

Lecture 12 advanced
Thur. 2.25.2016

Kepler Geometry of IHO (Isotropic Harmonic Oscillator) *Elliptical Orbits*
(Ch. 8 and Ch. 9 of Unit 1)

Kepler “laws” (Some that apply to all central (isotropic) $F(r)$ force fields)

Angular momentum invariance of IHO: $F(r)=-k\cdot r$ with $U(r)=k\cdot r^2/2$ (Derived here)

*Angular momentum invariance of **Coulomb**: $F(r)=-GMm/r^2$ with $U(r)=-GMm\cdot/r$ (Derived later)*

Total energy $E=KE+PE$ invariance of IHO: $F(r)=-k\cdot r$ (Derived here)

*Total energy $E=KE+PE$ invariance of **Coulomb**: $F(r)=-GMm/r^2$ (Derived later)*

A confusing introduction to Coriolis-centrifugal force geometry (Derived better later)

Introduction to dual matrix operator contact geometry (based on IHO orbits)

Quadratic form ellipse $\mathbf{r}\cdot\mathbf{Q}\cdot\mathbf{r}=1$ vs. inverse form ellipse $\mathbf{p}\cdot\mathbf{Q}^{-1}\cdot\mathbf{p}=1$

Duality norm relations ($\mathbf{r}\cdot\mathbf{p}=1$)

\mathbf{Q} -Ellipse tangents \mathbf{r}' normal to dual \mathbf{Q}^{-1} -ellipse position \mathbf{p} ($\mathbf{r}'\cdot\mathbf{p}=0=\mathbf{r}\cdot\mathbf{p}'$)

Operator geometric sequences and eigenvectors

Alternative scaling of matrix operator geometry

Vector calculus of tensor operation


Q: Where is this headed? A: Lagrangian-Hamiltonian duality

[Link \$\Rightarrow\$ BoxIt simulation of IHO orbits](#)

[Link \$\rightarrow\$ IHO orbital time rates of change](#)

[Link \$\rightarrow\$ IHO Exegesis Plot](#)

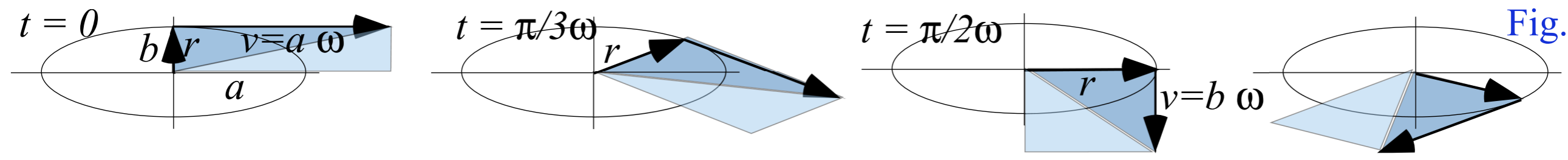
Kepler “laws” (Some that apply to all central (isotropic) $F(r)$ force fields)

-  *Angular momentum invariance of IHO: $F(r) = -k \cdot r$ with $U(r) = k \cdot r^2 / 2$ (Derived here)*
- Angular momentum invariance of Coulomb: $F(r) = -GMm/r^2$ with $U(r) = -GMm \cdot /r$ (Derived in Unit 5)*
- Total energy $E = KE + PE$ invariance of IHO: $F(r) = -k \cdot r$ (Derived here)*
- Total energy $E = KE + PE$ invariance of Coulomb: $F(r) = -GMm/r^2$ (Derived in Unit 5)*

Some Kepler's "laws" for central (isotropic) force $F(r)$

...and certainly apply to the IHO: $F(r) = -k \cdot r$ with $U(r) = k \cdot r^2 / 2$ (Recall from Lect.12 p.19: $k = G \frac{4\pi}{3} m \rho_{\oplus}$)

Unit 1
Fig. 9.8



1. Area of triangle $\Delta_r^v = \mathbf{r} \times \mathbf{v} / 2$ is constant

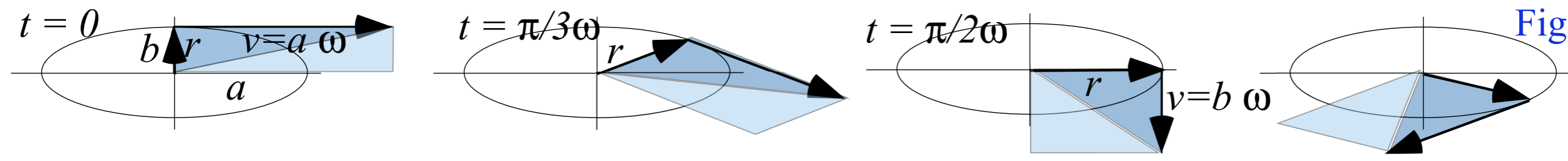
$$\mathbf{r} \times \mathbf{v} = r_x v_y - r_y v_x = a \cos \omega t \cdot (b \omega \cos \omega t) - b \sin \omega t \cdot (-a \omega \sin \omega t) = ab \cdot \omega (\cos^2 \omega t + \sin^2 \omega t) \quad \checkmark \text{ for IHO}$$

$$\begin{pmatrix} r_x \\ r_y \end{pmatrix} = \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} a \cos \omega t \\ b \sin \omega t \end{pmatrix} \quad \begin{pmatrix} v_x \\ v_y \end{pmatrix} = \begin{pmatrix} -a \omega \sin \omega t \\ b \omega \cos \omega t \end{pmatrix}$$

Some Kepler's "laws" that apply to any central (isotropic) force $F(r)$

...and certainly apply to the IHO: $F(r) = -k \cdot r$ with $U(r) = k \cdot r^2 / 2$ (Recall from Lect.12 p.19: $k = G \frac{4\pi}{3} m \rho_{\oplus}$)

Unit 1
Fig. 9.8



1. Area of triangle $\Delta_r^v = \mathbf{r} \times \mathbf{v} / 2$ is constant

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✓ for IHO

2. Angular momentum $\mathbf{L} = m \mathbf{r} \times \mathbf{v}$ is conserved

$$L = m |\mathbf{r} \times \mathbf{v}| = m (r_x v_y - r_y v_x) = m \cdot ab \cdot \omega$$

✓ for IHO

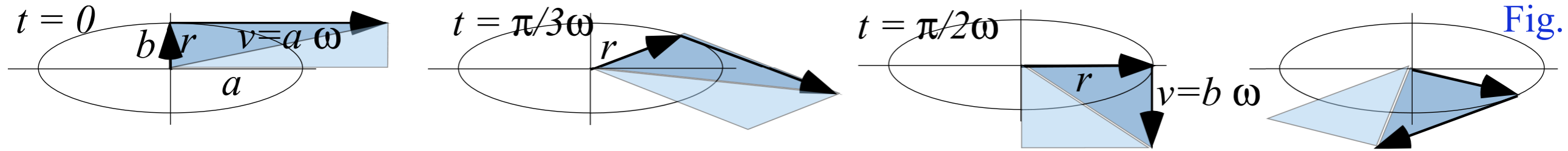
$$|\mathbf{r} \times \mathbf{v}| = r \cdot v \cdot \sin \Delta_r^v$$

$$|\mathbf{r} \cdot \mathbf{v}| = r \cdot v \cdot \cos \Delta_r^v$$

Some Kepler's "laws" that apply to any central (isotropic) force $F(r)$

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3. Equal area is swept by radius vector in each equal time interval T

$$A_T = \int_0^T \frac{\mathbf{r} \times d\mathbf{r}}{2} = \int_0^T \frac{\mathbf{r} \times \frac{d\mathbf{r}}{dt}}{2} dt = \int_0^T \frac{\mathbf{r} \times \mathbf{v}}{2} dt = \frac{L}{2m} \int_0^T dt = \frac{L}{2m} T$$

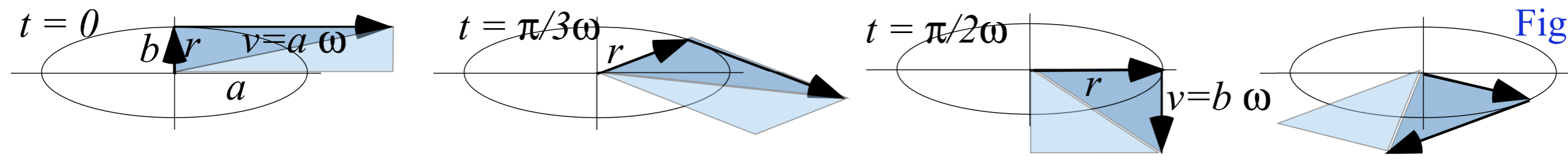
✓ for IHO

$$|\mathbf{r} \times d\mathbf{r}| = r \cdot dr \cdot \sin \angle_r$$

Some Kepler's "laws" that apply to any central (isotropic) force $F(r)$

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Unit 1
Fig. 9.8



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$$L = m \mathbf{r} \times \mathbf{v} = m (r_x v_y - r_y v_x) = m \cdot ab \cdot \omega = m \cdot ab \cdot \frac{2\pi}{\tau}$$

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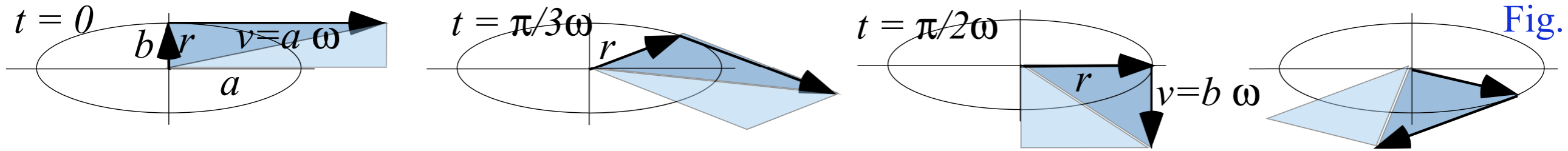
✓ for IHO

In one period: $\tau = \frac{1}{\nu} = \frac{2\pi}{\omega} = \frac{2mA_\tau}{L}$ the area is: $A_\tau = \frac{L\tau}{2m}$ (= $ab \cdot \pi$ for ellipse orbit)

Some Kepler's "laws" that apply to any central (isotropic) force $F(r)$

...and certainly apply to the IHO: $F(r) = -k \cdot r$ with $U(r) = k \cdot r^2 / 2$ (Recall from Lect.12 p.19: $k = G \frac{4\pi}{3} m \rho_{\oplus}$)

Unit 1
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In one period: $\tau = \frac{1}{\nu} = \frac{2\pi}{\omega} = \frac{2mA_\tau}{L}$ the area is: $A_\tau = \frac{L\tau}{2m}$ ($= ab \cdot \pi$ for ellipse orbit)

(Recall from Lecture 7: $\omega = \sqrt{k/m} = \sqrt{G\rho_{\oplus} 4\pi/3}$)

Some Kepler's "laws" for all central (isotropic) force $F(r)$ fields

Angular momentum invariance of IHO: $F(r) = -k \cdot r$ with $U(r) = k \cdot r^2 / 2$ (Derived here)

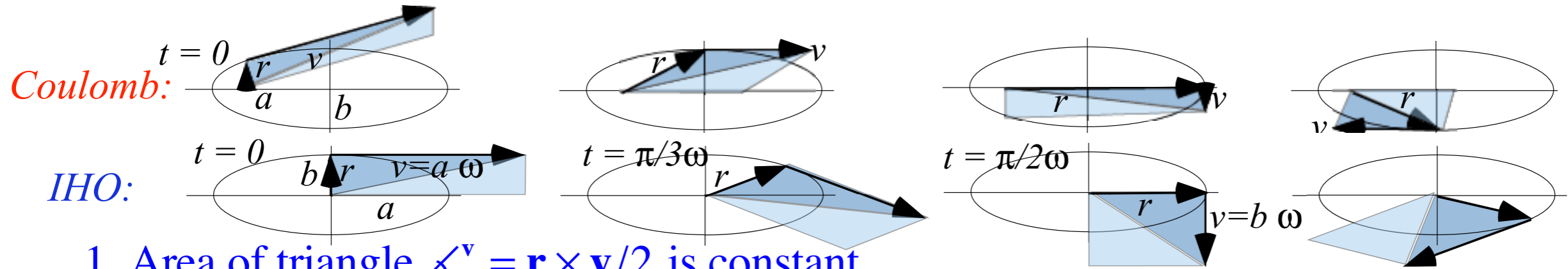
 *Angular momentum invariance of Coulomb: $F(r) = -GMm/r^2$ with $U(r) = -GMm \cdot /r$ (Derived in Unit 5)*

Total energy $E = KE + PE$ invariance of IHO: $F(r) = -k \cdot r$ (Derived here)

Total energy $E = KE + PE$ invariance of Coulomb: $F(r) = -GMm/r^2$ (Derived in Unit 5)

Some Kepler's "laws" that apply to any central (isotropic) force $F(r)$

Apply to IHO: $F(r) = -k \cdot r$ with $U(r) = k \cdot r^2 / 2$ and Coulomb: $F(r) = -GMm/r^2$ with $U(r) = -GMm \cdot / r$



1. Area of triangle $\Delta_r^v = \mathbf{r} \times \mathbf{v} / 2$ is constant

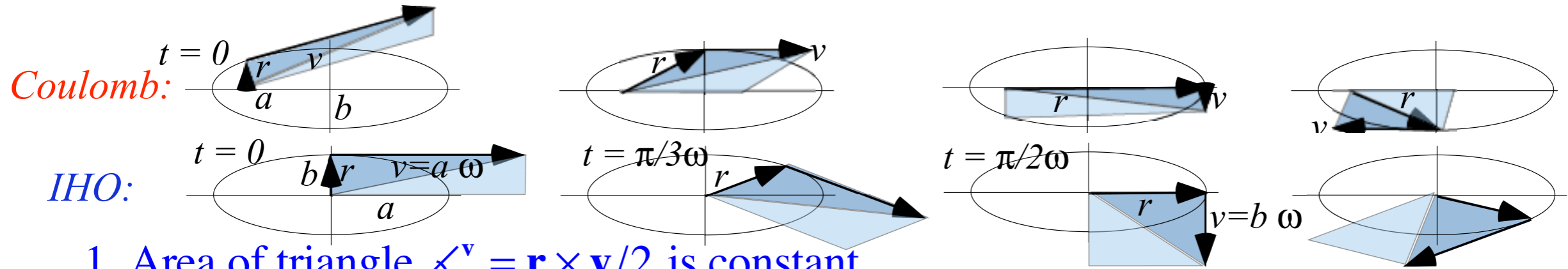
$$\mathbf{r} \times \mathbf{v} = r_x v_y - r_y v_x = \begin{cases} ab \cdot \sqrt{G\rho_{\oplus} 4\pi / 3} & \text{for IHO} \\ a^{-1/2} b \sqrt{GM_{\oplus}} & \text{for Coul.} \end{cases}$$

✓ for IHO

(Derived in Unit 5) ✓ for Coul.

Some Kepler's "laws" that apply to any central (isotropic) force $F(r)$

Apply to IHO: $F(r) = -k \cdot r$ with $U(r) = k \cdot r^2 / 2$ and Coulomb: $F(r) = -GMm/r^2$ with $U(r) = -GMm \cdot /r$



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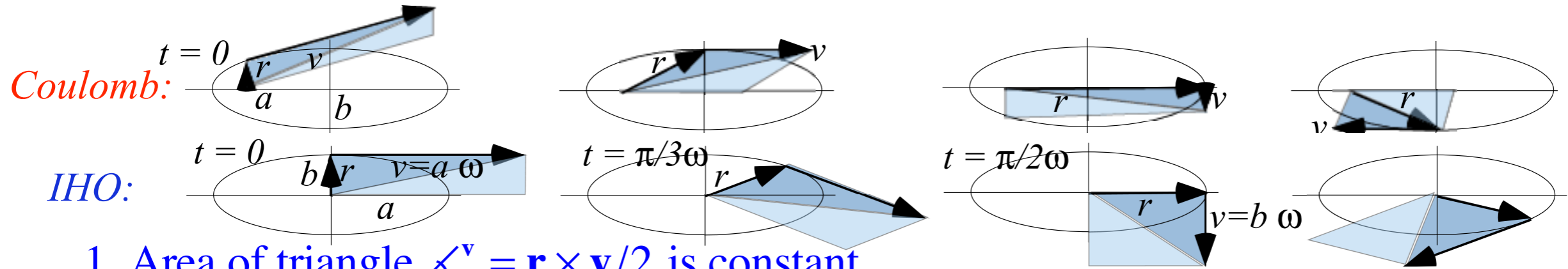
2. Angular momentum $L = m \mathbf{r} \times \mathbf{v}$ is conserved

$$L = m \mathbf{r} \times \mathbf{v} = m (r_x v_y - r_y v_x) = \begin{cases} m \cdot ab \cdot \sqrt{G\rho_{\oplus} 4\pi / 3} & \text{for IHO} \\ m \cdot a^{-1/2} b \sqrt{GM_{\oplus}} & \text{for Coul. (... in Unit 5)} \end{cases}$$

✓ for IHO
✓ for Coul.

Some Kepler's "laws" that apply to any central (isotropic) force $F(r)$

Apply to IHO: $F(r) = -k \cdot r$ with $U(r) = k \cdot r^2 / 2$ and Coulomb: $F(r) = -GMm/r^2$ with $U(r) = -GMm/r$



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3. Equal area is swept by radius vector in each equal time interval T

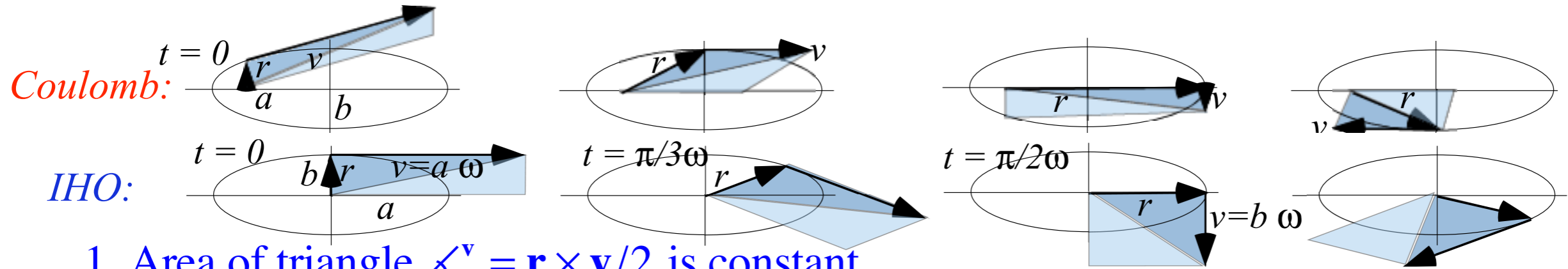
In one period:

$$\tau = \frac{1}{\nu} = \frac{2\pi}{\omega} = \frac{2mA_{\tau}}{L} = \frac{2m \cdot ab \cdot \pi}{L} = \begin{cases} \frac{2m \cdot ab \cdot \pi}{m \cdot ab \cdot \sqrt{G\rho_{\oplus} 4\pi / 3}} \\ \frac{2m \cdot ab \cdot \pi}{m \cdot a^{-1/2} b \sqrt{GM_{\oplus}}} \end{cases}$$

Applies to any central $F(r)$ (IHO and Coulomb)

Some Kepler's "laws" that apply to any central (isotropic) force $F(r)$

Apply to IHO: $F(r) = -k \cdot r$ with $U(r) = k \cdot r^2 / 2$ and Coulomb: $F(r) = -GMm/r^2$ with $U(r) = -GMm \cdot / r$



1. Area of triangle $\Delta_r^v = \mathbf{r} \times \mathbf{v} / 2$ is constant

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✓ for IHO
✓ for Coul.

2. Angular momentum $L = m \mathbf{r} \times \mathbf{v}$ is conserved

$$L = m \mathbf{r} \times \mathbf{v} = m (r_x v_y - r_y v_x) = \begin{cases} m \cdot ab \cdot \sqrt{G\rho_{\oplus} 4\pi / 3} & \text{for IHO} \\ m \cdot a^{-1/2} b \sqrt{GM_{\oplus}} & \text{for Coul. (... in Unit 5)} \end{cases}$$

✓ for IHO
✓ for Coul.

3. Equal area is swept by radius vector in each equal time interval T

In one period:

$$\tau = \frac{1}{\nu} = \frac{2\pi}{\omega} = \frac{2mA_{\tau}}{L} = \frac{2m \cdot ab \cdot \pi}{L}$$

Applies to any central $F(r)$

$$= \begin{cases} \frac{2m \cdot ab \cdot \pi}{m \cdot ab \cdot \sqrt{G\rho_{\oplus} 4\pi / 3}} = \frac{2\pi}{\sqrt{G\rho_{\oplus} 4\pi / 3}} & \text{for IHO} \\ \frac{2m \cdot ab \cdot \pi}{m \cdot a^{-1/2} b \sqrt{GM_{\oplus}}} = \frac{2\pi}{a^{-3/2} \sqrt{GM_{\oplus}}} & \text{for Coul.} \end{cases}$$

(not a function of a or b) that is ω_{IHO}
(not a function of b) that is ω_{Coul}

Some Kepler's "laws" for all central (isotropic) force $F(r)$ fields

Angular momentum invariance of IHO: $F(r)=-k\cdot r$ with $U(r)=k\cdot r^2/2$ (Derived here)

Angular momentum invariance of Coulomb: $F(r)=-GMm/r^2$ with $U(r)=-GMm\cdot/r$ (Derived in Unit 5)

 *Total energy $E=KE+PE$ invariance of IHO: $F(r)=-k\cdot r$ (Derived here)*

Total energy $E=KE+PE$ invariance of Coulomb: $F(r)=-GMm/r^2$ (Derived in Unit 5)

Kepler laws involve \mathcal{L} -momentum conservation in isotropic force $F(r)$

Now consider orbital energy conservation of the IHO: $F(r)=-k \cdot r$ with $U(r)=k \cdot r^2/2$

Total energy= $KE + PE$ is constant

$$\begin{aligned} KE + PE &= \frac{1}{2} \mathbf{v} \cdot \mathbf{M} \cdot \mathbf{v} + \frac{1}{2} \mathbf{r} \cdot \mathbf{K} \cdot \mathbf{r} \\ &= \frac{1}{2} \begin{pmatrix} v_x & v_y \end{pmatrix} \cdot \begin{pmatrix} m & 0 \\ 0 & m \end{pmatrix} \cdot \begin{pmatrix} v_x \\ v_y \end{pmatrix} + \begin{pmatrix} r_x & r_y \end{pmatrix} \cdot \begin{pmatrix} k & 0 \\ 0 & k \end{pmatrix} \cdot \begin{pmatrix} r_x \\ r_y \end{pmatrix} \\ &= \frac{1}{2} m v_x^2 + \frac{1}{2} m v_y^2 + \frac{1}{2} k r_x^2 + \frac{1}{2} k r_y^2 \\ &= \frac{1}{2} m (-a\omega \sin \omega t)^2 + \frac{1}{2} m (b\omega \cos \omega t)^2 + \frac{1}{2} k (a \cos \omega t)^2 + \frac{1}{2} k (b \sin \omega t)^2 \end{aligned}$$

$$\begin{pmatrix} v_x \\ v_y \end{pmatrix} = \begin{pmatrix} -a\omega \sin \omega t \\ b\omega \cos \omega t \end{pmatrix}$$

$$\begin{pmatrix} r_x \\ r_y \end{pmatrix} = \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} a \cos \omega t \\ b \sin \omega t \end{pmatrix}$$

Kepler laws involve \mathcal{L} -momentum conservation in isotropic force $F(r)$

Now consider orbital energy conservation of the IHO: $F(r) = -k \cdot r$ with $U(r) = k \cdot r^2 / 2$

Total IHO energy = $KE + PE$ is constant

$$\begin{aligned} KE + PE &= \frac{1}{2} \mathbf{v} \cdot \mathbf{M} \cdot \mathbf{v} + \frac{1}{2} \mathbf{r} \cdot \mathbf{K} \cdot \mathbf{r} \\ &= \frac{1}{2} \begin{pmatrix} v_x & v_y \end{pmatrix} \cdot \begin{pmatrix} m & 0 \\ 0 & m \end{pmatrix} \cdot \begin{pmatrix} v_x \\ v_y \end{pmatrix} + \begin{pmatrix} r_x & r_y \end{pmatrix} \cdot \begin{pmatrix} k & 0 \\ 0 & k \end{pmatrix} \cdot \begin{pmatrix} r_x \\ r_y \end{pmatrix} \\ &= \frac{1}{2} m v_x^2 + \frac{1}{2} m v_y^2 + \frac{1}{2} k r_x^2 + \frac{1}{2} k r_y^2 \\ &= \frac{1}{2} m (-a\omega \sin \omega t)^2 + \frac{1}{2} m (b\omega \cos \omega t)^2 + \frac{1}{2} k (a \cos \omega t)^2 + \frac{1}{2} k (b \sin \omega t)^2 \\ &= \frac{1}{2} m a^2 \omega^2 (\sin^2 \omega t) + \frac{1}{2} m b^2 \omega^2 (\cos^2 \omega t) + \frac{1}{2} k a^2 (\cos^2 \omega t) + \frac{1}{2} k b^2 (\sin^2 \omega t) \\ &= \frac{1}{2} m \omega^2 (a^2 + b^2) \quad \text{Given : } k = m\omega^2 \end{aligned}$$

Kepler laws involve Δ -momentum conservation in isotropic force $F(r)$

Now consider orbital energy conservation of the IHO: $F(r)=-k \cdot r$ with $U(r)=k \cdot r^2/2$

Total IHO energy= $KE + PE$ is constant

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 KE + PE &= \frac{1}{2} \mathbf{v} \cdot \mathbf{M} \cdot \mathbf{v} + \frac{1}{2} \mathbf{r} \cdot \mathbf{K} \cdot \mathbf{r} \\
 &= \frac{1}{2} \begin{pmatrix} v_x & v_y \end{pmatrix} \cdot \begin{pmatrix} m & 0 \\ 0 & m \end{pmatrix} \cdot \begin{pmatrix} v_x \\ v_y \end{pmatrix} + \begin{pmatrix} r_x & r_y \end{pmatrix} \cdot \begin{pmatrix} k & 0 \\ 0 & k \end{pmatrix} \cdot \begin{pmatrix} r_x \\ r_y \end{pmatrix} \\
 &= \frac{1}{2} m v_x^2 + \frac{1}{2} m v_y^2 + \frac{1}{2} k r_x^2 + \frac{1}{2} k r_y^2 \\
 &= \frac{1}{2} m (-a\omega \sin \omega t)^2 + \frac{1}{2} m (b\omega \cos \omega t)^2 + \frac{1}{2} k (a \cos \omega t)^2 + \frac{1}{2} k (b \sin \omega t)^2 \\
 &= \frac{1}{2} m a^2 \omega^2 (\sin^2 \omega t) + \frac{1}{2} m b^2 \omega^2 (\cos^2 \omega t) + \frac{1}{2} k a^2 (\cos^2 \omega t) + \frac{1}{2} k b^2 (\sin^2 \omega t) \\
 &= \frac{1}{2} m \omega^2 (a^2 + b^2) \quad \text{Given : } k = m\omega^2
 \end{aligned}$$

$$E = KE + PE = \frac{1}{2} m \omega^2 (a^2 + b^2) = \frac{1}{2} k (a^2 + b^2) \quad \text{since: } \omega = \sqrt{\frac{k}{m}} = \sqrt{G \rho_{\oplus} 4\pi / 3} \quad \text{or: } m\omega^2 = k$$

Some Kepler's "laws" for all central (isotropic) force $F(r)$ fields

Angular momentum invariance of IHO: $F(r) = -k \cdot r$ with $U(r) = k \cdot r^2 / 2$ (Derived here)

Angular momentum invariance of Coulomb: $F(r) = -GMm/r^2$ with $U(r) = -GMm \cdot /r$ (Derived in Unit 5)

Total energy $E = KE + PE$ invariance of IHO: $F(r) = -k \cdot r$ (Derived here)

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 KE + PE &= \frac{1}{2} \mathbf{v} \cdot \mathbf{M} \cdot \mathbf{v} + \frac{1}{2} \mathbf{r} \cdot \mathbf{K} \cdot \mathbf{r} \\
 &= \frac{1}{2} \begin{pmatrix} v_x & v_y \end{pmatrix} \cdot \begin{pmatrix} m & 0 \\ 0 & m \end{pmatrix} \cdot \begin{pmatrix} v_x \\ v_y \end{pmatrix} + \begin{pmatrix} r_x & r_y \end{pmatrix} \cdot \begin{pmatrix} k & 0 \\ 0 & k \end{pmatrix} \cdot \begin{pmatrix} r_x \\ r_y \end{pmatrix} \\
 &= \frac{1}{2} m v_x^2 + \frac{1}{2} m v_y^2 + \frac{1}{2} k r_x^2 + \frac{1}{2} k r_y^2 \\
 &= \frac{1}{2} m (-a\omega \sin \omega t)^2 + \frac{1}{2} m (b\omega \cos \omega t)^2 + \frac{1}{2} k (a \cos \omega t)^2 + \frac{1}{2} k (b \sin \omega t)^2 \\
 &= \frac{1}{2} m a^2 \omega^2 (\sin^2 \omega t) + \frac{1}{2} m b^2 \omega^2 (\cos^2 \omega t) + \frac{1}{2} k a^2 (\cos^2 \omega t) + \frac{1}{2} k b^2 (\sin^2 \omega t) \\
 &= \frac{1}{2} m \omega^2 (a^2 + b^2) \quad \text{Given: } k = m\omega^2
 \end{aligned}$$

$$E = KE + PE = \frac{1}{2} m \omega^2 (a^2 + b^2) = \frac{1}{2} k (a^2 + b^2) \quad \text{since: } \omega = \sqrt{\frac{k}{m}} = \sqrt{G\rho_{\oplus} 4\pi / 3} \quad \text{or: } m\omega^2 = k$$

We'll see that the Coul. orbits are simpler:

(like the period...not a function of b)

Kepler laws involve Δ -momentum conservation in isotropic force $F(r)$

Now consider orbital energy conservation of the IHO: $F(r)=-k \cdot r$ with $U(r)=k \cdot r^2/2$

Total IHO energy= $KE + PE$ is constant

$$\begin{aligned}
 KE + PE &= \frac{1}{2} \mathbf{v} \cdot \mathbf{M} \cdot \mathbf{v} + \frac{1}{2} \mathbf{r} \cdot \mathbf{K} \cdot \mathbf{r} \\
 &= \frac{1}{2} \begin{pmatrix} v_x & v_y \end{pmatrix} \cdot \begin{pmatrix} m & 0 \\ 0 & m \end{pmatrix} \cdot \begin{pmatrix} v_x \\ v_y \end{pmatrix} + \begin{pmatrix} r_x & r_y \end{pmatrix} \cdot \begin{pmatrix} k & 0 \\ 0 & k \end{pmatrix} \cdot \begin{pmatrix} r_x \\ r_y \end{pmatrix} \\
 &= \frac{1}{2} m v_x^2 + \frac{1}{2} m v_y^2 + \frac{1}{2} k r_x^2 + \frac{1}{2} k r_y^2 \\
 &= \frac{1}{2} m (-a\omega \sin \omega t)^2 + \frac{1}{2} m (b\omega \cos \omega t)^2 + \frac{1}{2} k (a \cos \omega t)^2 + \frac{1}{2} k (b \sin \omega t)^2 \\
 &= \frac{1}{2} m a^2 \omega^2 (\sin^2 \omega t) + \frac{1}{2} m b^2 \omega^2 (\cos^2 \omega t) + \frac{1}{2} k a^2 (\cos^2 \omega t) + \frac{1}{2} k b^2 (\sin^2 \omega t) \\
 &= \frac{1}{2} m \omega^2 (a^2 + b^2) \quad \text{Given : } k = m\omega^2
 \end{aligned}$$

$$E = KE + PE = \frac{1}{2} m \omega^2 (a^2 + b^2) = \frac{1}{2} k (a^2 + b^2) \quad \text{since: } \omega = \sqrt{\frac{k}{m}} = \sqrt{G\rho_{\oplus} 4\pi/3} \quad \text{or: } m\omega^2 = k$$

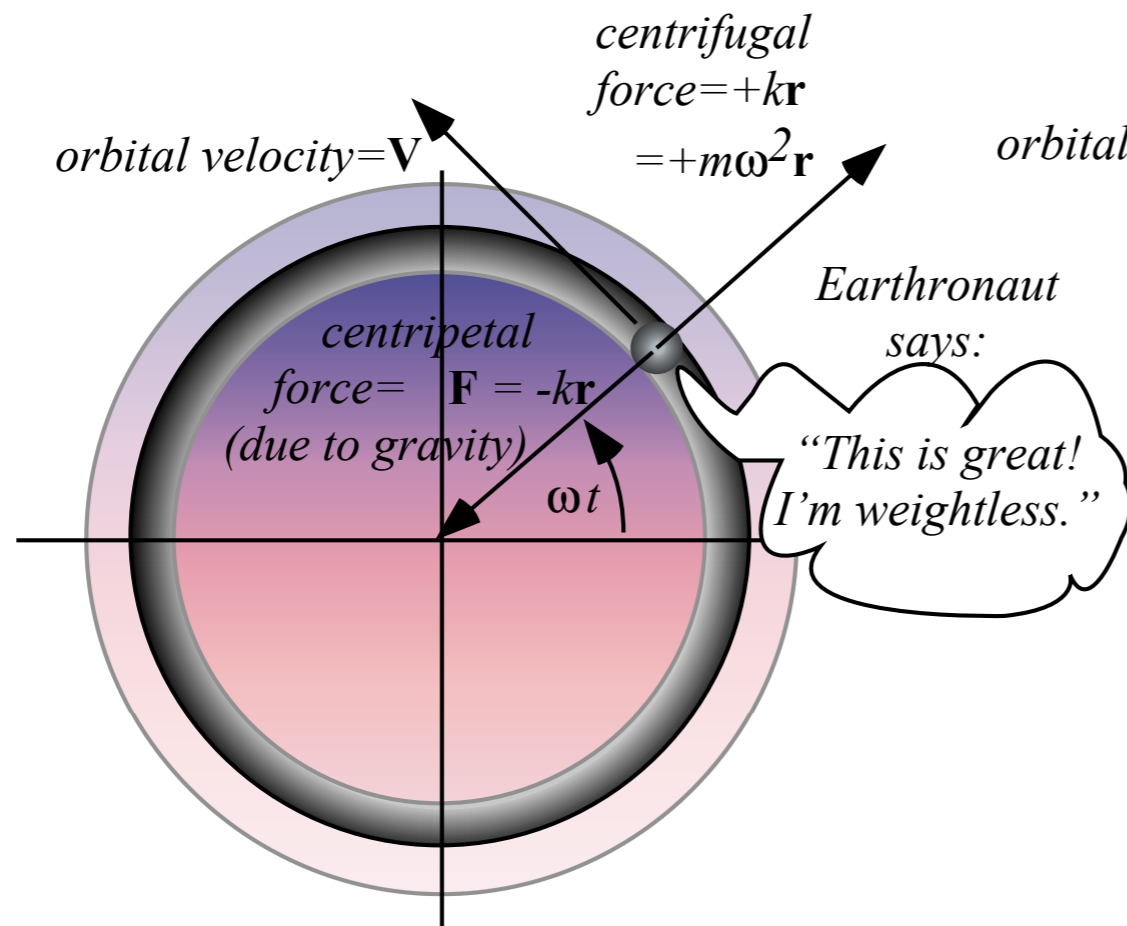
We'll see that the Coul. orbits are simpler:

(like the period...not a function of b)

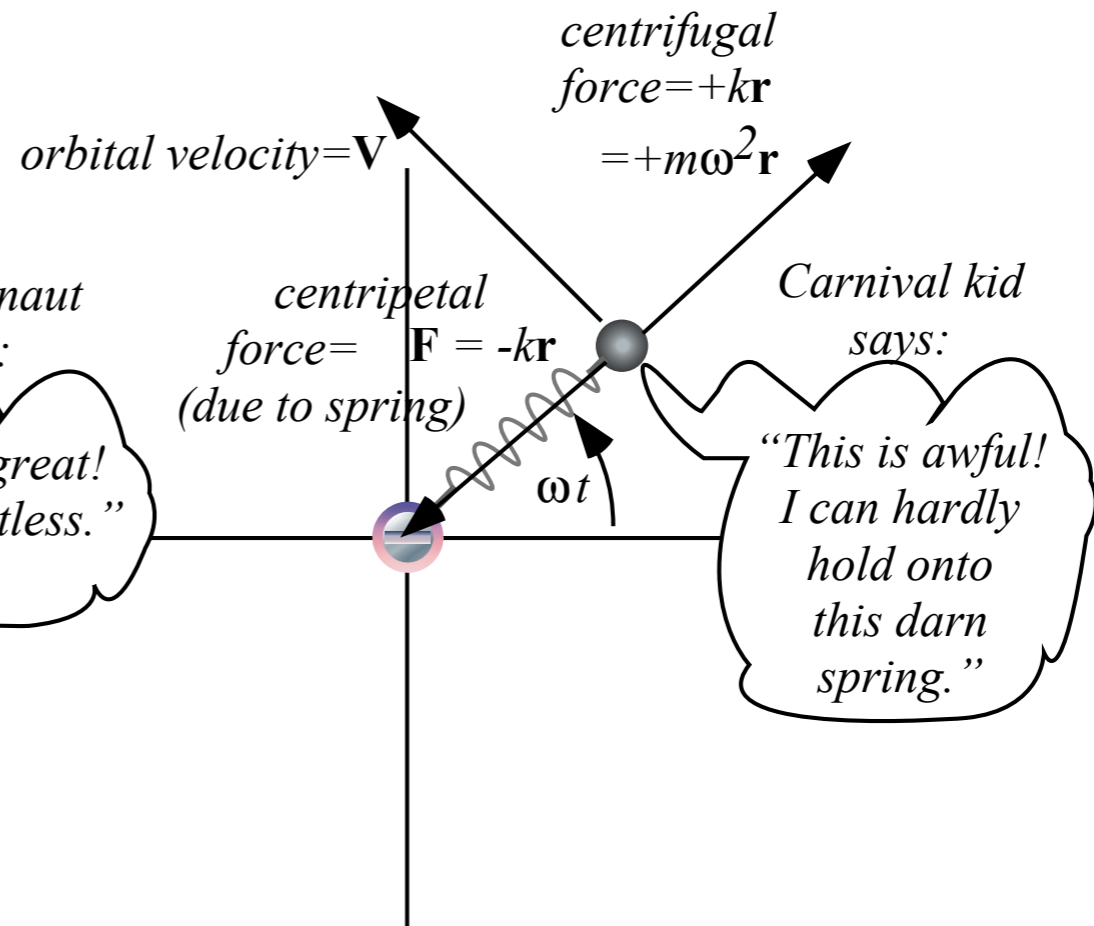
$$E = KE + PE = \frac{1}{2} m v_x^2 + \frac{1}{2} m v_y^2 - \frac{k}{r} = \frac{1}{2} m v_x^2 + \frac{1}{2} m v_y^2 - \frac{GM_{\oplus} m}{r} = -\frac{GM_{\oplus} m}{a}$$

 *A confusing introduction to Coriolis-centrifugal force geometry* (Derived better later)

(a) "Earthronaut" orbiting tunnel inside Earth

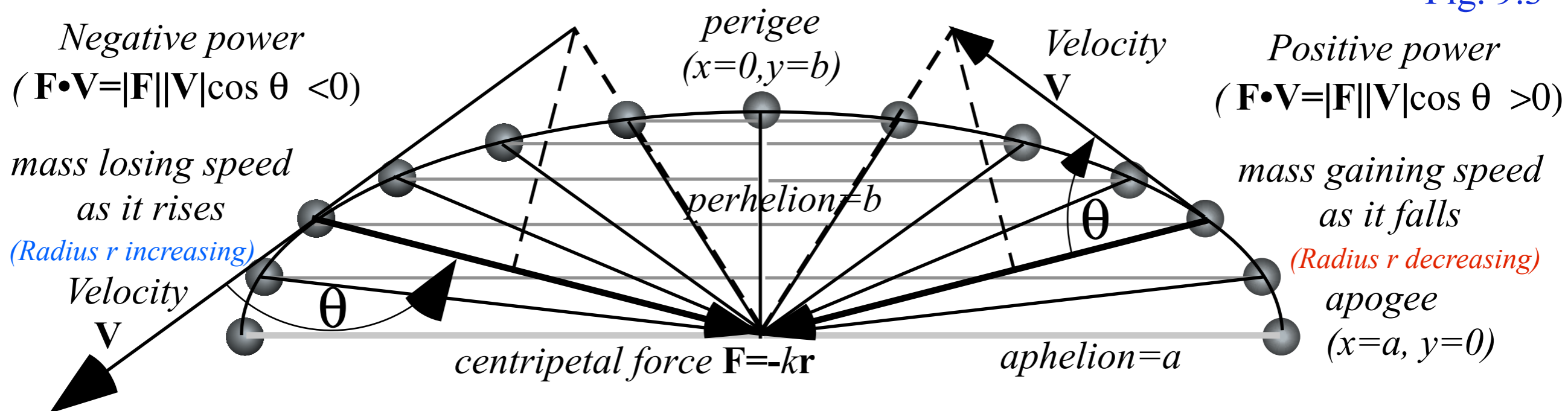


(b) "Carnival kid" orbiting in space attached to a spring



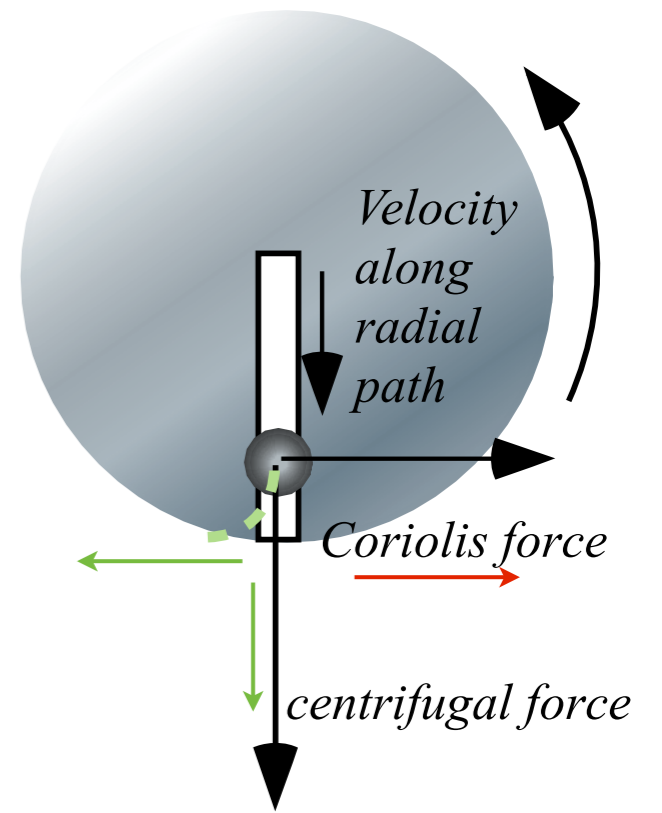
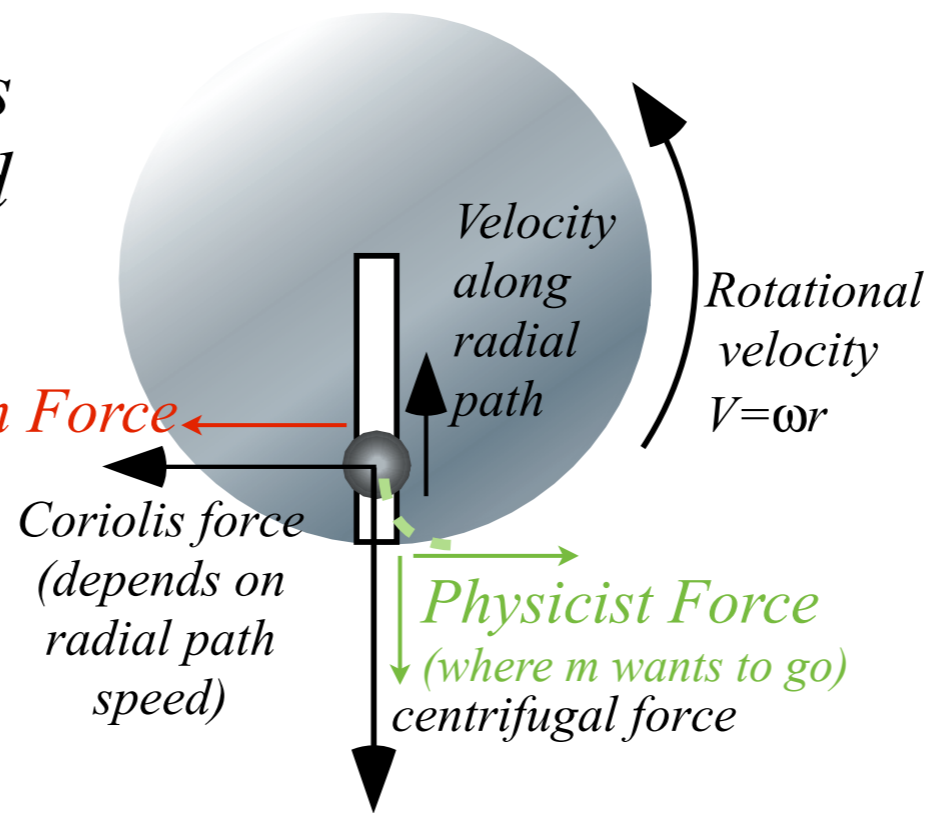
Unit 1
Fig. 9.2

Unit 1
Fig. 9.3



(a) Centrifugal and Coriolis Forces on Merry-Go-Round

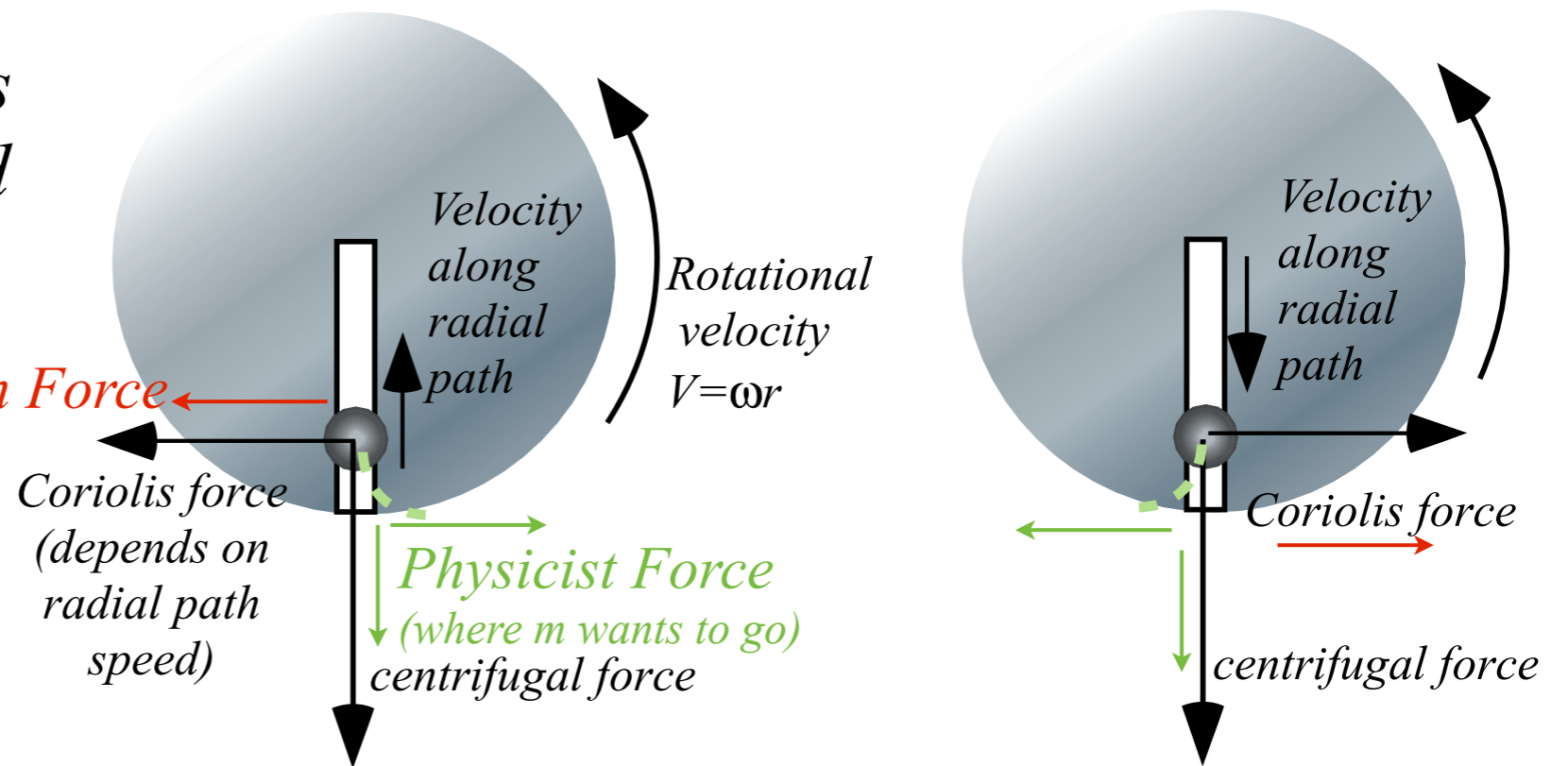
Mathematician Force
(to hold m back)
Constraint force
keeps m in radial slot



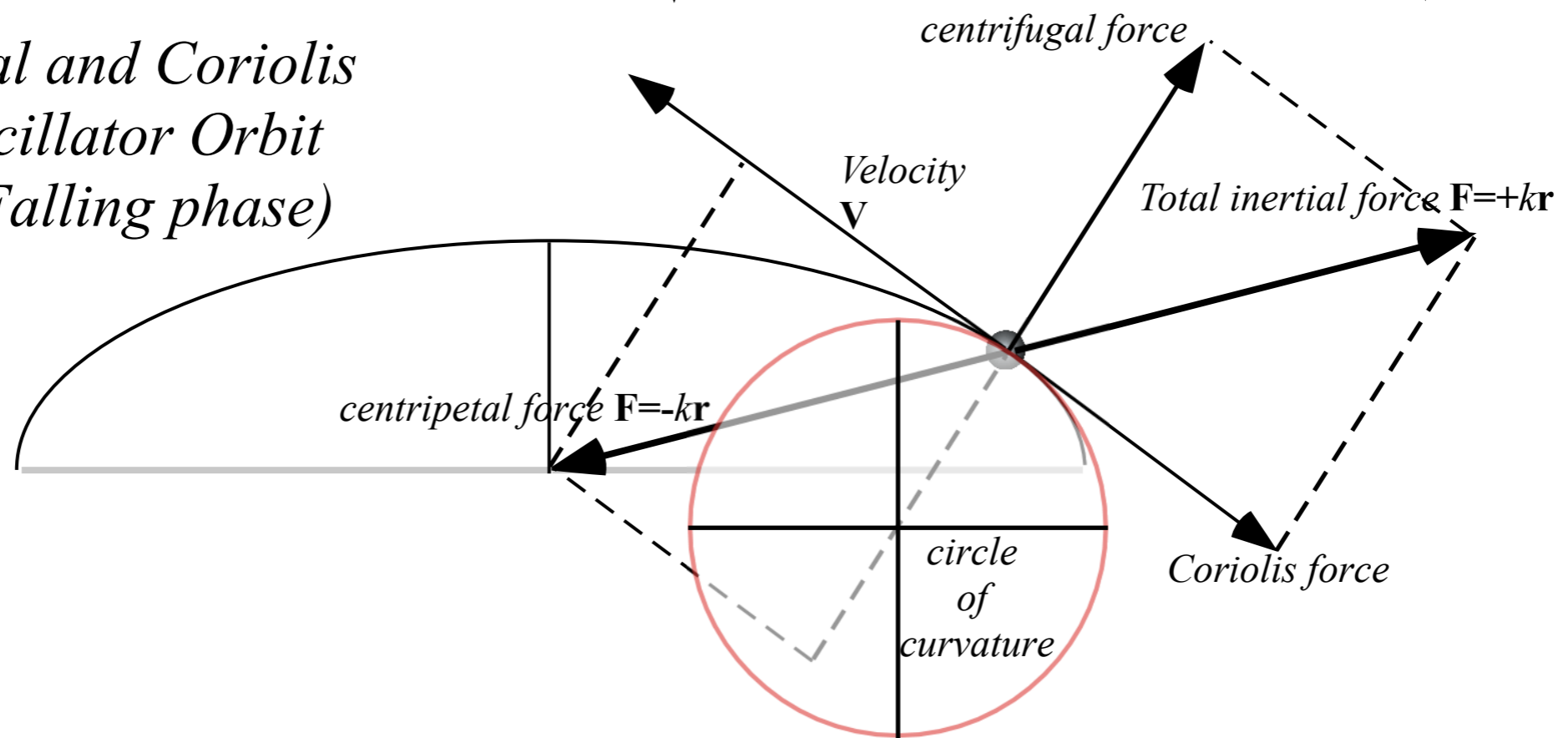
(a) Centrifugal and Coriolis Forces on Merry-Go-Round

Mathematician Force
(to hold m back)

Constraint force
keeps m in radial slot



(b) Centrifugal and Coriolis Forces on Oscillator Orbit (Falling phase)



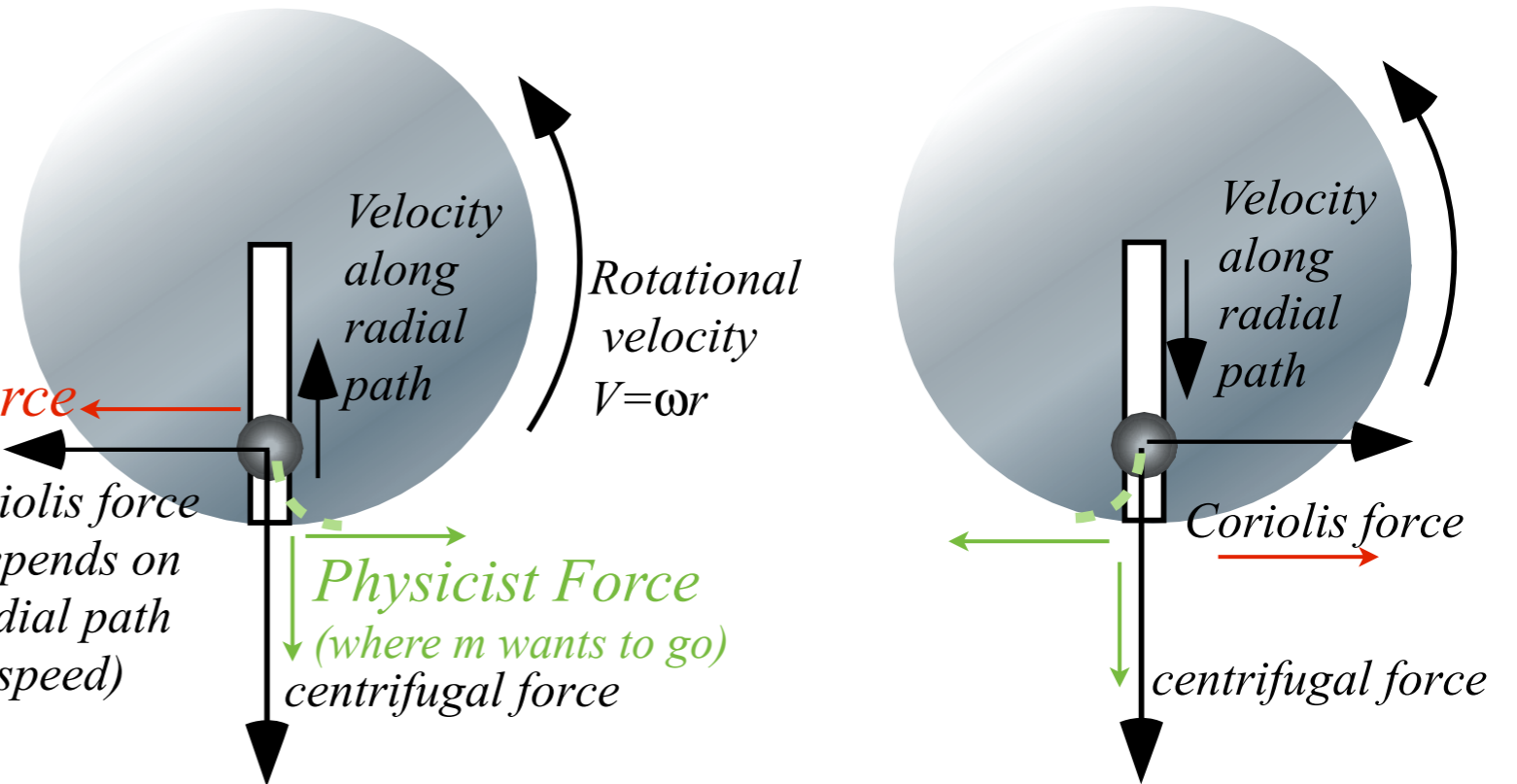
(a) Centrifugal and Coriolis Forces on Merry-Go-Round

Mathematician Force
(to hold m back)

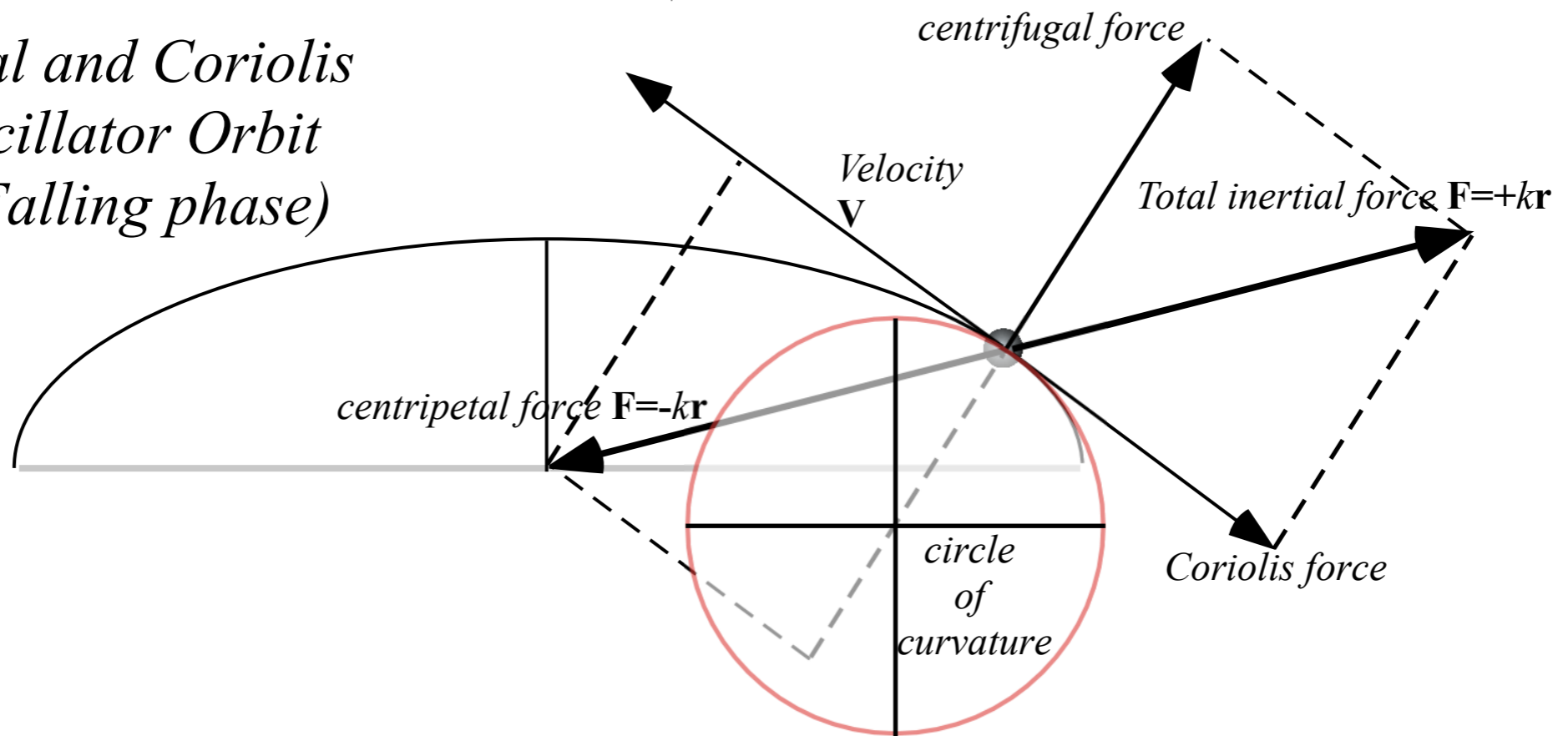
Constraint force
keeps m in radial slot

Coriolis force
(depends on
radial path
speed)

Physicist Force
(where m wants to go)
centrifugal force



(b) Centrifugal and Coriolis Forces on Oscillator Orbit (Falling phase)

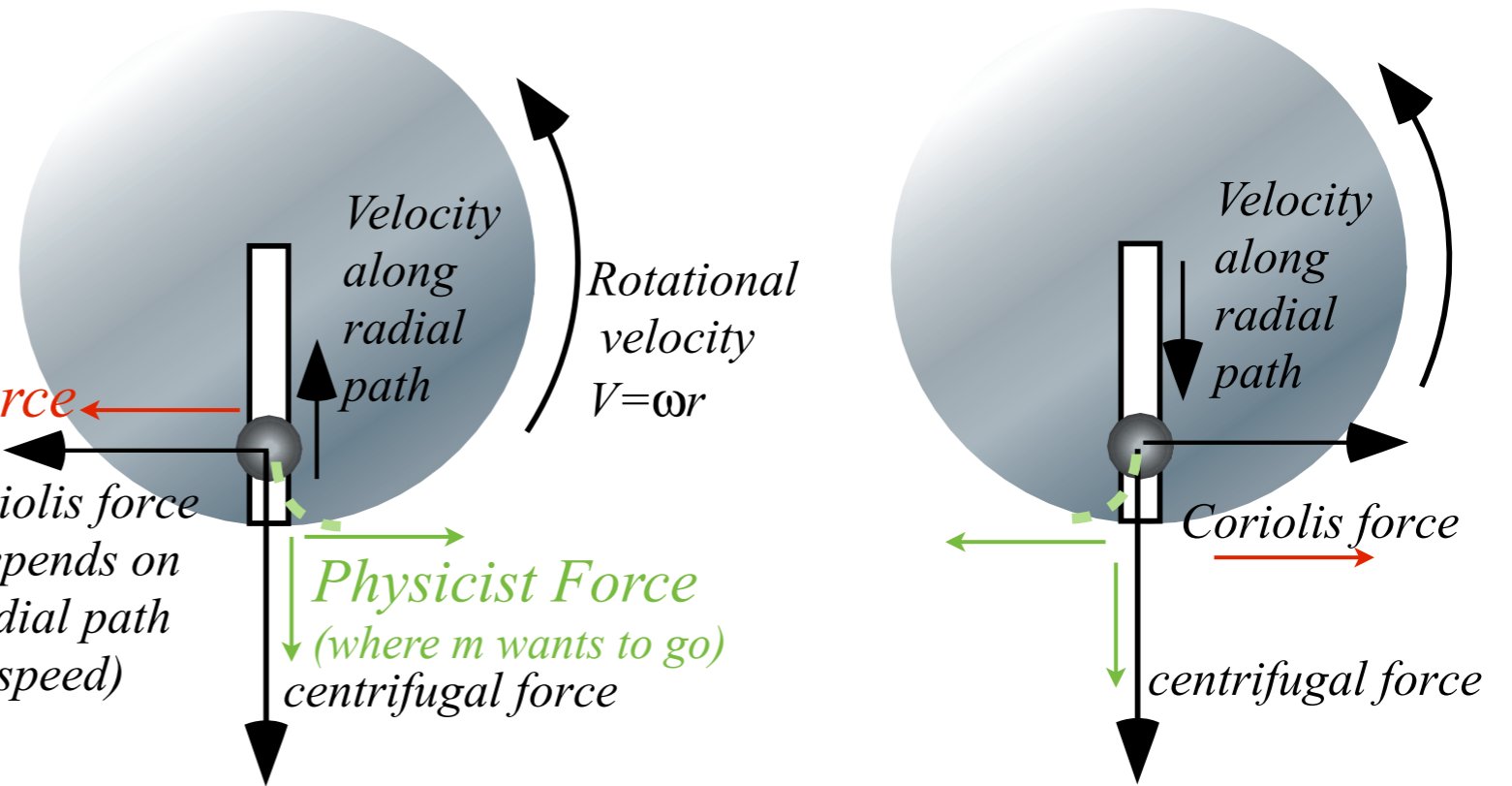


(a) Centrifugal and Coriolis Forces on Merry-Go-Round

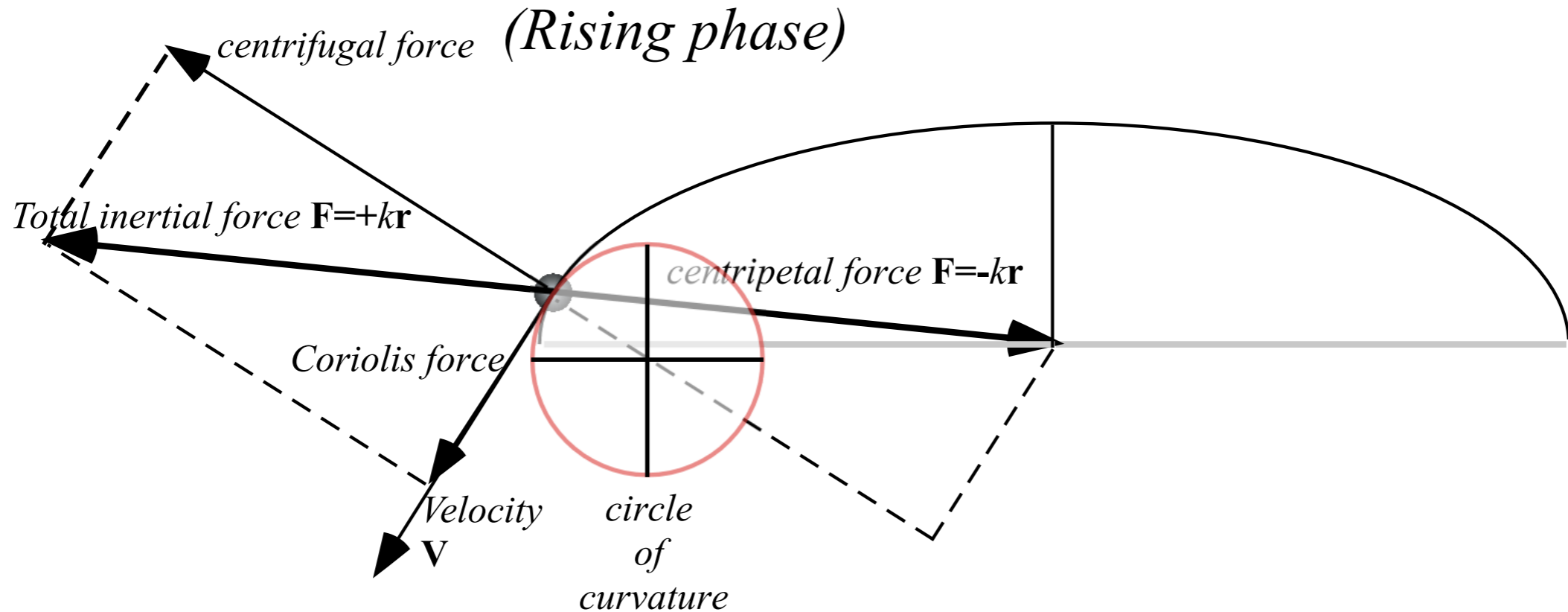
Mathematician Force
(to hold m back)
Constraint force
keeps m in radial slot

Coriolis force
(depends on
radial path
speed)

Physicist Force
(where m wants to go)
centrifugal force



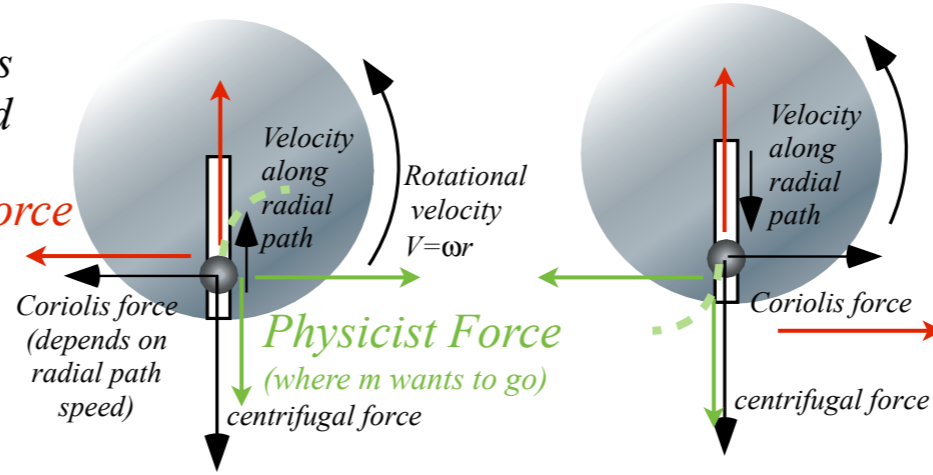
(c) Centrifugal and Coriolis Forces on Oscillator Orbit



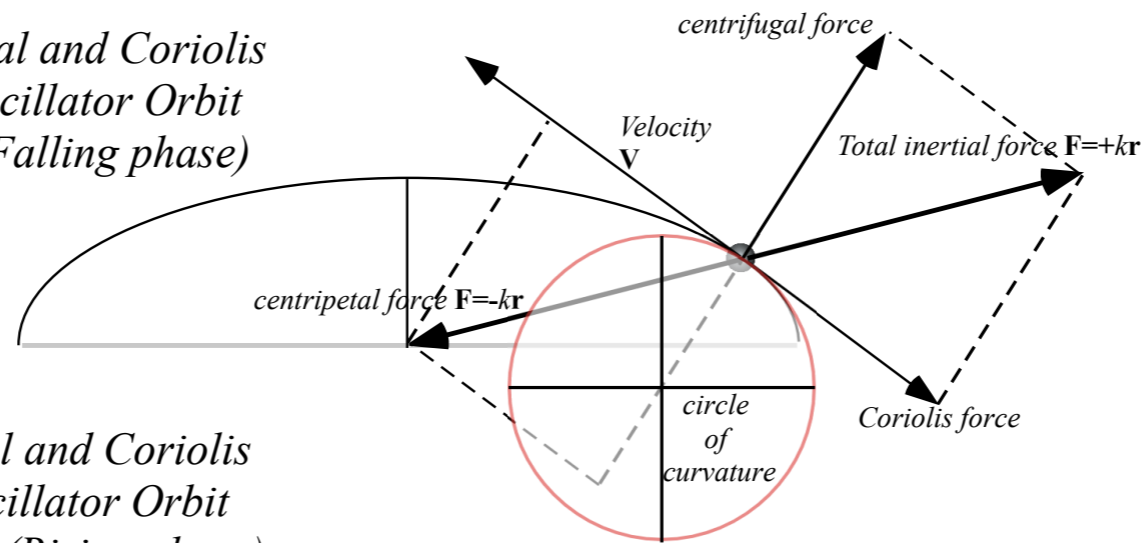
(a) Centrifugal and Coriolis Forces on Merry-Go-Round

Mathematician Force
(to hold m back)

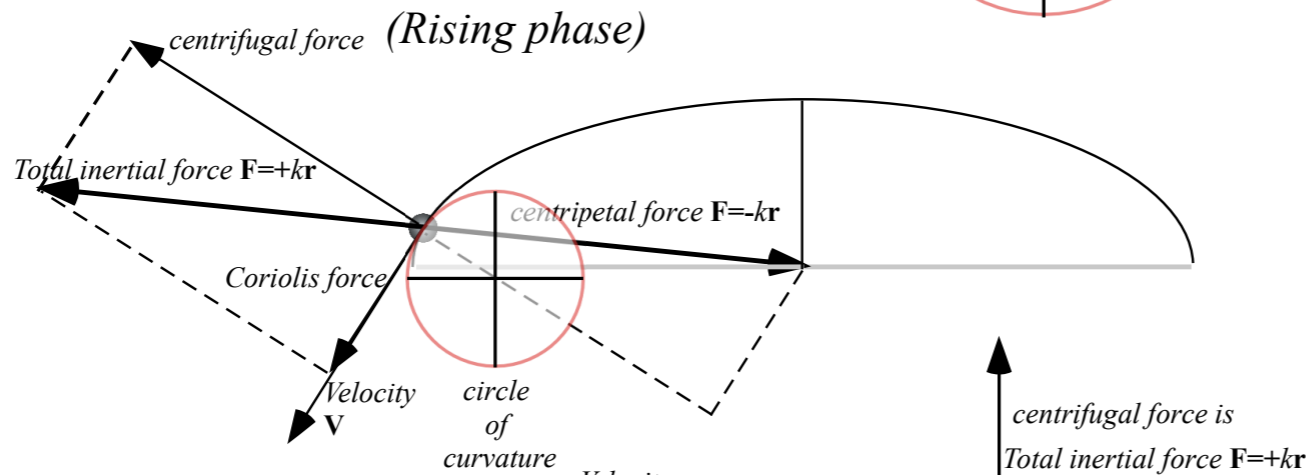
Constraint force
keeps m in radial slot



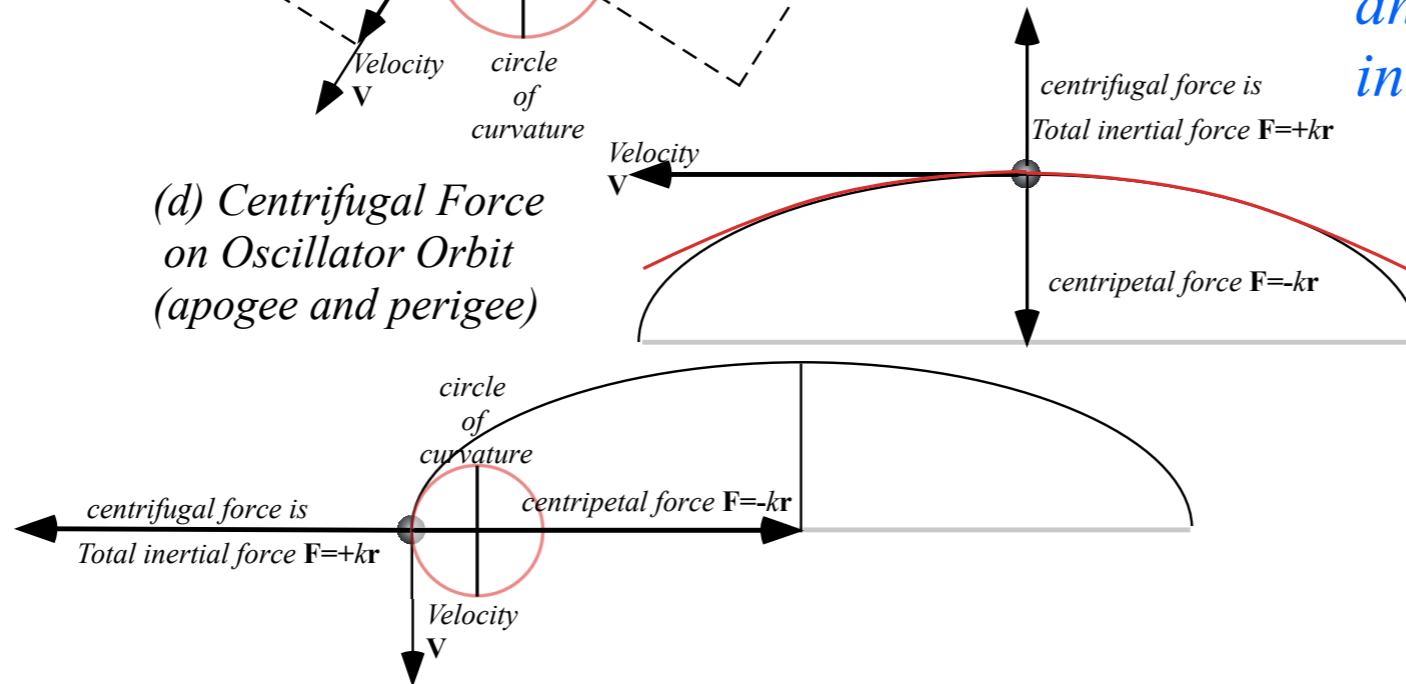
(b) Centrifugal and Coriolis Forces on Oscillator Orbit (Falling phase)



(c) Centrifugal and Coriolis Forces on Oscillator Orbit (Rising phase)



(d) Centrifugal Force on Oscillator Orbit (apogee and perigee)



Unit 1
Fig. 11.4
a-d

*Quite confusing?
Discussion of Coriolis forces will be done more elegantly and made more physically intuitive in Ch. 12 of Unit 1 and in Unit 6.*

→ *Introduction to dual matrix operator contact geometry (based on IHO orbits)*
Quadratic form ellipse $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = 1$ vs. inverse form ellipse $\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p} = 1$
Duality norm relations ($\mathbf{r} \cdot \mathbf{p} = 1$)
 \mathbf{Q} -Ellipse tangents \mathbf{r}' normal to dual \mathbf{Q}^{-1} -ellipse position \mathbf{p} ($\mathbf{r}' \cdot \mathbf{p} = 0 = \mathbf{r} \cdot \mathbf{p}'$)
Operator geometric sequences and eigenvectors
Alternative scaling of matrix operator geometry
Vector calculus of tensor operation

Quadratic forms and tangent contact geometry of their ellipses

A matrix Q that generates an ellipse by $\mathbf{r} \bullet Q \bullet \mathbf{r} = 1$ is called positive-definite (if $\mathbf{r} \bullet Q \bullet \mathbf{r}$ always > 0)

$$\begin{pmatrix} x & y \end{pmatrix} \bullet \overbrace{\begin{pmatrix} \frac{1}{a^2} & 0 \\ 0 & \frac{1}{b^2} \end{pmatrix}}^{\mathbf{r} \bullet Q \bullet \mathbf{r}} \bullet \begin{pmatrix} x \\ y \end{pmatrix} = 1 = \overbrace{\begin{pmatrix} x & y \end{pmatrix}}^{\mathbf{r}} \bullet \overbrace{\begin{pmatrix} \frac{x}{a^2} \\ \frac{y}{b^2} \end{pmatrix}}^{Q \bullet \mathbf{r}} = \frac{x^2}{a^2} + \frac{y^2}{b^2}$$

A inverse matrix Q^{-1} generates an ellipse by $\mathbf{p} \bullet Q^{-1} \bullet \mathbf{p} = 1$ called inverse or dual ellipse:

$$\begin{pmatrix} p_x & p_y \end{pmatrix} \bullet \overbrace{\begin{pmatrix} a^2 & 0 \\ 0 & b^2 \end{pmatrix}}^{\mathbf{p} \bullet Q^{-1} \bullet \mathbf{p}} \bullet \begin{pmatrix} p_x \\ p_y \end{pmatrix} = 1 = \overbrace{\begin{pmatrix} p_x & p_y \end{pmatrix}}^{\mathbf{p}} \bullet \overbrace{\begin{pmatrix} a^2 p_x \\ b^2 p_y \end{pmatrix}}^{Q^{-1} \bullet \mathbf{p}} = a^2 p_x^2 + b^2 p_y^2$$

Quadratic forms and tangent contact geometry of their ellipses

A matrix Q that generates an ellipse by $\mathbf{r} \bullet Q \bullet \mathbf{r} = 1$ is called positive-definite (if $\mathbf{r} \bullet Q \bullet \mathbf{r}$ always > 0)

$$\left(\begin{array}{cc} x & y \end{array} \right) \bullet \overbrace{\begin{pmatrix} \frac{1}{a^2} & 0 \\ 0 & \frac{1}{b^2} \end{pmatrix}}^{\mathbf{r} \bullet Q \bullet \mathbf{r}} \bullet \begin{pmatrix} x \\ y \end{pmatrix} = 1 = \overbrace{\begin{pmatrix} x & y \end{pmatrix}}^{\mathbf{r}} \bullet \overbrace{\begin{pmatrix} \frac{x}{a^2} \\ \frac{y}{b^2} \end{pmatrix}}^{Q \bullet \mathbf{r} = \mathbf{p}} = \frac{x^2}{a^2} + \frac{y^2}{b^2}$$

Defined mapping between ellipses

A inverse matrix Q^{-1} generates an ellipse by $\mathbf{p} \bullet Q^{-1} \bullet \mathbf{p} = 1$ called inverse or dual ellipse:

$$\left(\begin{array}{cc} p_x & p_y \end{array} \right) \bullet \overbrace{\begin{pmatrix} a^2 & 0 \\ 0 & b^2 \end{pmatrix}}^{\mathbf{p} \bullet Q^{-1} \bullet \mathbf{p}} \bullet \begin{pmatrix} p_x \\ p_y \end{pmatrix} = 1 = \overbrace{\begin{pmatrix} p_x & p_y \end{pmatrix}}^{\mathbf{p}} \bullet \overbrace{\begin{pmatrix} a^2 p_x \\ b^2 p_y \end{pmatrix}}^{Q^{-1} \bullet \mathbf{p} = \mathbf{r}} = a^2 p_x^2 + b^2 p_y^2$$

Introduction to dual matrix operator contact geometry (based on IHO orbits)

→ *Quadratic form ellipse $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = 1$ vs. inverse form ellipse $\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p} = 1$*

Duality norm relations ($\mathbf{r} \cdot \mathbf{p} = 1$)

\mathbf{Q} -Ellipse tangents \mathbf{r}' normal to dual \mathbf{Q}^{-1} -ellipse position \mathbf{p} ($\mathbf{r}' \cdot \mathbf{p} = 0 = \mathbf{r} \cdot \mathbf{p}'$)

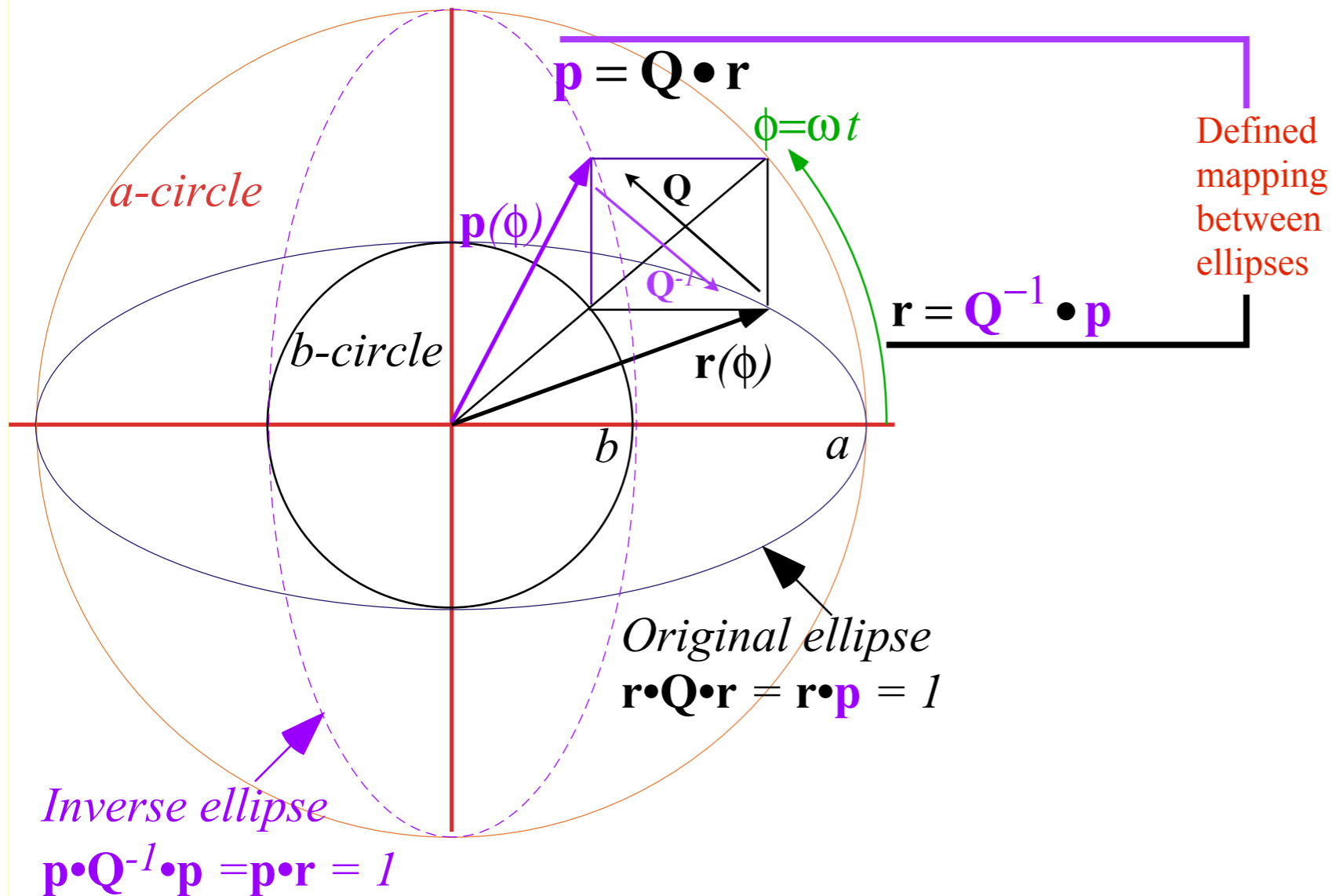
Operator geometric sequences and eigenvectors

Alternative scaling of matrix operator geometry

Vector calculus of tensor operation

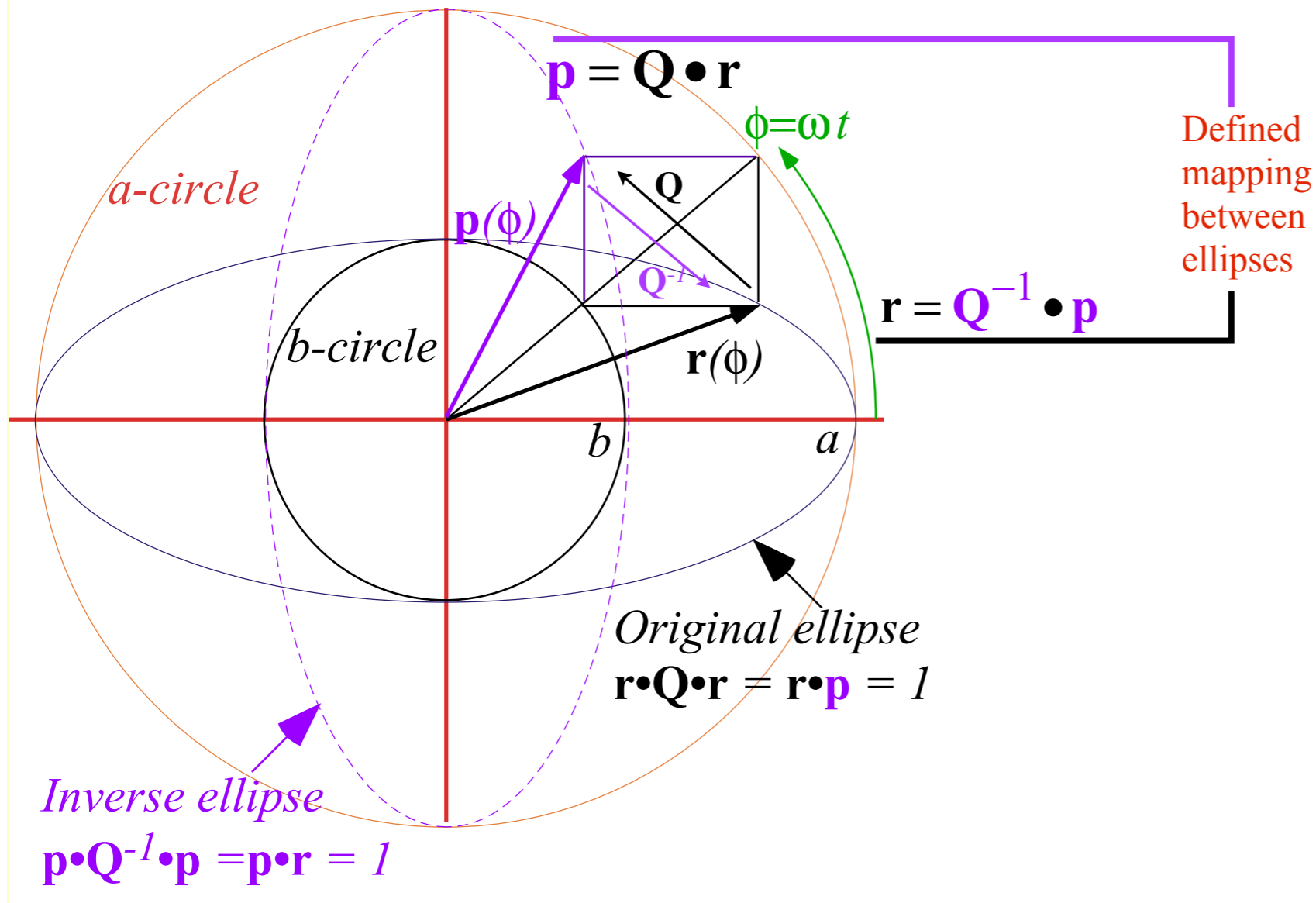
(a) Quadratic form ellipse and
Inverse quadratic form ellipse

based on
Unit 1
Fig. 11.6



(a) Quadratic form ellipse and
Inverse quadratic form ellipse

based on
Unit 1
Fig. 11.6



Here plot of \mathbf{p} -ellipse is re-scaled by scalefactor $S = a \cdot b$

\mathbf{p} -ellipse x -radius = $1/a$ plotted at: $S(1/a) = b$ ($=1$ for $a=2, b=1$)

\mathbf{p} -ellipse y -radius = $1/b$ plotted at: $S(1/b) = a$ ($=2$ for $a=2, b=1$)

Introduction to dual matrix operator geometry (based on IHO orbits)

Quadratic form ellipse $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = 1$ vs. inverse form ellipse $\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p} = 1$



Duality norm relations ($\mathbf{r} \cdot \mathbf{p} = 1$)

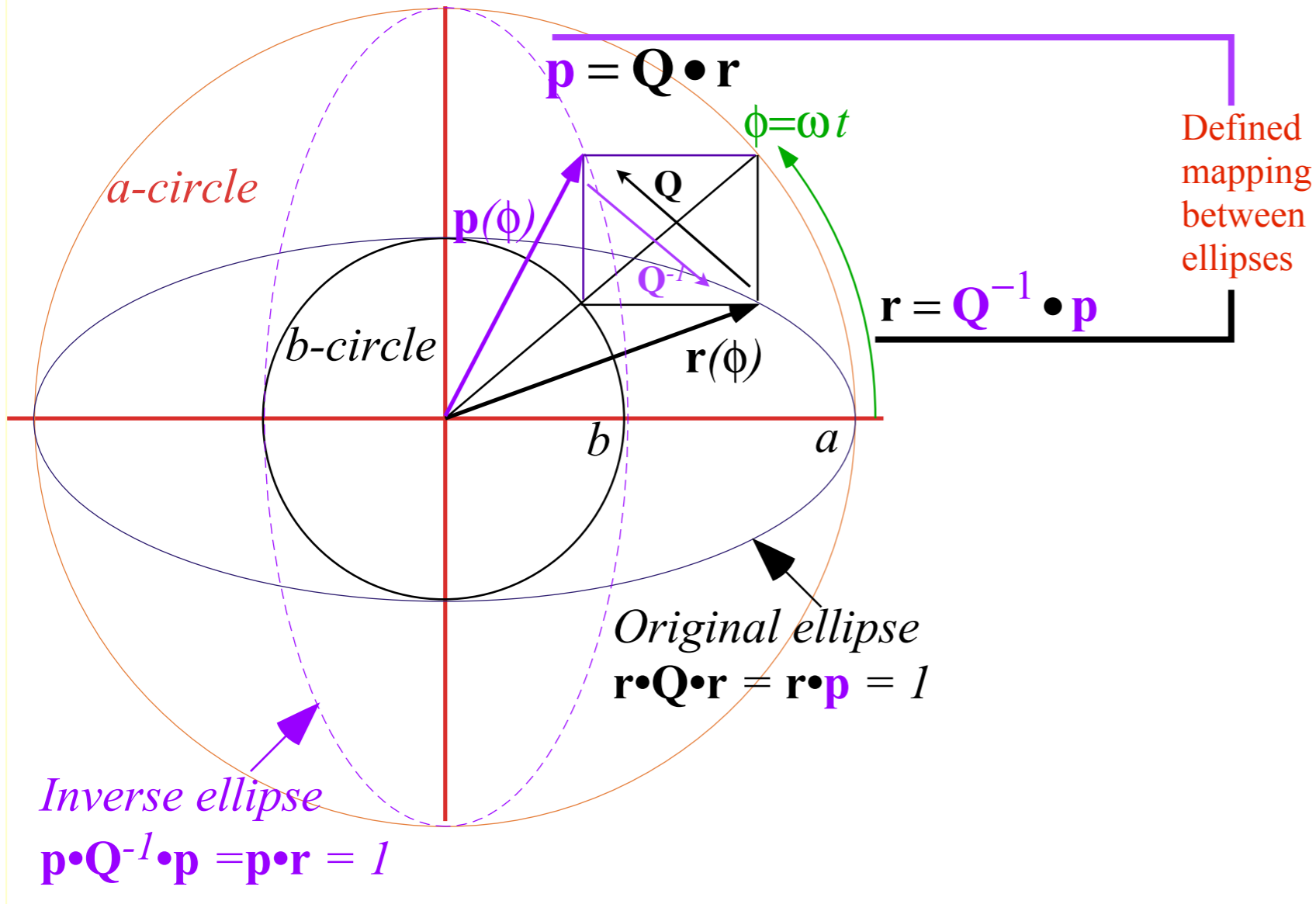
\mathbf{Q} -Ellipse tangents \mathbf{r}' normal to dual \mathbf{Q}^{-1} -ellipse position \mathbf{p} ($\mathbf{r}' \cdot \mathbf{p} = 0 = \mathbf{r} \cdot \mathbf{p}'$)

Operator geometric sequences and eigenvectors

Alternative scaling of matrix operator geometry

Vector calculus of tensor operation

(a) Quadratic form ellipse and
Inverse quadratic form ellipse



Quadratic form $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = 1$ has mutual duality relations with inverse form $\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p} = 1 = \mathbf{p} \cdot \mathbf{r}$

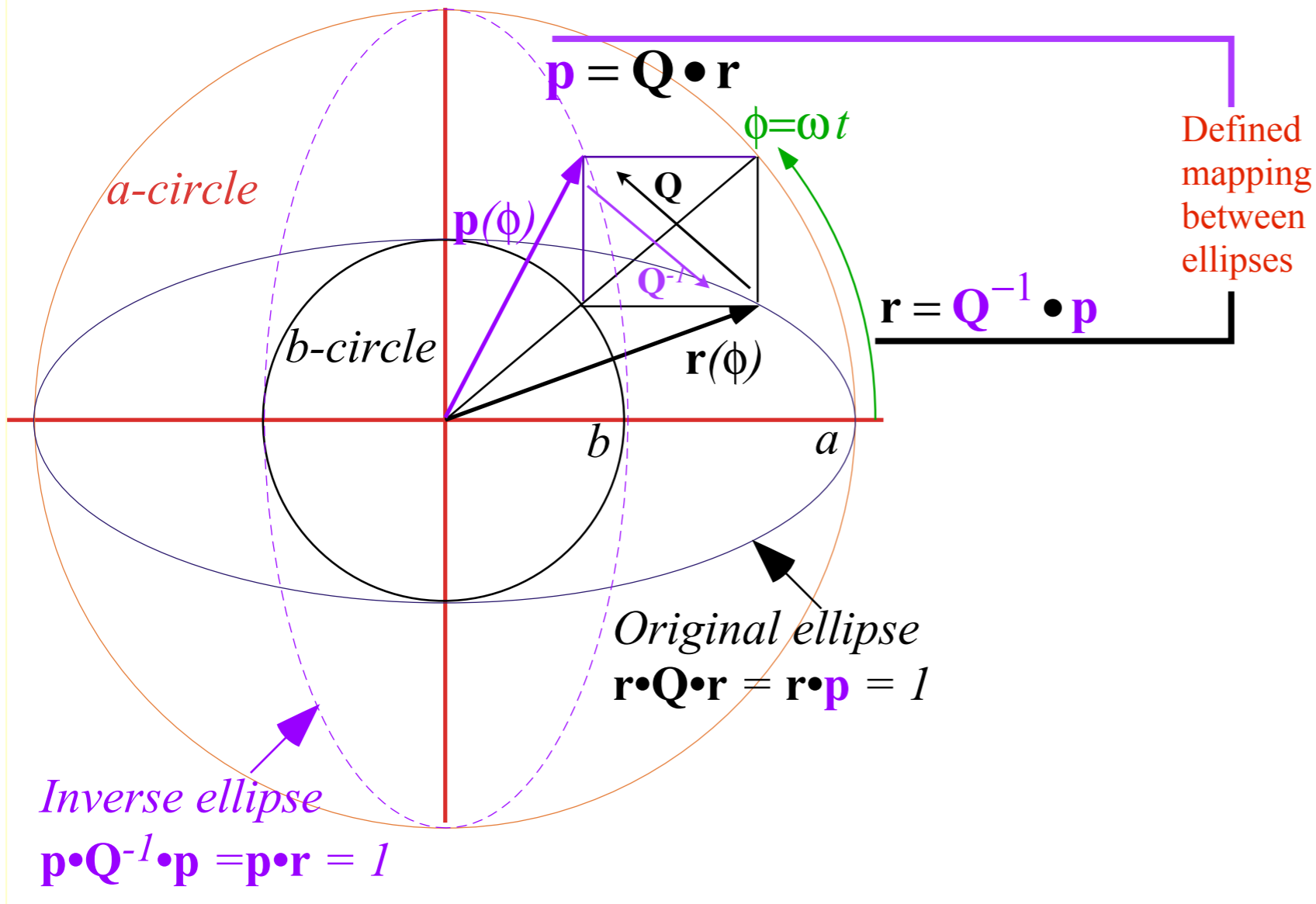
Here plot of \mathbf{p} -ellipse is re-scaled by scalefactor $S = a \cdot b$

\mathbf{p} -ellipse x -radius = $1/a$ plotted at: $S(1/a) = b$ ($=1$ for $a=2, b=1$)

\mathbf{p} -ellipse y -radius = $1/b$ plotted at: $S(1/b) = a$ ($=2$ for $a=2, b=1$)

(a) Quadratic form ellipse and
Inverse quadratic form ellipse

based on
Unit 1
Fig. 11.6



Quadratic form $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = 1$ has mutual duality relations with inverse form $\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p} = 1 = \mathbf{p} \cdot \mathbf{r}$

$$\mathbf{p} = \mathbf{Q} \cdot \mathbf{r} = \begin{pmatrix} \overbrace{1/a^2}^{\mathbf{Q}} & 0 \\ 0 & \overbrace{1/b^2}^{\mathbf{Q}} \end{pmatrix} \cdot \begin{pmatrix} \overbrace{x}^{\mathbf{r}} \\ \overbrace{y}^{\mathbf{r}} \end{pmatrix} = \begin{pmatrix} \overbrace{x/a^2}^{\mathbf{p}} \\ \overbrace{y/b^2}^{\mathbf{p}} \end{pmatrix} = \begin{pmatrix} (1/a)\cos\phi \\ (1/b)\sin\phi \end{pmatrix} \quad \text{where:} \quad \begin{matrix} x = r_x = a \cos\phi = a \cos\omega t \\ y = r_y = b \sin\phi = b \sin\omega t \end{matrix} \quad \text{so: } \boxed{\mathbf{p} \cdot \mathbf{r} = 1}$$

Here plot of \mathbf{p} -ellipse is re-scaled by scalefactor $S = a \cdot b$

\mathbf{p} -ellipse x -radius $= 1/a$ plotted at: $S(1/a) = b$ ($= 1$ for $a=2, b=1$)

\mathbf{p} -ellipse y -radius $= 1/b$ plotted at: $S(1/b) = a$ ($= 2$ for $a=2, b=1$)

[Link ⇒ BoxIt simulation of IHO orbits](#)
[Link → IHO orbital time rates of change](#)
[Link → IHO Exegesis Plot](#)

Introduction to dual matrix operator geometry (based on IHO orbits)

Quadratic form ellipse $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = 1$ vs. inverse form ellipse $\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p} = 1$

Duality norm relations ($\mathbf{r} \cdot \mathbf{p} = 1$)

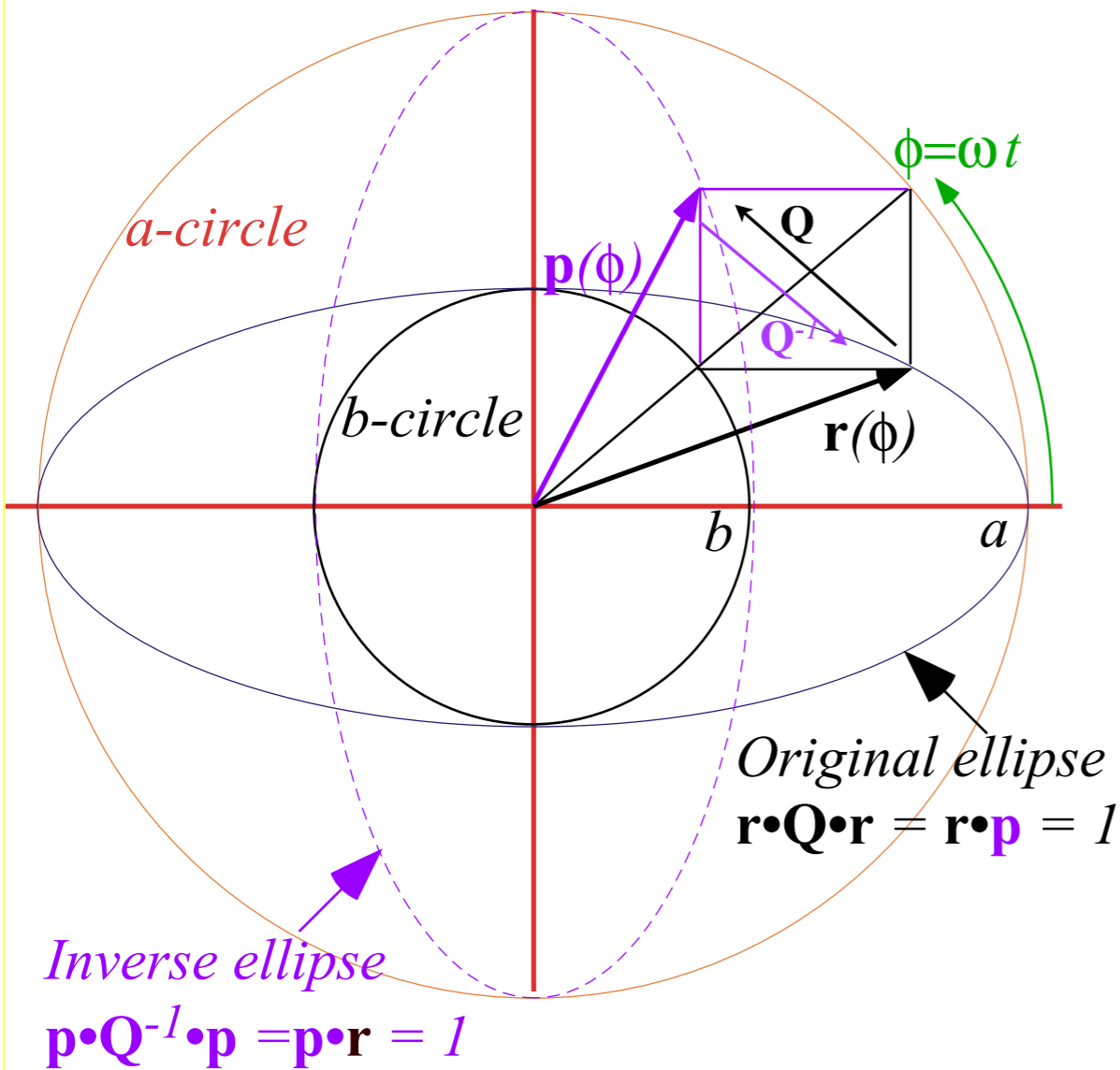
 *\mathbf{Q} -Ellipse tangents \mathbf{r}' normal to dual \mathbf{Q}^{-1} -ellipse position \mathbf{p} ($\mathbf{r}' \cdot \mathbf{p} = 0 = \mathbf{r} \cdot \mathbf{p}'$)*

Operator geometric sequences and eigenvectors

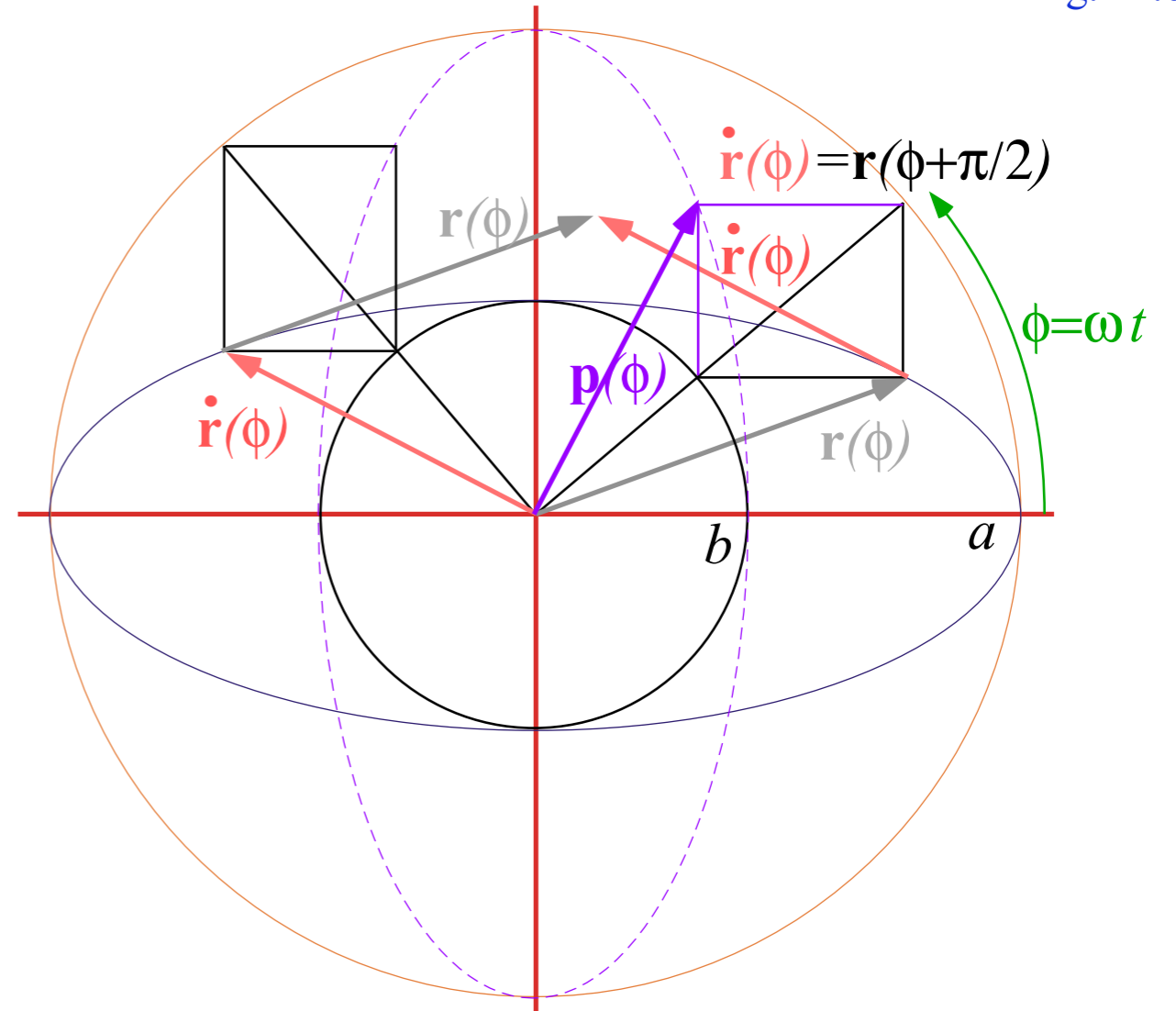
Alternative scaling of matrix operator geometry

Vector calculus of tensor operation

(a) Quadratic form ellipse and Inverse quadratic form ellipse



(b) Ellipse tangents



based on
Unit 1
Fig. 11.6

Quadratic form $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = 1$ has mutual duality relations with inverse form $\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p} = 1 = \mathbf{p} \cdot \mathbf{r}$

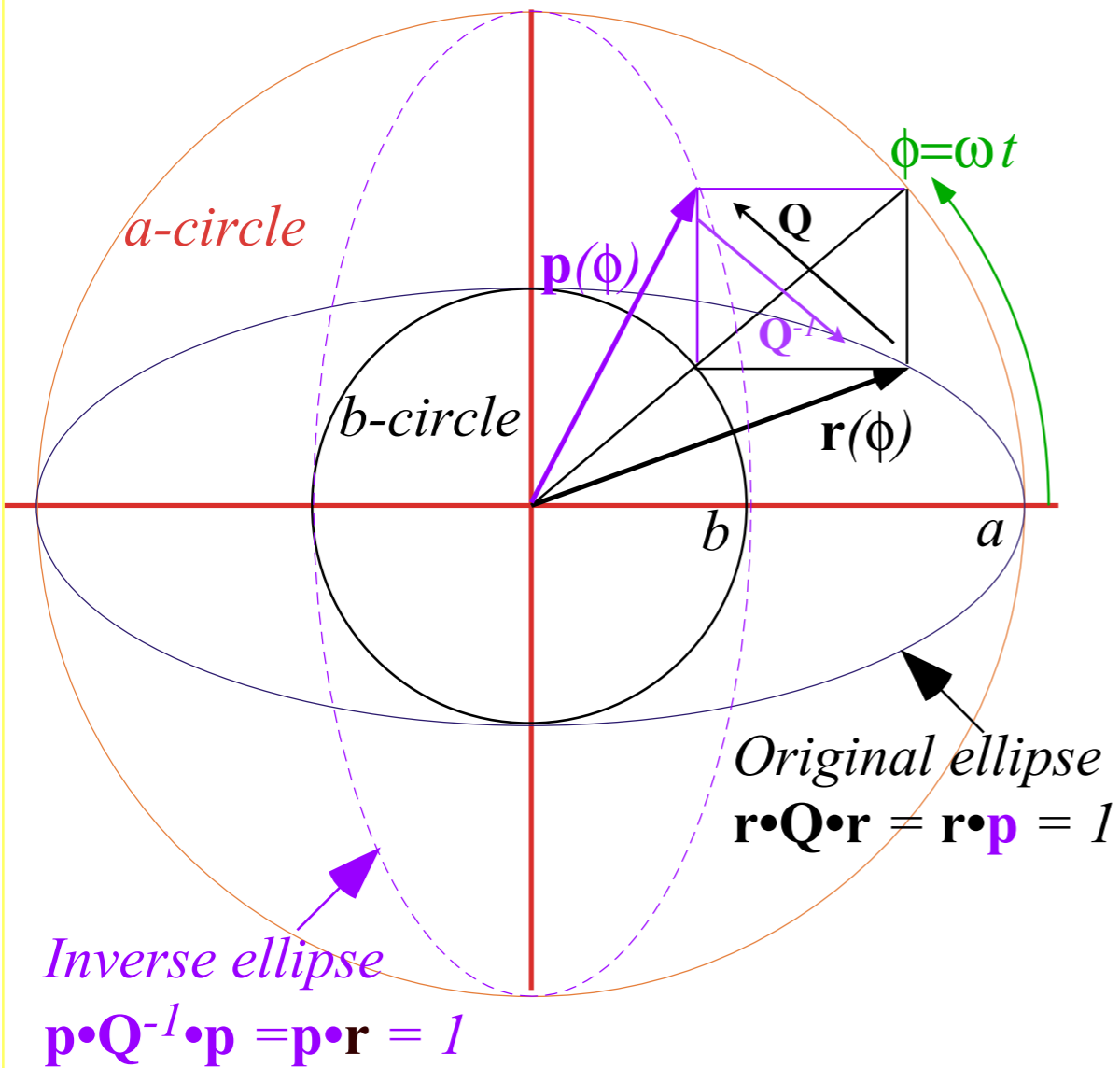
$$\mathbf{p} = \mathbf{Q} \cdot \mathbf{r} = \begin{pmatrix} 1/a^2 & 0 \\ 0 & 1/b^2 \end{pmatrix} \cdot \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a^2 \\ y/b^2 \end{pmatrix} = \begin{pmatrix} (1/a)\cos\phi \\ (1/b)\sin\phi \end{pmatrix} \quad \text{where:} \quad \begin{matrix} x = r_x = a \cos\phi = a \cos\omega t \\ y = r_y = b \sin\phi = b \sin\omega t \end{matrix} \quad \text{so: } \boxed{\mathbf{p} \cdot \mathbf{r} = 1}$$

Here plot of \mathbf{p} -ellipse is re-scaled by scalefactor $S = a \cdot b$

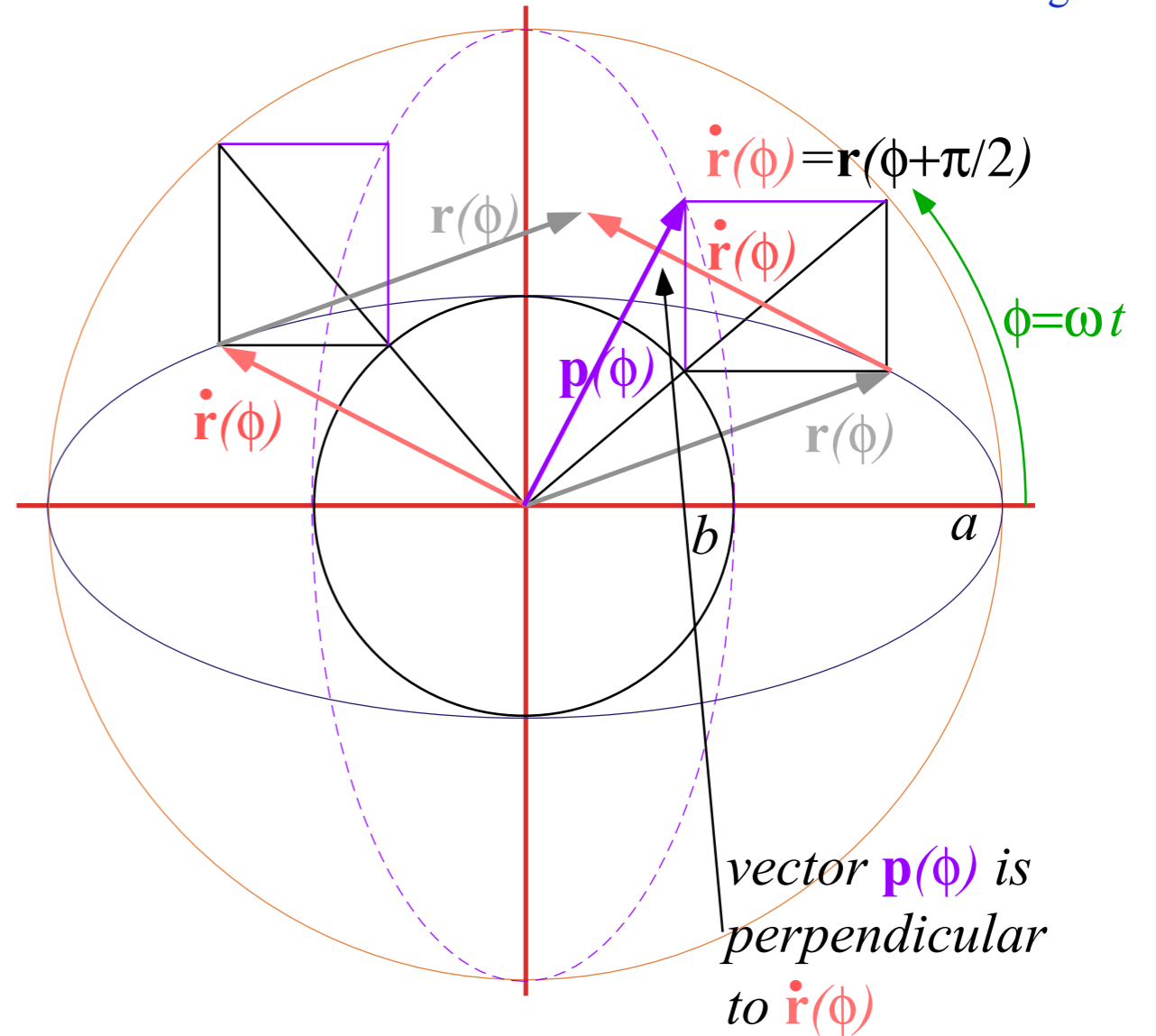
\mathbf{p} -ellipse x -radius $= 1/a$ plotted at: $S(1/a) = b$ ($= 1$ for $a=2, b=1$)

\mathbf{p} -ellipse y -radius $= 1/b$ plotted at: $S(1/b) = a$ ($= 2$ for $a=2, b=1$)

(a) Quadratic form ellipse and Inverse quadratic form ellipse



(b) Ellipse tangents



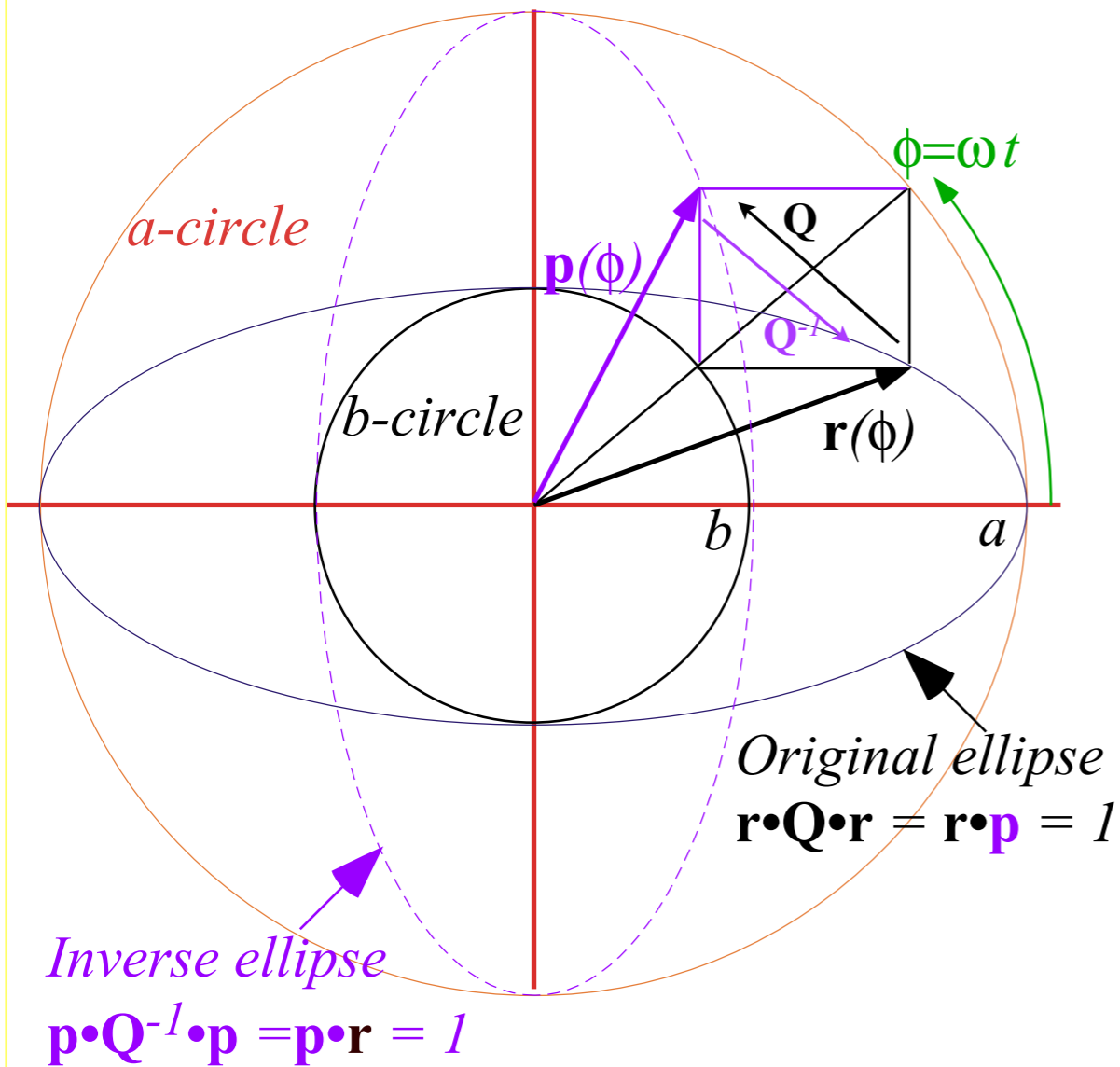
Quadratic form $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = 1$ has mutual duality relations with inverse form $\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p} = 1 = \mathbf{p} \cdot \mathbf{r}$

$$\mathbf{p} = \mathbf{Q} \cdot \mathbf{r} = \begin{pmatrix} 1/a^2 & 0 \\ 0 & 1/b^2 \end{pmatrix} \cdot \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a^2 \\ y/b^2 \end{pmatrix} = \begin{pmatrix} (1/a)\cos\phi \\ (1/b)\sin\phi \end{pmatrix} \quad \text{where:} \quad \begin{aligned} x = r_x &= a \cos\phi = a \cos\omega t \\ y = r_y &= b \sin\phi = b \sin\omega t \end{aligned} \quad \text{so: } \mathbf{p} \cdot \mathbf{r} = 1$$

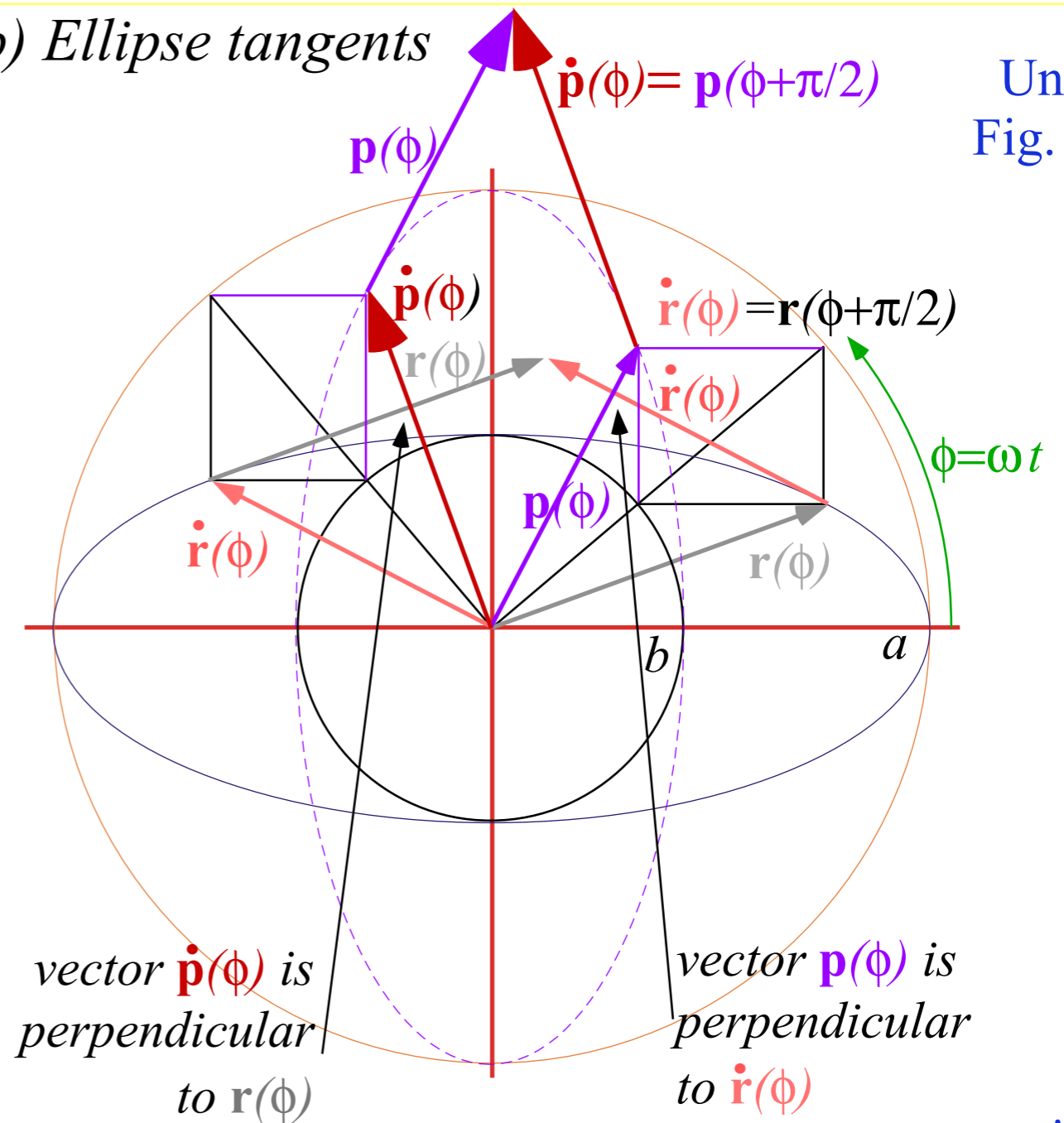
\mathbf{p} is perpendicular to velocity $\mathbf{v} = \dot{\mathbf{r}}$, a mutual orthogonality

$$\dot{\mathbf{r}} \cdot \mathbf{p} = 0 = \begin{pmatrix} \dot{r}_x & \dot{r}_y \end{pmatrix} \cdot \begin{pmatrix} p_x \\ p_y \end{pmatrix} = \begin{pmatrix} -a \sin\phi & b \cos\phi \end{pmatrix} \cdot \begin{pmatrix} (1/a)\cos\phi \\ (1/b)\sin\phi \end{pmatrix} \quad \text{where:} \quad \begin{aligned} \dot{r}_x &= -a \sin\phi \\ \dot{r}_y &= b \cos\phi \end{aligned} \quad \text{and:} \quad \begin{aligned} p_x &= (1/a)\cos\phi \\ p_y &= (1/b)\sin\phi \end{aligned}$$

(a) Quadratic form ellipse and Inverse quadratic form ellipse



(b) Ellipse tangents



Quadratic form $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = 1$ has mutual duality relations with inverse form $\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p} = 1$

unit mutual projection

$$\mathbf{p} = \mathbf{Q} \cdot \mathbf{r} = \begin{pmatrix} 1/a^2 & 0 \\ 0 & 1/b^2 \end{pmatrix} \cdot \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a^2 \\ y/b^2 \end{pmatrix} = \begin{pmatrix} (1/a)\cos\phi \\ (1/b)\sin\phi \end{pmatrix} \text{ where: } \begin{matrix} x = r_x = a \cos\phi = a \cos\omega t \\ y = r_y = b \sin\phi = b \sin\omega t \end{matrix} \text{ so: } \boxed{\mathbf{p} \cdot \mathbf{r} = 1}$$

\mathbf{p} is perpendicular to velocity $\mathbf{v} = \dot{\mathbf{r}}$, a mutual orthogonality. So is \mathbf{r} perpendicular to $\dot{\mathbf{p}}$: $\boxed{\dot{\mathbf{p}} \cdot \mathbf{r} = 0}$

$$\boxed{\dot{\mathbf{r}} \cdot \mathbf{p} = 0} = \begin{pmatrix} \dot{r}_x & \dot{r}_y \end{pmatrix} \cdot \begin{pmatrix} p_x \\ p_y \end{pmatrix} = \begin{pmatrix} -a \sin\phi & b \cos\phi \end{pmatrix} \cdot \begin{pmatrix} (1/a)\cos\phi \\ (1/b)\sin\phi \end{pmatrix} \text{ where: } \begin{matrix} \dot{r}_x = -a \sin\phi \\ \dot{r}_y = b \cos\phi \end{matrix} \text{ and: } \begin{matrix} p_x = (1/a)\cos\phi \\ p_y = (1/b)\sin\phi \end{matrix}$$

Introduction to dual matrix operator geometry (based on IHO orbits)

Quadratic form ellipse $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = 1$ vs. inverse form ellipse $\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p} = 1$

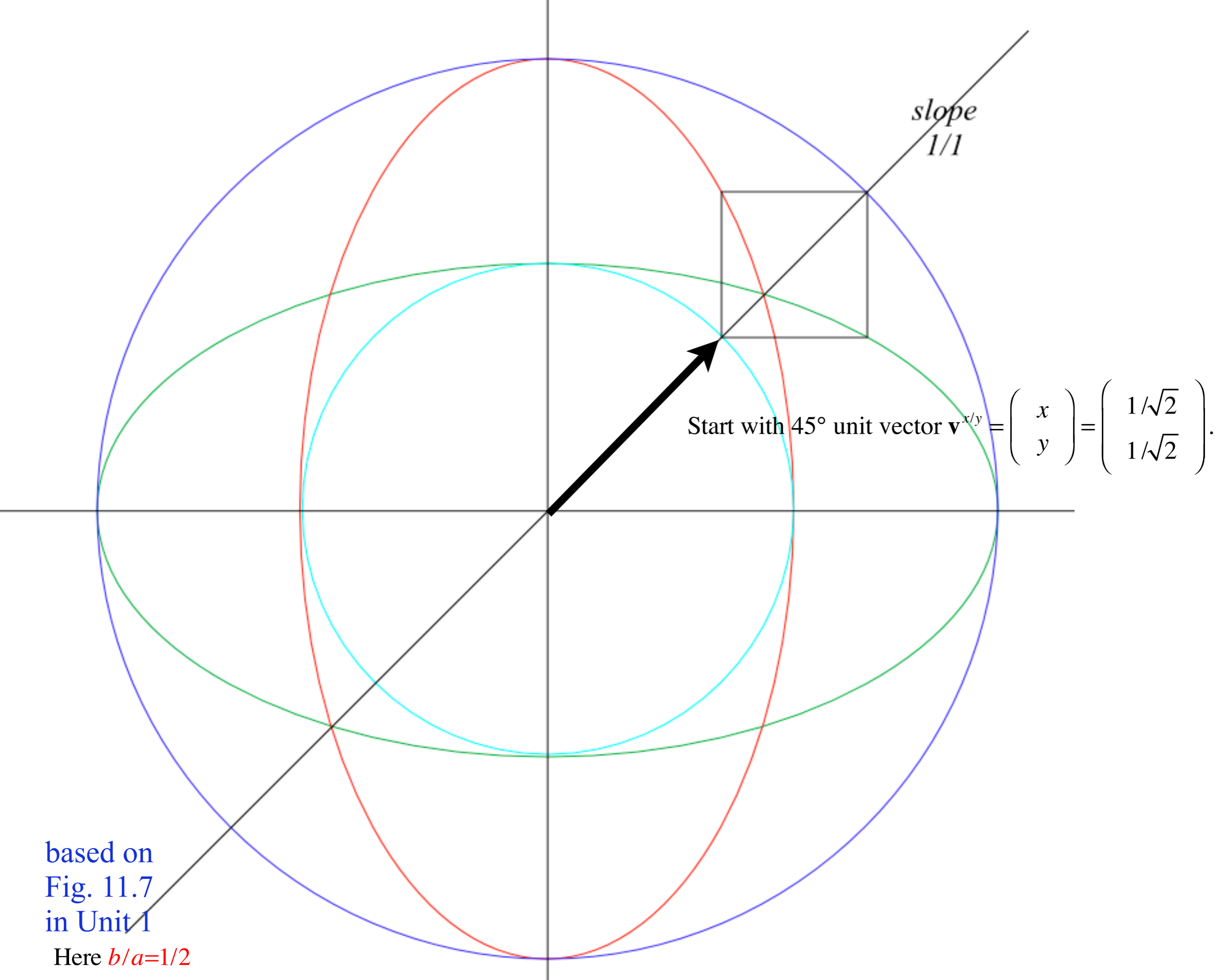
Duality norm relations ($\mathbf{r} \cdot \mathbf{p} = 1$)

\mathbf{Q} -Ellipse tangents \mathbf{r}' normal to dual \mathbf{Q}^{-1} -ellipse position \mathbf{p} ($\mathbf{r}' \cdot \mathbf{p} = 0 = \mathbf{r} \cdot \mathbf{p}'$)

 *Operator geometric sequences and eigenvectors*

Alternative scaling of matrix operator geometry

Vector calculus of tensor operation



based on
 Fig. 11.7
 in Unit 1

Here $b/a=1/2$

Diagonal \mathbf{R} -matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $a/b=2$.

$$\mathbf{R} \cdot \mathbf{v}^{x/y} = \begin{pmatrix} 1/a & 0 \\ 0 & 1/b \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a \\ y/b \end{pmatrix}$$

(Slope increases if $a > b$.)

Action of "sqrt-" matrix $R = \sqrt{Q}$

slope
 a/b

slope
 $1/1$

slope
 b/a

Action of "sqrt⁻¹-" matrix $R^{-1} = \sqrt{Q^{-1}}$

Diagonal \mathbf{R}^{-1} -matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor b/a .

$$\mathbf{R}^{-1} \cdot \mathbf{v}^{x/y} = \begin{pmatrix} a & 0 \\ 0 & b \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x \cdot a \\ y \cdot b \end{pmatrix}$$

(Slope decreases if $b < a$.)

based on
Fig. 11.7
in Unit 1

Here $b/a=1/2$

Diagonal \mathbf{R} -matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $a/b=2$.

$$\mathbf{R} \cdot \mathbf{v}^{x/y} = \begin{pmatrix} 1/a & 0 \\ 0 & 1/b \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a \\ y/b \end{pmatrix}$$

(It increases if $a > b$.)

Diagonal ($\mathbf{R}^2 = \mathbf{Q}$)-matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $a^2/b^2 = 4$.

$$\mathbf{Q} \cdot \mathbf{v}^{x/y} = \begin{pmatrix} 1/a^2 & 0 \\ 0 & 1/b^2 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a^2 \\ y/b^2 \end{pmatrix}$$

(It increases if $a > b$.)

Action of "sqrt-" matrix $R = \sqrt{Q}$

slope a^2/b^2

slope a/b

slope $1/1$

slope b/a

slope b^2/a^2

Action of "sqrt⁻¹-" matrix $R^{-1} = \sqrt{Q^{-1}}$

Diagonal \mathbf{R}^{-1} -matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $b/a=1/2$.

$$\mathbf{R}^{-1} \cdot \mathbf{v}^{x/y} = \begin{pmatrix} a & 0 \\ 0 & b \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x \cdot a \\ y \cdot b \end{pmatrix}$$

Diagonal ($\mathbf{R}^{-2} = \mathbf{Q}^{-1}$)-matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $b^2/a^2=1/4$.

$$\mathbf{Q}^{-1} \cdot \mathbf{v}^{x/y} = \begin{pmatrix} a^2 & 0 \\ 0 & b^2 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x \cdot a^2 \\ y \cdot b^2 \end{pmatrix}$$

based on
Fig. 11.7
in Unit 1

Here $b/a=1/2$

Diagonal \mathbf{R} -matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $a/b = 2$.

$$\mathbf{R} \cdot \mathbf{v}^{x/y} = \begin{pmatrix} 1/a & 0 \\ 0 & 1/b \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a \\ y/b \end{pmatrix}$$

(It increases if $a > b$.)

Diagonal ($\mathbf{R}^2 = \mathbf{Q}$)-matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $a^2/b^2 = 4$.

$$\mathbf{Q} \cdot \mathbf{v}^{x/y} = \begin{pmatrix} 1/a^2 & 0 \\ 0 & 1/b^2 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a^2 \\ y/b^2 \end{pmatrix}$$

(It increases if $a > b$.)

Either process can go on forever...

Diagonal ($\mathbf{R}^{2n} = \mathbf{Q}^n$)-matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $a^{2n}/b^{2n} = 4^n$.

Either process can go on forever...

Diagonal ($\mathbf{R}^{-2n} = \mathbf{Q}^{-n}$)-matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $b^{2n}/a^{2n} = 4^{-n}$.

based on
Fig. 11.7
in Unit 1

Here $b/a = 1/2$

slopeslope

a^3/b^3 a^2/b^2

slope
 $/a/b$

slope
 $1/1$

slope
 b/a

slope
 b^2/a^2

slope
 b^3/a^3

Diagonal \mathbf{R} -matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $a/b = 2$.

$$\mathbf{R} \cdot \mathbf{v}^{x/y} = \begin{pmatrix} 1/a & 0 \\ 0 & 1/b \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a \\ y/b \end{pmatrix}$$

(It increases if $a > b$.)

Diagonal ($\mathbf{R}^2 = \mathbf{Q}$)-matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $a^2/b^2 = 4$.

$$\mathbf{Q} \cdot \mathbf{v}^{x/y} = \begin{pmatrix} 1/a^2 & 0 \\ 0 & 1/b^2 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a^2 \\ y/b^2 \end{pmatrix}$$

(It increases if $a > b$.)

Either process can go on forever...

Diagonal ($\mathbf{R}^{2n} = \mathbf{Q}^n$)-matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $a^{2n}/b^{2n} = 4^n$.

...Finally, the result approaches **EIGENVECTOR** $|y\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$

of ∞ -slope which is "immune" to \mathbf{R} , \mathbf{Q} or \mathbf{Q}^n :

$$\mathbf{R}|y\rangle = (1/b)|y\rangle \quad \mathbf{Q}^n|y\rangle = (1/b^2)^n|y\rangle$$

Here $b/a = 1/2$

slopeslope

a^3/b^3 a^2/b^2

slope
 $/a/b$

slope
 $1/1$

EIGENVECTOR

$|y\rangle$

slope
 b/a

slope
 b^2/a^2

slope
 b^3/a^3

EIGENVECTOR

$|x\rangle$

Either process can go on forever...

Diagonal ($\mathbf{R}^{-2n} = \mathbf{Q}^{-n}$)-matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $b^{2n}/a^{2n} = 4^{-n}$.

...Finally, the result approaches **EIGENVECTOR** $|x\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$

of 0 -slope which is "immune" to \mathbf{R}^{-1} , \mathbf{Q}^{-1} or \mathbf{Q}^{-n} :

$$\mathbf{R}^{-1}|x\rangle = (a)|x\rangle \quad \mathbf{Q}^{-n}|x\rangle = (a^2)^n|x\rangle$$

Diagonal \mathbf{R} -matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $a/b=2$.

$$\mathbf{R} \cdot \mathbf{v}^{x/y} = \begin{pmatrix} 1/a & 0 \\ 0 & 1/b \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a \\ y/b \end{pmatrix}$$

(It increases if $a > b$.)

EIGENVECTOR

Diagonal ($\mathbf{R}^2=\mathbf{Q}$)-matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $a^2/b^2=4$.

$$\mathbf{Q} \cdot \mathbf{v}^{x/y} = \begin{pmatrix} 1/a^2 & 0 \\ 0 & 1/b^2 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a^2 \\ y/b^2 \end{pmatrix}$$

(It increases if $a > b$.)

EIGENVECTOR

Either process can go on forever...

Diagonal ($\mathbf{R}^{2n}=\mathbf{Q}^n$)-matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $a^{2n}/b^{2n}=4^n$.

...Finally, the result approaches **EIGENVECTOR** $|y\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$

of ∞ -slope which is "immune" to \mathbf{R} , \mathbf{Q} or \mathbf{Q}^n :

$$\mathbf{R}|y\rangle = (1/b)|y\rangle$$

$$\mathbf{Q}^n|y\rangle = (1/b^2)^n|y\rangle$$

Eigenvalues

Eigensolution Relations

Either process can go on forever...

Diagonal ($\mathbf{R}^{-2n}=\mathbf{Q}^{-n}$)-matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $b^{2n}/a^{2n}=4^{-n}$.

...Finally, the result approaches **EIGENVECTOR** $|x\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$

of 0 -slope which is "immune" to \mathbf{R}^{-1} , \mathbf{Q}^{-1} or \mathbf{Q}^{-n} :

$$\mathbf{R}^{-1}|x\rangle = (a)|x\rangle$$

$$\mathbf{Q}^{-n}|x\rangle = (a^2)^n|x\rangle$$

Eigenvalues

Introduction to dual matrix operator geometry (based on IHO orbits)

Quadratic form ellipse $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = 1$ vs. inverse form ellipse $\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p} = 1$

Duality norm relations ($\mathbf{r} \cdot \mathbf{p} = 1$)

\mathbf{Q} -Ellipse tangents \mathbf{r}' normal to dual \mathbf{Q}^{-1} -ellipse position \mathbf{p} ($\mathbf{r}' \cdot \mathbf{p} = 0 = \mathbf{r} \cdot \mathbf{p}'$)

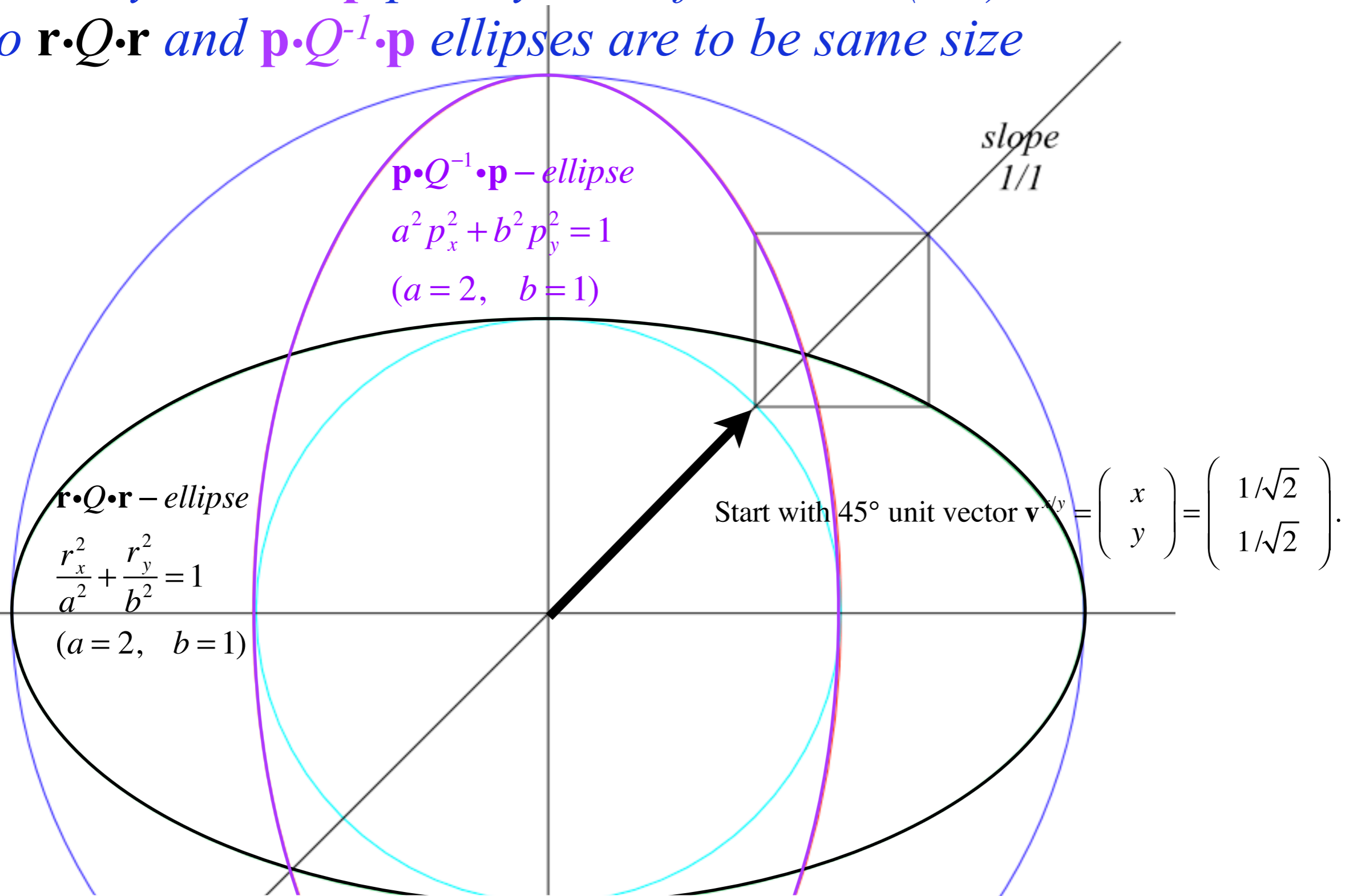
Operator geometric sequences and eigenvectors



Alternative scaling of matrix operator geometry

Vector calculus of tensor operation

You may rescale \mathbf{p} -plot by scale factor $S=(a \cdot b)$
 so $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r}$ and $\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p}$ ellipses are to be same size

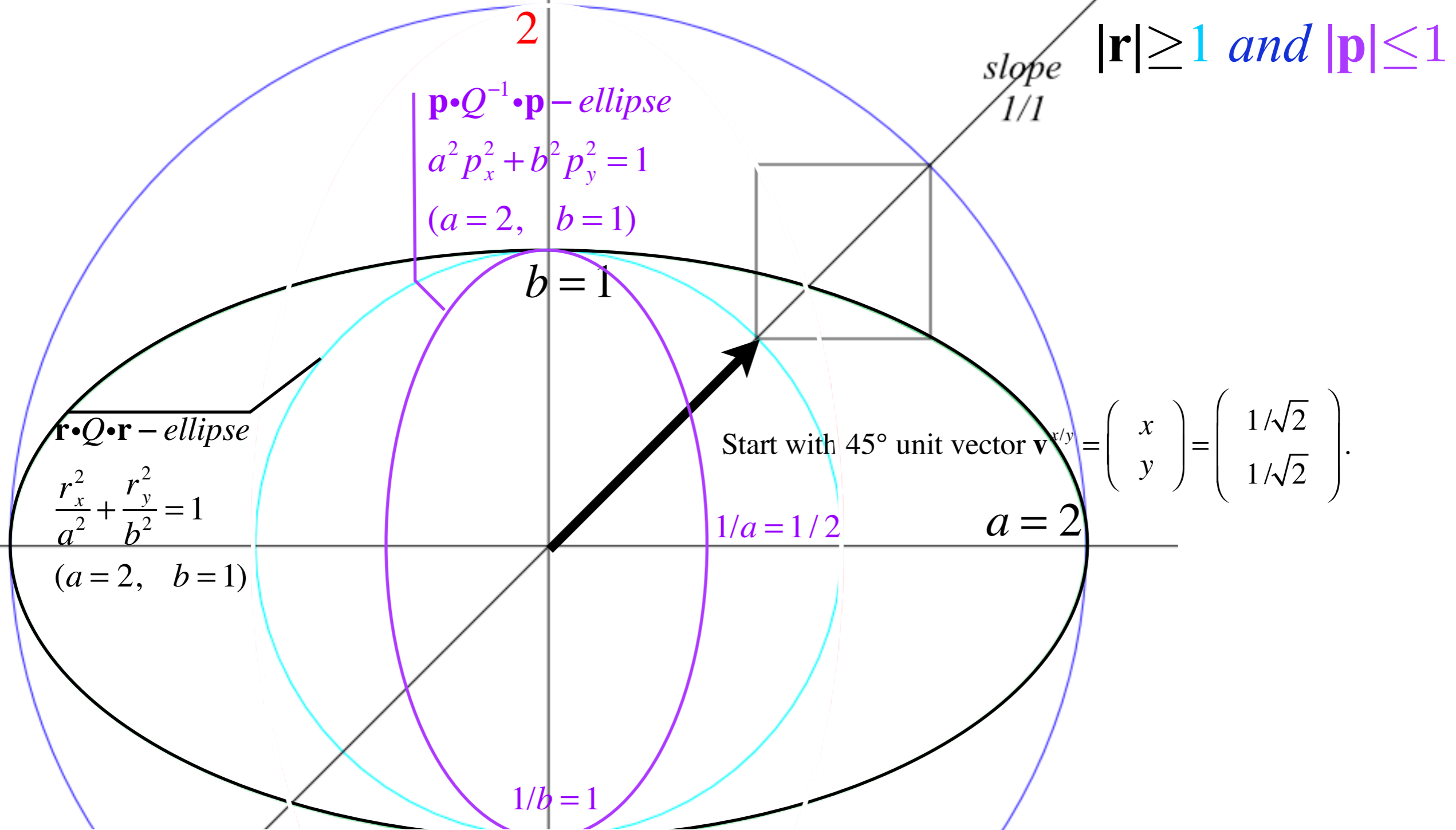


Here plot of \mathbf{p} -ellipse is re-scaled by scalefactor $S=a \cdot b$

\mathbf{p} -ellipse x -radius= $1/a$ plotted at: $S(1/a)=b$ ($=1$ for $a=2, b=1$)

\mathbf{p} -ellipse y -radius= $1/b$ plotted at: $S(1/b)=a$ ($=2$ for $a=2, b=1$)

..or else rescale **p**-plot by scale factor $S=b$
 to separate $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r}$ and $\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p}$ ellipses into different regions



Here plot of **p**-ellipse is re-scaled by scalefactor $S=b$

p-ellipse x -radius= $1/a$ plotted at: $S(1/a)=b/a$ ($=1/2$ for $a=2, b=1$)

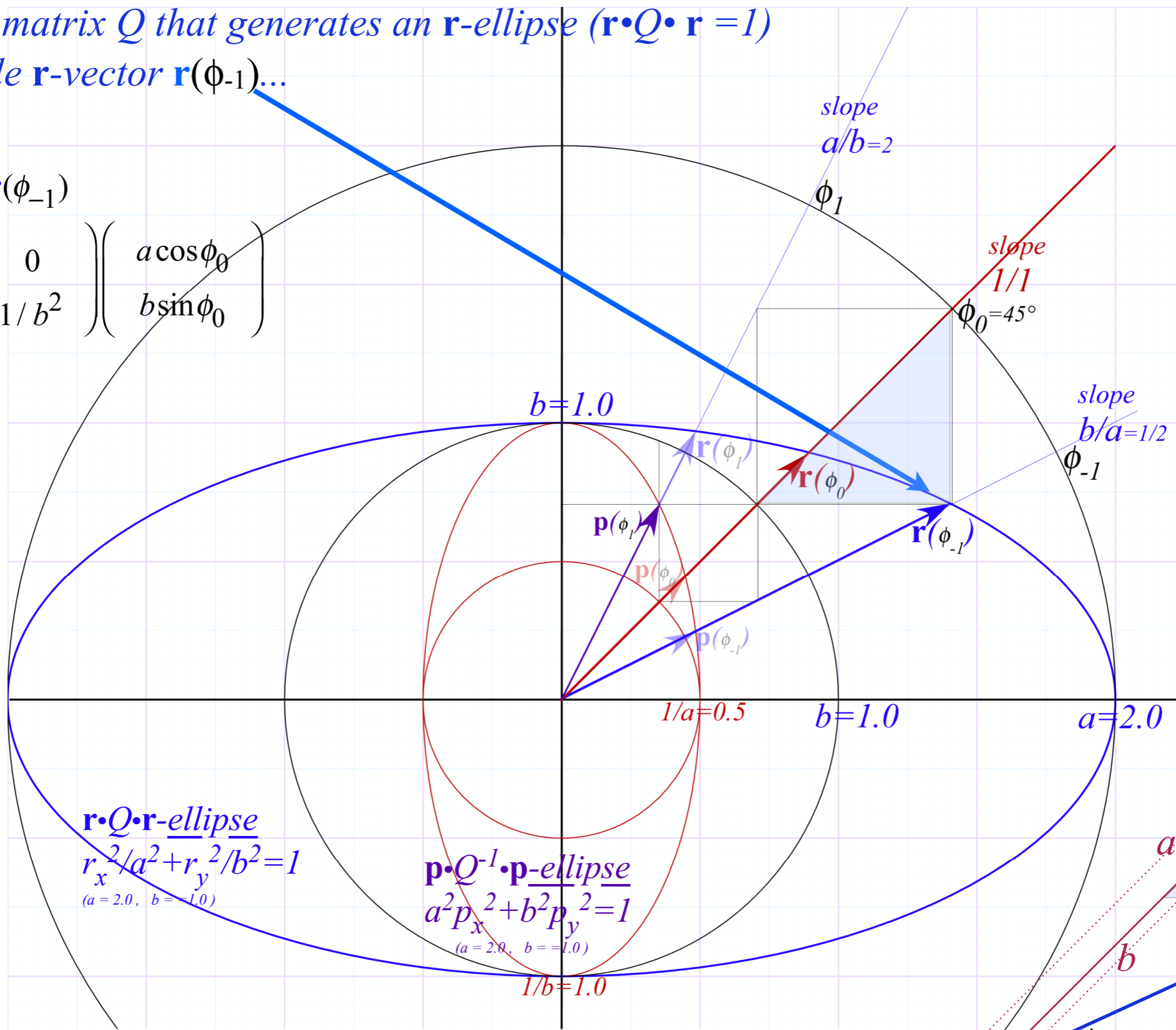
p-ellipse y -radius= $1/b$ plotted at: $S(1/b)=1$

Action of matrix Q that generates an \mathbf{r} -ellipse ($\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = 1$)

on a single \mathbf{r} -vector $\mathbf{r}(\phi_{-1}) \dots$

$$\mathbf{p}(\phi_1) = \mathbf{Q} \cdot \mathbf{r}(\phi_{-1})$$

$$= \begin{pmatrix} 1/a^2 & 0 \\ 0 & 1/b^2 \end{pmatrix} \begin{pmatrix} a \cos \phi_0 \\ b \sin \phi_0 \end{pmatrix}$$



Variation of Fig. 11.7 in Unit 1

$$\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} \text{-ellipse}$$

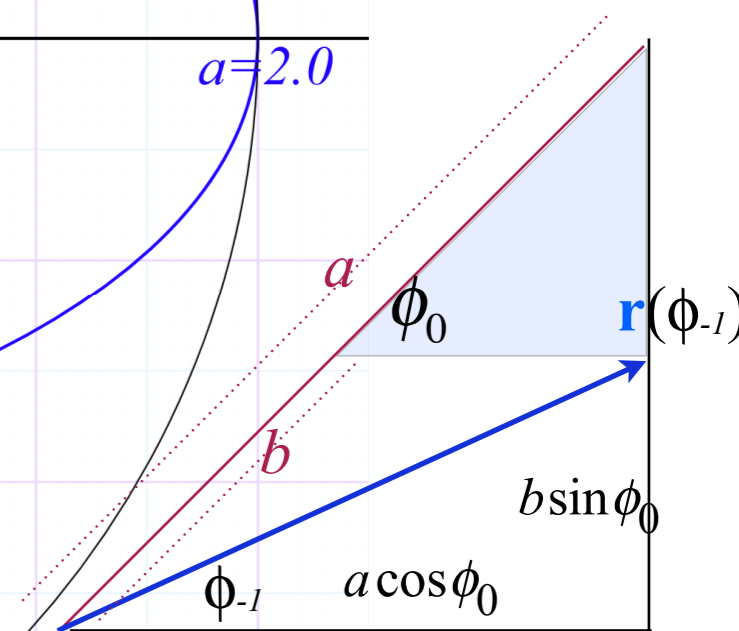
$$\frac{r_x^2}{a^2} + \frac{r_y^2}{b^2} = 1$$

(a = 2.0, b = 1.0)

$$\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p} \text{-ellipse}$$

$$a^2 p_x^2 + b^2 p_y^2 = 1$$

(a = 2.0, b = 1.0)



Here plot of \mathbf{p} -ellipse is re-scaled by scalefactor $S=b$

\mathbf{p} -ellipse x -radius= $1/a$ plotted at: $S(1/a)=b/a$ (=1/2 for $a=2, b=1$)

\mathbf{p} -ellipse y -radius= $1/b$ plotted at: $S(1/b)=1$

Action of matrix Q that generates an \mathbf{r} -ellipse ($\mathbf{r} \cdot Q \cdot \mathbf{r} = 1$)

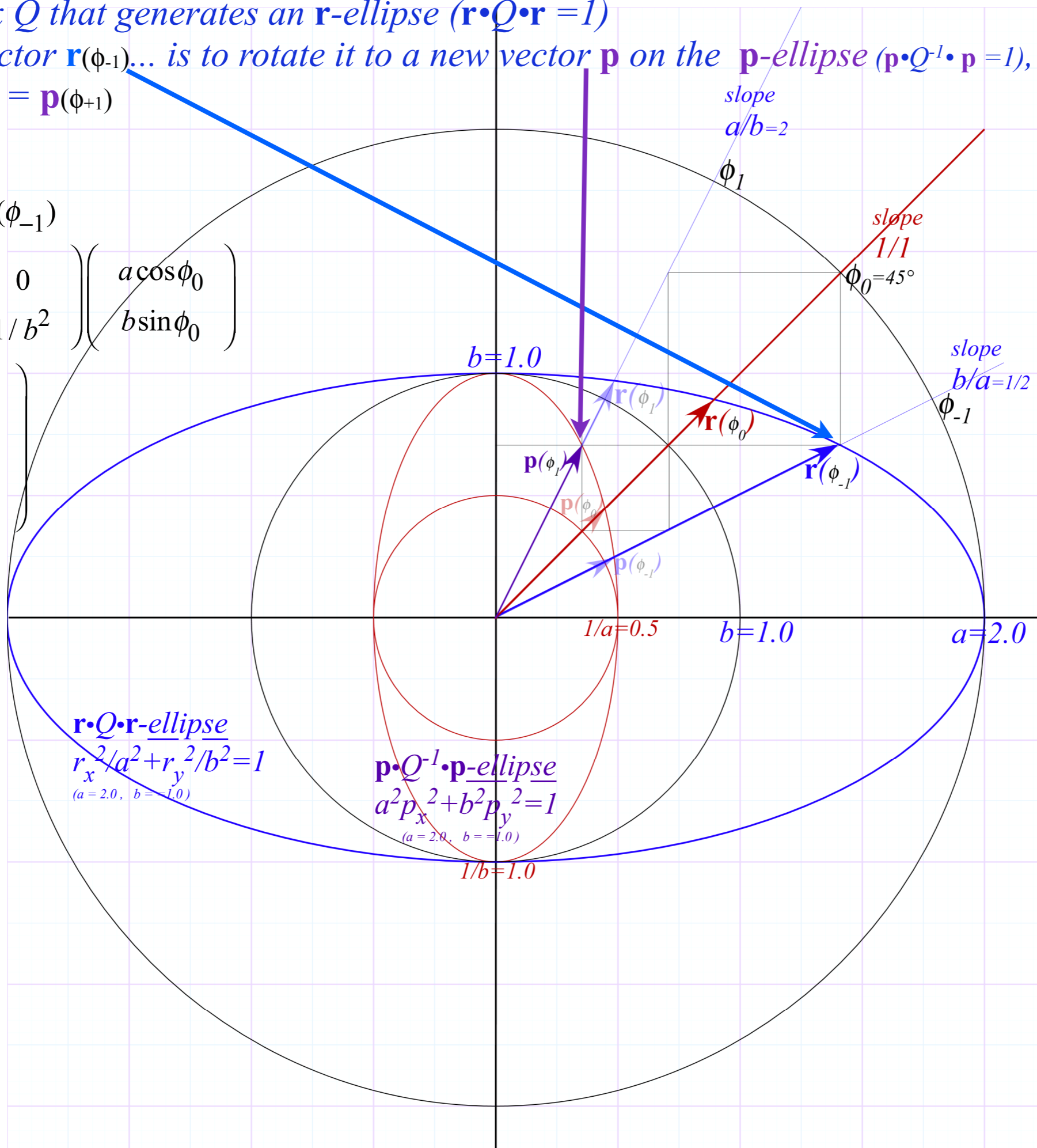
on a single \mathbf{r} -vector $\mathbf{r}(\phi_{-1})$... is to rotate it to a new vector \mathbf{p} on the \mathbf{p} -ellipse ($\mathbf{p} \cdot Q^{-1} \cdot \mathbf{p} = 1$), that is, $Q \cdot \mathbf{r}(\phi_{-1}) = \mathbf{p}(\phi_{+1})$

$$\mathbf{p}(\phi_1) = Q \cdot \mathbf{r}(\phi_{-1})$$

$$= \begin{pmatrix} 1/a^2 & 0 \\ 0 & 1/b^2 \end{pmatrix} \begin{pmatrix} a \cos \phi_0 \\ b \sin \phi_0 \end{pmatrix}$$

$$= \begin{pmatrix} \frac{1}{a} \cos \phi_0 \\ \frac{1}{b} \sin \phi_0 \end{pmatrix}$$

$$= \begin{pmatrix} \frac{1}{2} & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & \frac{1}{2} \end{pmatrix}$$



Variation of Fig. 11.7 in Unit 1

Action of matrix Q that generates an \mathbf{r} -ellipse ($\mathbf{r} \cdot Q \cdot \mathbf{r} = 1$)

on a single \mathbf{r} -vector $\mathbf{r}(\phi_{-1})$... is to rotate it to a new vector \mathbf{p} on the \mathbf{p} -ellipse ($\mathbf{p} \cdot Q^{-1} \cdot \mathbf{p} = 1$), that is, $Q \cdot \mathbf{r}(\phi_{-1}) = \mathbf{p}(\phi_{+1})$

$$\mathbf{p}(\phi_1) = Q \cdot \mathbf{r}(\phi_{-1})$$

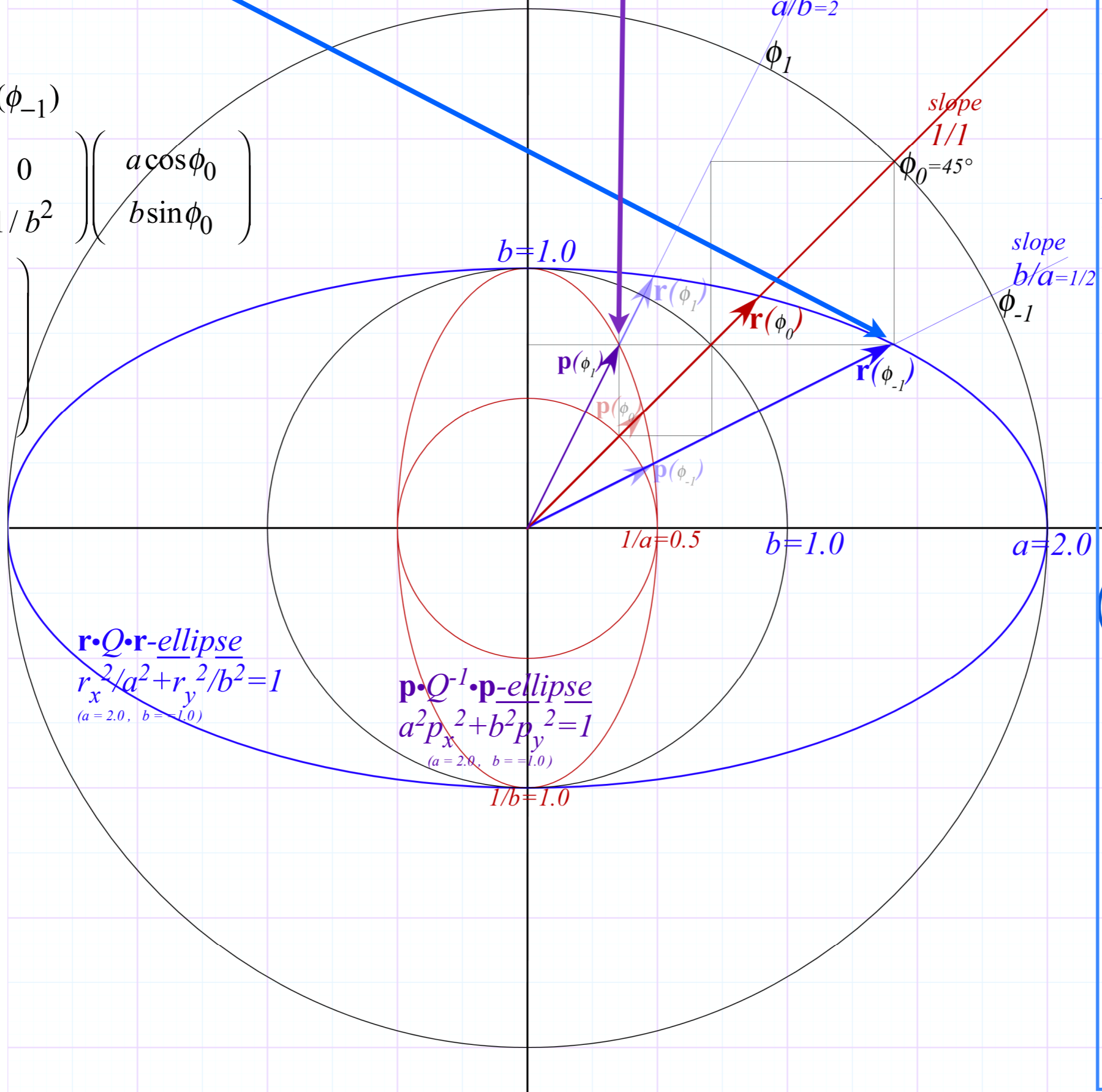
$$= \begin{pmatrix} 1/a^2 & 0 \\ 0 & 1/b^2 \end{pmatrix} \begin{pmatrix} a \cos \phi_0 \\ b \sin \phi_0 \end{pmatrix}$$

$$= \begin{pmatrix} \frac{1}{a} \cos \phi_0 \\ \frac{1}{b} \sin \phi_0 \end{pmatrix}$$

$$= \begin{pmatrix} \frac{1}{2} & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \end{pmatrix}$$

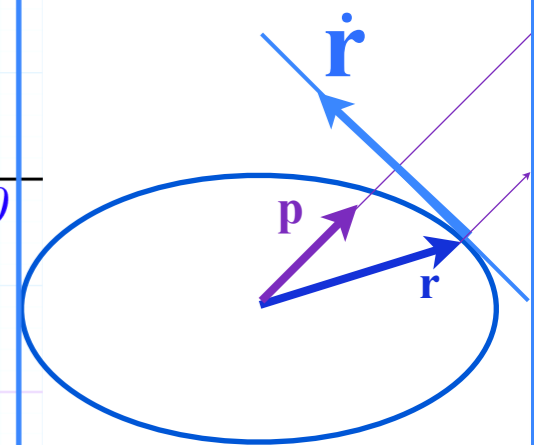
$\mathbf{r} \cdot Q \cdot \mathbf{r}$ -ellipse
 $r_x^2/a^2 + r_y^2/b^2 = 1$
 ($a = 2.0, b = 1.0$)

$\mathbf{p} \cdot Q^{-1} \cdot \mathbf{p}$ -ellipse
 $a^2 p_x^2 + b^2 p_y^2 = 1$
 ($a = 2.0, b = 1.0$)



Key points of matrix geometry:

Matrix Q maps any vector \mathbf{r} to a new vector \mathbf{p} normal to the tangent $\dot{\mathbf{r}}$ to its $\mathbf{r} \cdot Q \cdot \mathbf{r}$ -ellipse.



Variation of Fig. 11.7 in Unit 1

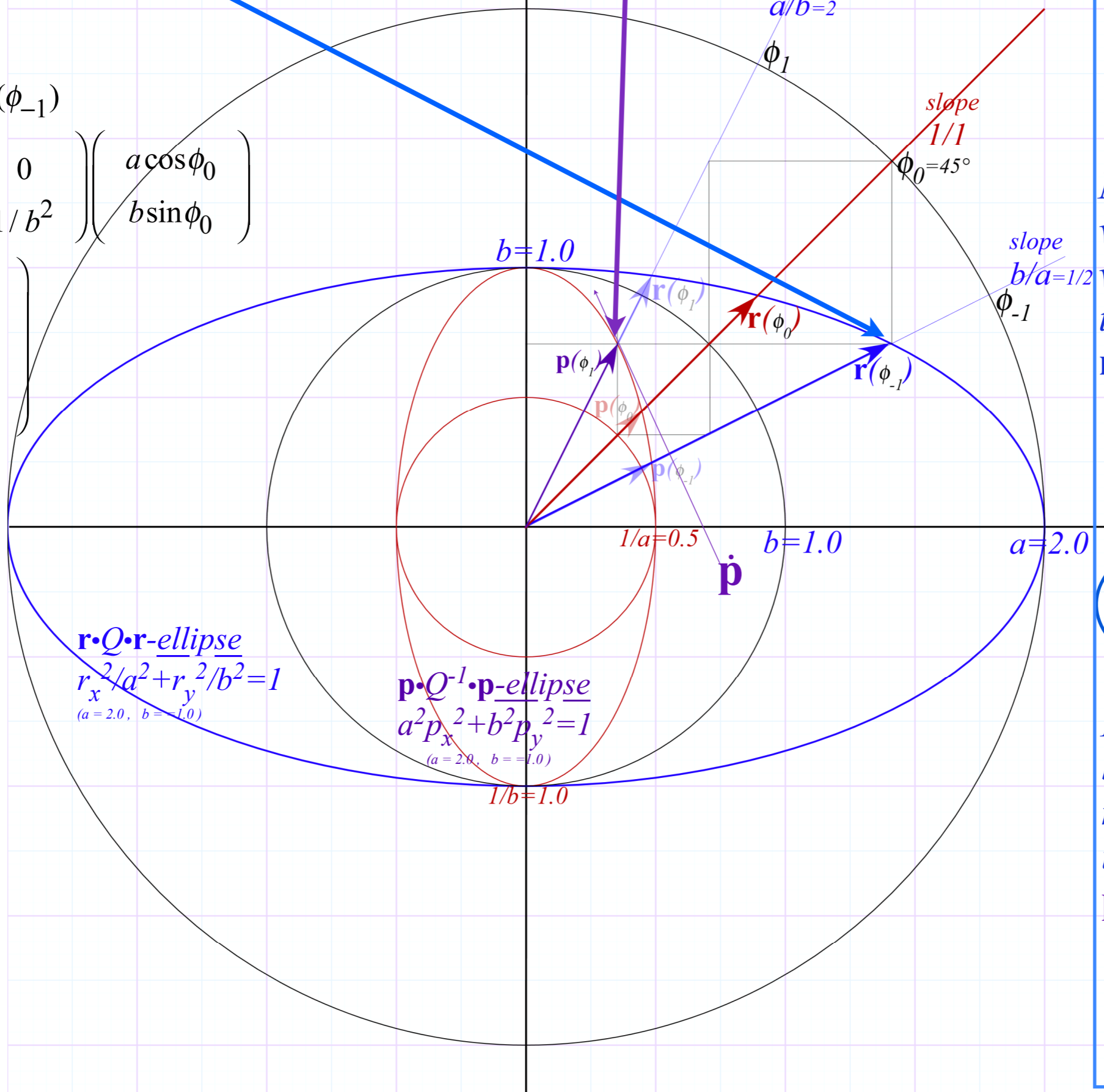
Action of matrix Q that generates an \mathbf{r} -ellipse ($\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = 1$) on a single \mathbf{r} -vector $\mathbf{r}(\phi_{-1})$... is to rotate it to a new vector \mathbf{p} on the \mathbf{p} -ellipse ($\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p} = 1$), that is, $\mathbf{Q} \cdot \mathbf{r}(\phi_{-1}) = \mathbf{p}(\phi_{+1})$

$$\mathbf{p}(\phi_1) = \mathbf{Q} \cdot \mathbf{r}(\phi_{-1})$$

$$= \begin{pmatrix} 1/a^2 & 0 \\ 0 & 1/b^2 \end{pmatrix} \begin{pmatrix} a \cos \phi_0 \\ b \sin \phi_0 \end{pmatrix}$$

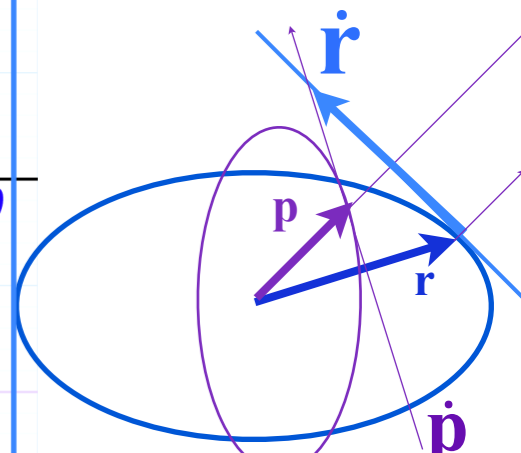
$$= \begin{pmatrix} \frac{1}{a} \cos \phi_0 \\ \frac{1}{b} \sin \phi_0 \end{pmatrix}$$

$$= \begin{pmatrix} \frac{1}{2} & \frac{1}{\sqrt{2}} \\ \frac{1}{1} & \frac{1}{\sqrt{2}} \end{pmatrix}$$



Key points of matrix geometry:

Matrix Q maps any vector \mathbf{r} to a new vector \mathbf{p} normal to the tangent $\dot{\mathbf{r}}$ to its $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r}$ -ellipse.



Matrix Q^{-1} maps \mathbf{p} back to \mathbf{r} that is normal to the tangent $\dot{\mathbf{p}}$ to its $\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p}$ -ellipse.

Variation of Fig. 11.7 in Unit 1

Introduction to dual matrix operator geometry (based on IHO orbits)

Quadratic form ellipse $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = 1$ vs. inverse form ellipse $\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p} = 1$

Duality norm relations ($\mathbf{r} \cdot \mathbf{p} = 1$)

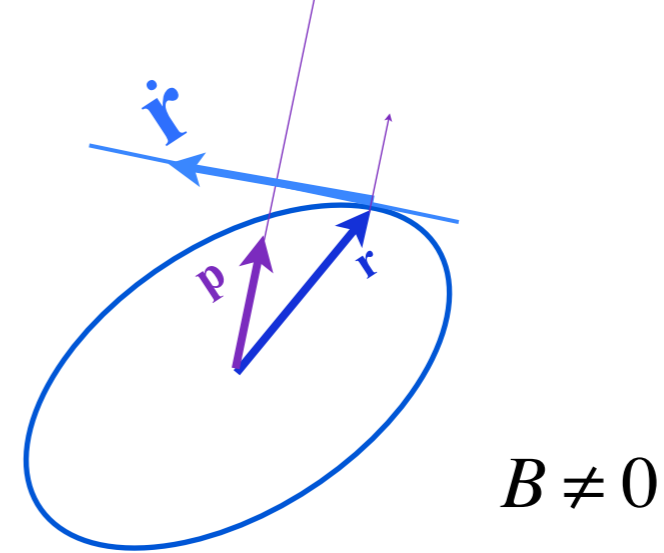
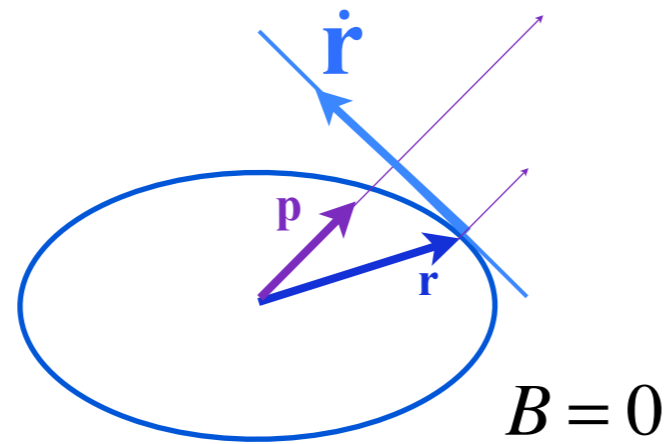
\mathbf{Q} -Ellipse tangents \mathbf{r}' normal to dual \mathbf{Q}^{-1} -ellipse position \mathbf{p} ($\mathbf{r}' \cdot \mathbf{p} = 0 = \mathbf{r} \cdot \mathbf{p}'$)

Operator geometric sequences and eigenvectors

Alternative scaling of matrix operator geometry



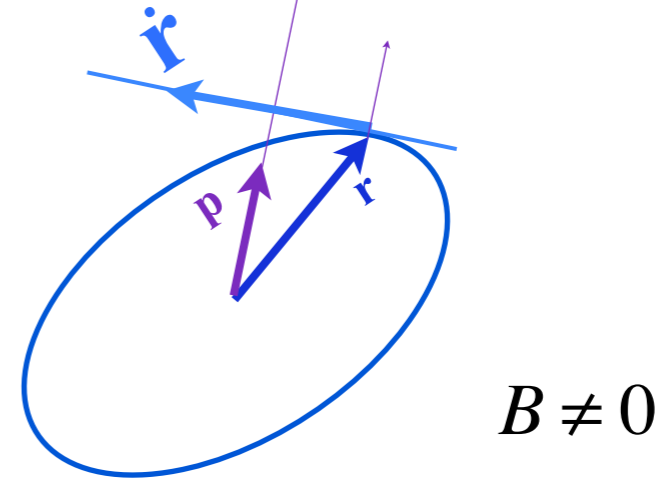
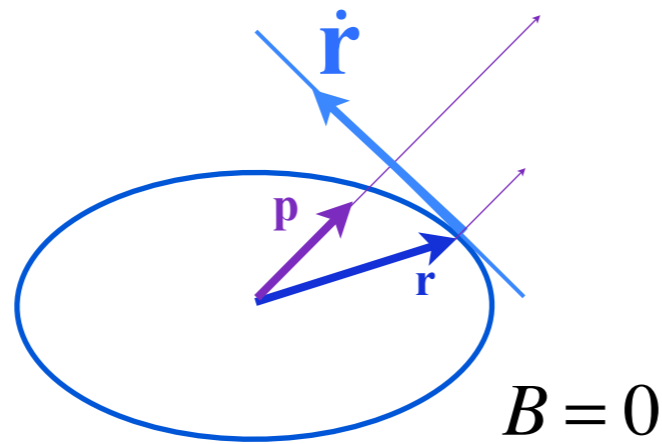
Vector calculus of tensor operation



Derive matrix “normal-to-ellipse” geometry by vector calculus:

Let matrix $Q = \begin{pmatrix} A & B \\ B & D \end{pmatrix}$

define the ellipse $1 = \mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = \begin{pmatrix} x & y \end{pmatrix} \cdot \begin{pmatrix} A & B \\ B & D \end{pmatrix} \cdot \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x & y \end{pmatrix} \cdot \begin{pmatrix} A \cdot x + B \cdot y \\ B \cdot x + D \cdot y \end{pmatrix} = A \cdot x^2 + 2B \cdot xy + D \cdot y^2 = 1$



Derive matrix “normal-to-ellipse” geometry by vector calculus:

Let matrix $Q = \begin{pmatrix} A & B \\ B & D \end{pmatrix}$

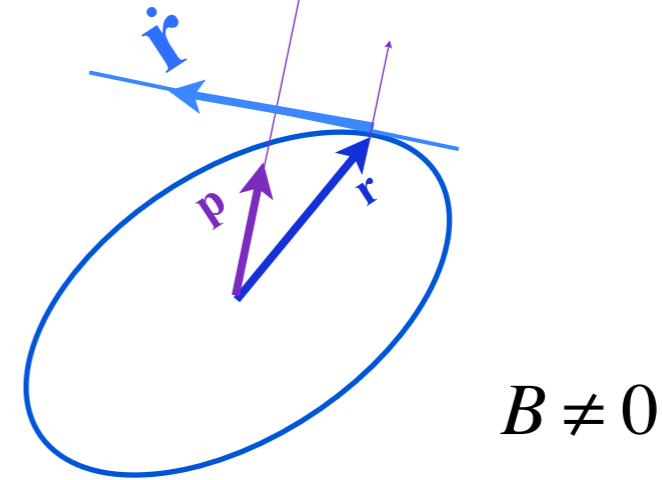
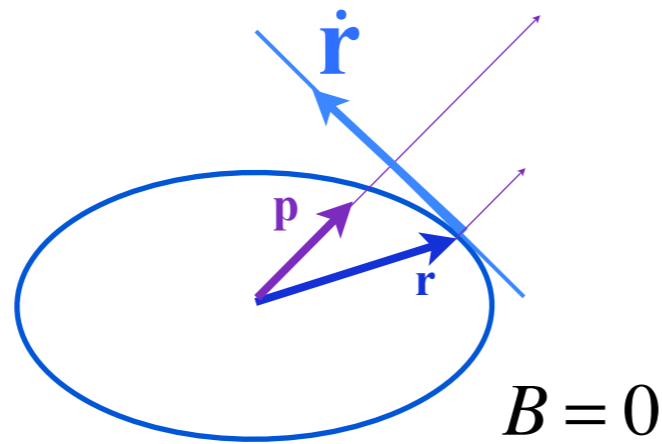
define the ellipse $1 = \mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = \begin{pmatrix} x & y \end{pmatrix} \cdot \begin{pmatrix} A & B \\ B & D \end{pmatrix} \cdot \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x & y \end{pmatrix} \cdot \begin{pmatrix} A \cdot x + B \cdot y \\ B \cdot x + D \cdot y \end{pmatrix} = A \cdot x^2 + 2B \cdot xy + D \cdot y^2 = 1$

Compare operation by Q on vector \mathbf{r} with vector derivative or gradient of $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r}$

$$\begin{pmatrix} A & B \\ B & D \end{pmatrix} \cdot \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} A \cdot x + B \cdot y \\ B \cdot x + D \cdot y \end{pmatrix}$$

$$\frac{\partial}{\partial \mathbf{r}} (\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r}) = \nabla (\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r})$$

$$\begin{pmatrix} \frac{\partial}{\partial x} \\ \frac{\partial}{\partial y} \end{pmatrix} (A \cdot x^2 + 2B \cdot xy + D \cdot y^2) = \begin{pmatrix} 2A \cdot x + 2B \cdot y \\ 2B \cdot x + 2D \cdot y \end{pmatrix}$$



Derive matrix “normal-to-ellipse” geometry by vector calculus:

Let matrix $Q = \begin{pmatrix} A & B \\ B & D \end{pmatrix}$

define the ellipse $1 = \mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = \begin{pmatrix} x & y \end{pmatrix} \cdot \begin{pmatrix} A & B \\ B & D \end{pmatrix} \cdot \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x & y \end{pmatrix} \cdot \begin{pmatrix} A \cdot x + B \cdot y \\ B \cdot x + D \cdot y \end{pmatrix} = A \cdot x^2 + 2B \cdot xy + D \cdot y^2 = 1$

Compare operation by Q on vector \mathbf{r} with vector derivative or gradient of $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r}$

$$\begin{pmatrix} A & B \\ B & D \end{pmatrix} \cdot \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} A \cdot x + B \cdot y \\ B \cdot x + D \cdot y \end{pmatrix}$$

$$\frac{\partial}{\partial \mathbf{r}} (\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r}) = \nabla (\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r})$$

$$\begin{pmatrix} \frac{\partial}{\partial x} \\ \frac{\partial}{\partial y} \end{pmatrix} (A \cdot x^2 + 2B \cdot xy + D \cdot y^2) = \begin{pmatrix} 2A \cdot x + 2B \cdot y \\ 2B \cdot x + 2D \cdot y \end{pmatrix}$$

Very simple result:

$$\frac{\partial}{\partial \mathbf{r}} \left(\frac{\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r}}{2} \right) = \nabla \left(\frac{\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r}}{2} \right) = \mathbf{Q} \cdot \mathbf{r}$$

Introduction to dual matrix operator geometry (based on IHO orbits)

Quadratic form ellipse $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = 1$ vs. inverse form ellipse $\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p} = 1$

Duality norm relations ($\mathbf{r} \cdot \mathbf{p} = 1$)

\mathbf{Q} -Ellipse tangents \mathbf{r}' normal to dual \mathbf{Q}^{-1} -ellipse position \mathbf{p} ($\mathbf{r}' \cdot \mathbf{p} = 0 = \mathbf{r} \cdot \mathbf{p}'$)

(Still more) Operator geometric sequences and eigenvectors

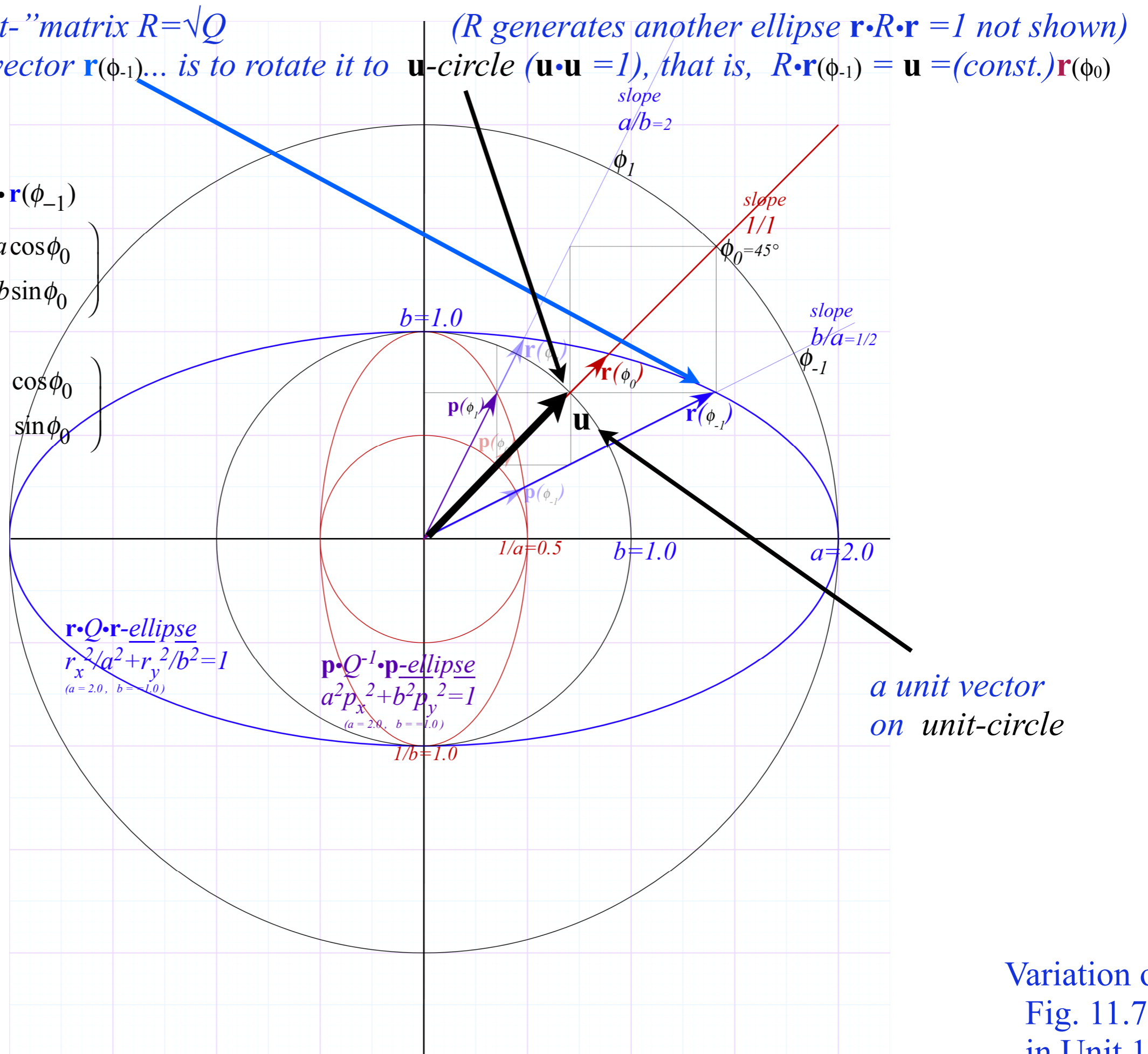
Alternative scaling of matrix operator geometry

Vector calculus of tensor operation



Action of "sqrt-"matrix $R=\sqrt{Q}$ (R generates another ellipse $\mathbf{r}\cdot R\cdot\mathbf{r} = 1$ not shown) on a single \mathbf{r} -vector $\mathbf{r}(\phi_{-1})\dots$ is to rotate it to \mathbf{u} -circle ($\mathbf{u}\cdot\mathbf{u} = 1$), that is, $R\cdot\mathbf{r}(\phi_{-1}) = \mathbf{u} = (\text{const.})\mathbf{r}(\phi_0)$

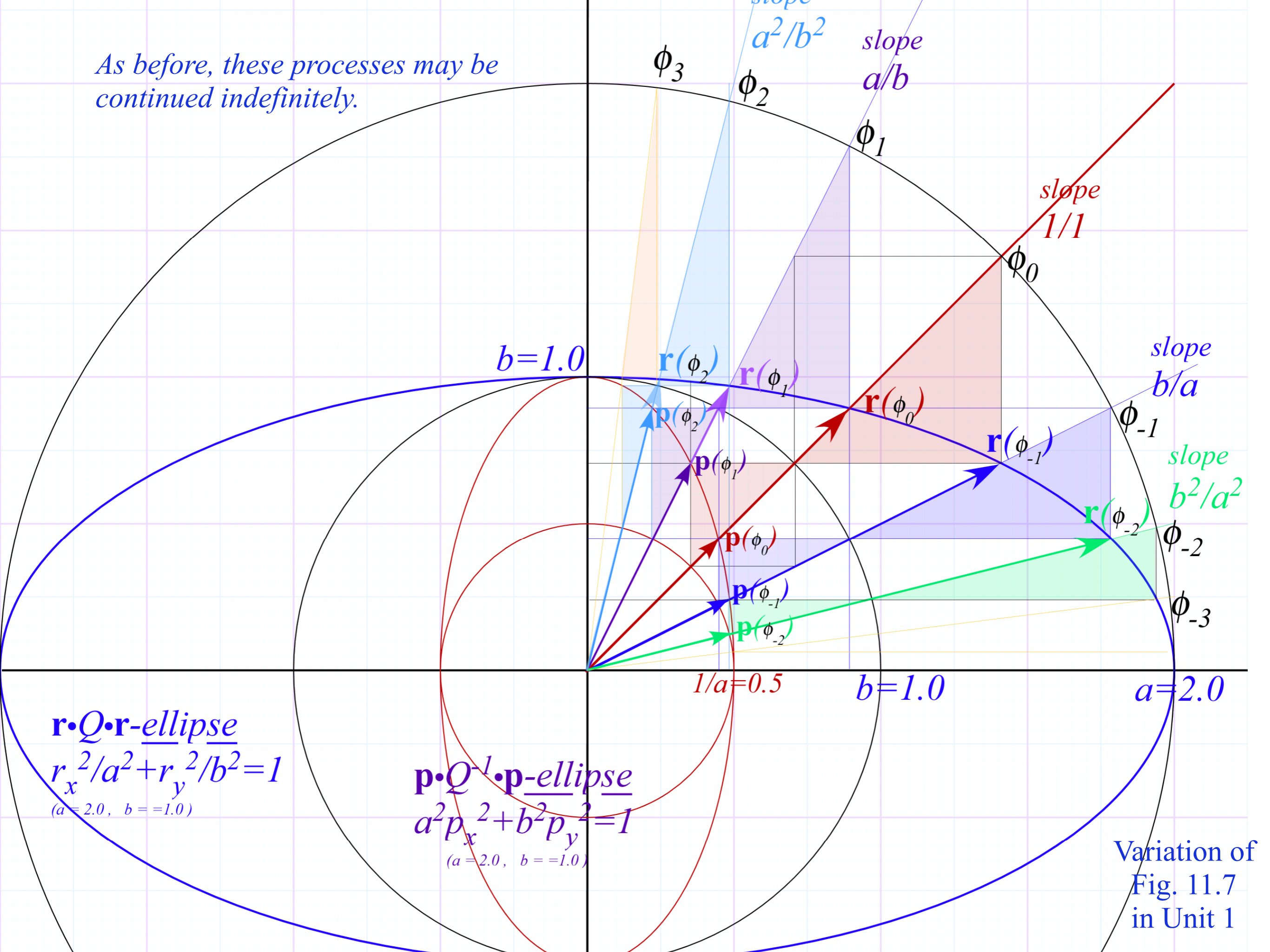
$$\begin{aligned} \mathbf{u} &= \sqrt{\mathbf{Q}} \cdot \mathbf{r}(\phi_{-1}) = \mathbf{R} \cdot \mathbf{r}(\phi_{-1}) \\ &= \begin{pmatrix} 1/a & 0 \\ 0 & 1/b \end{pmatrix} \begin{pmatrix} a \cos \phi_0 \\ b \sin \phi_0 \end{pmatrix} \\ &= \begin{pmatrix} \frac{1}{a} a \cos \phi_0 \\ \frac{1}{b} b \sin \phi_0 \end{pmatrix} = \begin{pmatrix} \cos \phi_0 \\ \sin \phi_0 \end{pmatrix} \\ &= \begin{pmatrix} \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} \end{pmatrix} \end{aligned}$$



a unit vector on unit-circle

Variation of Fig. 11.7 in Unit 1

As before, these processes may be continued indefinitely.

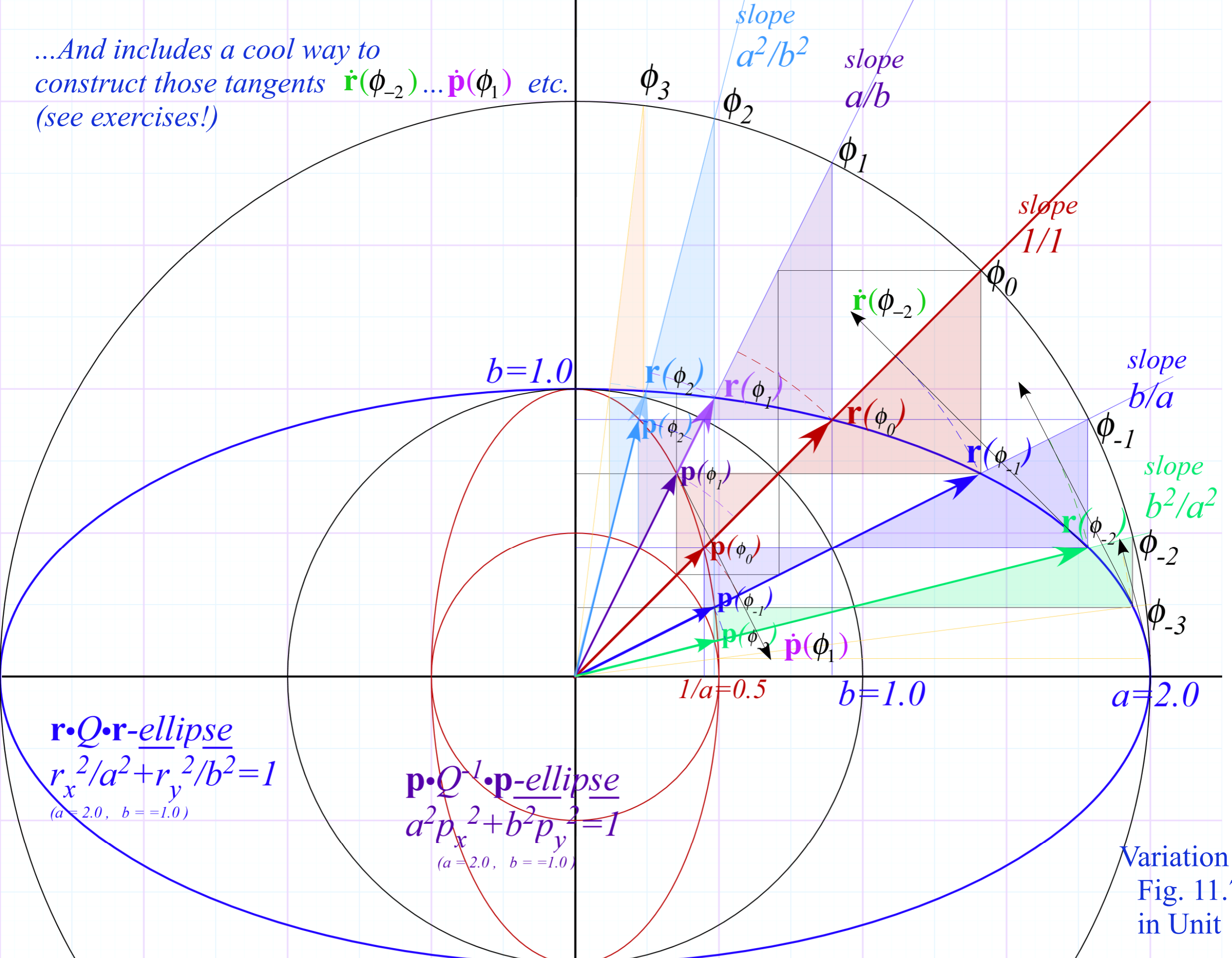


$\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r}$ -ellipse
 $r_x^2/a^2 + r_y^2/b^2 = 1$
 ($a=2.0, b=1.0$)

$\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p}$ -ellipse
 $a^2 p_x^2 + b^2 p_y^2 = 1$
 ($a=2.0, b=1.0$)

Variation of Fig. 11.7 in Unit 1

...And includes a cool way to construct those tangents $\mathbf{r}(\phi_{-2}) \dots \mathbf{p}(\phi_1)$ etc. (see exercises!)



$\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r}$ -ellipse

$$r_x^2/a^2 + r_y^2/b^2 = 1$$

($a=2.0, b=1.0$)

$\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p}$ -ellipse

$$a^2 p_x^2 + b^2 p_y^2 = 1$$

($a=2.0, b=1.0$)

Variation of Fig. 11.7 in Unit 1

*Q: Where is this headed?
Preview of Lecture 9*

A: Lagrangian-Hamiltonian duality

The R and Q matrix transformations are like the mechanics rescaling matrices $\sqrt{\mathbf{M}}$ and \mathbf{M} :

Like $Q=R^2$:

$$\mathbf{M} = \begin{pmatrix} m_1 & 0 \\ 0 & m_2 \end{pmatrix} = \mathbf{R}^2$$

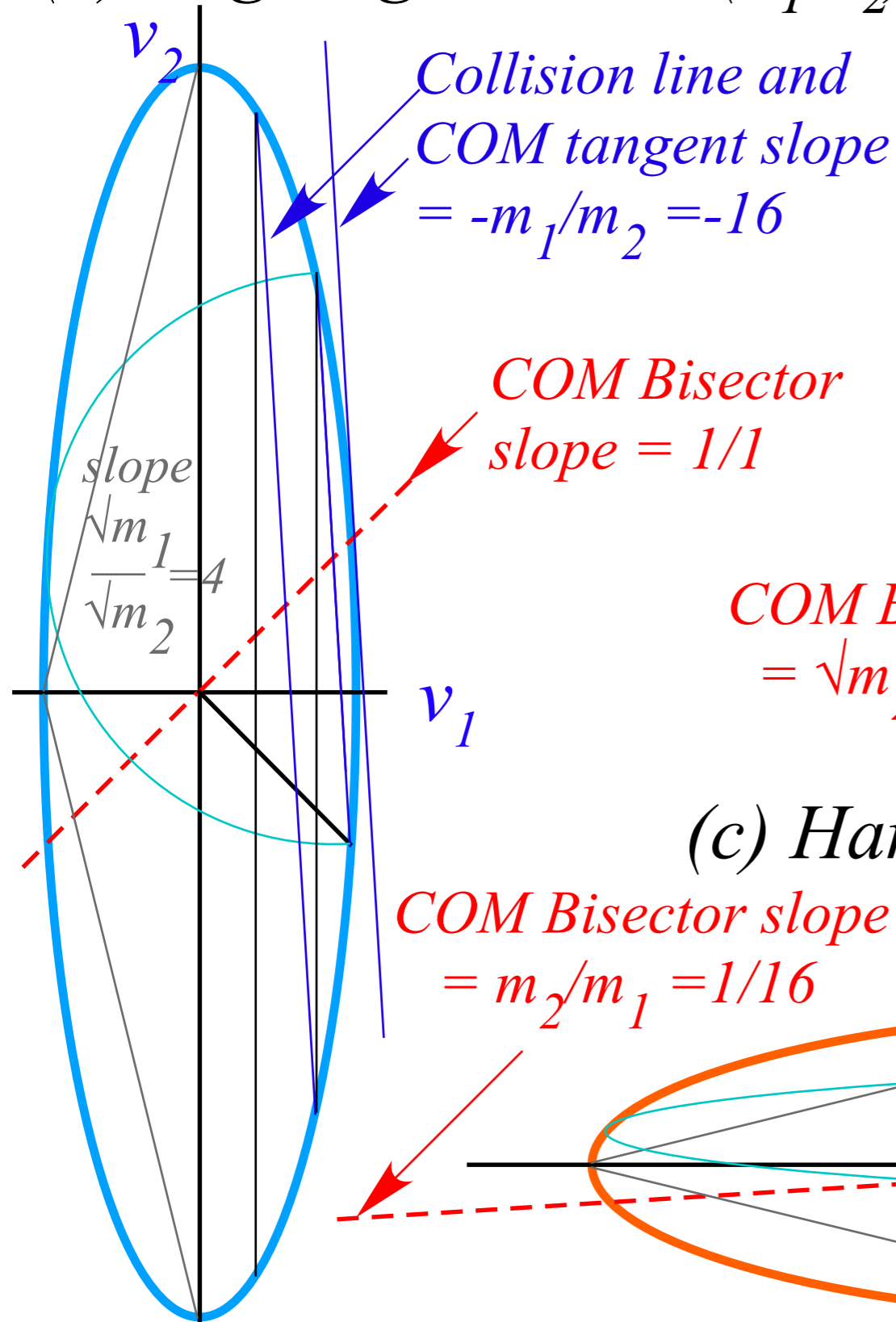
Like $\sqrt{Q}=R$:

$$\sqrt{\mathbf{M}} = \begin{pmatrix} \sqrt{m_1} & 0 \\ 0 & \sqrt{m_2} \end{pmatrix} = \mathbf{R}$$

Like $Q^{-1}=R^{-2}$:

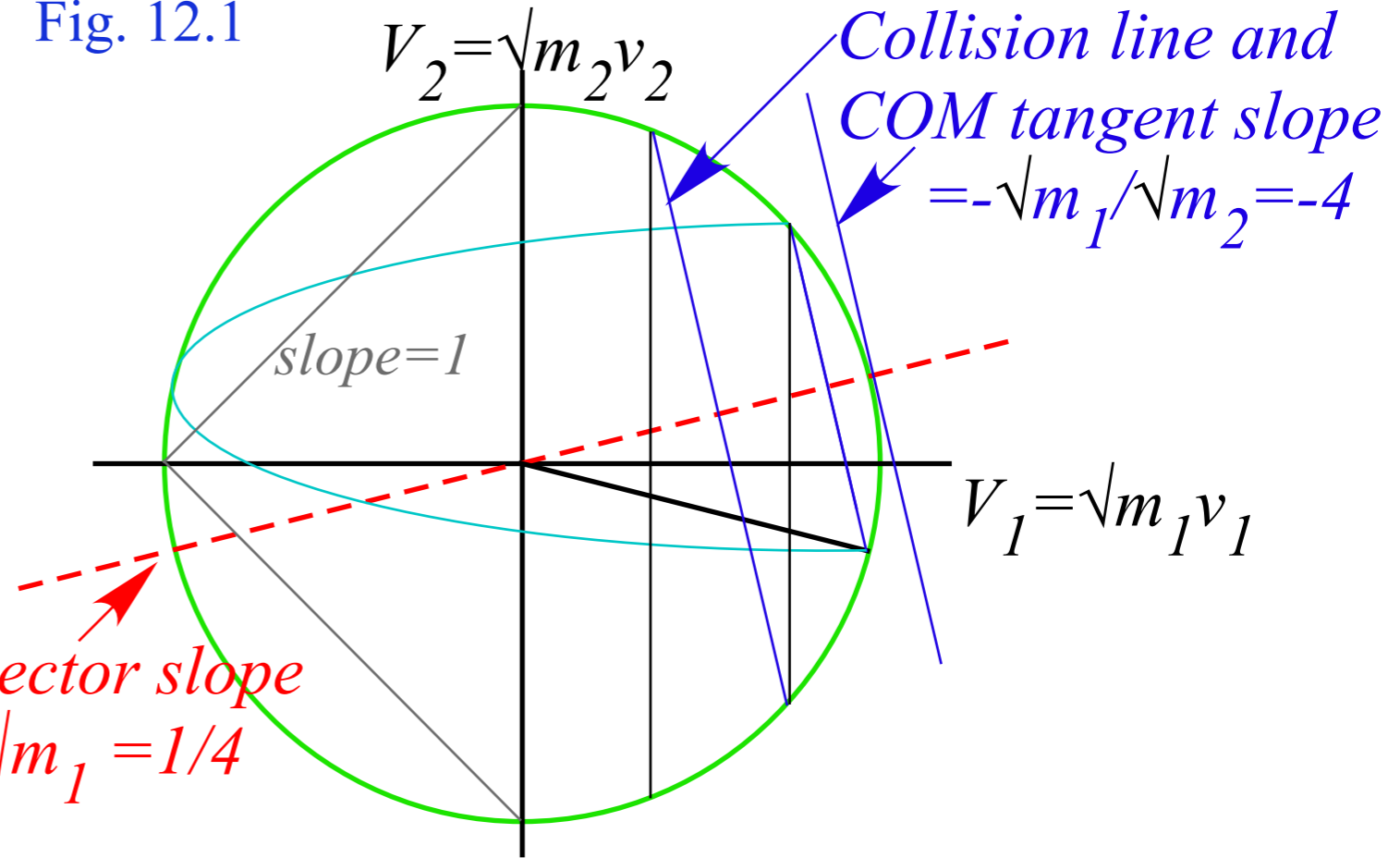
$$\mathbf{M}^{-1} = \begin{pmatrix} 1/m_1 & 0 \\ 0 & 1/m_2 \end{pmatrix} = \mathbf{R}^{-2}$$

(a) Lagrangian $L = L(v_1, v_2)$

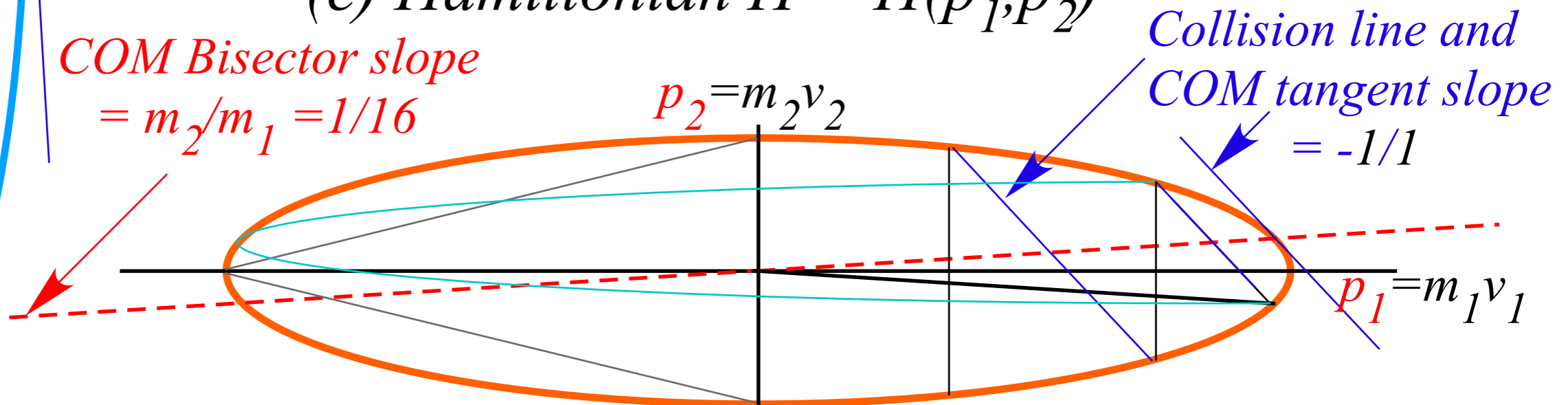


Unit 1
Fig. 12.1

(b) Estrangian $E = E(V_1, V_2)$

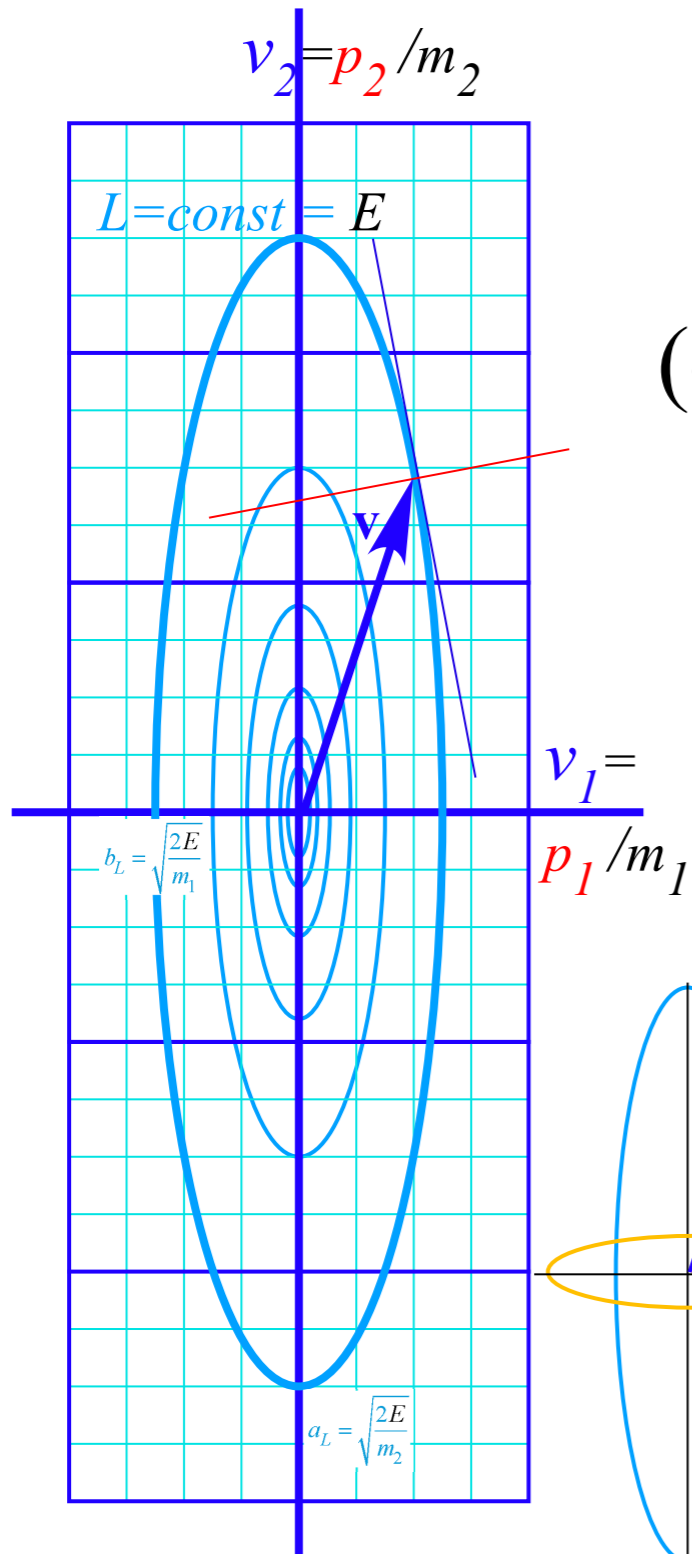


(c) Hamiltonian $H = H(p_1, p_2)$

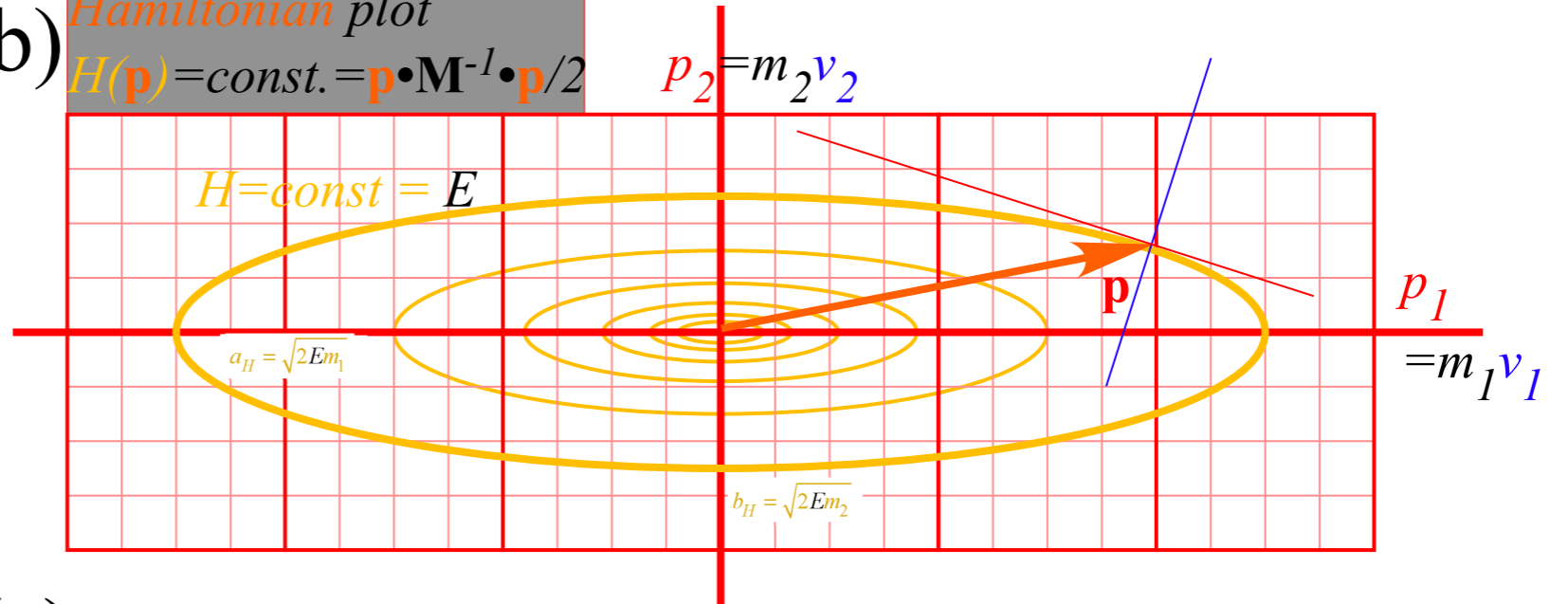


Unit 1
Fig. 12.2

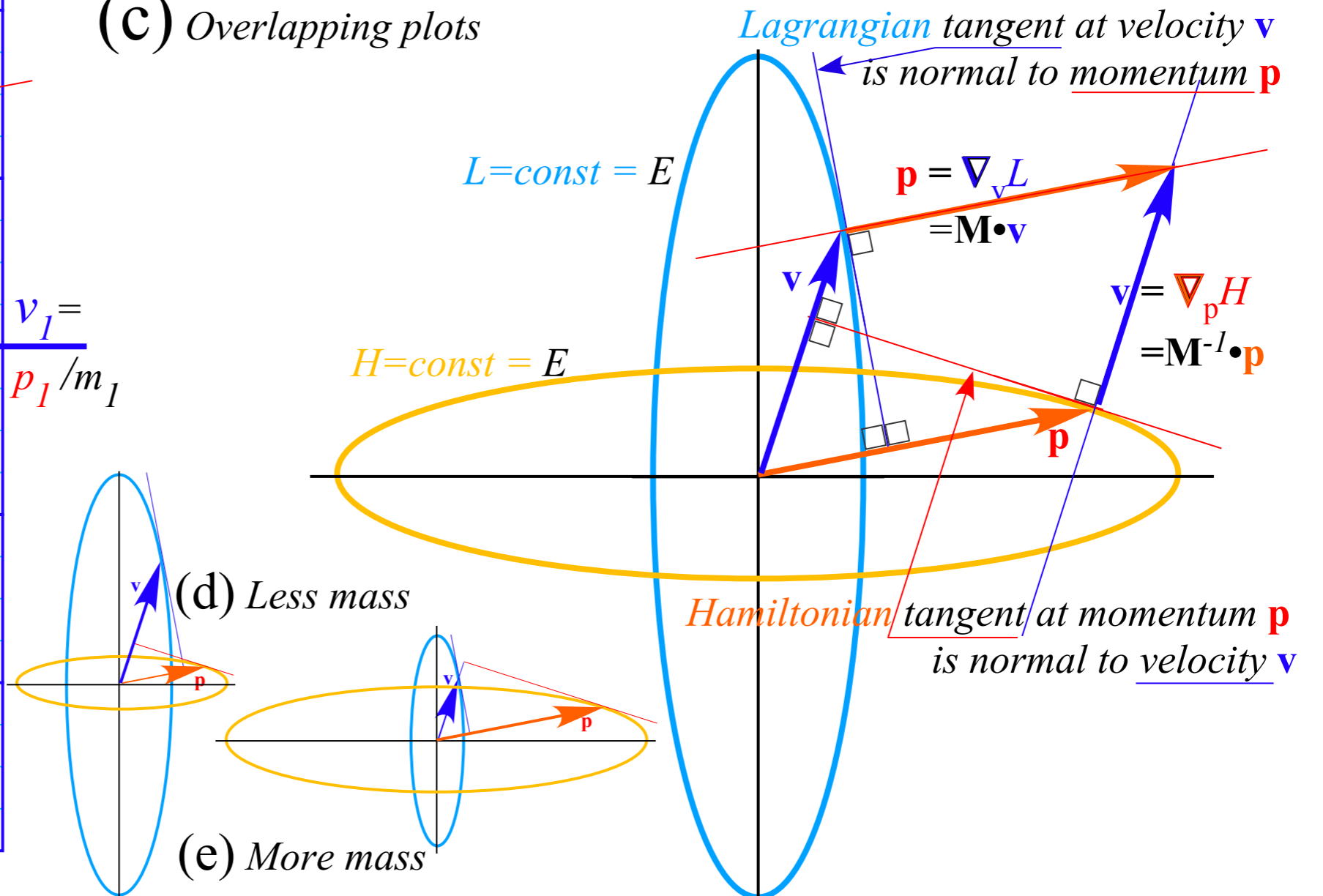
(a) *Lagrangian plot*
 $L(\mathbf{v}) = \text{const.} = \mathbf{v} \cdot \mathbf{M} \cdot \mathbf{v} / 2$



(b) *Hamiltonian plot*
 $H(\mathbf{p}) = \text{const.} = \mathbf{p} \cdot \mathbf{M}^{-1} \cdot \mathbf{p} / 2$



(c) *Overlapping plots*



(d) *Less mass*

(e) *More mass*