

Classical, semi-classical, and quantum dynamics of uni-axial and multi-axial floppy rotors

*William G. Harter and Horace Crogman,
Department of Physics, University of Arkansas
Fayetteville, AR 72701*

(1st Half)

*Ways to geometrically visualize **single-rotor** quantum states and dynamics*

(2nd Half)

*Ways to begin visualizing **compound-rotor** quantum states and dynamics*

(Next talk)

*Ways to begin computing **compound-rotor** states...*

Simple Rigid Rotor Hamiltonian...

$$\mathbf{H} = A\mathbf{J}_x^2 + B\mathbf{J}_y^2 + C\mathbf{J}_z^2 + \dots$$

...and its *multi-pole expansion*...

$$\left(\frac{A+B+C}{3}\right)(\mathbf{J}_x^2 + \mathbf{J}_y^2 + \mathbf{J}_z^2) + \left(\frac{2C-A-B}{6}\right)(2\mathbf{J}_z^2 - \mathbf{J}_x^2 - \mathbf{J}_y^2) + \left(\frac{A-B}{2}\right)(\mathbf{J}_x^2 - \mathbf{J}_y^2)$$

Spherical Top
(A=B=C)
 $\mathbf{H} = B\mathbf{J}^2$
 $\mathbf{T}_0^{(0)} = \mathbf{J}^2$

Symmetric Top
(A=B≠C)
 $\mathbf{H} = B\mathbf{J}^2 + (C-B)(2/3)\mathbf{T}_0^{(2)}$
 $2\mathbf{T}_0^{(2)}$

Asymmetric Top
(A≠B≠C)
 $\sqrt{\frac{2}{3}}(\mathbf{T}_2^{(2)} - \mathbf{T}_{-2}^{(2)})$

$$\mathbf{H} = B\mathbf{J}^2 + (2C-A-B)/3 \mathbf{T}_0^{(2)} + (A-B)/\sqrt{6} (\mathbf{T}_2^{(2)} - \mathbf{T}_{-2}^{(2)})$$

Some Approaches for Treating Rotor Hamiltonians

(Q) Quantum: Find H-matrix rep and diagonalize by computer

$$\left\langle \begin{matrix} J' \\ K' \end{matrix} \left| \mathbf{T}_0^{(0)} \right| \begin{matrix} J \\ K \end{matrix} \right\rangle = \delta_{K'K} J(J+1)$$

(But, is there life after diagonalization?!?)

$$\left\langle \begin{matrix} J' \\ K' \end{matrix} \left| \mathbf{T}_0^{(2)} \right| \begin{matrix} J \\ K \end{matrix} \right\rangle = C_{0KK'}^{2J J'} \langle J' \| 2 \| J \rangle$$

$$\left\langle \begin{matrix} J' \\ K' \end{matrix} \left| \mathbf{T}_q^{(2)} \right| \begin{matrix} J \\ K \end{matrix} \right\rangle = C_{qKK'}^{2J J'} \langle J' \| 2 \| J \rangle$$

(P) Classical RES Plot: Rotational Energy (RE) surfaces and/or H-phase paths

$$\left\langle \mathbf{T}_0^{(0)} \right\rangle = c Y_0^0 = J(J+1)$$

(tensor operator \mathbf{T}_q^k is replaced by spherical harmonic $Y_q^k[\beta, \gamma]$)

$$\left\langle 2\mathbf{T}_0^{(2)} \right\rangle = c Y_0^2 = J(J+1)(3 \cos^2 \beta - 1)$$

$$\sqrt{\frac{2}{3}} \left\langle \left(\mathbf{T}_2^{(2)} - \mathbf{T}_{-2}^{(2)} \right) \right\rangle = c \left(Y_2^2 - Y_{-2}^2 \right) = J(J+1) \left(\sin^2 \beta \cos 2\gamma \right)$$

(S) Semiclassical: Some of both

Some Approaches for Treating Rotor Hamiltonians (contd)

(P) Classical RE Plot: Rotational Energy (RE) surfaces and/or H-phase paths

$$\langle \mathbf{T}_0^{(0)} \rangle = cY_0^0 = J(J+1)$$

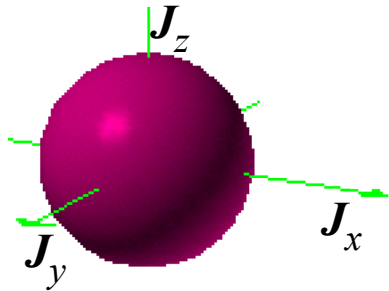
(But, there IS life before AND after diagonalization!)

$$\langle 2\mathbf{T}_0^{(2)} \rangle = cY_0^2 = J(J+1)(3\cos^2 \beta - 1)$$

$$\sqrt{\frac{2}{3}} \langle (\mathbf{T}_2^{(2)} - \mathbf{T}_{-2}^{(2)}) \rangle = c(Y_2^2 - Y_{-2}^2) = J(J+1)(\sin^2 \beta \cos 2\gamma)$$

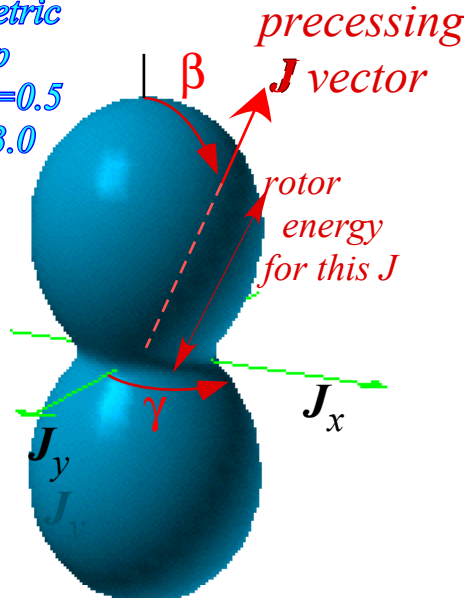
**Spherical
Top**

$A=B=C=1.0$

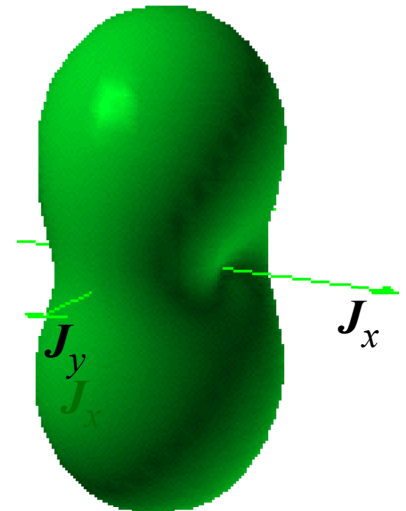


RE surface
Energy plotted radially
vs.
direction of J-vector
|for fixed magnitude |J

**Symmetric
Top**
 $A=B=0.5$
 $C=3.0$



**Asymmetric
Top**
 $A=0.5$
 $B=1.5$
 $C=3.0$

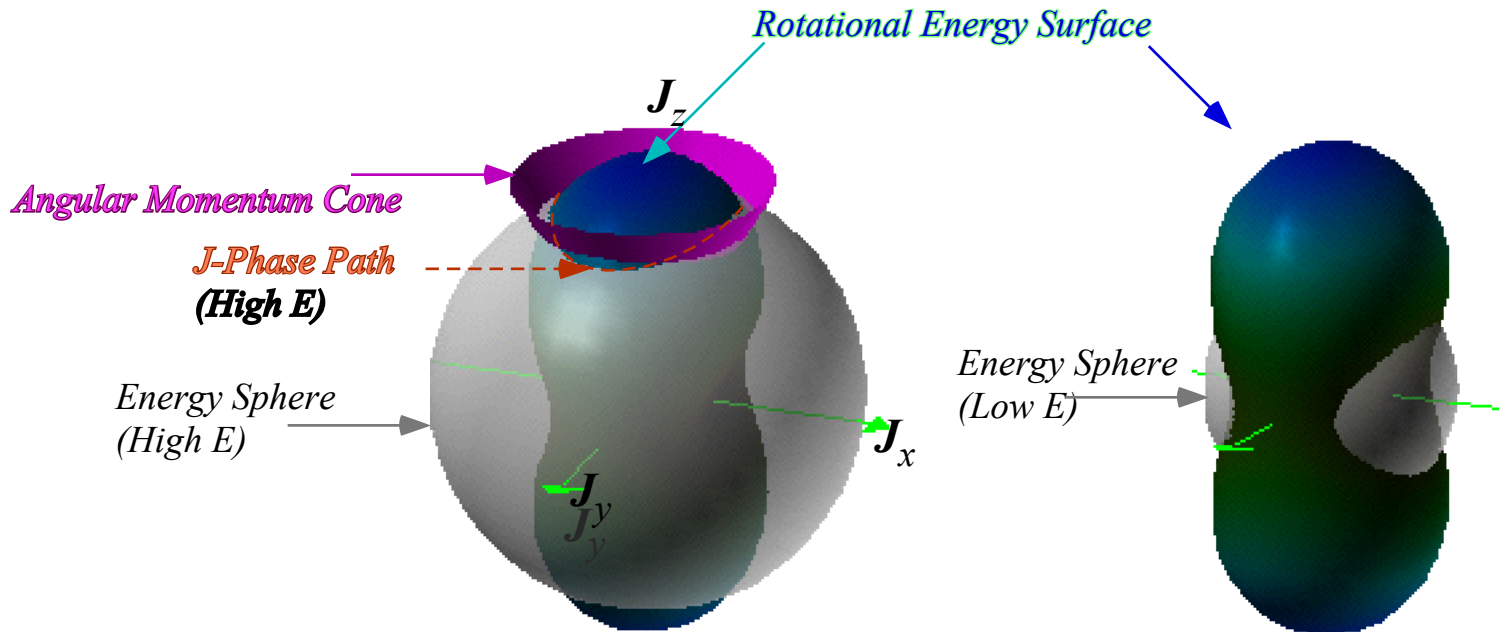


(S) Semiclassical Analysis

Uses

J-Phase Paths (Intersection(s) of RE Surface and Energy Sphere)
and

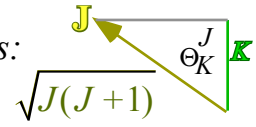
Quantum angular momentum cones



$$\mathbf{J}^2 \left| \begin{matrix} J \\ K \end{matrix} \right\rangle = J(J+1) \left| \begin{matrix} J \\ K \end{matrix} \right\rangle$$

$$\mathbf{J}_z \left| \begin{matrix} J \\ K \end{matrix} \right\rangle = K \left| \begin{matrix} J \\ K \end{matrix} \right\rangle$$

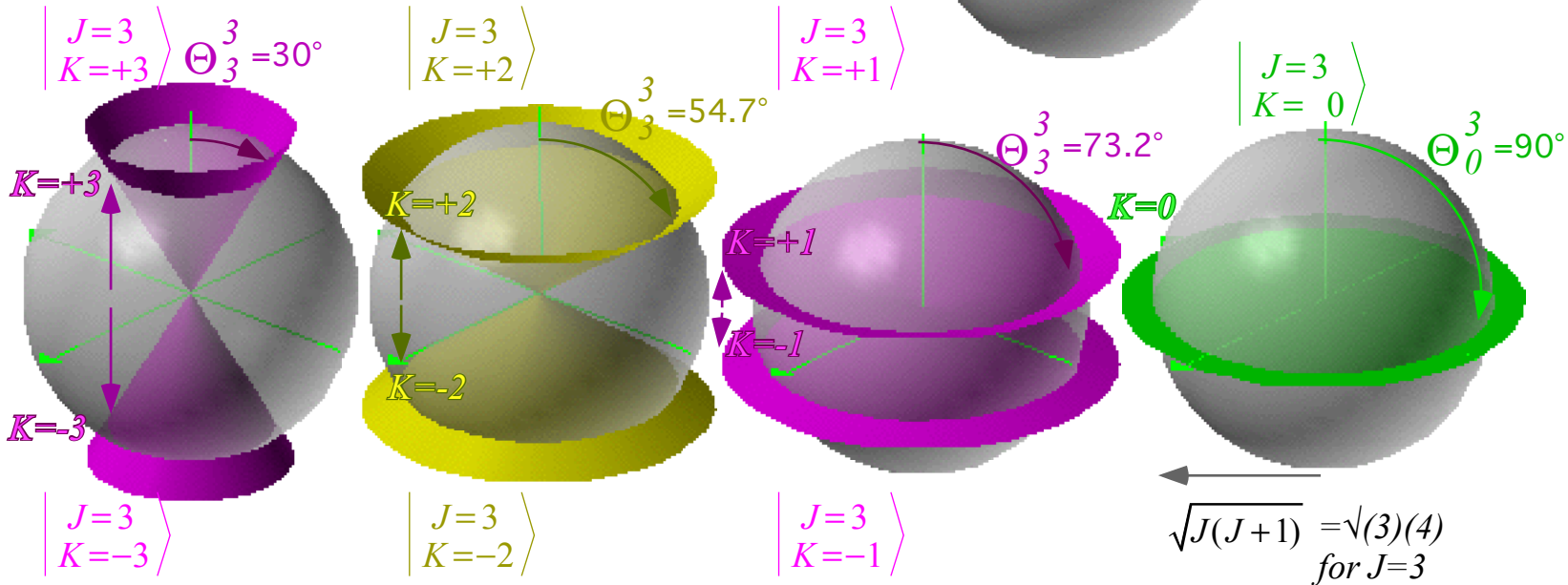
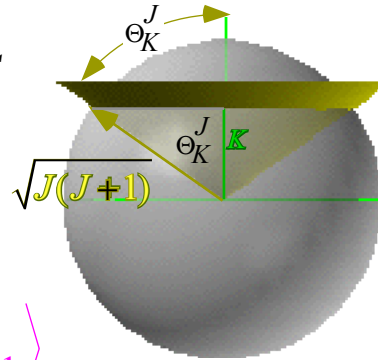
Interpreted "Literally" is:



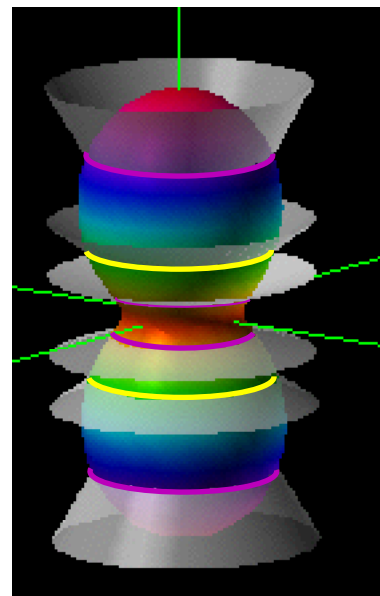
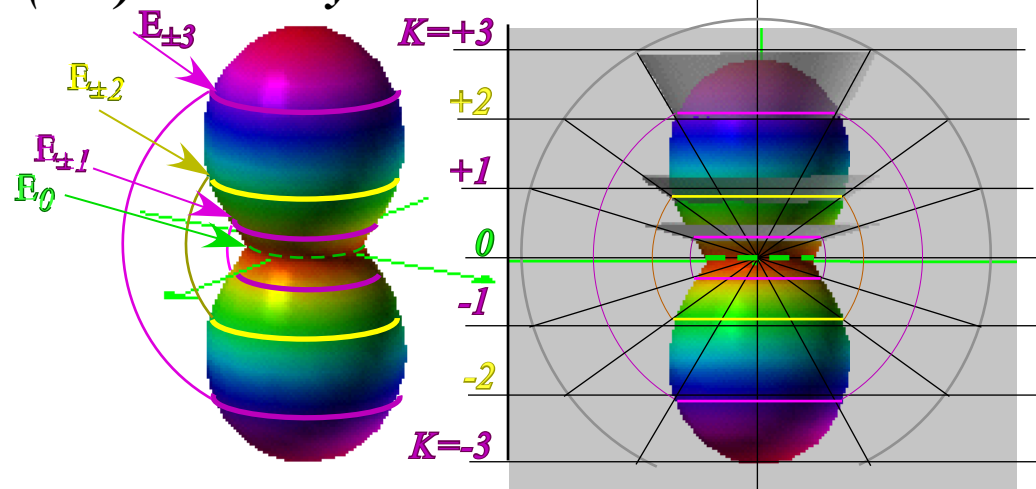
$\sqrt{J(J+1)}$
 $\sim J+1/2$

Quantum Angular Cone Uncertainty Angles

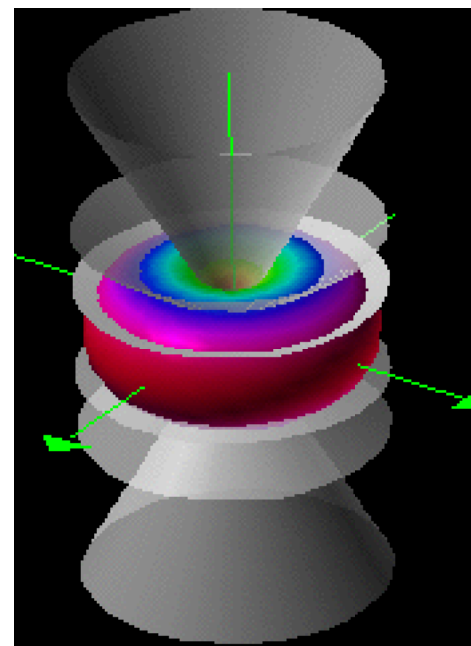
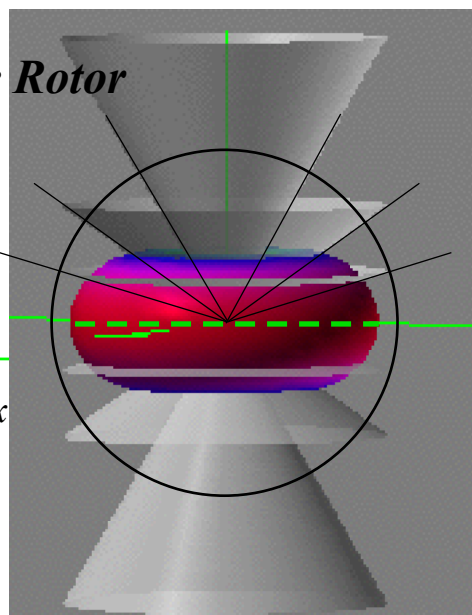
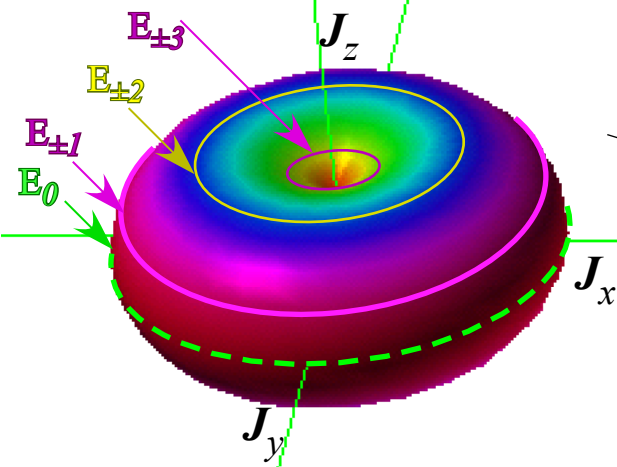
$$\cos \Theta_K^J = \frac{K}{\sqrt{J(J+1)}}$$

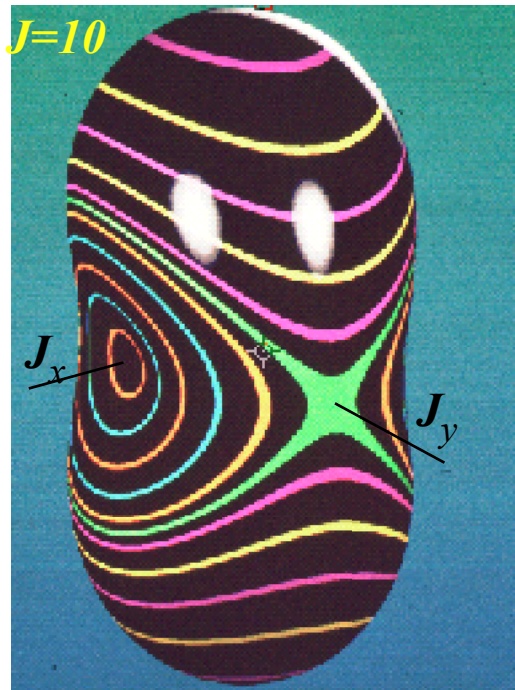
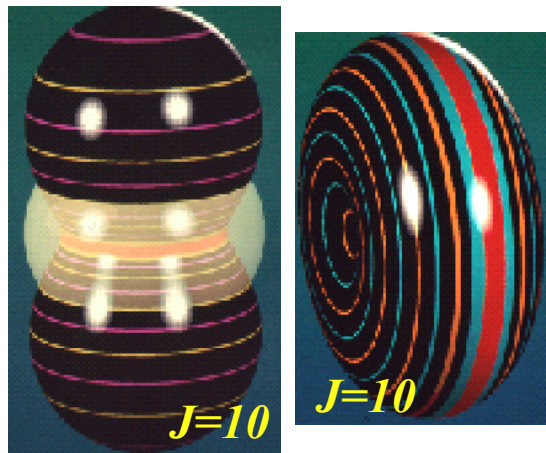


**(S) Semiclassical J-Phase Paths for
(J=3) Prolate Symmetric Rotor**



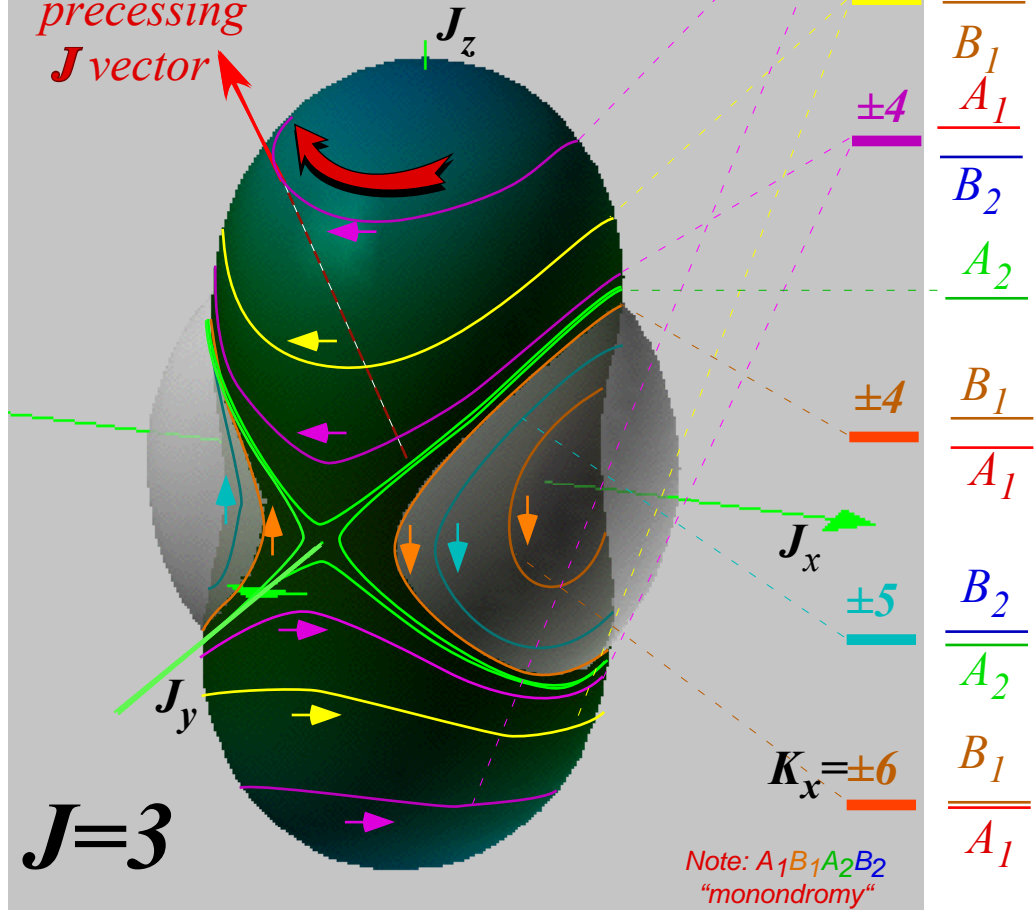
(J=3) Oblate Symmetric Rotor



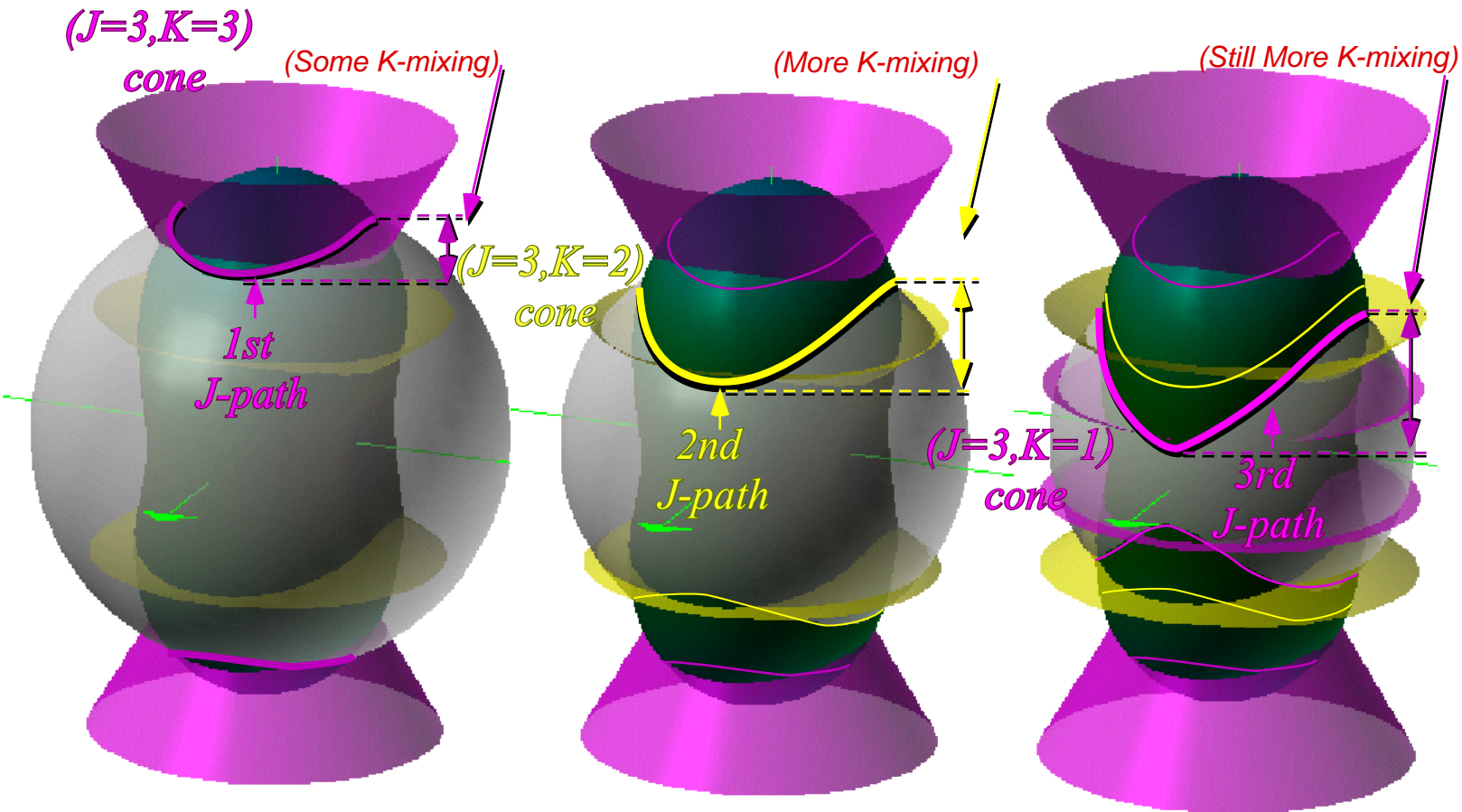


*Asymmetric Top Eigensolutions
Related to RE Surface
and semi-classical \mathbf{J} -phase paths*

*precessing
 \mathbf{J} vector*



Asymmetric Top quantum \mathbf{J} phase paths deviate from (J,K) - cones at low J and K
(This indicates more K -mixing in eigenstates)



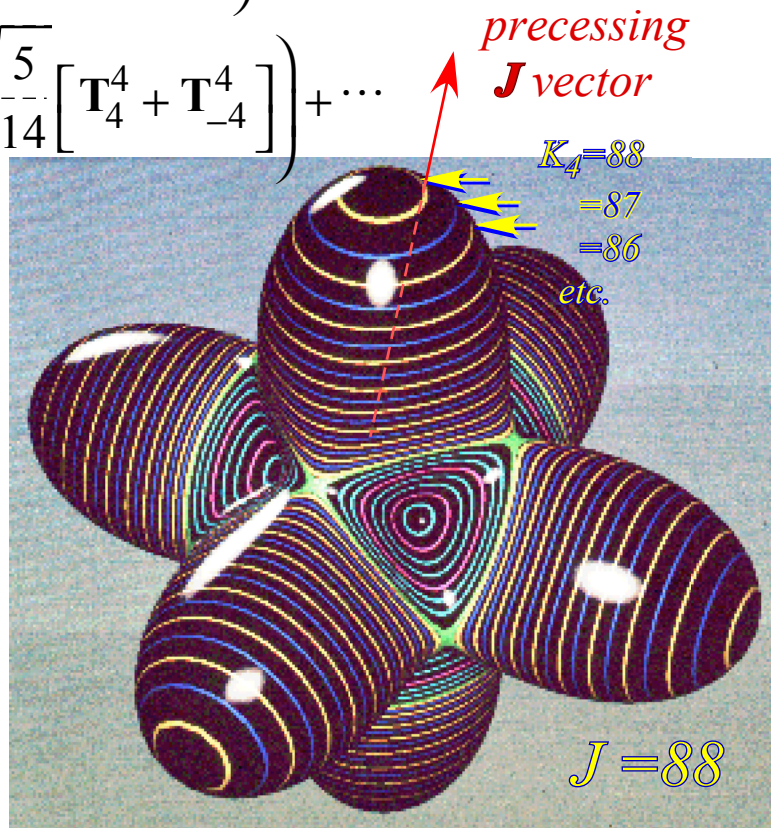
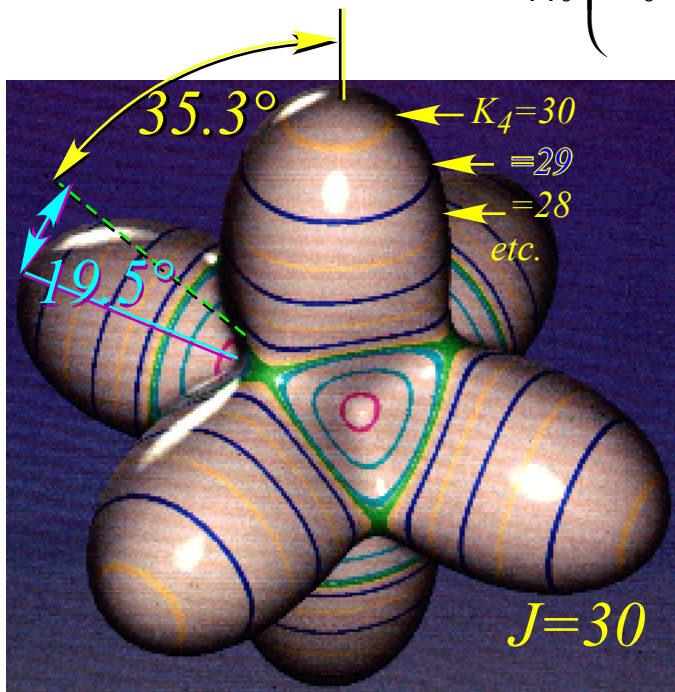
Semi Rigid Rotor Hamiltonian: Centrifugal and Coriolis terms...

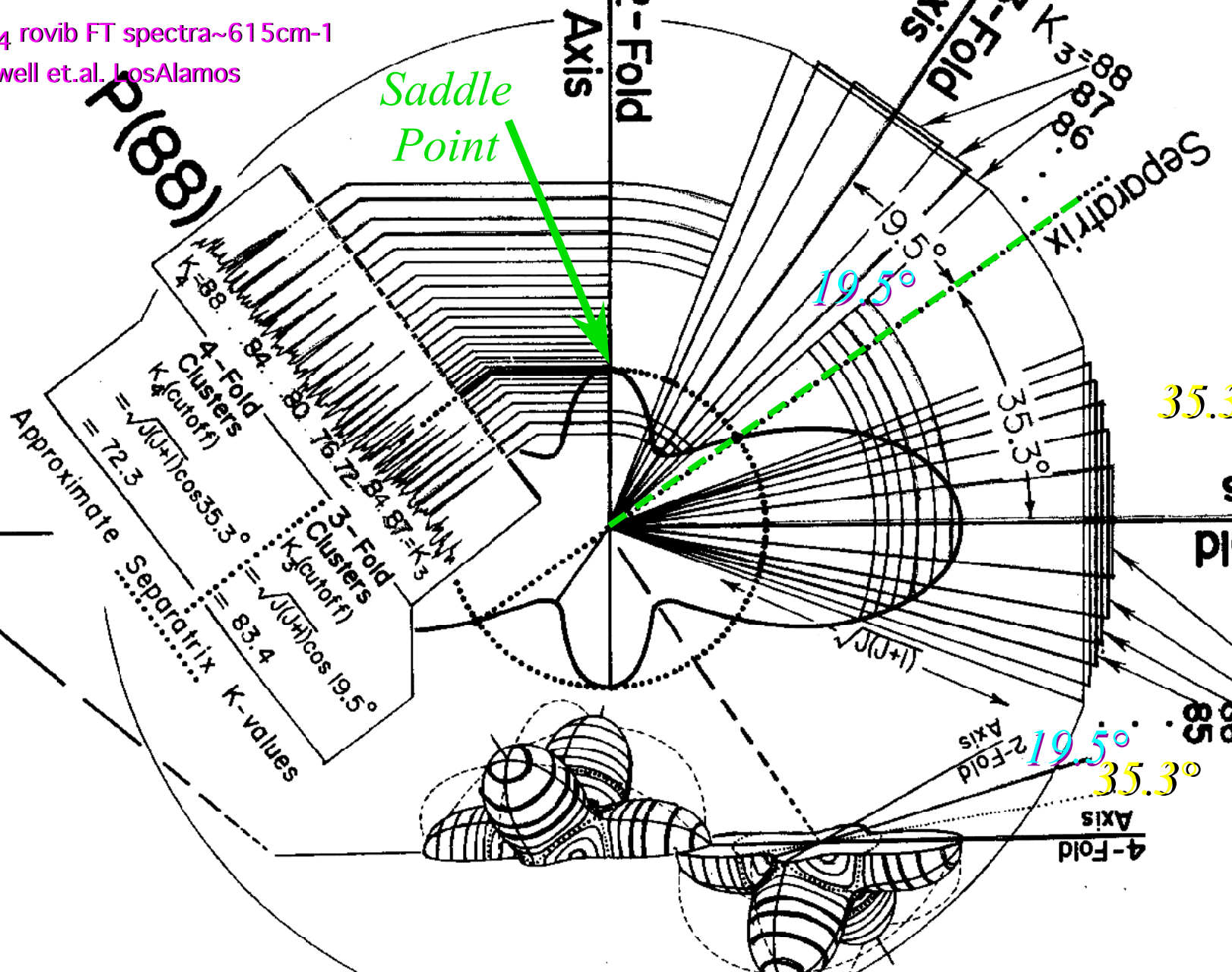
$$\mathbf{H} = A\mathbf{J}_x^2 + B\mathbf{J}_y^2 + C\mathbf{J}_z^2 + t_{xxxx}\mathbf{J}_x^4 + t_{xyxy}\mathbf{J}_x^2\mathbf{J}_y^2 + \dots$$

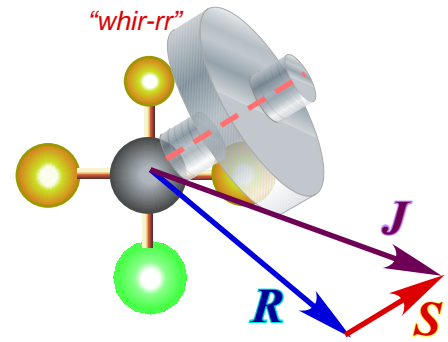
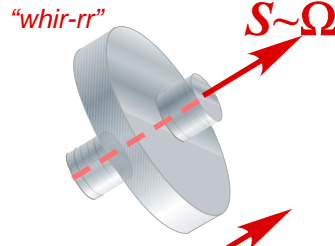
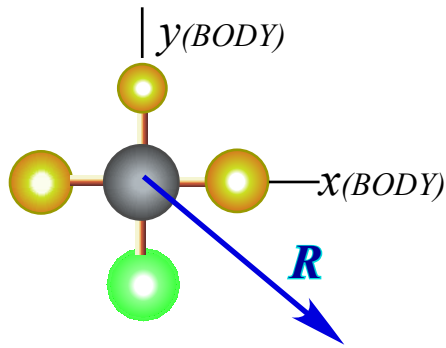
Semi Rigid O_h or T_d Spherical Top: (Hecht Hamiltonian 1960)

$$\mathbf{H} = B\left(\mathbf{J}_x^2 + \mathbf{J}_y^2 + \mathbf{J}_z^2\right) + t_{440}\left(\mathbf{J}_x^4 + \mathbf{J}_y^4 + \mathbf{J}_z^4 - \frac{3}{5}J^4\right) + \dots$$

$$= B\mathbf{J}^2 + t_{440}\left(\mathbf{T}_0^4 + \sqrt{\frac{5}{14}}\left[\mathbf{T}_4^4 + \mathbf{T}_{-4}^4\right]\right) + \dots$$







Rotor R PLUS "Gyro" Spin S EQUALS Compound Rotor $J=R+S$

Compound Rotor Hamiltonian: Rigid rotor with body-fixed "gyro"...

$$H = AR_x^2 + BR_y^2 + CR_z^2 + \dots + (\text{coupling or constraint}) + \dots + B_S S \cdot S$$

(discussed in next talk)

Here *constraint is rigid* so body components (S_x, S_y, S_z) are fixed ("slippery gyro")

ANALOGY: $p^2/2M$ becomes: $(p-eA)^2/2M$ in an em field

Let: $R = J - S$ and consider non-constant terms (ignore gyro S terms that are constant)

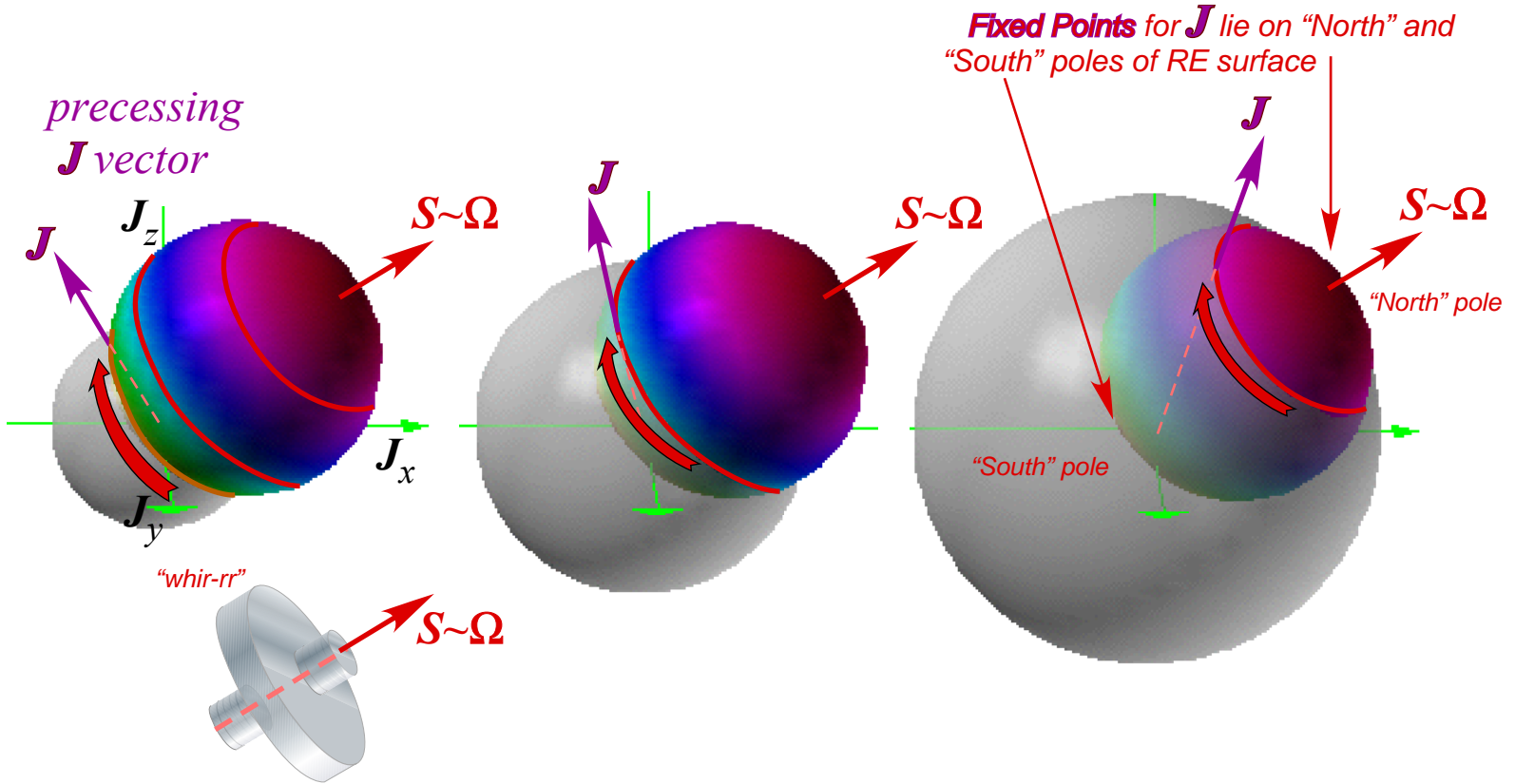
$$H = A(J_x - S_x)^2 + B(J_y - S_y)^2 + C(J_z - S_z)^2 + \dots + 0 \text{ (for constraint)} + \dots + (\text{constant BS terms})$$

$$H = AJ_x^2 + BJ_y^2 + CJ_z^2 + \dots - 2AJ_x S_x - 2BJ_y S_y - 2CJ_z S_z + \dots + (\text{more constant terms})$$

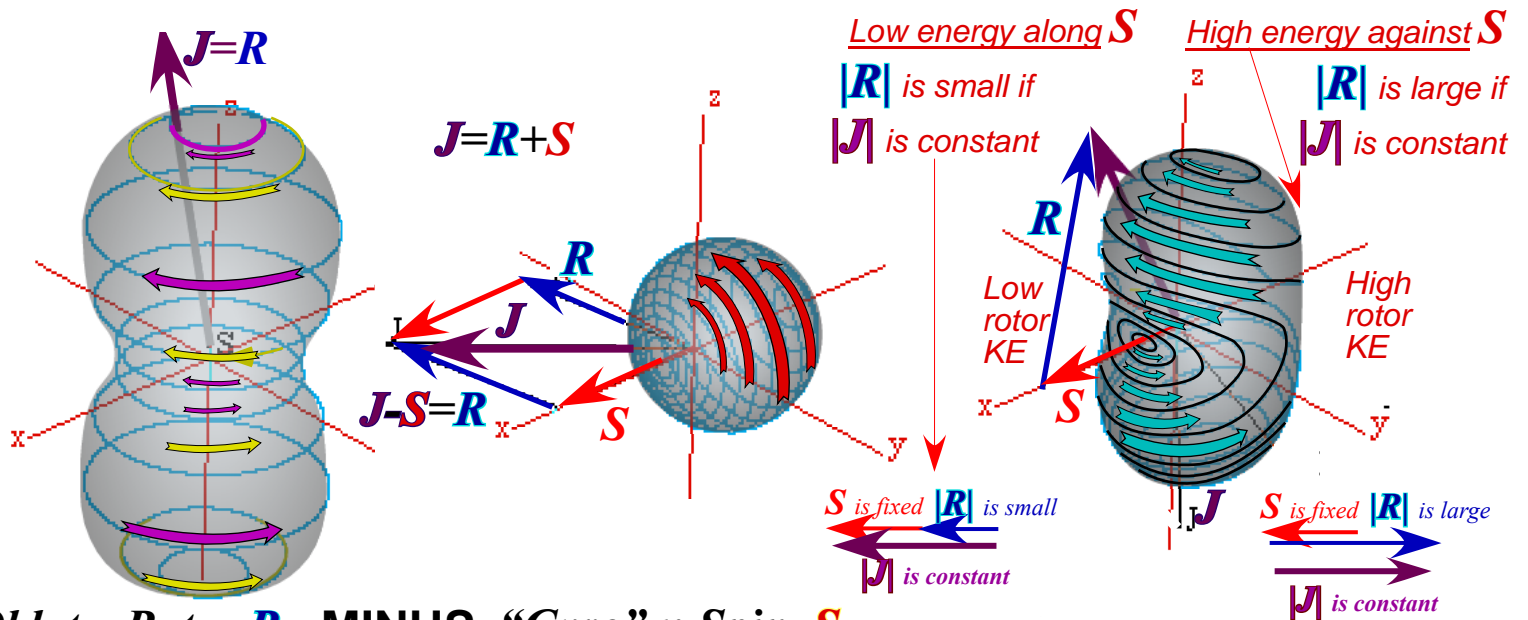
↑ ↑ ↑

"Coriolis effect" subtracts linear or 1st-order J_m or T^1_m terms for gyro-rotor H

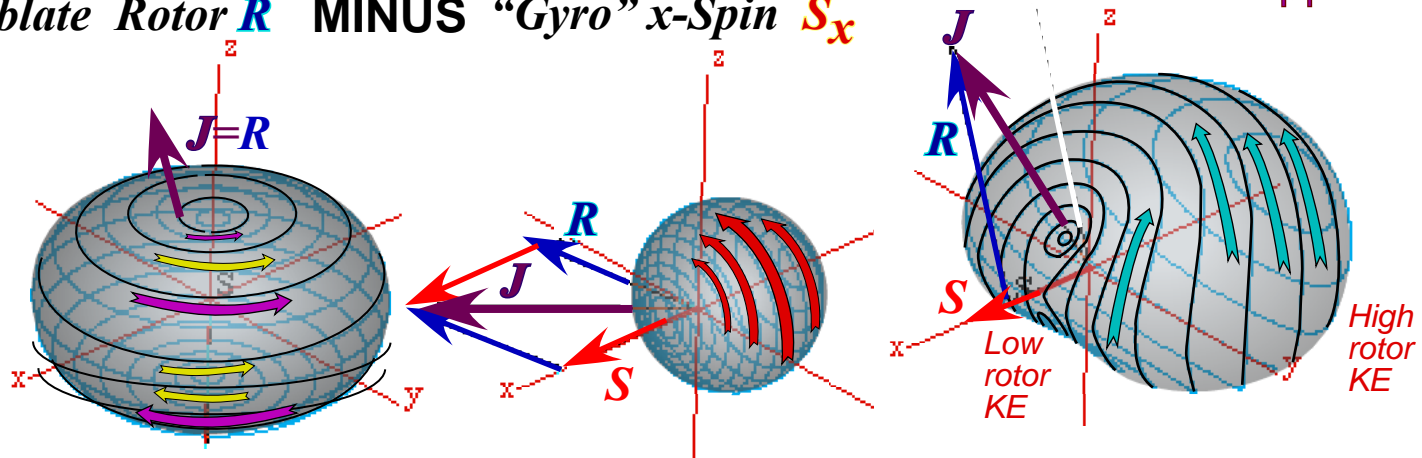
RE Surface for 1st-order \mathbf{J}_m or \mathbf{T}_m^1 term is a sphere displaced in J -direction
 Energy sphere intersections are concentric circular precession paths
 All paths precess with the same sense around gyro S -vector



Prolate Rotor R MINUS "Gyro" x -Spin S_x



Oblate Rotor R MINUS "Gyro" x -Spin S_x



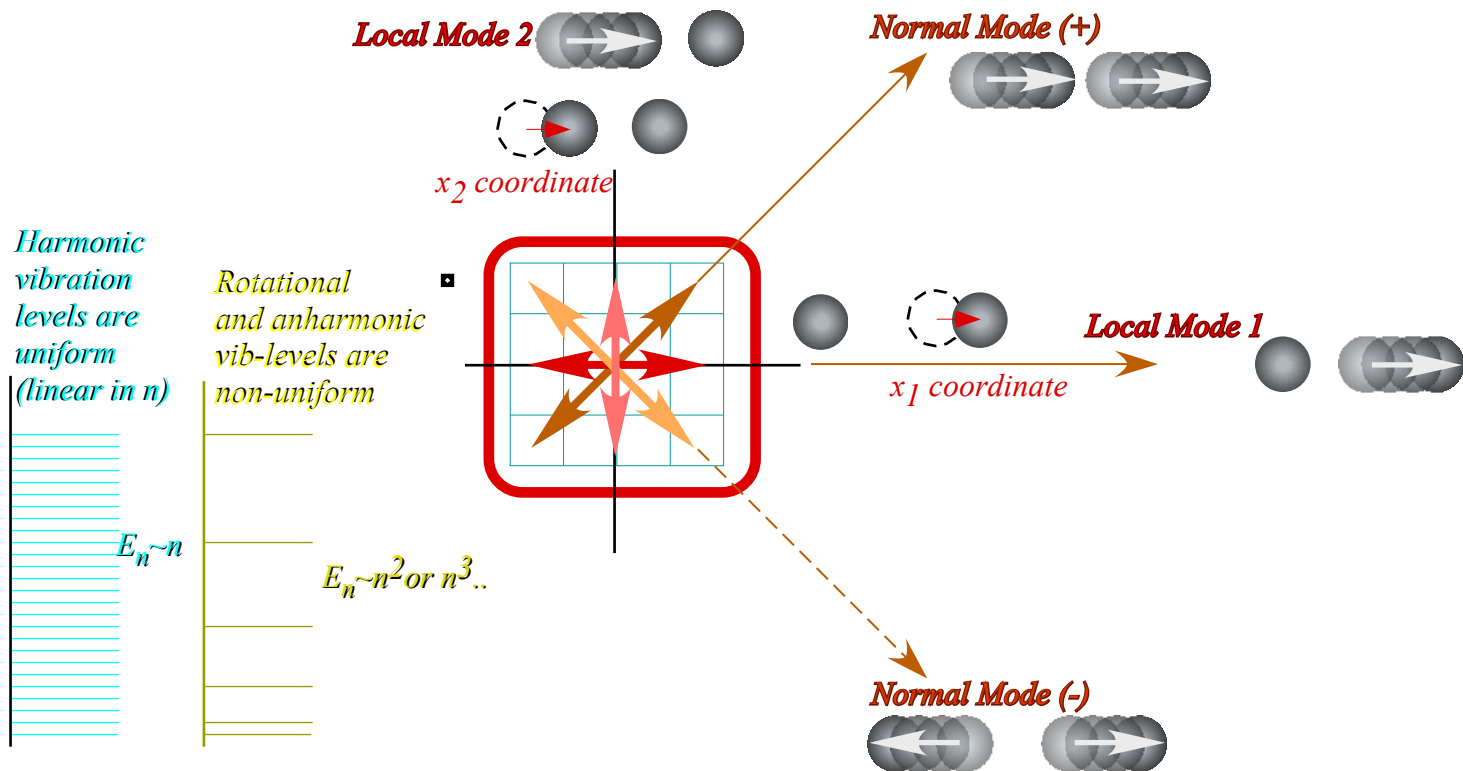
Recall Hamiltonian for 2D vibration has a (quasi-)spin theory, too

$$\mathbf{H} = \omega_0 \mathbf{1} + \Omega \left(\mathbf{a}_{(+)}^\dagger \mathbf{a}_{(+)} - \mathbf{a}_{(-)}^\dagger \mathbf{a}_{(-)} \right) / 2 + \dots + (\text{anharmonic } \mathbf{a}_\mu^\dagger \mathbf{a}_\nu \mathbf{a}_\lambda^\dagger \mathbf{a}_\kappa \text{ terms})$$

$$= \omega_0 \mathbf{1} + \Omega \mathbf{J}_x + \dots + B \mathbf{J}_x^2 + C \mathbf{J}_y^2 + A \mathbf{J}_z^2 + \dots + a_{xy} \mathbf{J}_x \mathbf{J}_y +$$

1st-order \mathbf{J}_m or \mathbf{T}_m^1 term
is *harmonic* part of \mathbf{H}

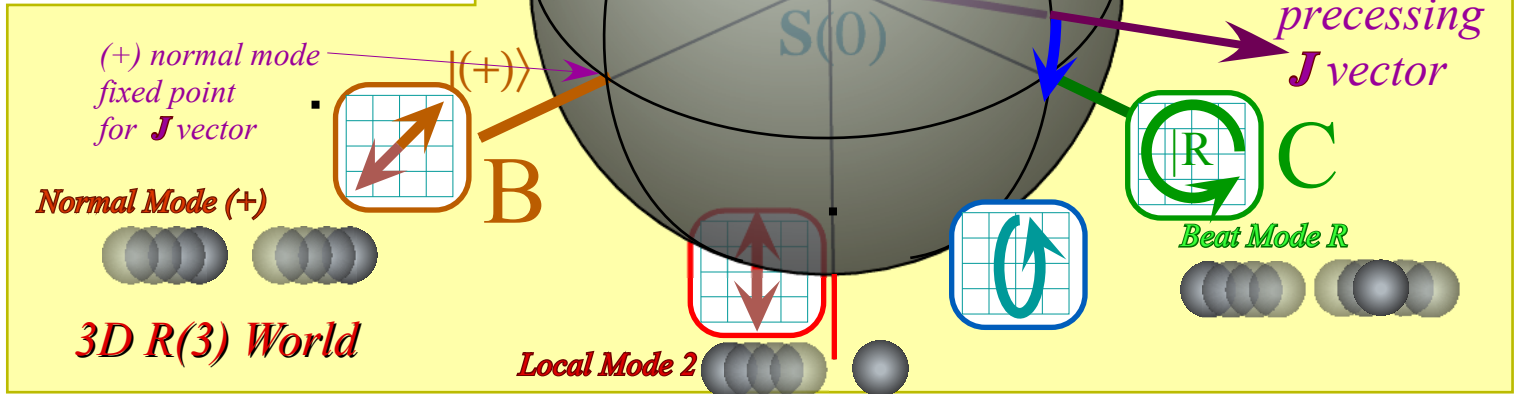
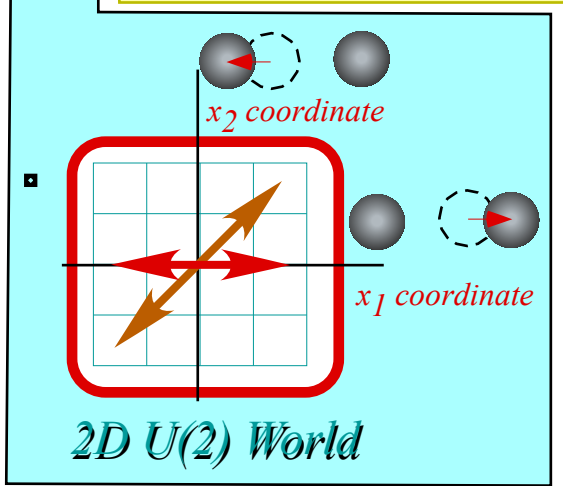
Higher-order \mathbf{J}_m or \mathbf{T}_m^1 terms
are *anharmonic* parts of \mathbf{H}



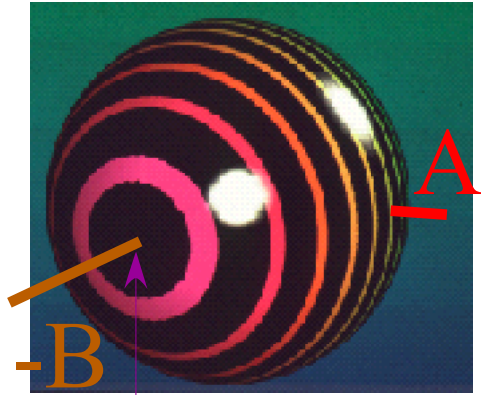
(contd) 2D vibration are related to 3D rotation of “quasi-spin” \mathbf{J}

$$H = \omega_0 \mathbf{1} + \Omega \left(\mathbf{a}_{(+)}^\dagger \mathbf{a}_{(+)} - \mathbf{a}_{(-)}^\dagger \mathbf{a}_{(-)} \right) / 2 + \dots + (\text{anharmonic } \mathbf{a}_\mu^\dagger \mathbf{a}_\nu \mathbf{a}_\lambda^\dagger \mathbf{a}_\kappa \text{ terms})$$

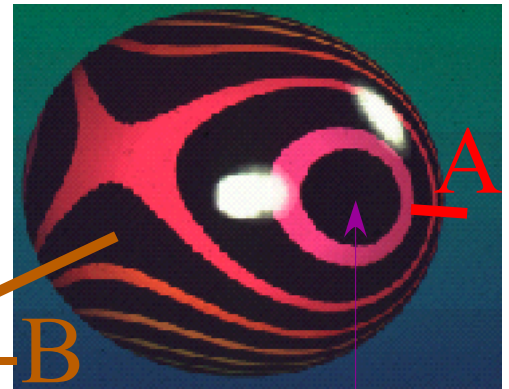
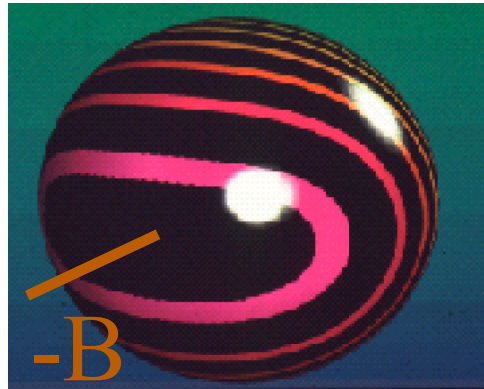
$$H = \omega_0 \mathbf{1} + \Omega \mathbf{J}_x + \dots + B \mathbf{J}_x^2 + C \mathbf{J}_y^2 + A \mathbf{J}_z^2 + \dots + a_{xy} \mathbf{J}_x \mathbf{J}_y + \dots$$



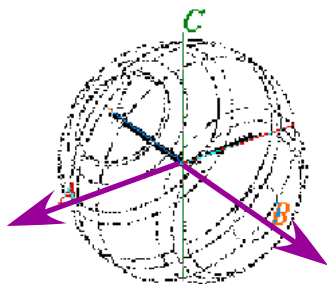
For higher J values, anharmonic terms grow to make stable local modes



(+) normal mode
fixed point
for \mathbf{J} vector

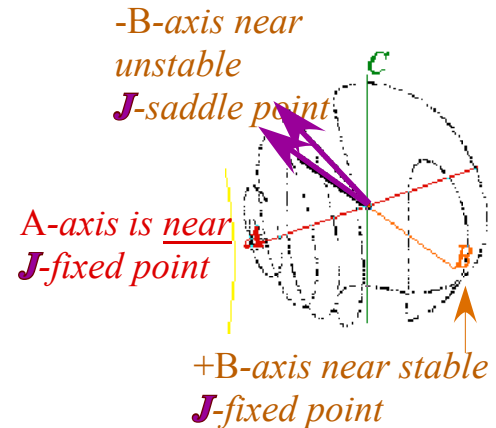


(1) local mode
fixed point
for \mathbf{J} vector

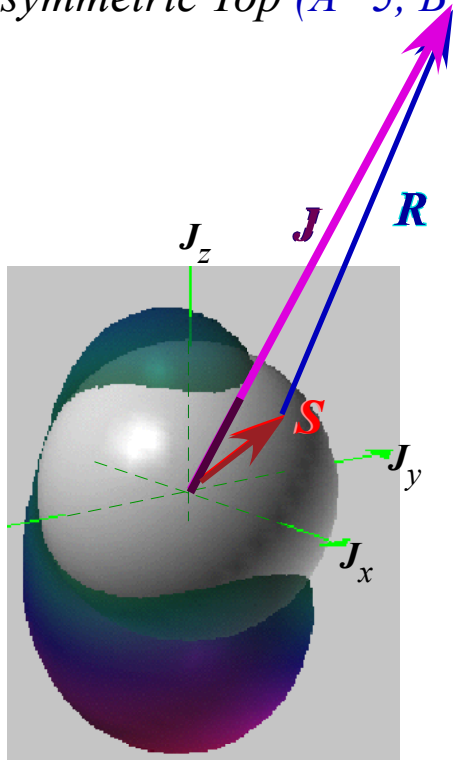


$\pm B$ -axes are
 \mathbf{J} -fixed points

A -axis is NOT
 \mathbf{J} -fixed point

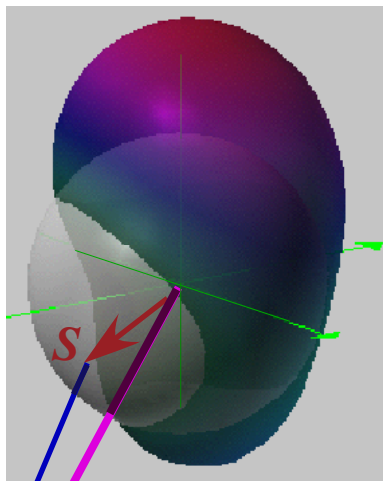


Spin gyro $S=(1,1,1)$ attached to
Asymmetric Top ($A=5, B=10, C=15$)

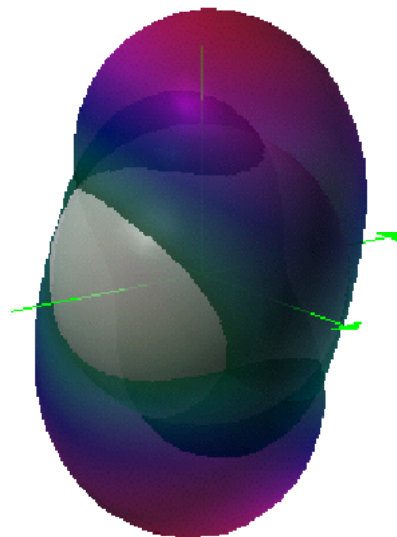


“Sherman” (The shark)
First appeared in a
1992 JCP article by
Hougan, Kleiner, and Ortigoso

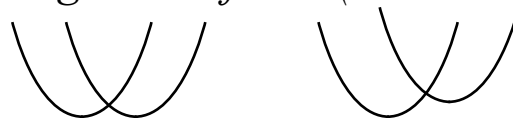
Time reversed
gyro $-S=(-1,-1,-1)$



The two together



Crossing RE surfaces
analogous to
Crossing PE surfaces (Jahn-Teller)



(Pre) Conclusion

Rotational Energy (RE) surfaces: Past & Future

Two or more RE's beg for an interaction.

Here it's the coupling we "turned off" into a constraint.

$$\langle \mathbf{H} \rangle = \begin{pmatrix} \text{Spin-up RE}(\beta, \gamma) & \text{Coupling}(\beta, \gamma) \\ \text{Coupling}(\beta, \gamma)^* & \text{Spin-down RE}(\beta, \gamma) \end{pmatrix}$$

Base RE surfaces are eigensolutions of this matrix.

Combination RES depends on eigenvector chosen.

This opens worlds of interesting mechanics. (Both QM and CM)

Two (or more) surfaces imply an infinity of surfaces "between" them.

Intermediate surfaces not unique for each energy

("Tide" rises and falls, saddles open and close. Result: Chaotic trajectory)

(Recall that Born told Otto Stern that his spin experiment wouldn't show quantization.)

(Final) Conclusion

Rotational Energy (RE) surfaces help analyze rotor dynamics as do Potential Energy (PE) surfaces for vibration.

PE surfaces based on vibrational coordinates.

RE surfaces based on rovibrational phase space.

Can approximate quantum levels and spectra and also mixing and transitions.

RES have a variety of complementary surfaces:

Angular Velocity surfaces: (AVS) $\boldsymbol{\omega} \cdot \mathbf{I} \cdot \boldsymbol{\omega} = 2E$ (Poinsot ellipsoid)

Angular Momentum surfaces: (AMS) $\mathbf{J} \cdot \mathbf{I}^{-1} \cdot \mathbf{J} = 2E$ (Landau ellipsoid)

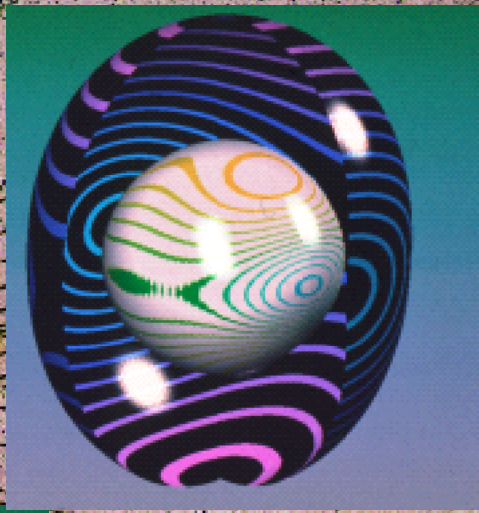
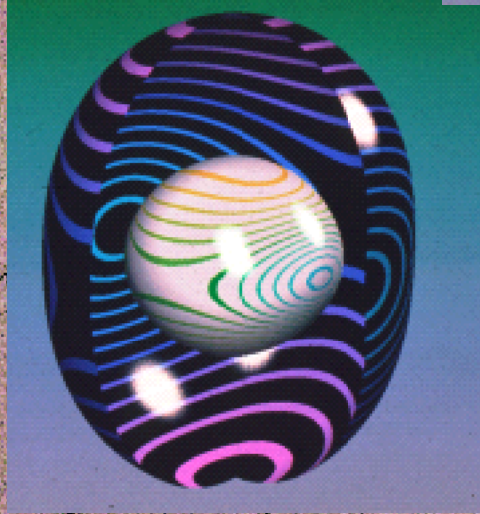
PE surfaces used since beginning of QM (Born 1926)

RE surfaces first used in 1976.

J = 10.5 Eigenvalues of Spin-Rotor

500.0 1000.0 1500.0

-1000.0



(A₁ B₁ A₂ B₂) clusters

R =

0.0 10.0 20.0 30.0 40.0 50.0 60.0 70.0 80.0 90.0

D_{XX} With D_{YY} = 2D_{XX} and D_{ZZ} = 3D_{XX}