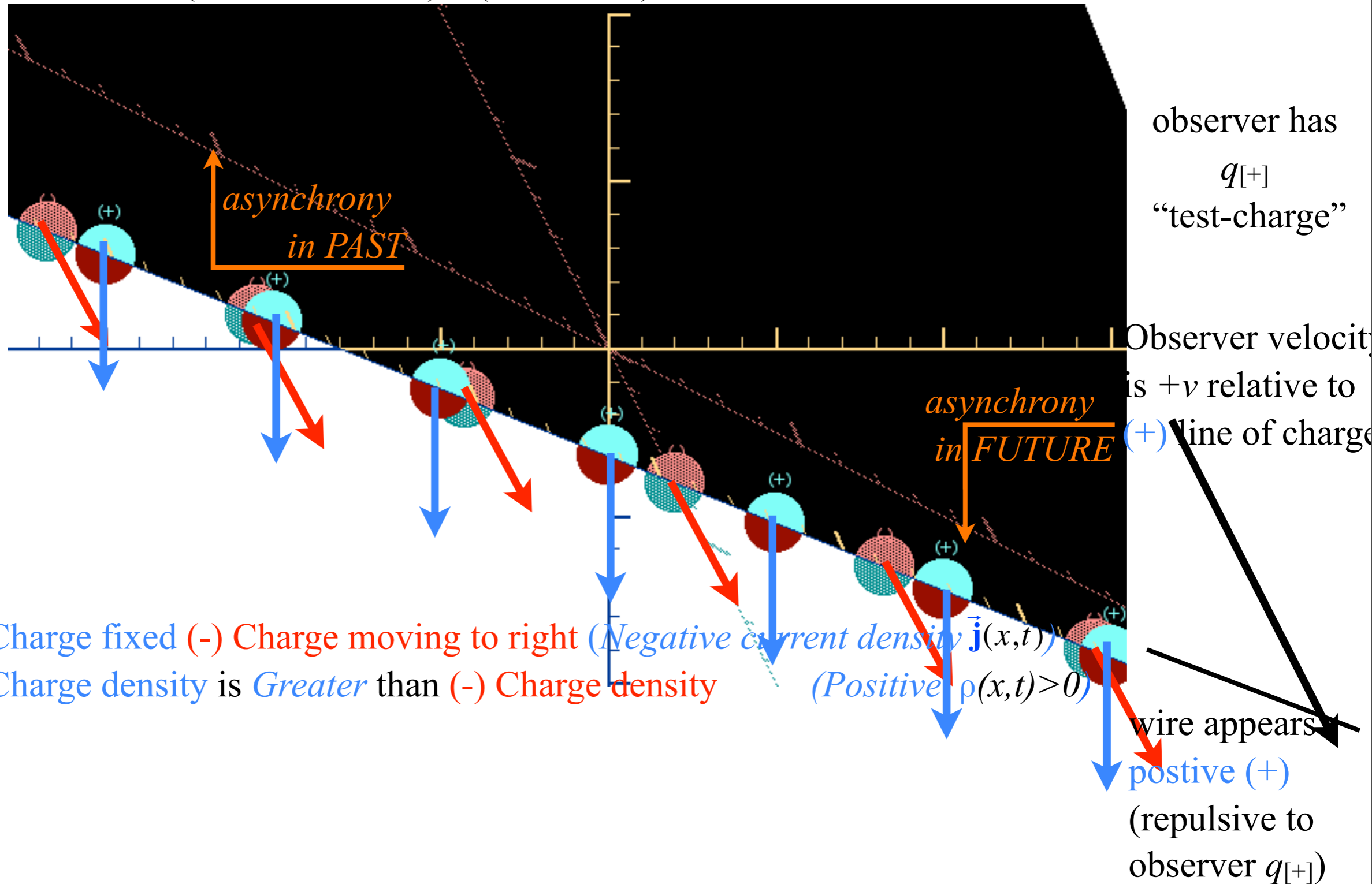
wire appears positive (+) (repulsive to observer $q_{[+]}$)

Relativistic effects on charge, current, and Maxwell Fields

Current density changes by Lorentz *asynchrony*

Asynchrony due to off-diagonal $\sinh \rho$ (a 1st-order effect)

in Lorentz transform:
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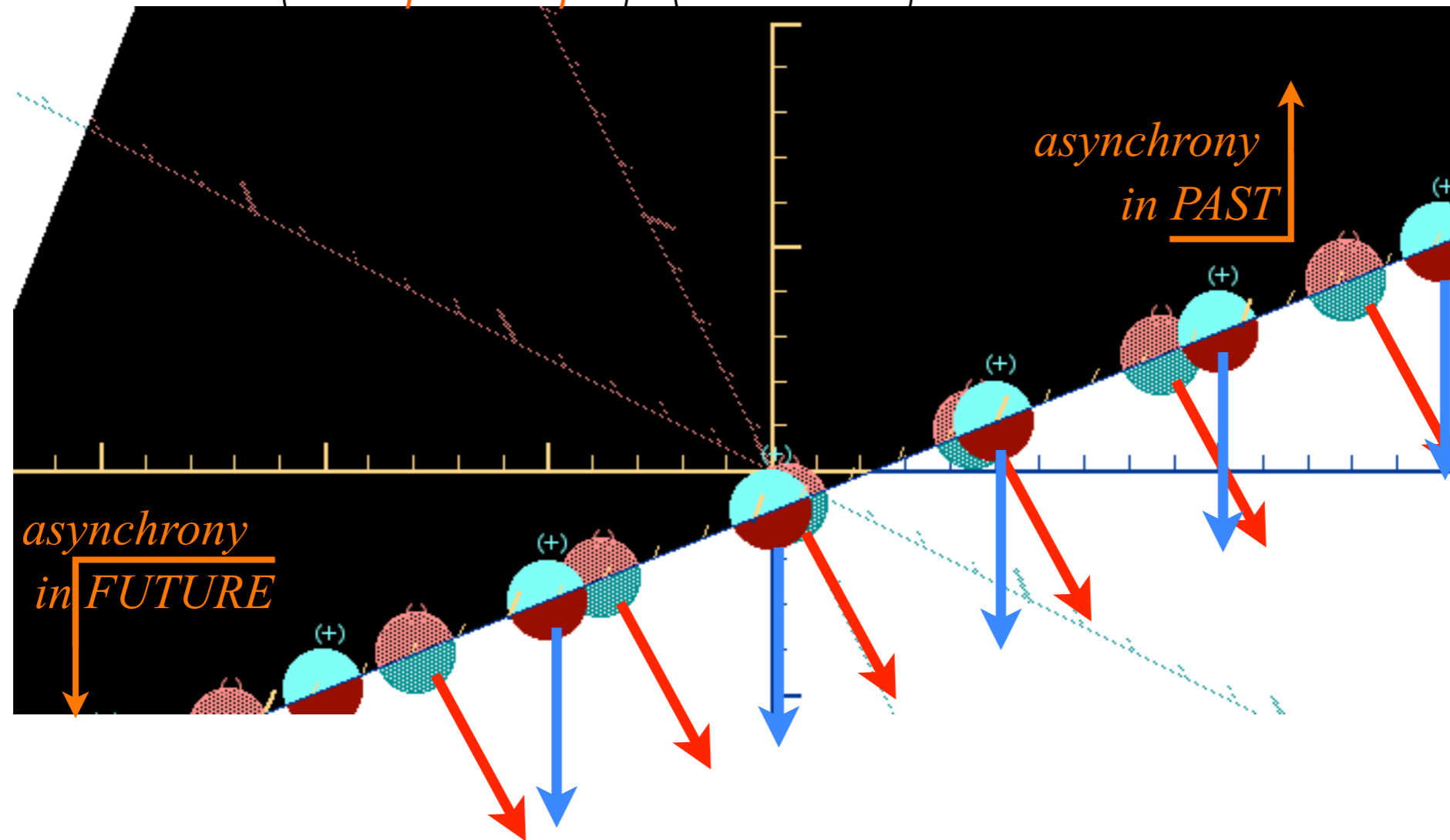


Relativistic effects on charge, current, and Maxwell Fields

Current density changes by Lorentz *asynchrony*

Asynchrony due to off-diagonal $\sinh \rho$ (a 1st-order effect)

in Lorentz transform:
$$\begin{pmatrix} \cosh \rho & \sinh \rho \\ \sinh \rho & \cosh \rho \end{pmatrix} \sim \begin{pmatrix} 1 & v/c \\ v/c & 1 \end{pmatrix}$$



observer has
 $q_{[+]}$
“test-charge”

Observer velocity
is $-v$ relative to
(+) line of charge

wire appears
negative (-)
(attractive to
observer $q_{[+]}$)

(+) Charge fixed (-) Charge moving to right (Negative current density $\vec{j}(x,t)$)

(+) Charge density is *Less* than (-) Charge density (Negative $\rho(x,t) < 0$)

Relativistic effects on charge, current, and Maxwell Fields

Current density changes by Lorentz *asynchrony*

Asynchrony due to off-diagonal $\sinh \rho$ (a 1st-order effect)

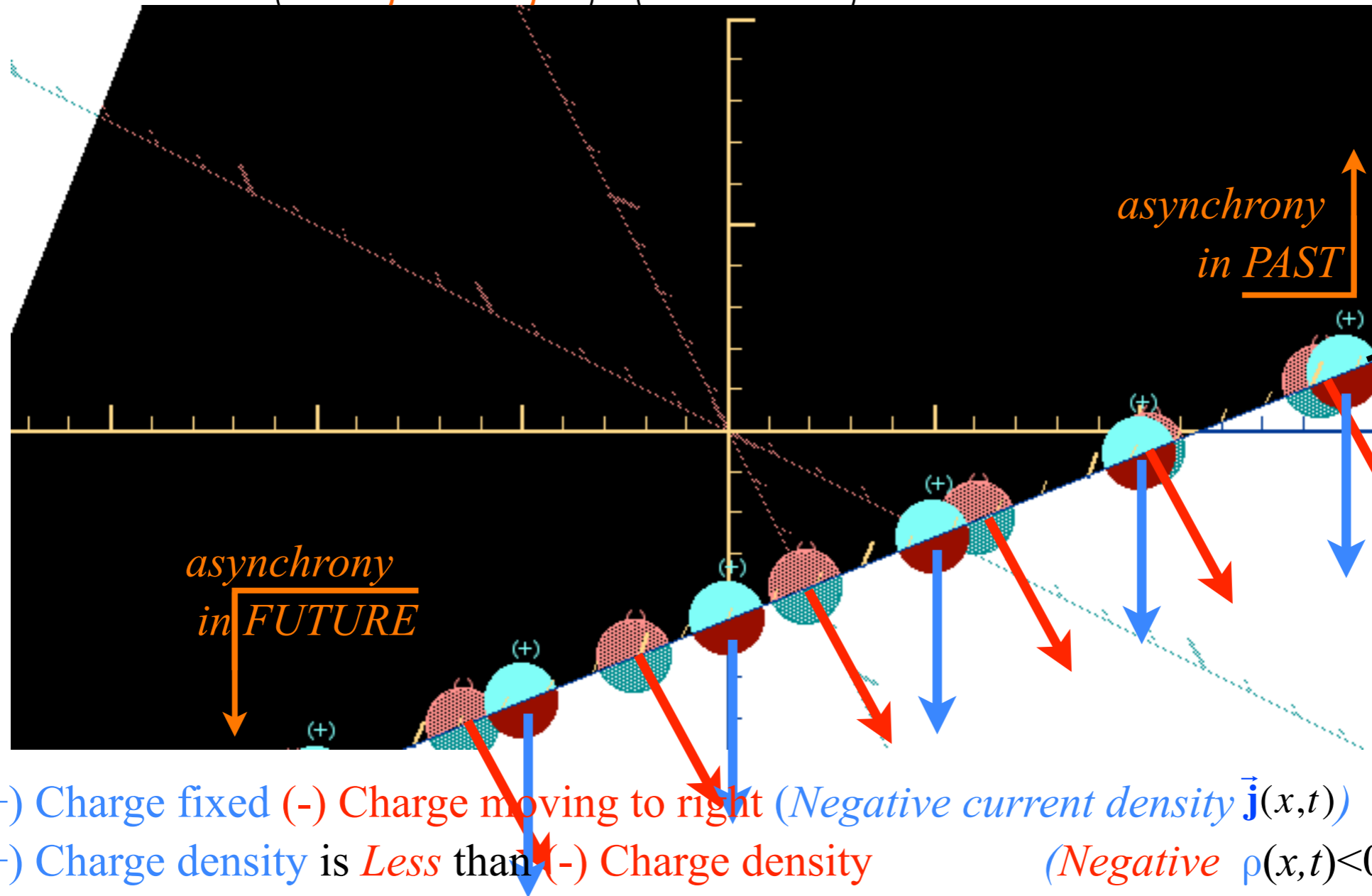
in Lorentz transform:
$$\begin{pmatrix} \cosh \rho & \sinh \rho \\ \sinh \rho & \cosh \rho \end{pmatrix} \sim \begin{pmatrix} 1 & v/c \\ v/c & 1 \end{pmatrix}$$

observer has

$q_{[+]}$

“test-charge”

Observer velocity is $-v$ relative to $(+)$ line of charge



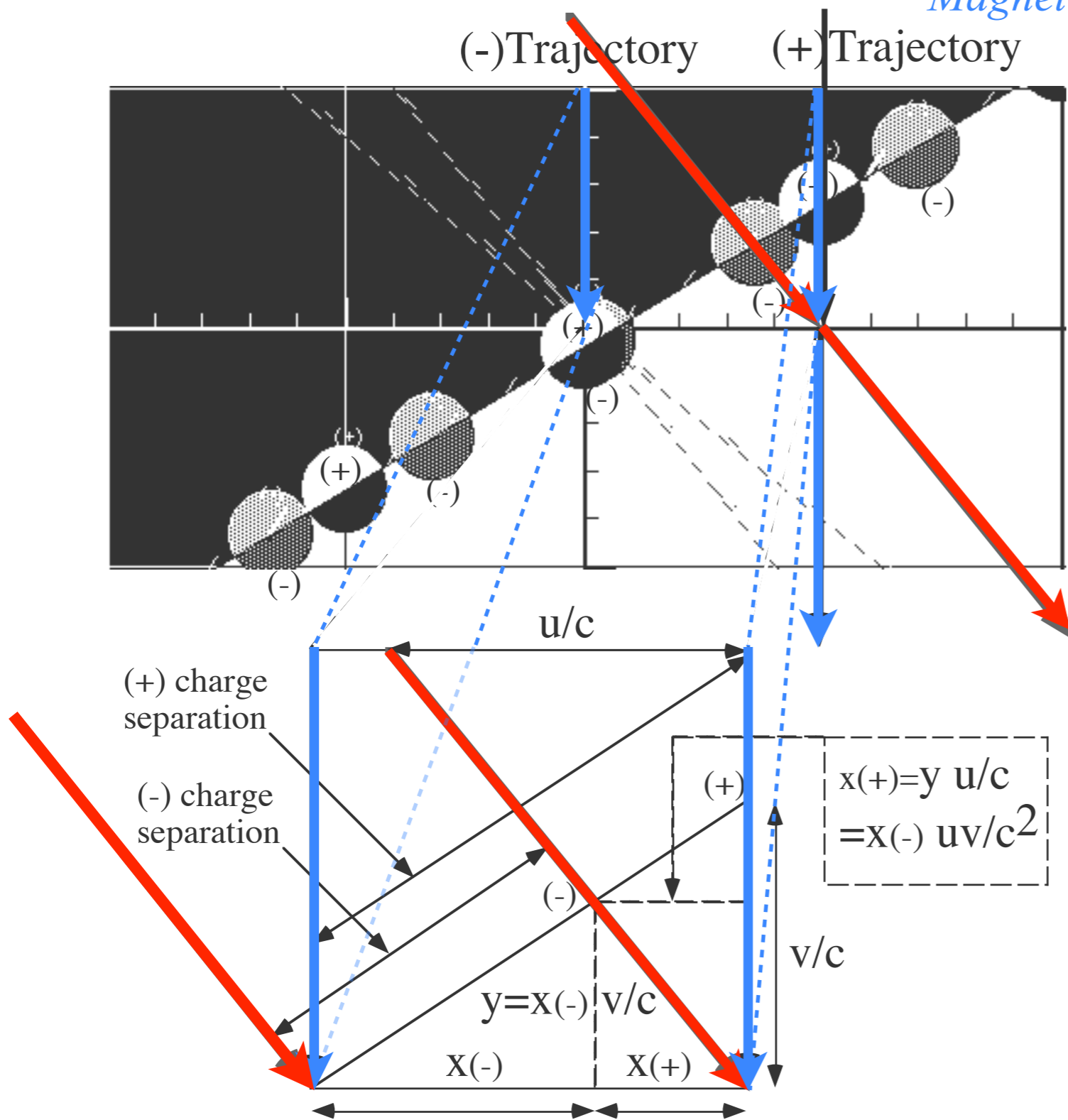
(+) Charge fixed (-) Charge moving to right (Negative current density $\vec{j}(x,t)$)

(+) Charge density is *Less* than (-) Charge density (Negative $\rho(x,t) < 0$)

Relativistic effects on charge, current, and Maxwell Fields

*Current density changes by Lorentz **asynchrony***

 *Magnetic B-field is relativistic $\sinh\rho$ 1st order-effect*



$$\frac{\rho(-)}{\rho(+)} = \frac{(+)\text{ charge separation}}{(-)\text{ charge separation}} = \frac{x(+)+x(-)}{x(-)}$$

$$\frac{\rho(-)}{\rho(+)} = \frac{x(+)}{x(-)} + 1 = \frac{uv}{c^2} + 1$$

$$\rho(+)-\rho(-) = \rho(+)\left(1 - \frac{\rho(-)}{\rho(+)}\right) = -\frac{uv}{c^2}\rho(+)$$

Unit square: $(u/c) / 1 = x(+)/y$
 $(v/c) / 1 = y/x(-)$

The electric force field \mathbf{E} of a charged line varies inversely with radius. The Gauss formula for force in mks units :

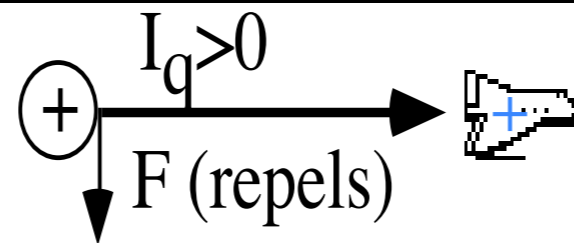
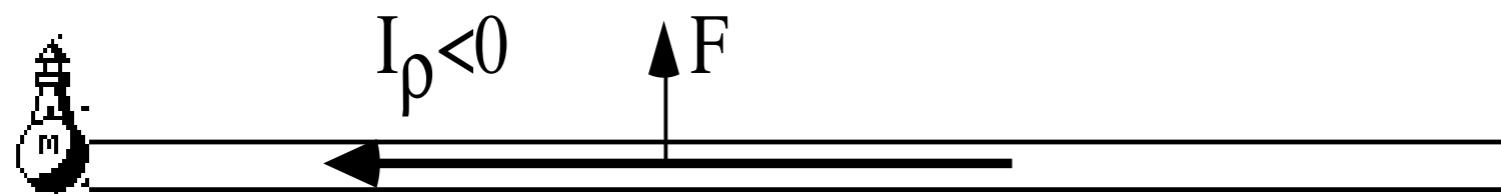
$$F = qE = q \left[\frac{1}{4\pi\epsilon_0} \frac{2\rho}{r} \right], \quad \text{where: } \frac{1}{4\pi\epsilon_0} = 9 \times 10^9 \frac{N \cdot m^2}{\text{Coul.}}$$

$$F = qE = q \left[\frac{1}{4\pi\epsilon_0} \frac{2}{r} \left(-\frac{uv}{c^2} \rho(+)\right) \right] = -\frac{2qv\rho(+)}{4\pi\epsilon_0 c^2 r} = -2 \times 10^{-7} \frac{I_q I_\rho}{r}$$

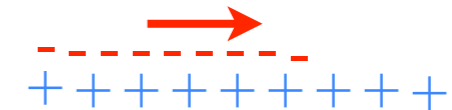
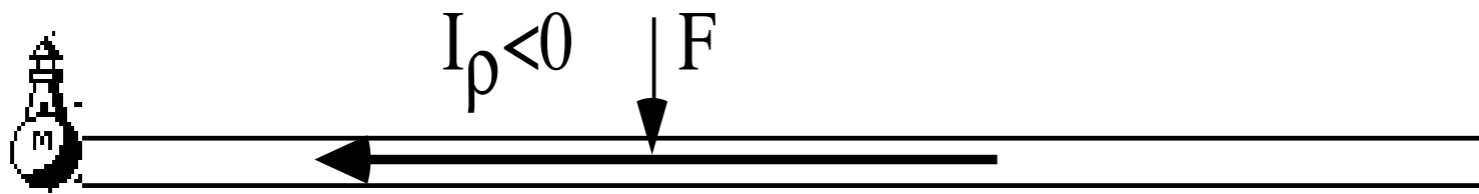
$$1/4\pi\epsilon_0 = 9 \cdot 10^9$$

$$c^2 = 9 \cdot 10^{16}$$

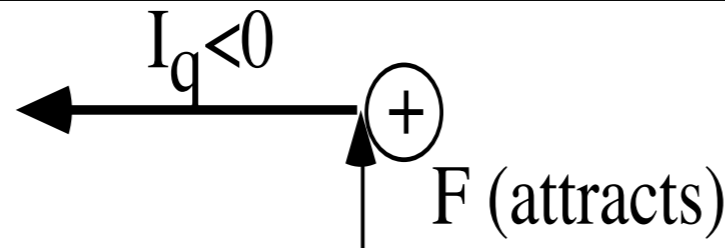
$$1/(4\pi\epsilon_0 c^2) = 10^{-7}$$



I see excess (+) charge up there. Yuk!



I see excess (-) charge up there. Yum!



The electric force field \mathbf{E} of a charged line varies inversely with radius. The Gauss formula for force in mks units :

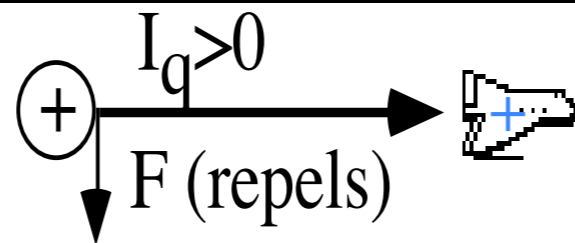
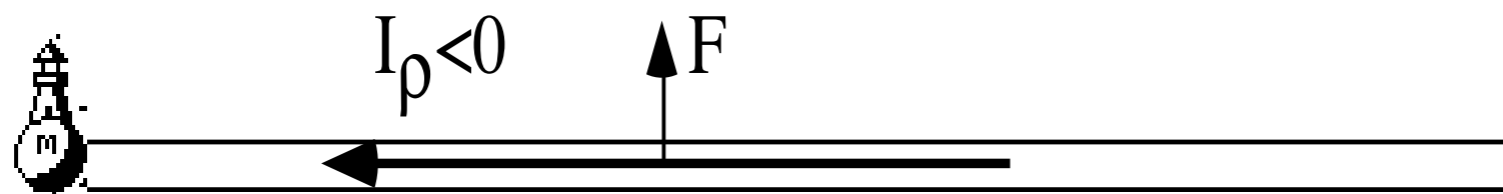
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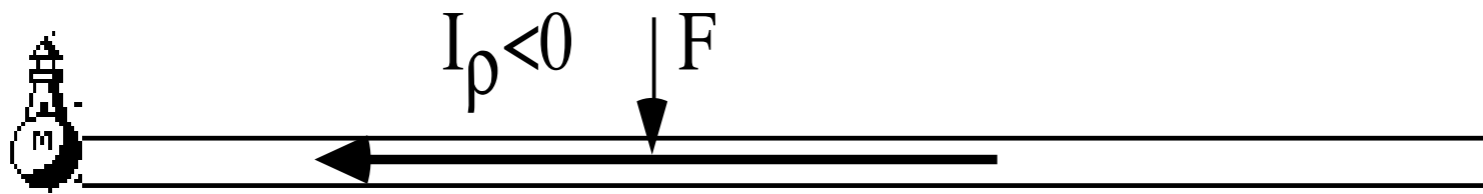
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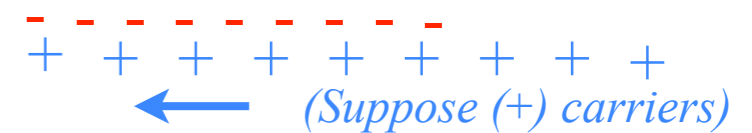
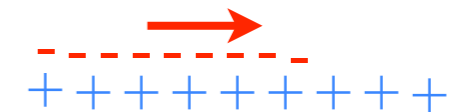
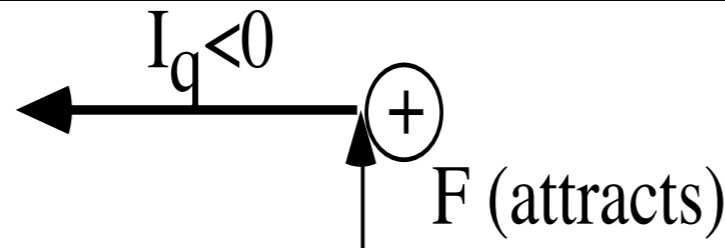
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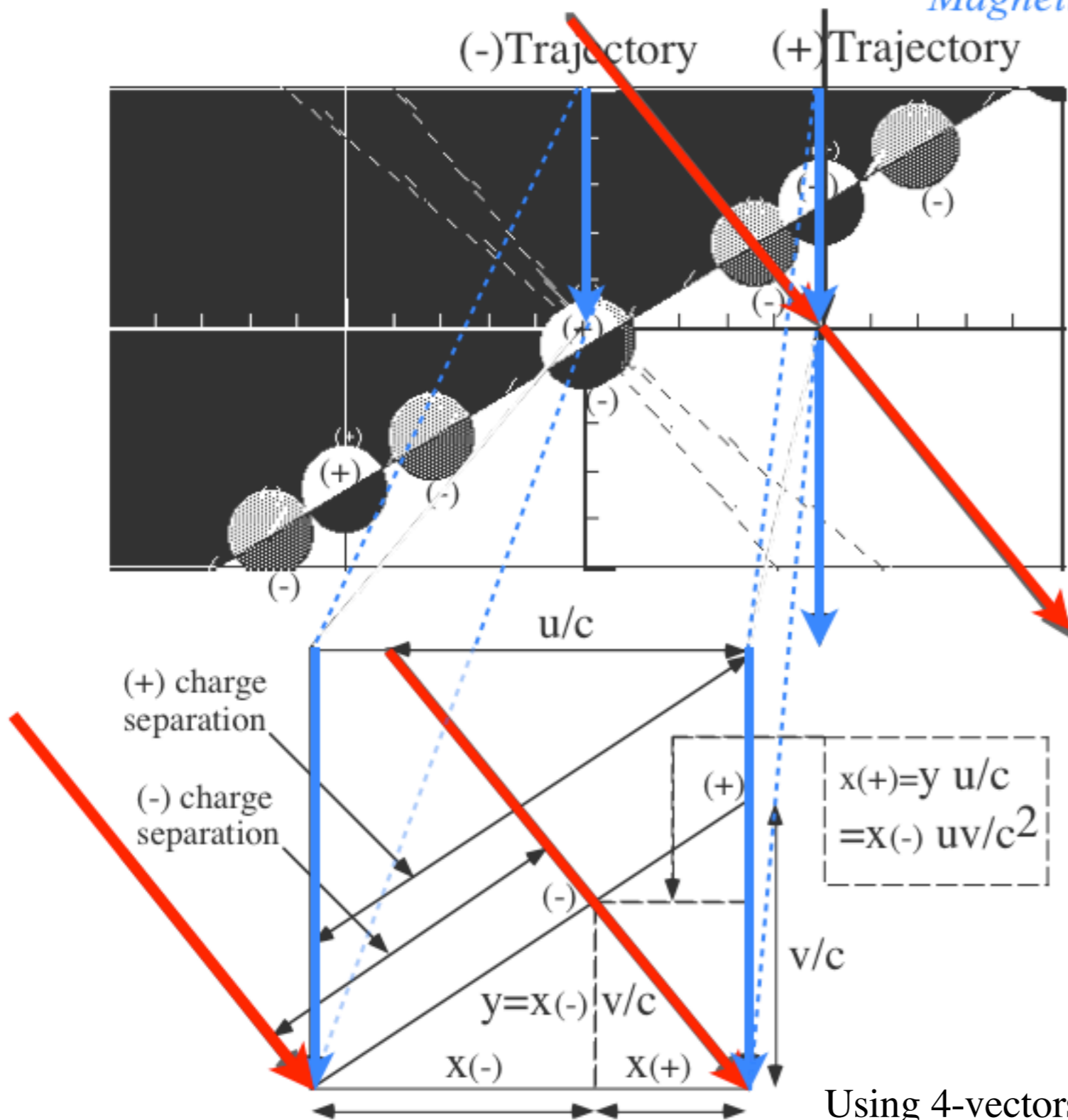


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Unit square: $(u/c) / 1 = x(+)/y$
 $(v/c) / 1 = y/x(-)$

Using 4-vectors to EL Transform (charge-current)=($c\rho, \mathbf{j}$)

$$\begin{pmatrix} c\rho' \\ j_{x'} \\ j_{y'} \\ j_{z'} \end{pmatrix} = \begin{pmatrix} \cosh \rho & \sinh \rho & \cdot & \cdot \\ \sinh \rho & \cosh \rho & \cdot & \cdot \\ \cdot & \cdot & 1 & \cdot \\ \cdot & \cdot & \cdot & 1 \end{pmatrix} \begin{pmatrix} c\rho \\ j_x \\ j_y \\ j_z \end{pmatrix}$$