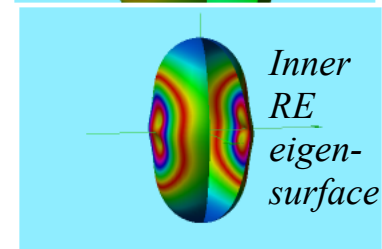
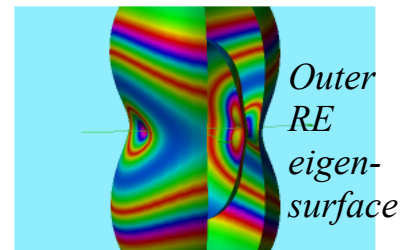
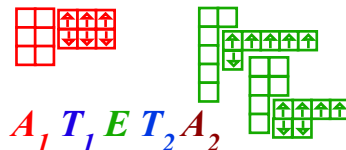
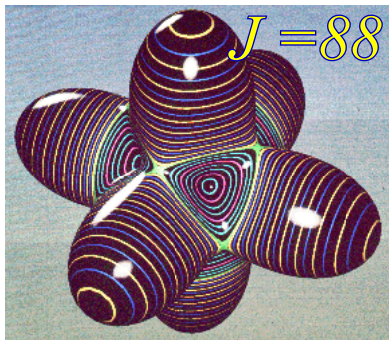


Dynamics of localized angular momentum and multi-surface rotational energy anisotropy: internal-rotor molecules and spin symmetry conversion effects

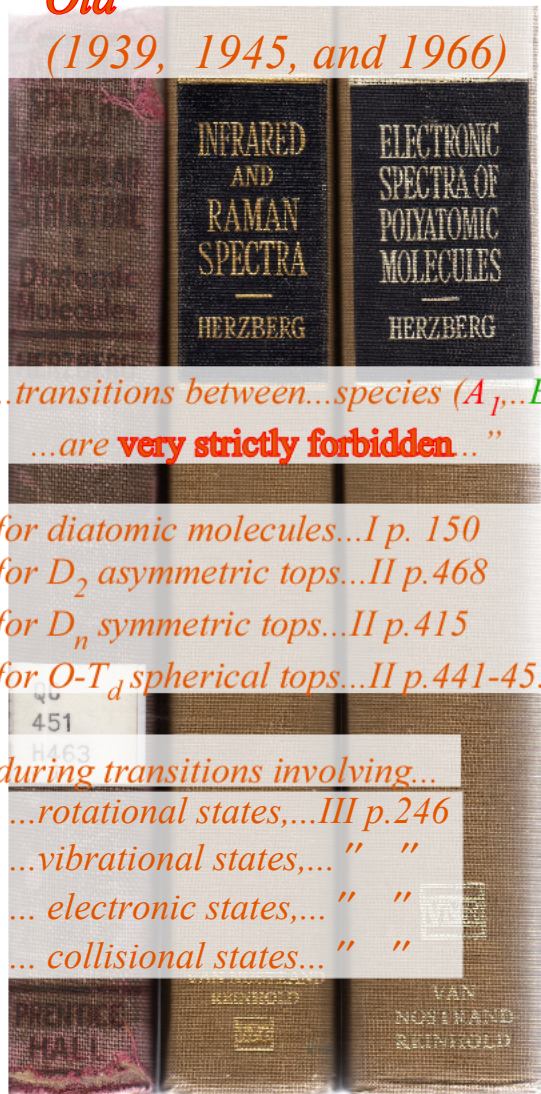
*William G. Harter and Justin C. Mitchell
Department of Physics, University of Arkansas
Fayetteville, AR 72701
wharter@uark.edu*



CONSERVATION OF ROVIBRONIC SPECIES - Two Views:

Old

(1939, 1945, and 1966)



versus

New (1978-2005)

www.sciencemag.org SCIENCE VOL 310 23 DECEMBER 2005

CHEMISTRY

Nuclear Spin Conversion in Molecules

Jon T. Hougen and Takeshi Oka

Molecules with identical nuclei having nonzero spin can exist in different states called nuclear spin modifications by most researchers and nuclear spin isomers by some. Once prepared in a

as initially shown by Bonhoeffer and Harteck in 1929 (3). Once prepared, a *para*-H₂ sample can be preserved for months.

[review of C₂H₄ study: Sun, Takagi, Matsushima, Science 310, 1938(2005)]

“...transitions between...species ($A_1, \dots, E, \dots, T_2, \dots$)
...are **very strictly forbidden**...”

Strictly versus **NOT!**
Conservation and preservation?

...for diatomic molecules...I p. 150
...for D₂ asymmetric tops...II p.468
...for D_n symmetric tops...II p.415
...for O-T_d spherical tops...II p.441-453

No Way! versus **WAY!**
Conversion, perversion or transition?

...during transitions involving...
...rotational states,...III p.246
...vibrational states,... " "
... electronic states,... " "
... collisional states... " "



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Strictly versus **NOT!**
Conservation and preservation?

To **conserve** vs. To **convert**
To **preserve** vs. To **pervert**

No Way! versus **WAY!**
Conversion, perversion or transition?

perversion
Widespread and extreme mixing of species reported in CF₄, SiF₄ and SF₆:

Ch. Borde, Phys. Rev. A20,254(1978)(expt.)
Harter, Phys. Rev. A24,192 (1981)(theory)

HOW CONSERVED IS ROVIBRONIC-SPIN SYMMETRY?

What preserves it? versus *What mixes it up?*

No Way!

WAY!

and...

What is it?

SPIN SYMMETRY correlation has a new name...



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What preserves it? versus *What mixes it up?*

No Way!

WAY!

and...

What is it?

SPIN SYMMETRY correlation has a new name...

*it's now called **ENTANGLEMENT!***



Herzberg's terms:

"..Overall ...symmetry..."

Better terms:

..Under-all ... or internal symmetry...spin frame..... "Bare" rotor

(From an overall "Coupled" state we SUBTRACT vibronic "Activity" to get underlying "Bare" rotor.)

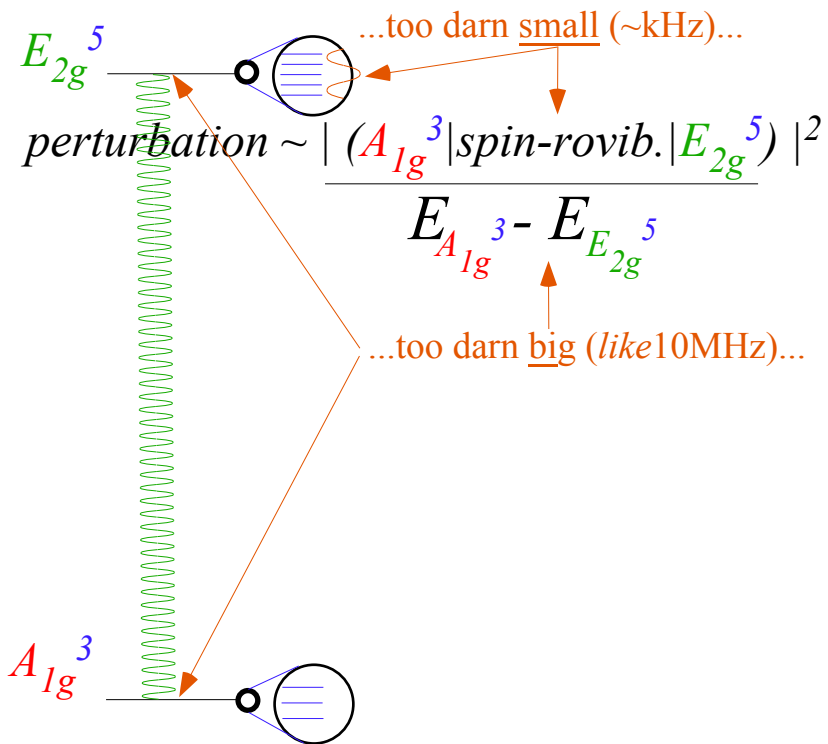
HOW CONSERVED IS ROVIBRONIC-SPIN SYMMETRY?

A_{2u}^1

What preserves it? versus *What messes it up?*

No Way!

...because nuclear moments...
...are so very slight..."



HOW CONSERVED IS ROVIBRONIC-SPIN SYMMETRY?

What preserves it? versus **What mixes it up?**

A_{2u}^1

No Way!

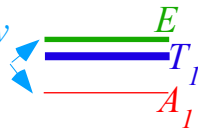
“...because nuclear moments...
...are so very slight...”

...too darn small (~kHz)...

$$\text{perturbation} \sim \frac{|(A_{1g}^3 | \text{spin-rovib.} | E_{2g}^5)|^2}{E_{A_{1g}^3} - E_{E_{2g}^5}}$$

...too darn big (like 10MHz)...

...exponentially
tiny!
(like 10^{-50} Hz)



RE Superhyperfine transitions

Hyperfine effects may rule! $A_1 T_1 E T_2 A_2$ get seriously mixed up.

Harter, Patterson, and daPaixao, Rev. Mod. Phys. 50, 37(1978)

Harter and Patterson, Phys. Rev. A19, 2277(1979) (CF₄)

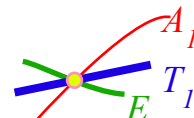
Harter, Phys. Rev. A24, 192-262(1981) (SF₆)

WAY!

...because levels of different species
are forced together by angular wave
localization or “level-clustering” or
(rarely) by “accidental” degeneracy.

“Accidental” degeneracy

Lea, Leask & Wolf JPCSol.23,1381(1962)



Level-clustering

Dorney and Watson JMS 42,135(1972)

Harter and Patterson PRL38,224(1977)

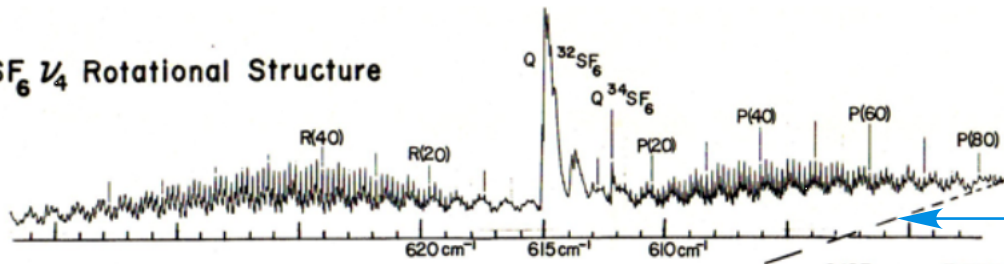
JCP 66,4872(1977)

RE Surface precession vs. tunneling

Harter and Patterson JMP 20,1453(1979)

JCP 80,4241(1984)

(a) SF₆ 1/4 Rotational Structure

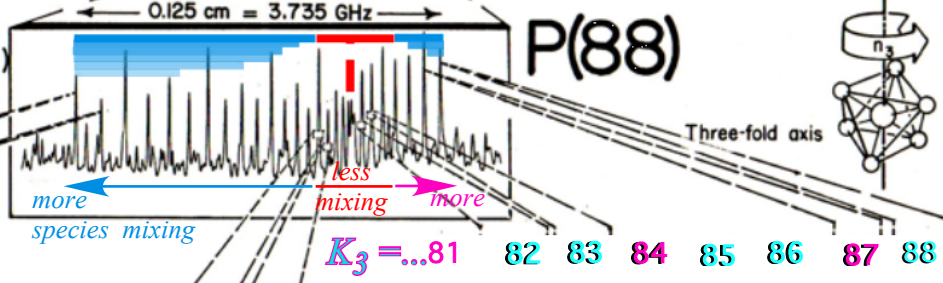


Primary AET species mixing
 increases with distance from
 "separatrix"

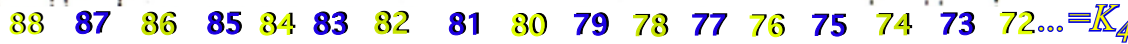
(b) P(88) Fine Structure (Rotational anisotropy effects)



Four fold axis



(c) Superfine Structure (Rotational axis tunneling)

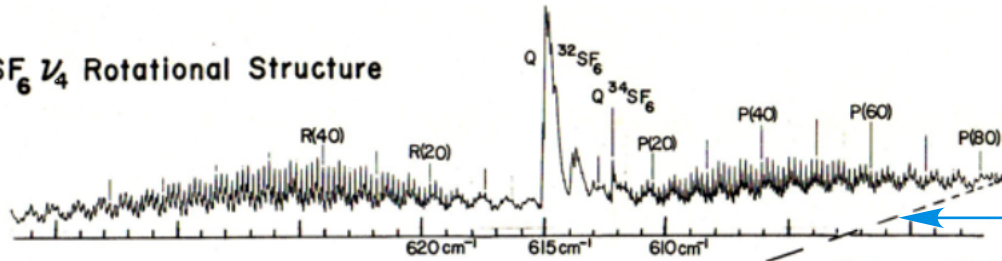


Internal 3-fold axial quanta
 label C₃-CLUSTERS

Internal "body frame" 4-fold axial quanta
 label C₄-CLUSTERS of lines and/or levels



(a) SF₆ 1/4 Rotational Structure



FT IR and Laser Diode Spectra
K.C. Kim, W.B. Person, D. Seitz, and B.J. Krohn
J. Mol. Spectrosc. 76, 322 (1979).

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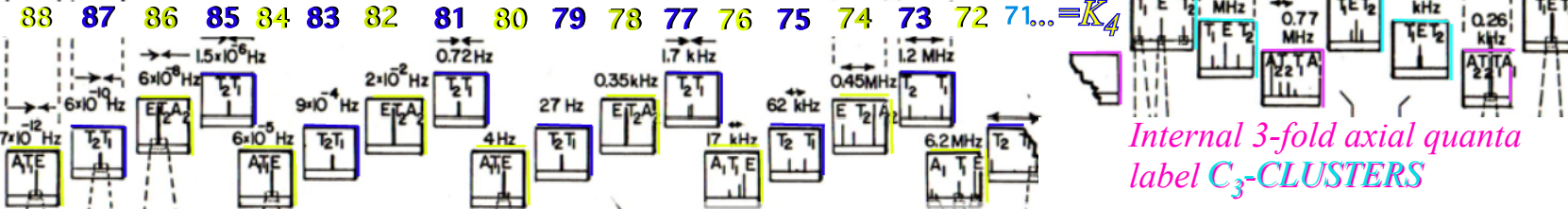
(b) P(88) Fine Structure (Rotational anisotropy effects)



Four fold axis



(c) Superfine Structure (Rotational axis tunneling)



Internal 3-fold axial quanta
label C₃-CLUSTERS

more species mixed

CASE 1 Unmixed
primary A₁ T₁ E T₂ A₂ species

mixed

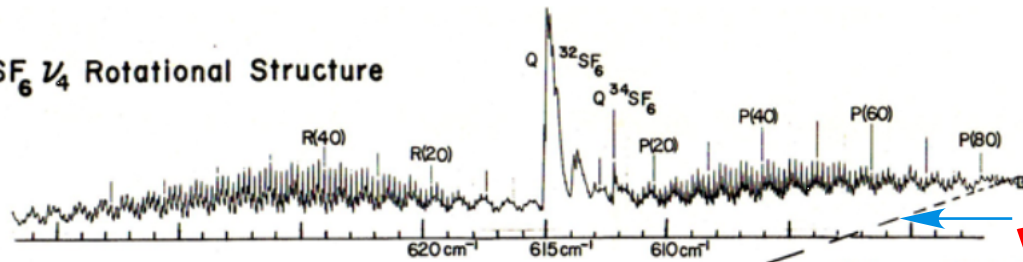
Internal "body frame" 4-fold axial quanta
label C₄-CLUSTERS of lines and/or levels

CASE 2₄ Extreme mixing

CASE 2₃-
Major mixing

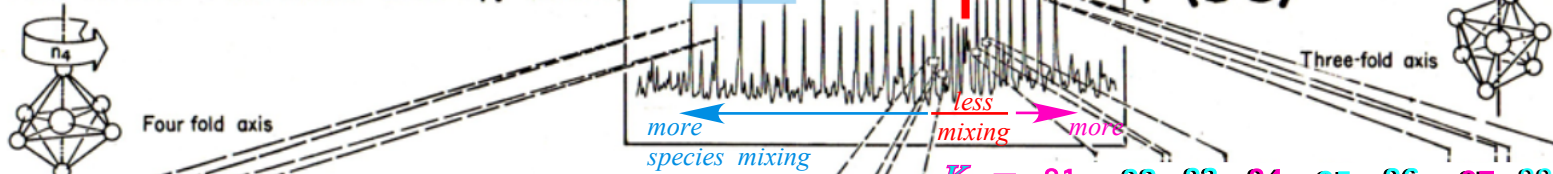


(a) SF₆ ¼ Rotational Structure

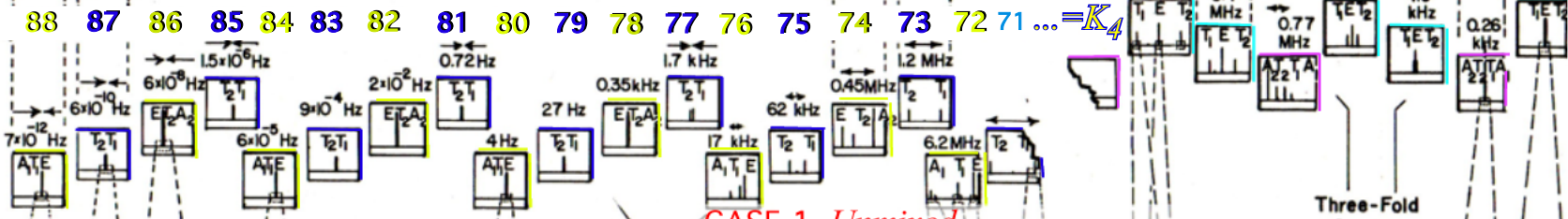


Primary AET species mixing
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"separatrix"

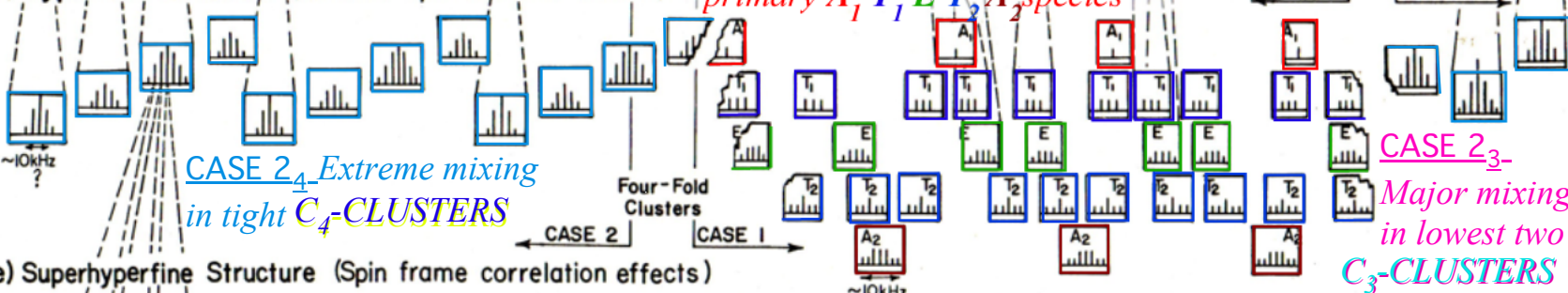
(b) P(88) Fine Structure (Rotational anisotropy effects)



(c) Superfine Structure (Rotational axis tunneling)



(d) Hyperfine Structure (Nuclear spin-rotation effects)



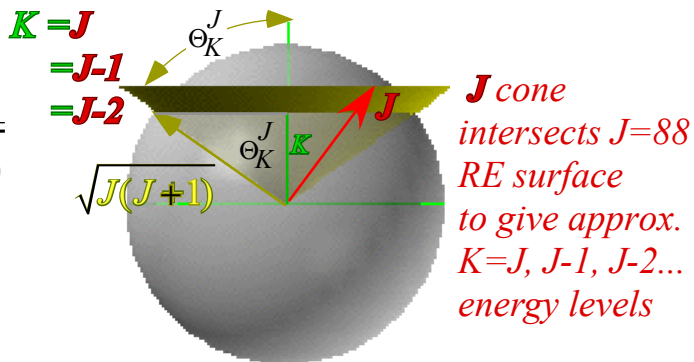
(e) Superhyperfine Structure (Spin frame correlation effects)



(Next page: approximate theory)

Finding Hamiltonian Eigensolutions by Geometry
 using
Uncertainty Cone Angles

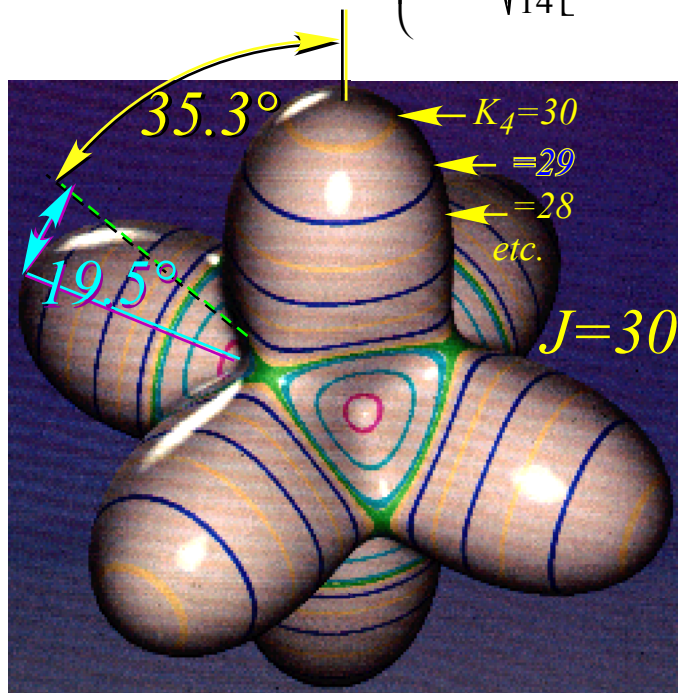
$$\cos \Theta_K^J = \frac{K}{\sqrt{J(J+1)}}$$



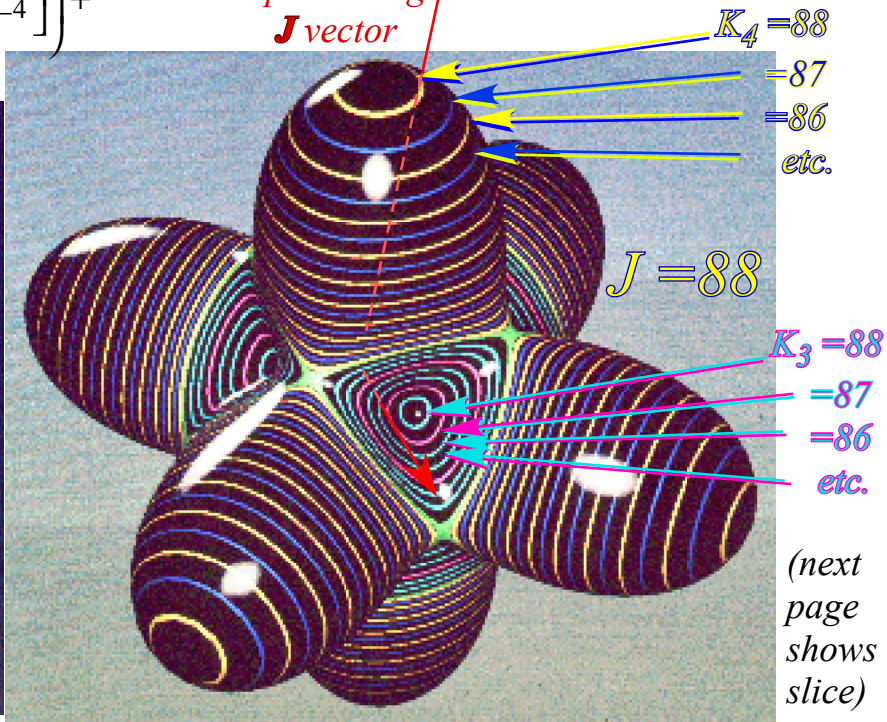
O_h or T_d Spherical Top: (Hecht Ro-vib Hamiltonian 1960)

$$H = B(\mathbf{J}_x^2 + \mathbf{J}_y^2 + \mathbf{J}_z^2) + t_{440} \left(\mathbf{J}_x^4 + \mathbf{J}_y^4 + \mathbf{J}_z^4 - \frac{3}{5} J^4 \right) + \dots$$

$$= BJ^2 + t_{440} \left(\mathbf{T}_0^4 + \sqrt{\frac{5}{14}} [\mathbf{T}_4^4 + \mathbf{T}_{-4}^4] \right) + \dots$$

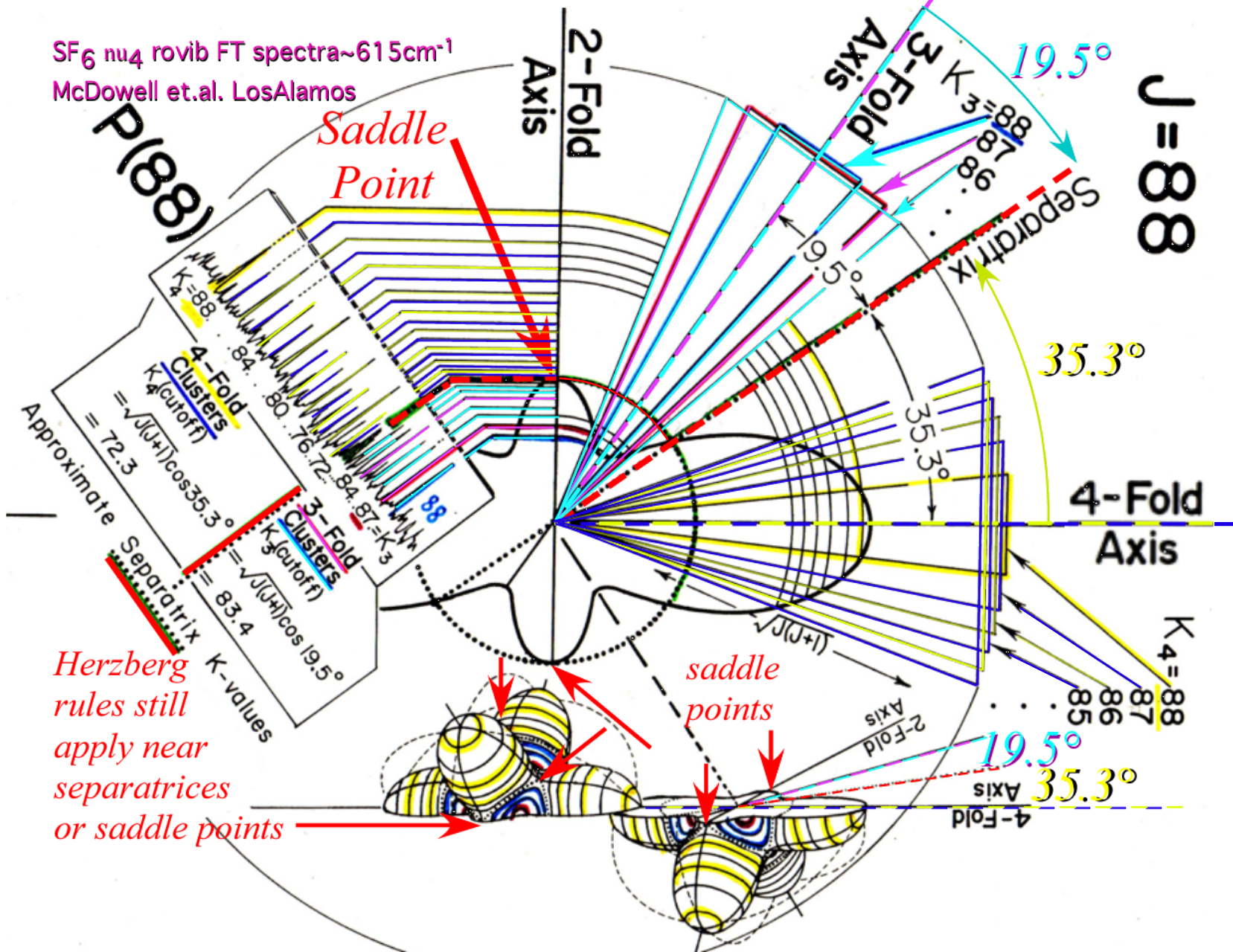


RE Surface topo-lines track precessing J vector semi-classical

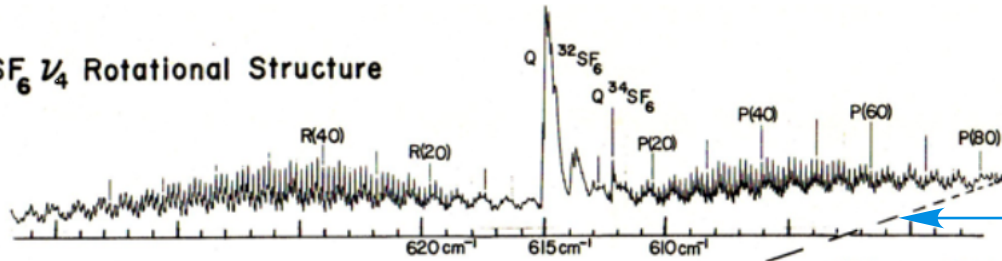


SF₆ nu₄ rovib FT spectra ~615cm⁻¹

McDowell et.al. LosAlamos



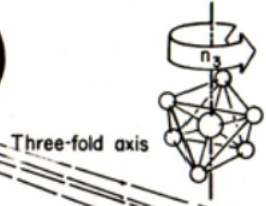
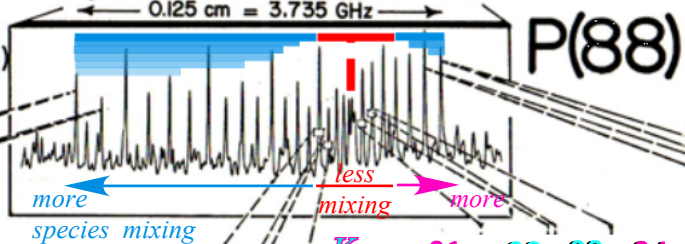
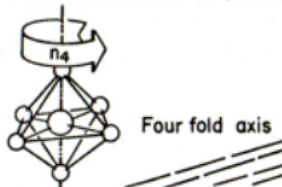
(a) SF₆ 1/4 Rotational Structure



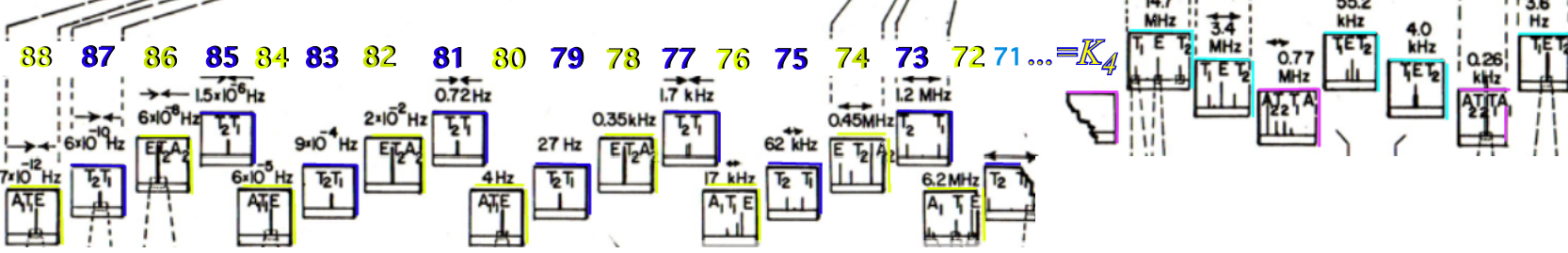
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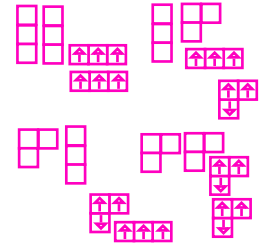
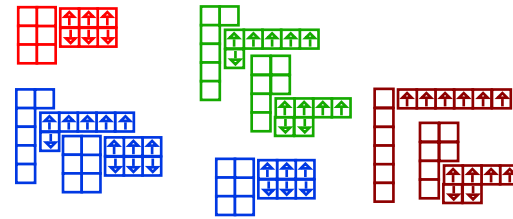
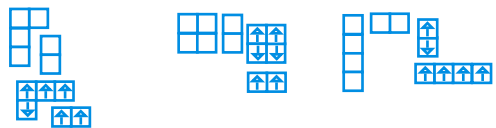
(c) Superfine Structure (Rotational axis tunneling)



CASE 2₄
Broken 4 + 2 tableau state description

CASE 1 Unmixed
primary A₁ T₁ E T₂ A₂ species
(Whole 6-box tableaus)

CASE 2₃
Broken 3 + 3 Tableaus



Spin-rovib ENTANGLEMENT symmetry
might be controllable!

Multiple-RE surfaces: Using semi-classical geometry...

Can we describe internal-rotor molecules and their spin symmetry?

Can we describe hyperfine spin dynamics?

The Simplest Cases:

Rigid top with one body fixed “Gyro” (one spin-1/2, one CH₃, ...)



Multiple-RE surfaces: Using semi-classical geometry...

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Rigid top with one body fixed "Gyro" (one spin-1/2, one CH₃, ...)

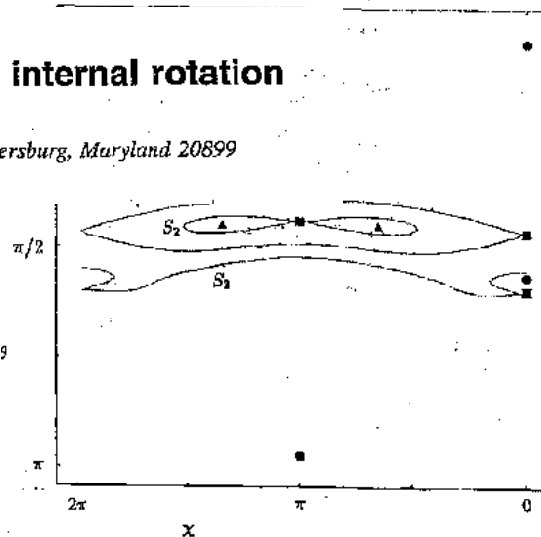
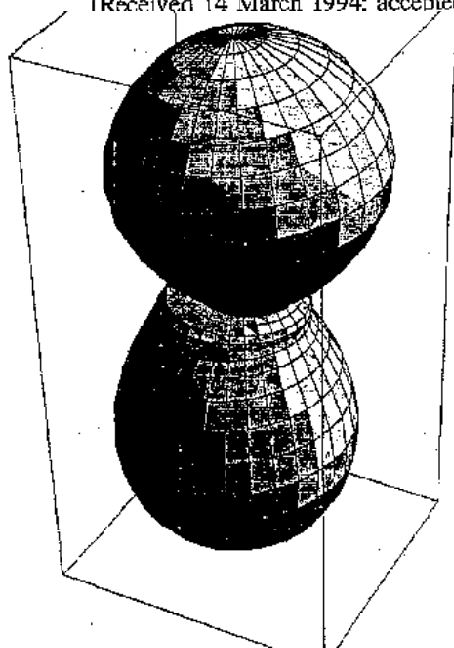
J. Chem. Phys. 101, 2710 (1994)

Rotational energy surfaces of molecules exhibiting internal rotation

Juan Ortigoso^{a)} and Jon T. Hougen

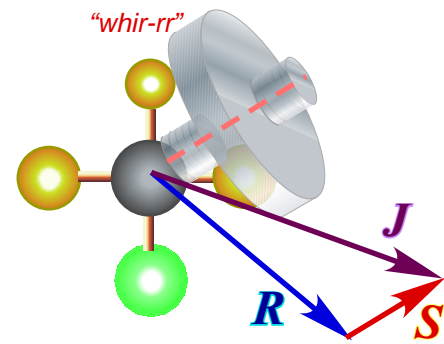
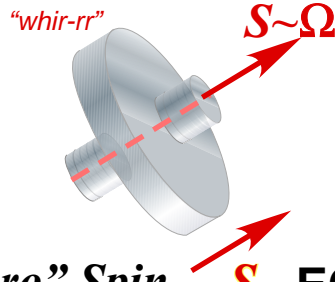
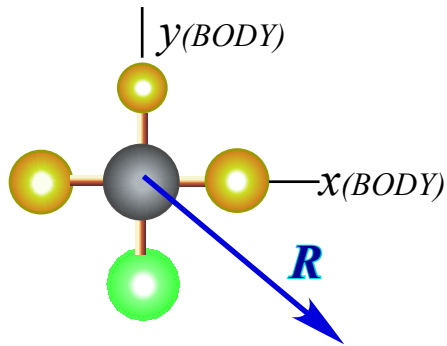
Molecular Physics Division, National Institute of Standards and Technology, Gaithersburg, Maryland 20899

(Received 14 March 1994; accepted 28 April 1994)



*One of the first Applications of
Multiple RES introduced in Comp.Phys.Rpt. 8,319(1988)*

*Problem: Mathematica graphic engines were not terrific!
(..and Los Alamos graphics was too \$\$expensive\$\$)*



Rotor R PLUS "Gyro" Spin S EQUALS Compound Rotor $J=R+S$

Compound Rotor Hamiltonian: Rigid rotor with body-fixed "gyro"...

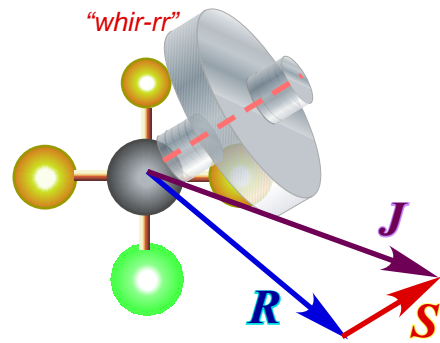
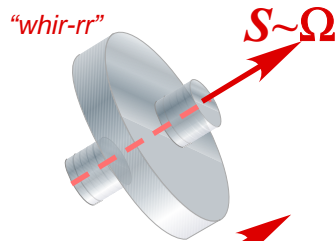
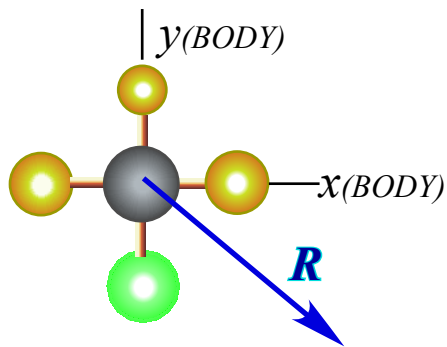
In general, this term is the difficult part...

$$H = AR_x^2 + BR_y^2 + CR_z^2 + \dots + (\text{coupling or constraint}) + \dots + B_S S \cdot S$$

Zero-Interaction Potential 'Proximation (ZIPP)

...but suppose it's zero!
Constraints do no work.





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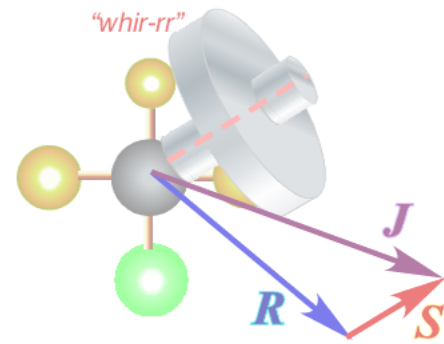
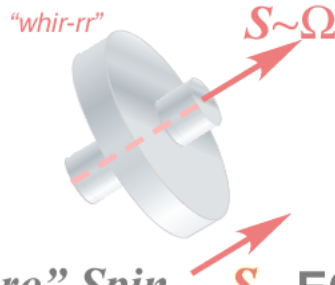
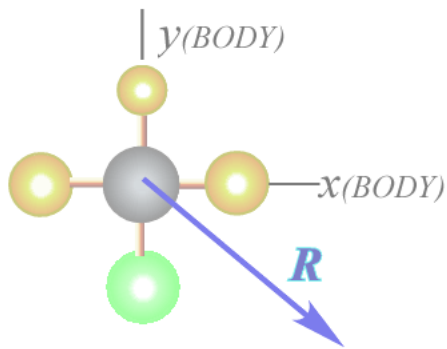
Zero-Interaction Potential 'Proximation (ZIPP)

*...but suppose it's zero!
Constraints do no work.*

Let: $R = J - S$ and consider non-constant terms (ignore gyro S terms that are constant)

$$H = A(J_x - S_x)^2 + B(J_y - S_y)^2 + C(J_z - S_z)^2 + \dots + 0 \text{ (for constraint)} + \dots + (\text{constant BS terms})$$





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Compound Rotor Hamiltonian: Rigid rotor with body-fixed "gyro"...

In general, this term is the difficult part...

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Zero-Interaction Potential 'Proximation (ZIPP)

Constraints do no work.

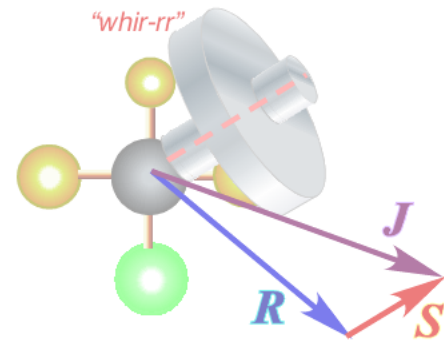
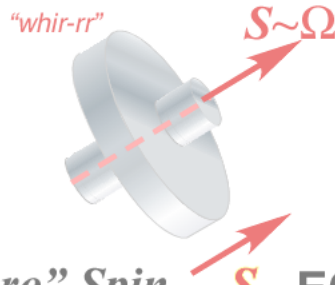
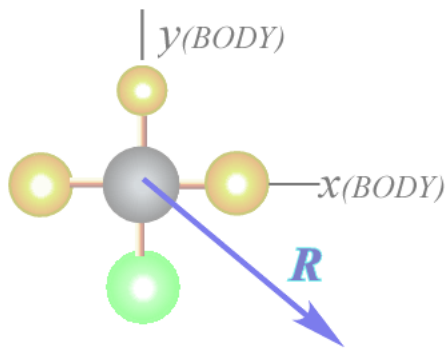
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$$H = A(J_x - S_x)^2 + B(J_y - S_y)^2 + C(J_z - S_z)^2 + \dots + 0 \text{ (for constraint)} + \dots + (\text{constant BS terms})$$

$$H = AJ_x^2 + BJ_y^2 + CJ_z^2 + \dots - 2AJ_xS_x - 2BJ_yS_y - 2CJ_zS_z + \dots + (\text{more constant terms})$$

"Coriolis effect" subtracts linear or 1st-order \mathbf{J}_m or \mathbf{T}_m^1 terms for gyro-rotor H





Rotor **R** PLUS “Gyro” Spin **S** EQUALS Compound Rotor **J=R+S**

Compound Rotor Hamiltonian: Rigid rotor with body-fixed “gyro”...

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Zero-Interaction Potential ‘Proximation (ZIPP)

Constraints do no work.

Let: $\mathbf{R} = \mathbf{J} - \mathbf{S}$ and consider non-constant terms (ignore gyro \mathbf{S} terms that are constant)

$$H = A(\mathbf{J}_x - \mathbf{S}_x)^2 + B(\mathbf{J}_y - \mathbf{S}_y)^2 + C(\mathbf{J}_z - \mathbf{S}_z)^2 + \dots + 0 \text{ (for constraint)} + \dots + (\text{constant BS terms})$$

$$H = A\mathbf{J}_x^2 + B\mathbf{J}_y^2 + C\mathbf{J}_z^2 + \dots - 2A\mathbf{J}_x\mathbf{S}_x - 2B\mathbf{J}_y\mathbf{S}_y - 2C\mathbf{J}_z\mathbf{S}_z + \dots + (\text{more constant terms})$$

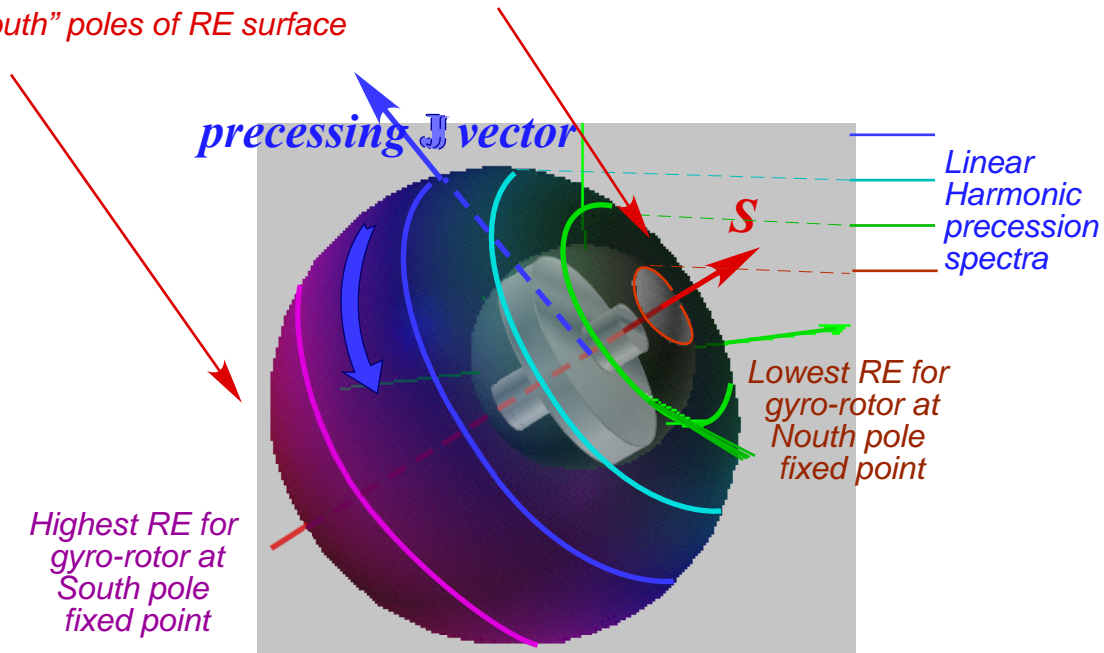
“Coriolis effect“ subtracts linear or 1st-order \mathbf{J}_m or \mathbf{T}_m^1 terms for gyro-rotor H

BR^2 to $B(\mathbf{J}-\mathbf{S})^2$ is analogous to $p^2/2M$ to $(\mathbf{p}-e\mathbf{A})^2/2M$ gauge-transformation
 ... $\mathbf{J} \cdot \mathbf{S}$ is analogous to $e\mathbf{p} \cdot \mathbf{A}$

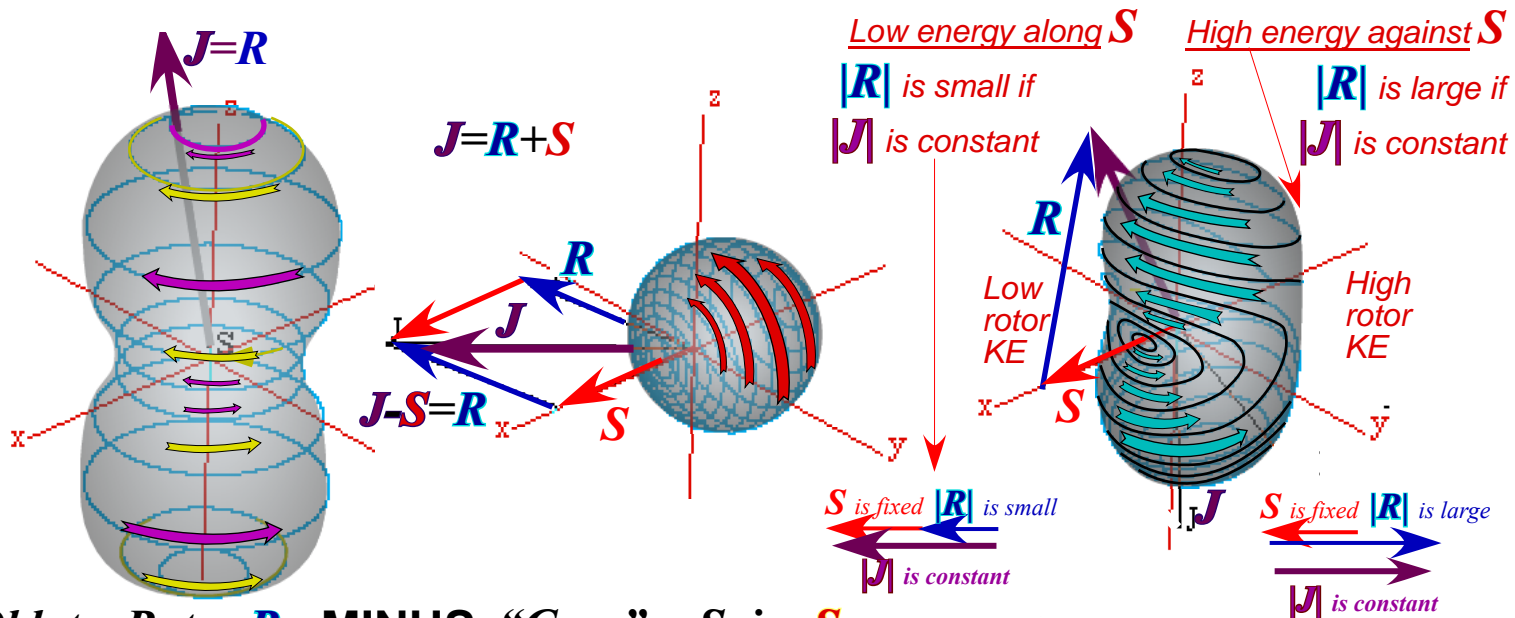


RE Surface for 1st-order \mathbf{J}_m or \mathbf{T}_m^1 term is a cardioid displaced in J -direction
Energy sphere intersections are concentric circular precession paths
All paths precess with the same sense around gyro S -vector

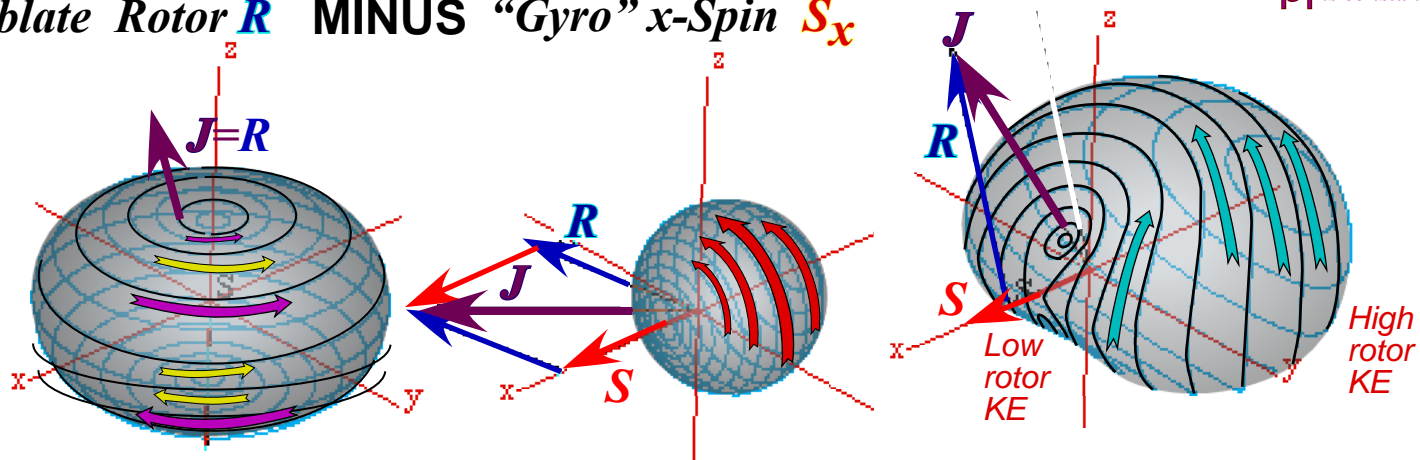
Fixed Points for \mathbf{J} lie on “North” and “South” poles of RE surface



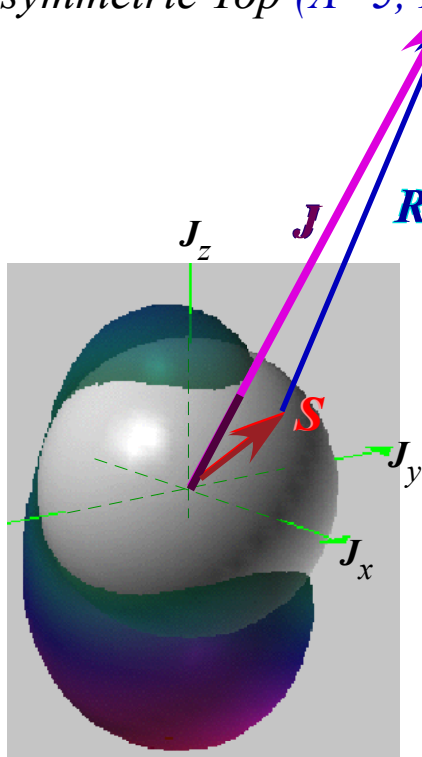
Prolate Rotor R MINUS "Gyro" x -Spin S_x



Oblate Rotor R MINUS "Gyro" x -Spin S_x

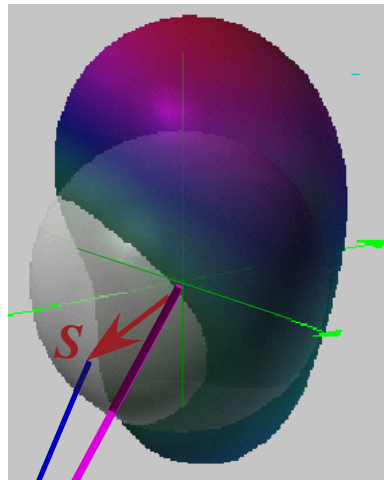


Spin gyro $S=(1,1,1)$ attached (ZIPPed) to
Asymmetric Top ($A=5, B=10, C=15$)

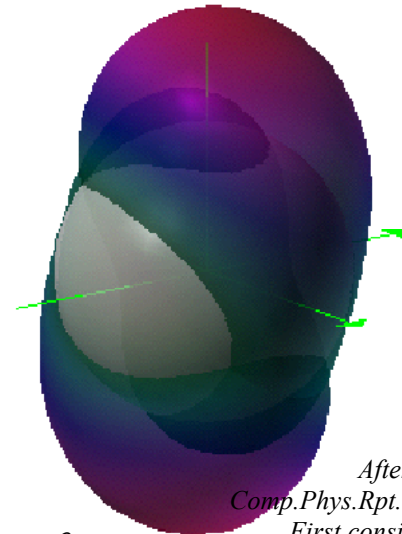


“Sherman” (The shark)

Time reversed
gyro $-S=(-1,-1,-1)$

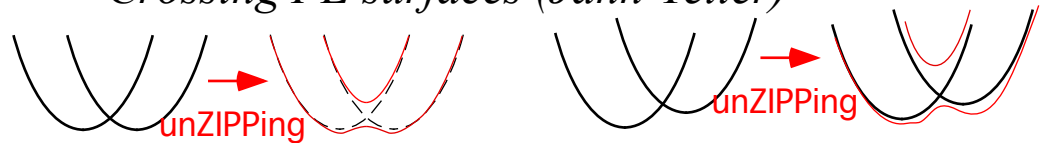


The two together



After
Comp.Phys.Rpt. 8,319(1988)
First considered in
1992 JCP article by
Hougan, Kleiner, and Ortigoso

Crossing RE surfaces
analogous to
Crossing PE surfaces (Jahn-Teller)



Two or more RE's beg to be **unZIPPED**. $\langle \mathbf{H} \rangle = \begin{pmatrix} \text{Spin-up RE}(\beta, \gamma) & \text{Coupling}(\beta, \gamma) \\ \text{Coupling}(\beta, \gamma)^* & \text{Spin-down RE}(\beta, \gamma) \end{pmatrix}$
 Base RE surfaces are eigenvalues of matrix.

Classical RE

$$H = AJ_x^2 + BJ_y^2 + CJ_z^2 + \dots - 2AJ_x S_x - 2BJ_y S_y - 2CJ_z S_z + \dots + (\text{more constant terms})$$

Semi-Classical Spin-1/2 RE $\sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \mathbf{1} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$ makes matrix

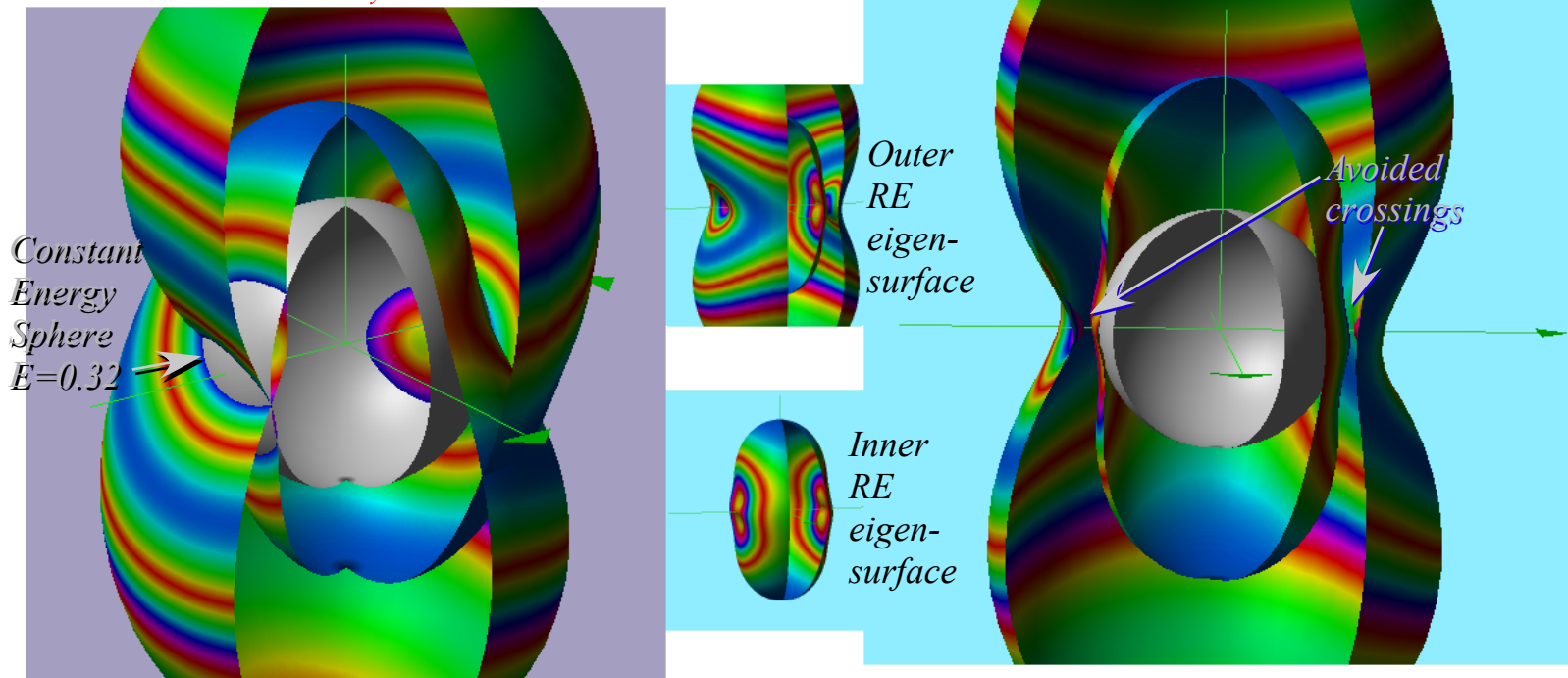
$$\mathbf{H} = (AJ_x^2 + BJ_y^2 + CJ_z^2)\mathbf{1} \dots - AJ_x s_x \sigma_x - BJ_y s_y \sigma_y - CJ_z s_z \sigma_z + \dots + \mathbf{1} (\text{more constant terms})$$

Classical **ZIP** $A=0.2, B=0.8, C=1.4$

$s_x=0.0, s_y=0.1, s_z=0.2$

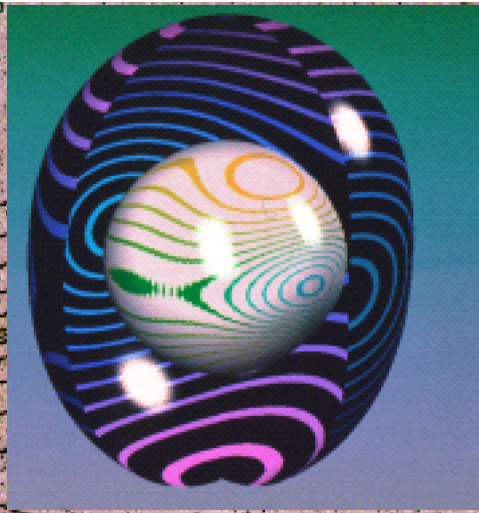
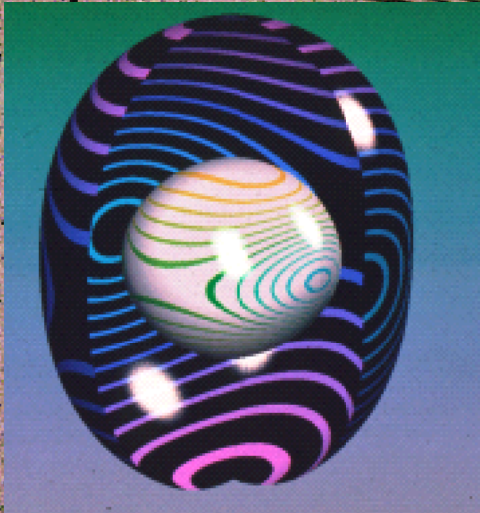
Semi-Classical spin-1/2 unZIP $A=0.2, B=0.8, C=1.4$

$s_x=0.0, s_y=0.1, s_z=0.2$



J= 10.5 Eigenvalues of Spin- Rotor

500.0 1000.0 1500.0



$(A_1 B_1 A_2 B_2)$ clusters

$(R=21/2) \times (l=1/2)$ Diagonalization $A=0.2, B=0.4, C=0.6$
 varying $D_{xx}=s_x, D_{yy}=s_y=2D_{xx}, D_{zz}=s_z=3D_{xx}$

-1000 0.0 1000 20.0 30.0 40.0 50.0 60.0 70.0 80.0 90.0
 D_{xx} With $D_{yy}=2D_{xx}$ and $D_{zz}=3D_{xx}$

R =

Rotational Energy (RE) surfaces: Future?

*Two or more RE's beg for full interaction **unZIPed** for higher spin quanta*

$$\langle \mathbf{H} \rangle = \begin{pmatrix} \text{Spin-up RE}(\beta, \gamma) & \text{Coupling}(\beta, \gamma) \\ \text{Coupling}(\beta, \gamma)^* & \text{Spin-down RE}(\beta, \gamma) \end{pmatrix}$$

Base RE surfaces are eigensolutions of such matrices.

Combination RES depends on eigenvector chosen.

Interesting mechanics and dynamics. (Both QM and CM)

Two (or more) surfaces imply an infinity of surfaces “between” them.

Intermediate surfaces not unique for each energy

(“Tide” rises and falls, saddles open and close. Result: Chaotic trajectory and opportunity for sensitive control schemes.)

Conclusion

Rotational Energy (RE) surfaces help analyze rotor dynamics as do Potential Energy (PE) surfaces for vibration.

PE surfaces based on vibrational coordinates.

RE surfaces based on rovibrational phase space.

Can approximate quantum levels and spectra and also mixing and transitions.

RES have a variety of complementary surfaces:

Angular Velocity surfaces: (AVS) $\boldsymbol{\omega} \cdot \mathbf{I} \cdot \boldsymbol{\omega} = 2E$ (Poincot ellipsoid)

Angular Momentum surfaces: (AMS) $\mathbf{J} \cdot \mathbf{I}^{-1} \cdot \mathbf{J} = 2E$ (Landau ellipsoid)

PE surfaces used since beginning of QM (Born 1926)

RE surfaces first used in 1976.